Hamiltonian Analysis of the Arbitrary-Dimensional Hilbert-Palatini Theory

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Outline

- Motivation
- 2 Introduction
- 3 Hilbert-Palatini theory
- 4 Conclusions & Outlook

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- Riemann (curvature) tensor $R_{\mu\nu\rho}{}^{\sigma} = 2\partial_{[\mu}\Gamma_{\nu]\rho}{}^{\sigma} + 2\Gamma_{[\mu|\rho}{}^{\lambda}\Gamma_{|\nu]\lambda}{}^{\sigma}$ Square brackets denote antisymmetrization: $A_{[\mu\nu]} = \frac{1}{2} (A_{\mu\nu} - A_{\nu\mu})$

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Action of the gravitational field (Einstein-Hilbert action):

$$S_{\mathsf{EH}} = \int d^{D+1} x \cdot \sqrt{-g} R$$

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Note: $\omega_{\mu I}^{J}$ is **not** related to $\Gamma_{\mu\nu}^{\rho}$ or e^{μ}_{I}

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Hilbert-Palatini action:

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Introduction: Comparison of GR with HP theory

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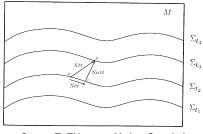
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⇒ Matter-free GR and HP theory are equivalent on-shell

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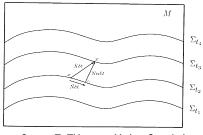
Conclusions & Outlook



Source: T. Thiemann. Modern Canonical

Quantum General Relativity (2007)

 \leftarrow Spacetime split Assumption: Spacetime topology $M\cong \mathbb{R}\times \sigma$ Foliation of M into hypersurfaces $\Sigma_t\cong \sigma$ of constant time t



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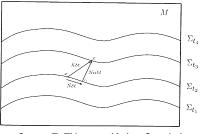
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Parametrization with

N . . . lapse function

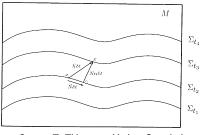
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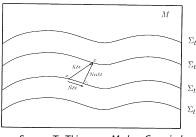
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Pullback all spatial quantities on σ :



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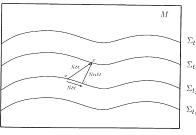
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$$N^{\mu} \rightarrow N^{a}, \; e_{\mu}{}^{I} \rightarrow e_{t}{}^{I}, e_{a}{}^{I},$$

$$\omega_{\mu I}^{J} \rightarrow \omega_{tI}^{J}, \omega_{aI}^{J}$$
, etc.

$$N = -\frac{1}{e^{t_i}e^{t_i}}$$
, $N^a = -\frac{e^t_ie^{al}}{e^t_ie^{t_i}}$

Canonical variables:

Coordinates	$e_t^{\ \prime}$	$ e_a '$	ω_t^{IJ}	ω_a^{IJ}
Momenta	$\Pi^t{}_I = 0$	$\Pi^a{}_I = 0$	$\Pi^t_{IJ} = 0$	$\Pi^{a}_{IJ} = 2ee^{[t}_{I}e^{a]}_{J}$

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Incorporate e_t^I and e_a^I into N, N^a , Π^a_{II}

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Get rid of additional freedom by introducing simplicity constraint:

$$\mathcal{S}^{ab}{}_{IJKL} := \Pi^a{}_{[IJ} \Pi^b{}_{KL]}$$

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HP-Hamiltonian:

$$H = \int d^D \mathbf{x} \cdot \left(\mathcal{N} \mathcal{H} + \mathcal{N}^a \mathcal{H}_a + \Lambda^{IJ} \mathcal{G}_{IJ} + c_{ab}{}^{IJKL} \mathcal{S}^{ab}{}_{IJKL} \right)$$

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$$\begin{split} \left\{\mathcal{B}^{ab}{}_{IJKL}\left[b_{ab}{}^{IJKL}\right], \mathcal{S}^{cd}{}_{MNOP}\left[c_{cd}{}^{MNOP}\right]\right\} =: \mathcal{F}^{abcd}{}_{IJKLMNOP}\left[b_{ab}{}^{IJKL}c_{cd}{}^{MNOP}\right] \\ \left\{\mathcal{B}^{ab}{}_{IJKL}\left[b_{ab}{}^{IJKL}\right], \mathcal{H}\left[\mathcal{N}\right]\right\} =: \Sigma^{ab}{}_{IJKL}\left[b_{ab}{}^{IJKL}\right] \end{split}$$

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Assume
$$\exists \left(\mathcal{F}^{-1}\right)_{abcd}^{IJKLMNOP}$$
 such that

$$\left(\mathcal{F}^{-1}\right)_{abcd}^{\textit{IJKLMNOP}}\mathcal{F}^{abef}_{\textit{IJKLQRST}} = \delta^{(ef)}_{(cd)}\delta^{[\textit{MNOP}]}_{[\textit{QRST}]}$$

$$\begin{split} \left\{ \mathcal{B}^{ab}{}_{IJKL} \left[b_{ab}{}^{IJKL} \right], \mathcal{S}^{cd}{}_{MNOP} \left[c_{cd}{}^{MNOP} \right] \right\} =: \mathcal{F}^{abcd}{}_{IJKLMNOP} \left[b_{ab}{}^{IJKL} c_{cd}{}^{MNOP} \right] \\ \left\{ \mathcal{B}^{ab}{}_{IJKL} \left[b_{ab}{}^{IJKL} \right], \mathcal{H} \left[\mathcal{N} \right] \right\} =: \Sigma^{ab}{}_{IJKL} \left[b_{ab}{}^{IJKL} \right] \end{split}$$

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 $\Rightarrow \mathcal{B}^{ab}{}_{IJKL}$ now Poisson-commutes with Hamiltonian

Still, $\mathcal{S}^{ab}{}_{IJKL}$ and $\mathcal{B}^{ab}{}_{IJKL}$ are second-class contraints



Proof of $\exists \left(\mathcal{F}^{-1}\right)_{abcd}^{\mathit{IJKLMNOP}}$

Proof of $\exists~ \left(\mathcal{F}^{-1}\right)_{abcd}^{\quad IJKLMNOP}$ has not been found

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Problem: For D > 3, S^{ab}_{IJKL} 's are not independent

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However, there is an indication that no further constraints exist

① Spacetime split in internal space: $I \rightarrow 0, i$

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Hilbert-Palatini theory: Reduction to ADM theory

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No further constraints needed to obtain ADM

Outline

- Motivation
- 2 Introduction
- 3 Hilbert-Palatini theory
- Conclusions & Outlook

Conclusions

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Conclusions

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Outlook

Conclusions

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Outlook

Investigations of other GR reformulations



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Outlook

- Investigations of other GR reformulations
- Maybe LQG "predicts" 3 + 1-dimensional spacetime?



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