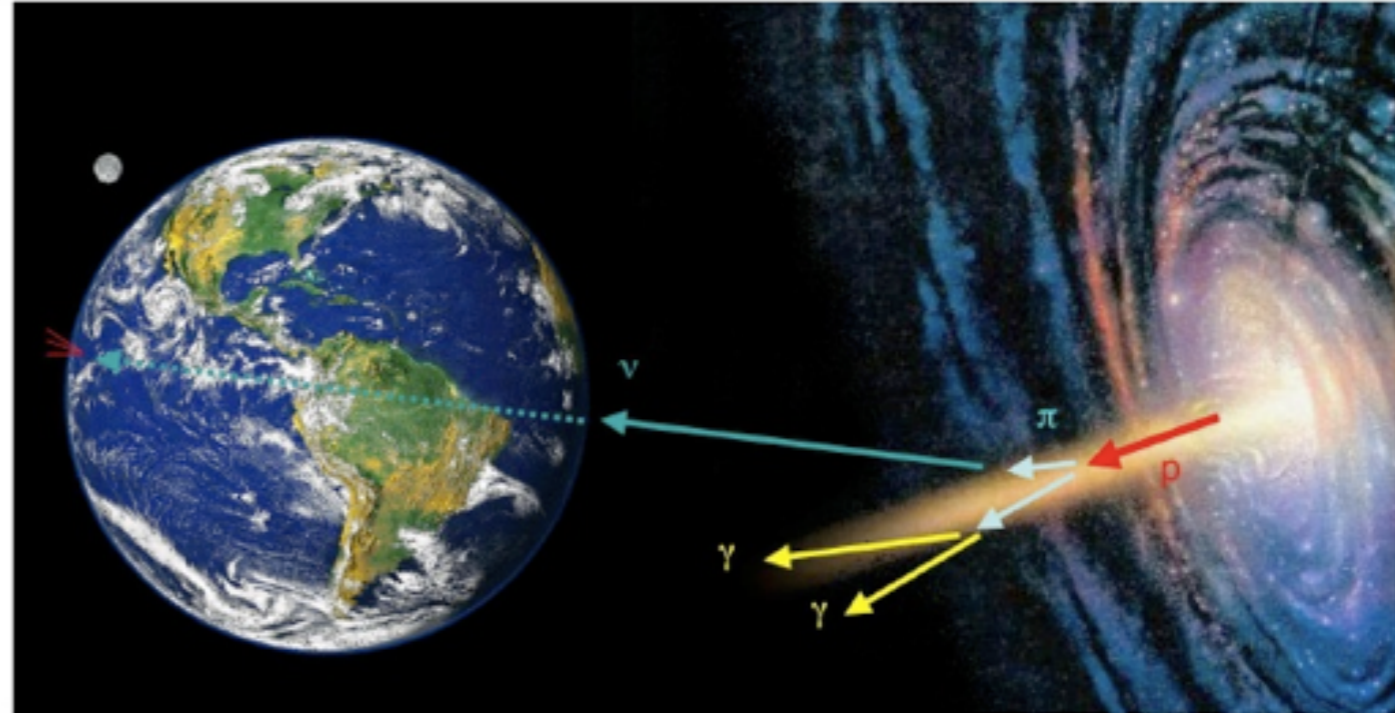
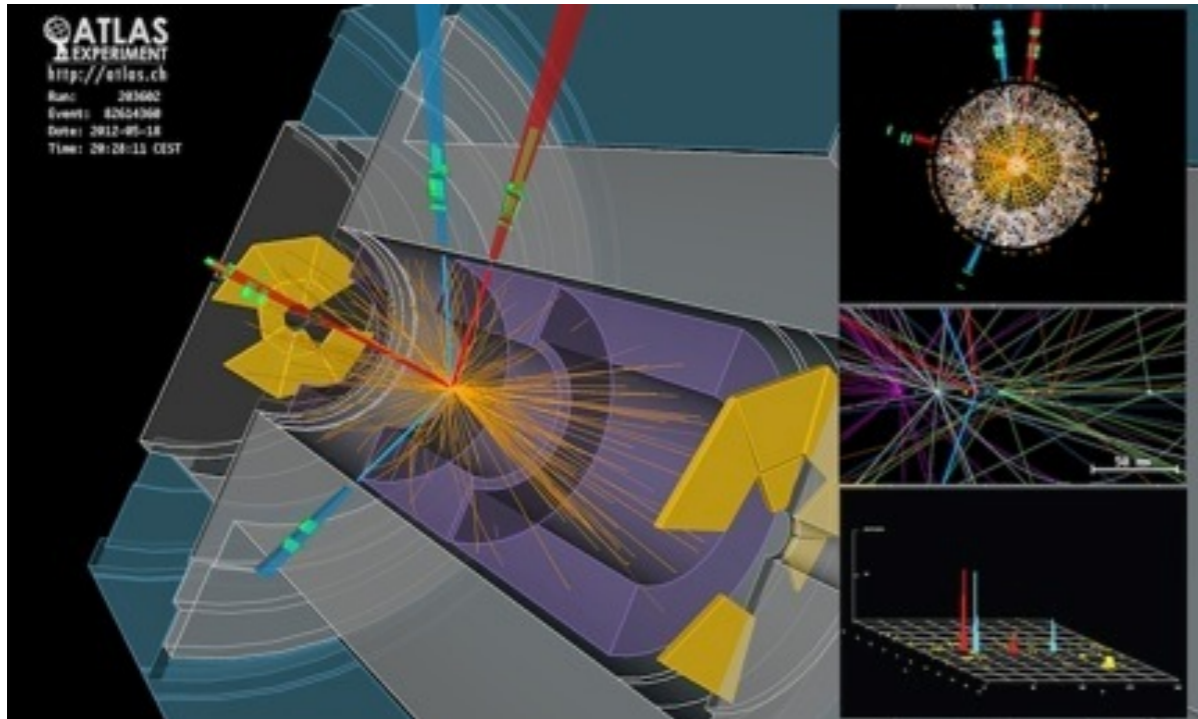


Particle Physics at Colliders and in the High Energy Universe



13. New Physics at the Energy Frontier

27.01.2019



Prüfungstermine

- Wir haben eine (kleine) Auswahl an möglichen Prüfungsterminen:
 - Mo, 10.02.
 - **Di, 11.02. (nachmittags)**
 - Mo, 2.03.
 - **Mo, 6.04. (vormittags)**

Overview

- Shortcomings of the Standard Model and Motivations for New Physics
- Features of Grand Unified Theories
- Supersymmetry
- LHC Searches for
 - SUSY
 - Dark Matter
 - Extra Dimensions
- A Broad Look at Current LHC Limits

Introduction: Beyond the Standard Model at LHC

- A significant fraction of all analyses performed at the LHC search for phenomena beyond the standard model - typically classified in:
 - **Supersymmetry**: Searching for indications for a concrete, popular and well-motivated extension of the Standard Model
 - **Exotics**: More generic new phenomena searches - often also motivated by theoretical ideas, but also very general searches for deviations from Standard Model expectations

Impossible to cover this in a single lecture - here:

A focus on Supersymmetry as an example to illustrate motivations and properties of BSM theories

A few examples of other phenomena and search strategies

Reminder: Limitations of the Standard Model

- The Standard Model with:
 - Fundamental fermions (3 pairs of quarks, 3 pairs of leptons)
 - Fundamental interactions through gauge fields, manifested through
 - W, Z, γ (electroweak - $SU(2) \times U(1)$)
 - gluons (strong - $SU(3)$)

successfully describes all HEP experiments and observations.

BUT:

- it has conceptual problems
- it is incomplete: Fails to describe astrophysical / cosmological observations.

Limitations: Conceptual Problems

- Too many free parameters: ~ 18 masses, couplings, mixing angles
- No unification of electroweak and strong interaction
- No inclusion of quantum gravity
- Family replication: Why 3 families of fundamental fermions?
- Hierarchy problem: “Fine tuning” of precise cancellation of radiative corrections
- Why $1/3$ - charges of quarks - or: What ensures exactly equal charge of protons and electrons?

possible solution

GUT; $E \sim 10^{16}$ GeV

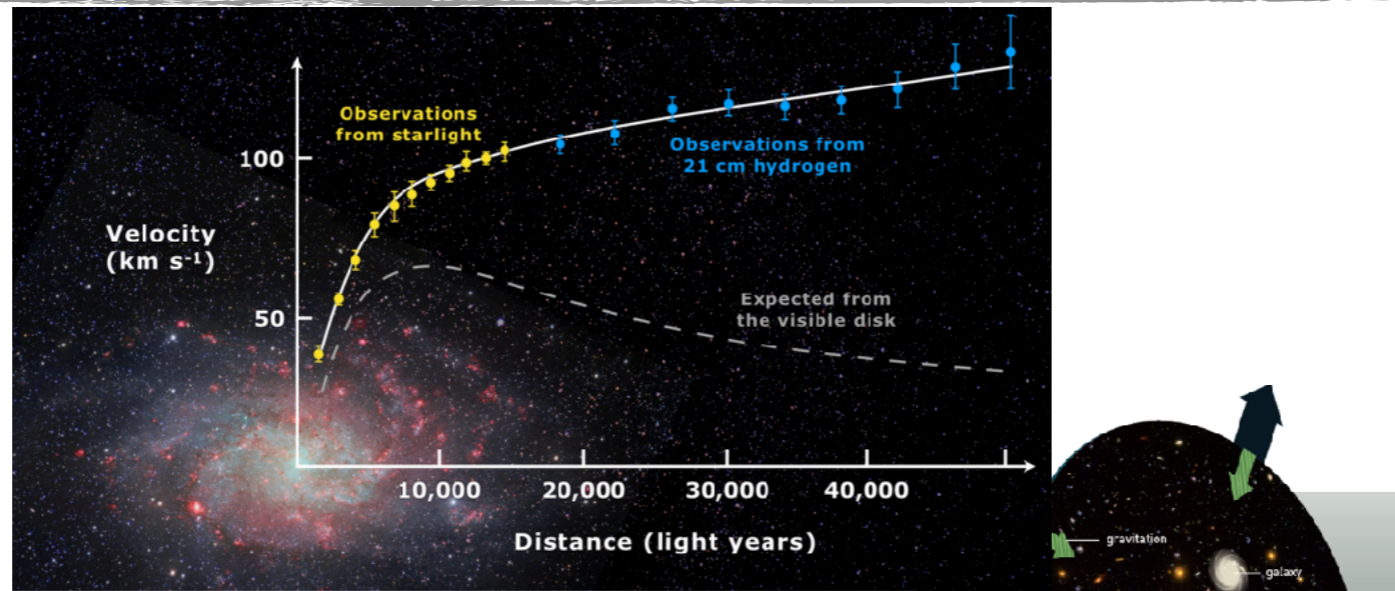
TOE; $E \sim 10^{19}$ GeV

SUSY, Extra dimensions, ...
 $E \sim 10^3$ GeV

GUT; $E \sim 10^{16}$ GeV

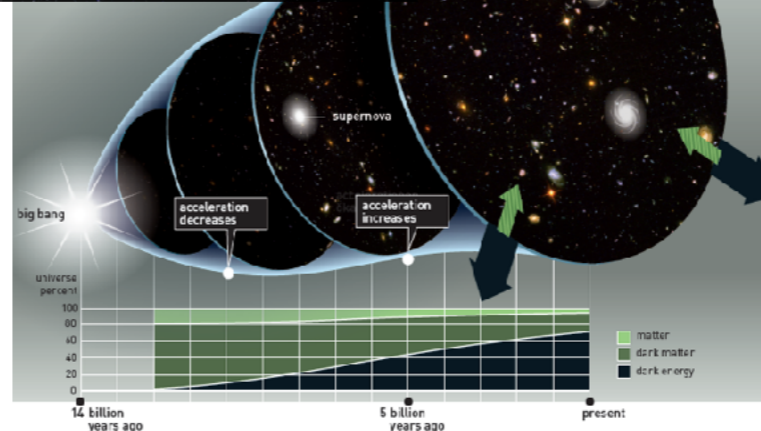
Limitations: Observations

- Dark Matter

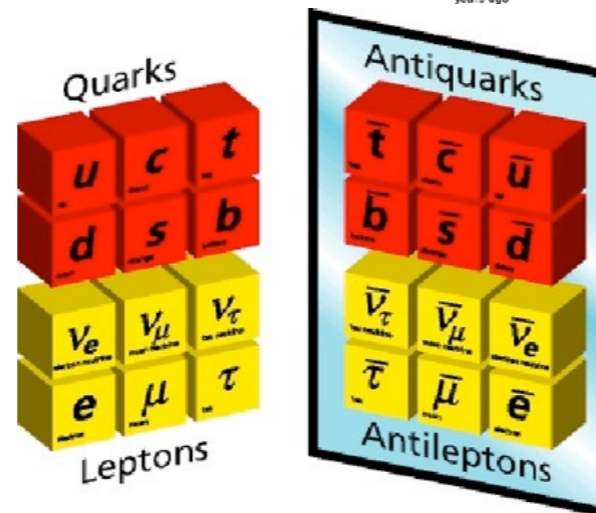


New particles?

- Dark Energy



- Baryon asymmetry



New interactions?

- Neutrino masses

Ideas Beyond: Grand Unified Theories



- The simplest symmetry that contains $U(1)$, $SU(2)$ and $SU(3)$: $SU(5)$ (Georgi, Glashow 1974)

Ideas Beyond: Grand Unified Theories

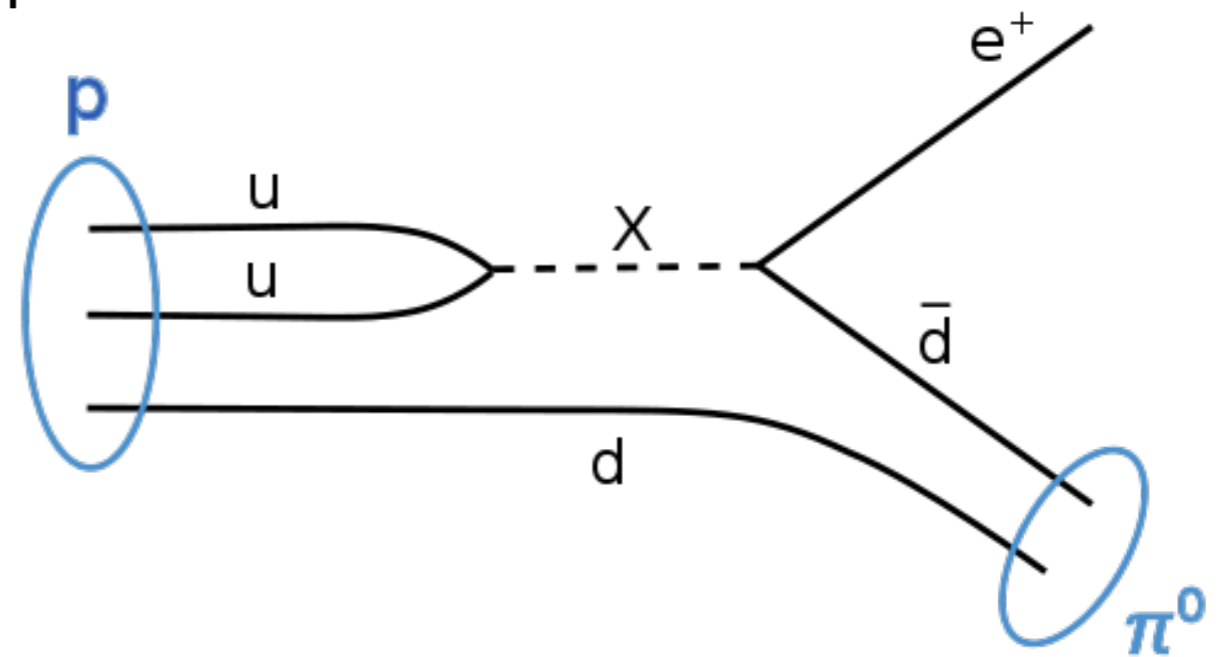
The particle structure in SU(5) GUT:

- Multiplets of (known) leptons and quarks, which can be transformed by the exchange of new heavy bosons (“leptoquarks”) X, Y with $-1/3$ and $-1/4$ charge

a direct consequence: Proton decay via $p \rightarrow \pi^0 e^+$

$$\tau_p \sim \frac{M_X^4}{\alpha_{GUT}^2 M_p^5} \sim 10^{30 \pm 1} \text{ yr}$$

for $M_X \sim 10^{15} \text{ GeV}$



already excluded by SuperKamiokande: Standard SUSY-GUT excluded.

more next week!

Ideas Beyond: Grand Unified Theories

- Electric charge is one of the generators of SU(5)
 - Quantisation of charge follows from exchange rules
 - Sum of all charges in each fermion multiplet = 0
(e.g. each family: neutrino, lepton, 3 x up-type quark, 3 x down-type quark)

=> Explains 1/3 charges of quarks by existence of 3 colors

=> Guarantees equal charge of proton and electron

Additional consequences:

- Small, but finite neutrino mass
- Existence of heavy magnetic monopoles

Coupling Constants: Unification?

- For GUTs: unification of running coupling constants?

$$\alpha_1(M_X) = \alpha_2(M_X) = \alpha_3(M_X) \quad \text{with:} \quad \alpha_1 = 8 \alpha_{\text{em}}/3 = 8(e^2/4\pi)/3 ;$$

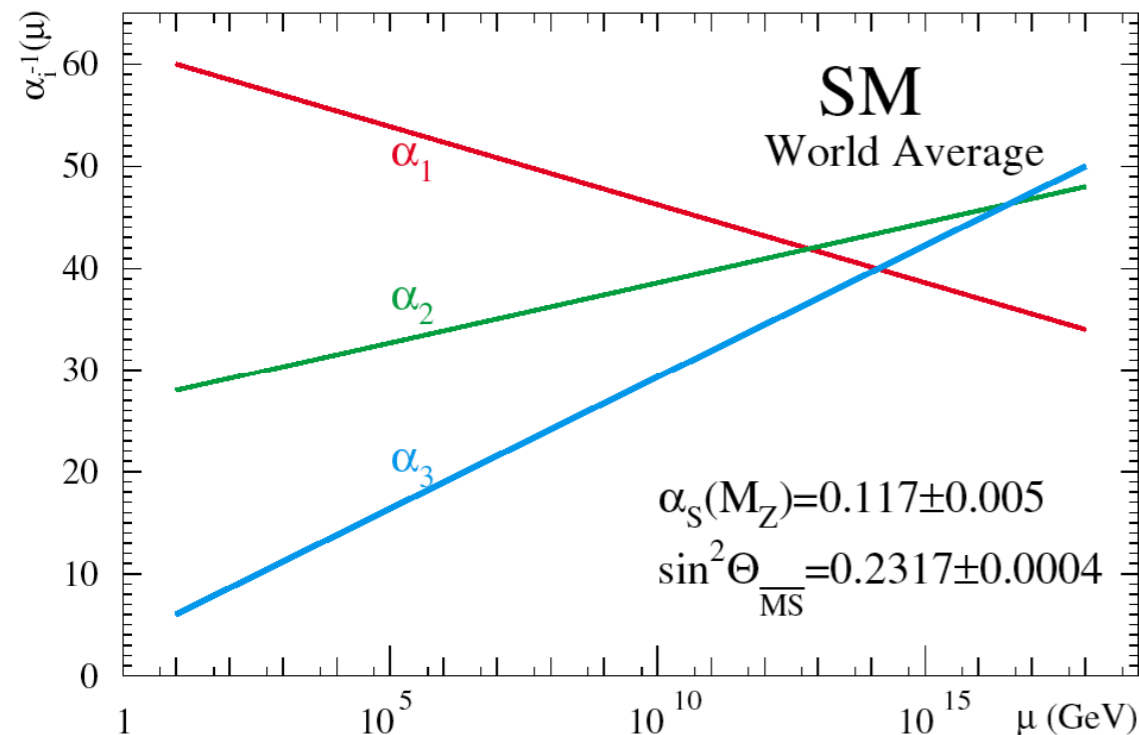
$$\alpha_2 = g^2/4\pi; \quad (g = e / \sin\theta_w)$$

$$\alpha_3 = \alpha_s$$

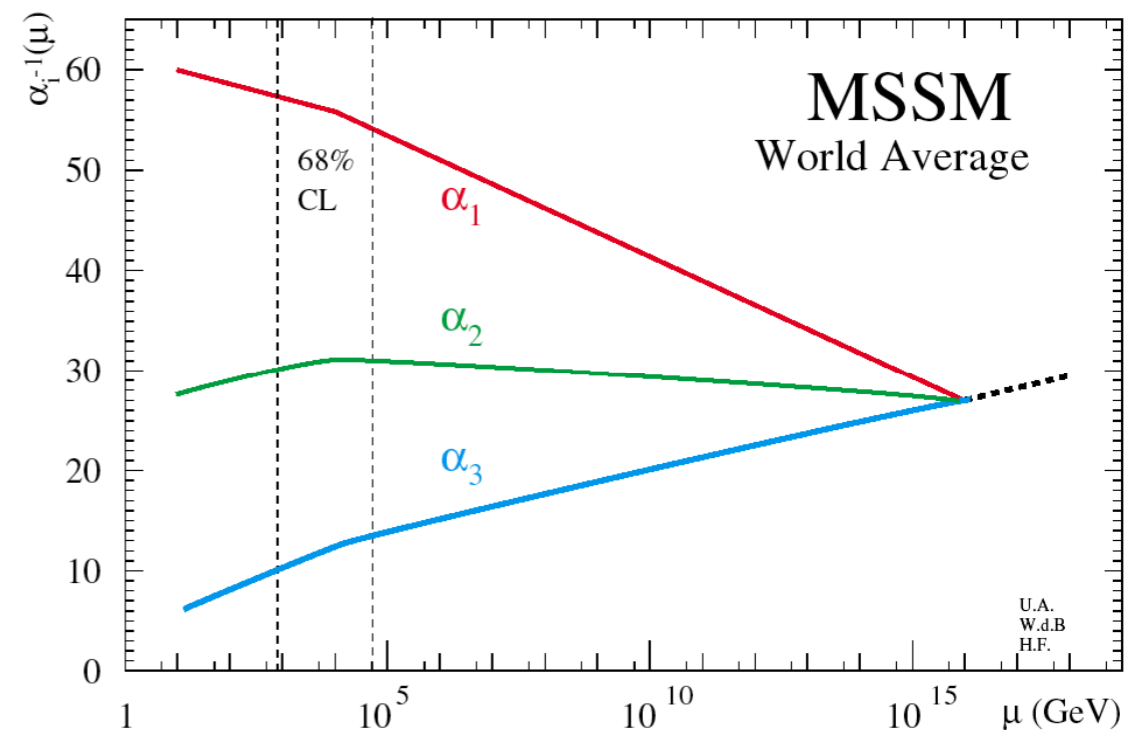
energy dependence: $\alpha(q^2) = \frac{\alpha(\mu^2)}{1 - \beta_0 \alpha(\mu^2) \ln(q^2 / \mu^2)}$; mit $-\beta_0 = \frac{11N_c - 4N_f}{12\pi}$

$N_c = 0, 2, 3$ für U(1), SU(2), SU(3),

$N_f = 3$ (Number of fermion generations)

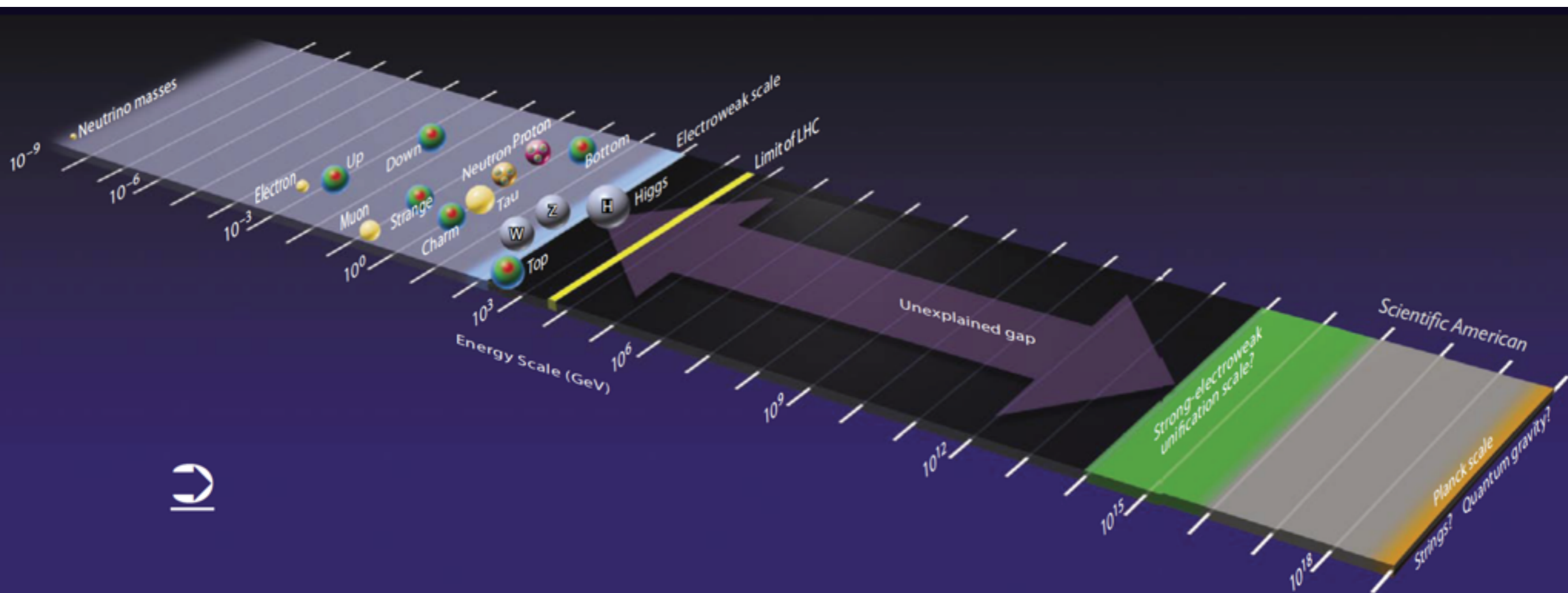


additional SUSY particles in loops:



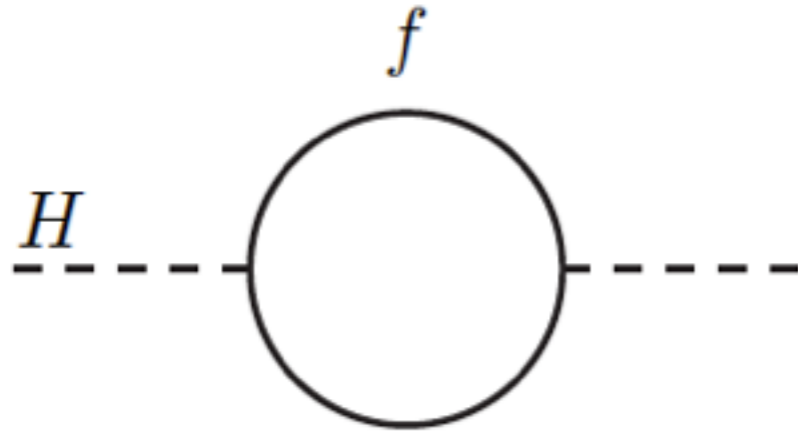
The Hierarchy Problem: A Closer Look

- The problem: Two mass scales: The Planck scale ($\sim 10^{19}$ GeV) and the electroweak scale ($\sim 10^2$ GeV) - separated by 17 orders of magnitude!
 - The consequence: Gravitation is much weaker than all other interactions
 - In the Standard Model: Higgs-Mass of 125 GeV: How is this stabilized?



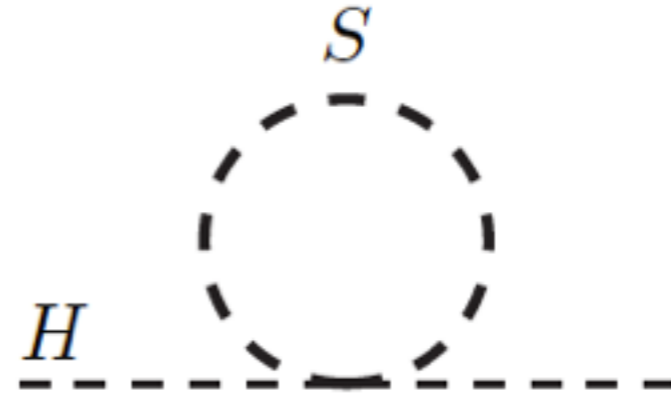
The Hierarchy Problem: Quantum Corrections

Quantum corrections to the Higgs mass via particle loops



Contribution of fermion loops

$$\Delta m_H^2 = -\frac{|\lambda_f|^2}{8\pi^2} \Lambda_{UV}^2 + \dots$$



contribution of boson loops

$$\Delta m_H^2 = \frac{\lambda_S}{16\pi^2} \left[\Lambda_{UV}^2 - 2m_S^2 \ln(\Lambda_{UV}/m_S) + \dots \right]$$

Corrections depend on the scale until which they have to be taken into account - The natural UV scale: M_P , results in enormous corrections due to the Λ^2 behavior

In the SM: Largest contribution from the top quark: Strongest coupling to the Higgs field, $\lambda_f \sim 1$

Also so far unknown, heavy particles contribute, the heaviest particle which couples to the Higgs field dominates!

⇒ Requires unnatural fine tuning to save a small Higgs mass

A popular Idea: Supersymmetry

- The strategy for a solution is suggested by the correction terms:

$$\Delta m_H^2 = -\frac{|\lambda_f|^2}{8\pi^2} \Lambda_{UV}^2 + \dots \quad \Delta m_H^2 = \frac{\lambda_S}{16\pi^2} \left[\Lambda_{UV}^2 - 2m_S^2 \ln(\Lambda_{UV}/m_S) + \dots \right]$$

Contributions of fermions and bosons have opposite sign!

- Cancellation of these contributions is automatic, if there is a symmetry between bosons and fermions, a so-called **Supersymmetry**
- A SUSY - gauge transformation transforms bosons into fermions and vice versa:

$$Q|\text{Boson}\rangle = |\text{Fermion}\rangle, \quad Q|\text{Fermion}\rangle = |\text{Boson}\rangle$$

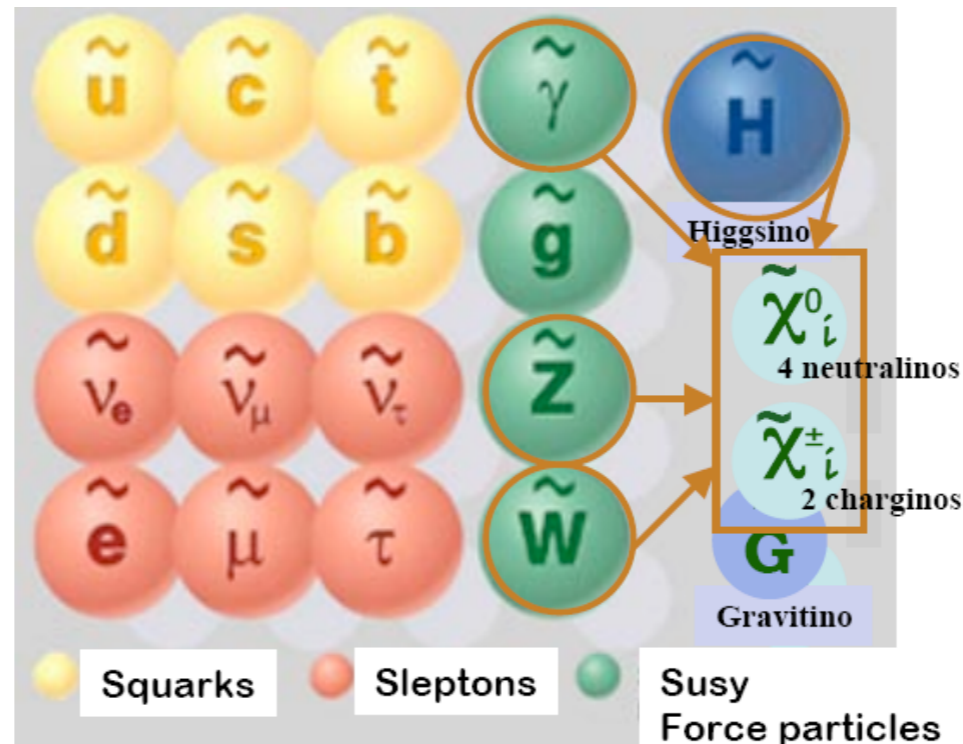
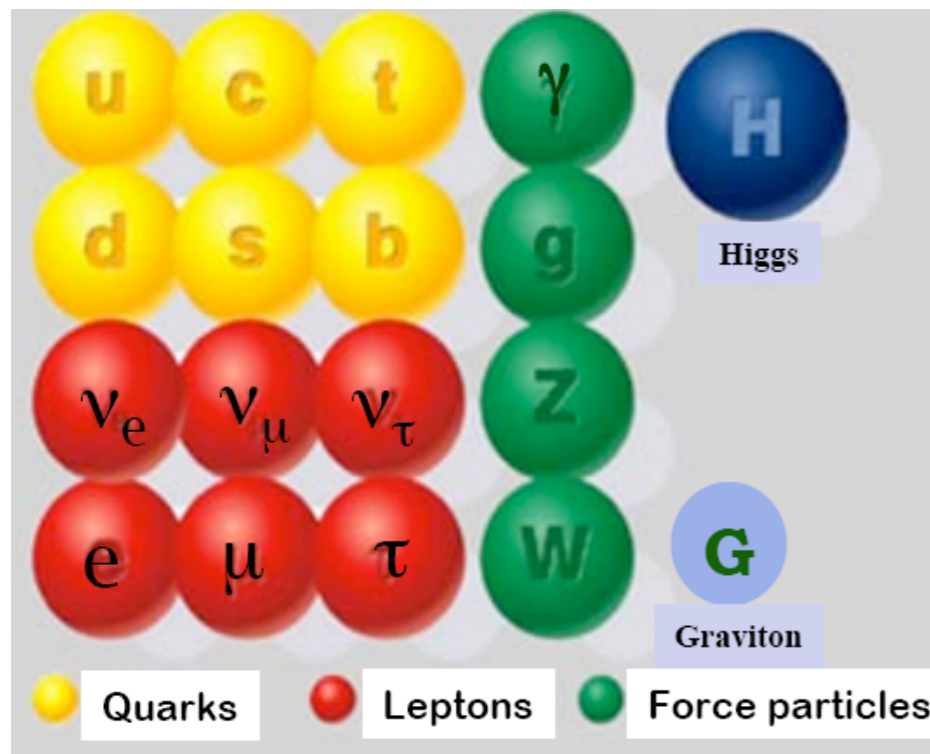
- Consequences of Supersymmetry:
 - Each fermion in the SM has a bosonic “superpartner” with 1/2 different spin, analogous for every boson
 - The partners are arranged in so-called “super-multiplets”
 - For an exact symmetry, the masses of particles and their superpartners are identical

Supersymmetry: Particles & Forces

- Each SM particle gets a supersymmetric partner

| Teilchen | Spin | S-Teilchen | Spin |
|-----------------|------|--------------------------|------|
| Quark Q | 1/2 | Squark \tilde{Q} | 0 |
| Lepton l | 1/2 | Slepton \tilde{l} | 0 |
| Photon γ | 1 | Photino $\tilde{\gamma}$ | 1/2 |
| Gluon g | 1 | Gluino \tilde{g} | 1/2 |
| W^\pm | 1 | Wino \tilde{W}^\pm | 1/2 |
| Z^0 | 1 | Zino \tilde{Z}^0 | 1/2 |

Higgs structure gets more complex: 2 complex doublets, results in 5 physical Higgs fields - Gauginos and Higgsinos mix to form Charginos and Neutralinos



Supersymmetry: Models & Phenomenology

- 128 free parameters to describe masses and couplings of all SUSY particles
 - reduced to a few by adding additional assumptions on breaking mechanism etc.
- New conserved quantity: **R - Parity** - a multiplicative quantity

$$R = (-1)^{3(B-L)+2S} \quad (\text{B/L: baryon, lepton number, S: spin})$$

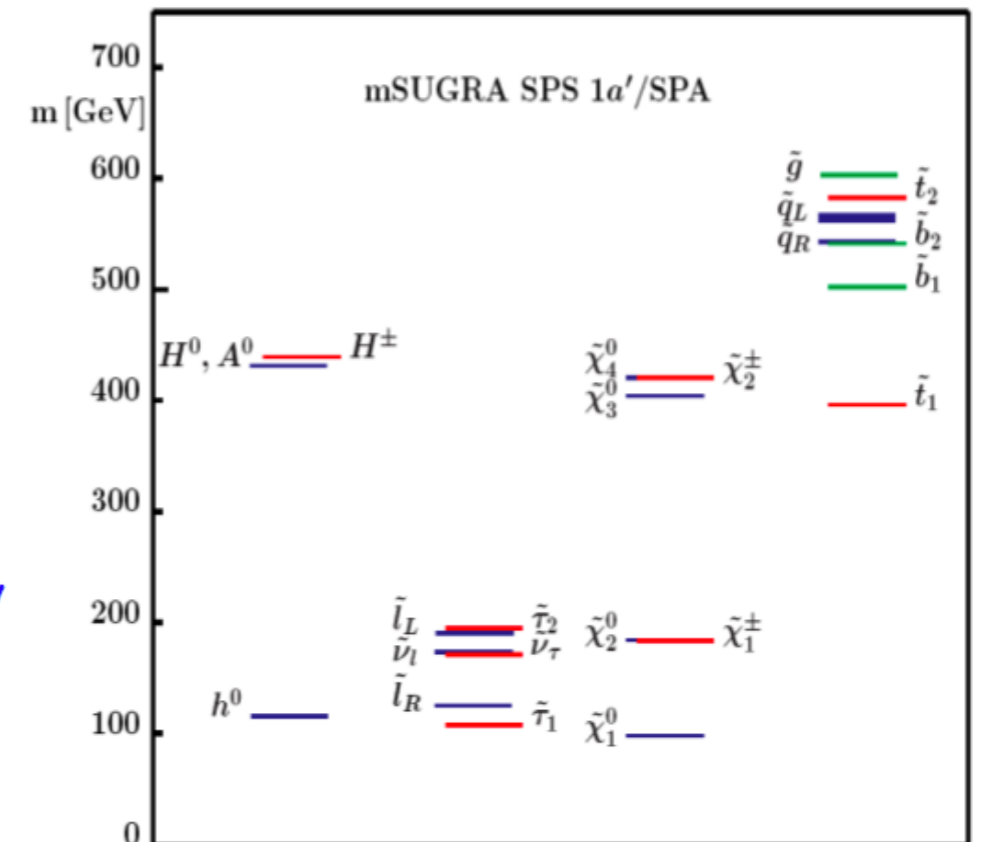
R = 1 for “normal” particles, = -1 for SUSY particles

R parity conservation implies that the lightest SUSY particle (LSP) has to be stable => a good Dark Matter candidate!

N.B.: Also models with R - Parity violation exist...

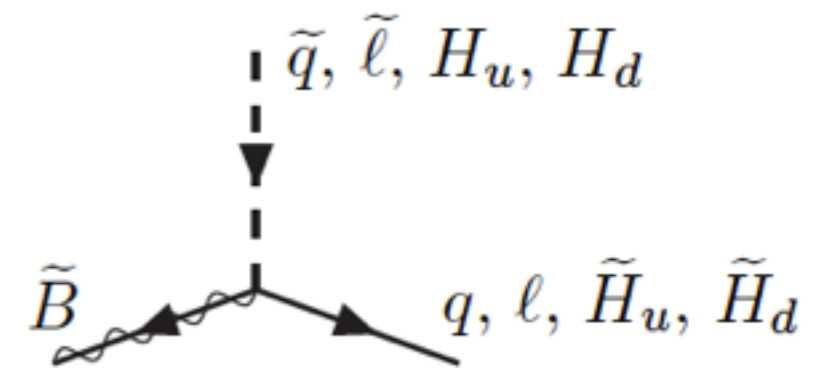
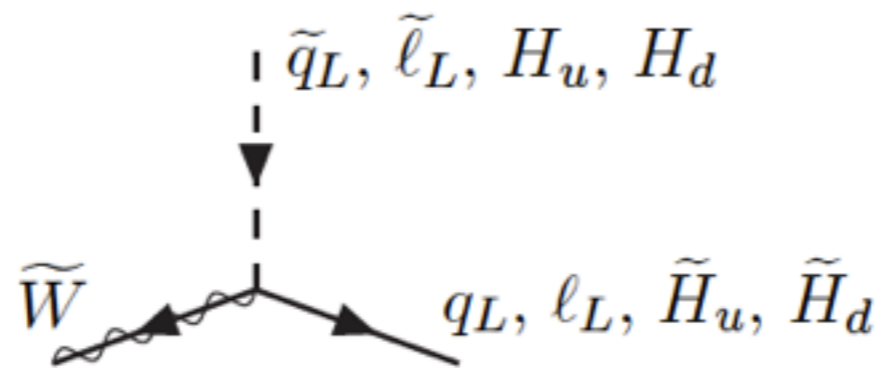
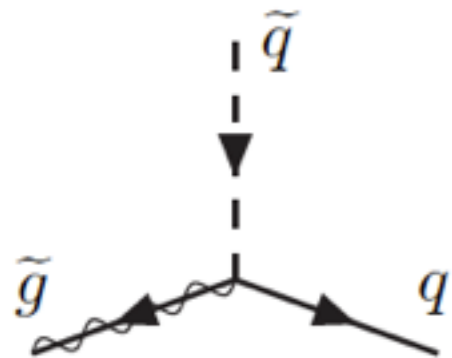
one example of a SUSY particle spectrum
 ... excluded by LHC after a few month of running

- $m_0 = 100 \text{ GeV}$
- $m_{\frac{1}{2}} = 250 \text{ GeV}$
- $A_0 = -100 \text{ GeV}$
- $\tan\beta = 10$
- $\text{sgn}(\mu) = +1$



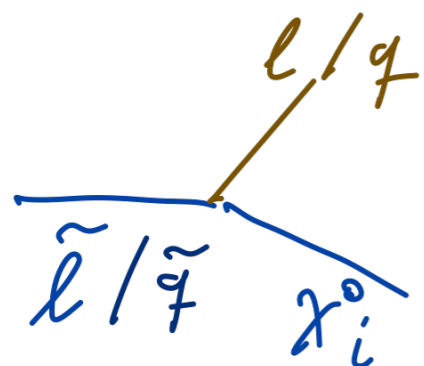
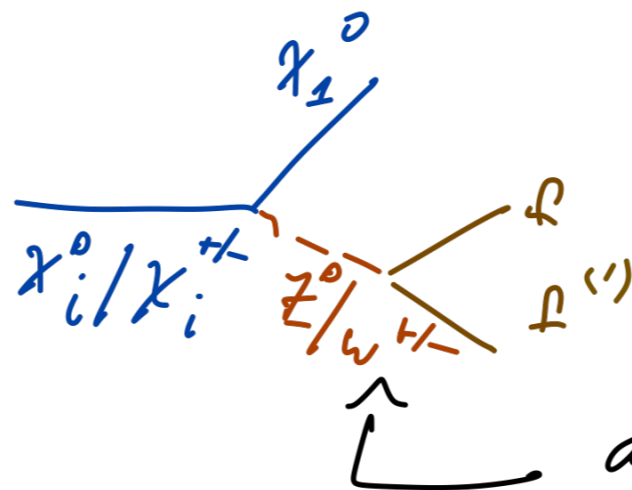
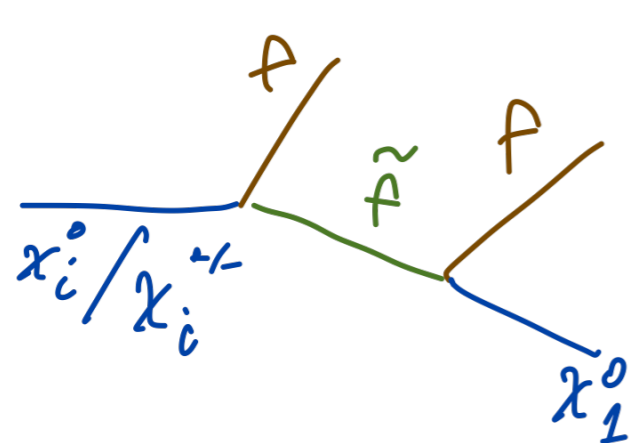
Interaction of SUSY Particles

- SUSY particles interact just as SM particles according to their quantum numbers
 - Right-sfermions (handedness here refers to the SM partners, since the sfermions have spin 0) do not carry weak isospin and therefore do not couple to W bosons
- Coupling to SUSY gauge bosons:

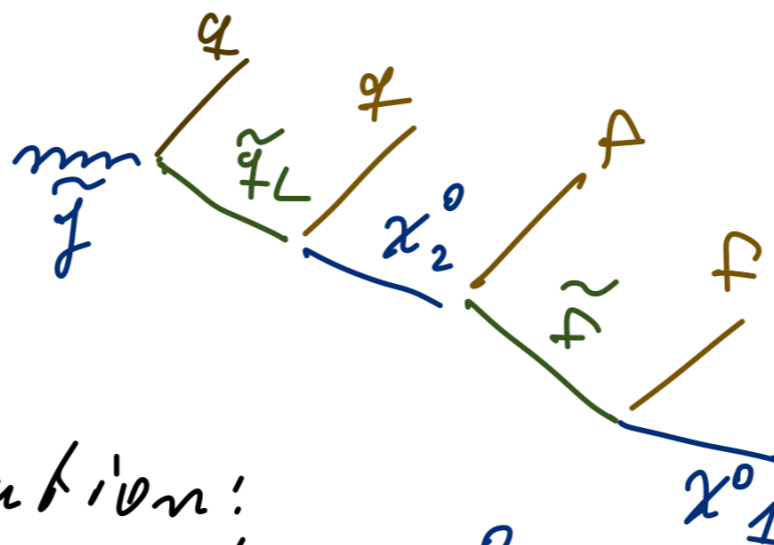
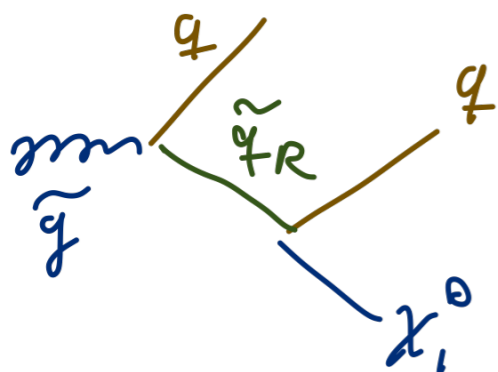


Decay of SUSY Particles

- Depends on spectrum: Ordering of masses - in general:



← also to χ_i^{H-}
 when changing from $\nu \leftrightarrow l^{+-} / q \leftrightarrow q'$

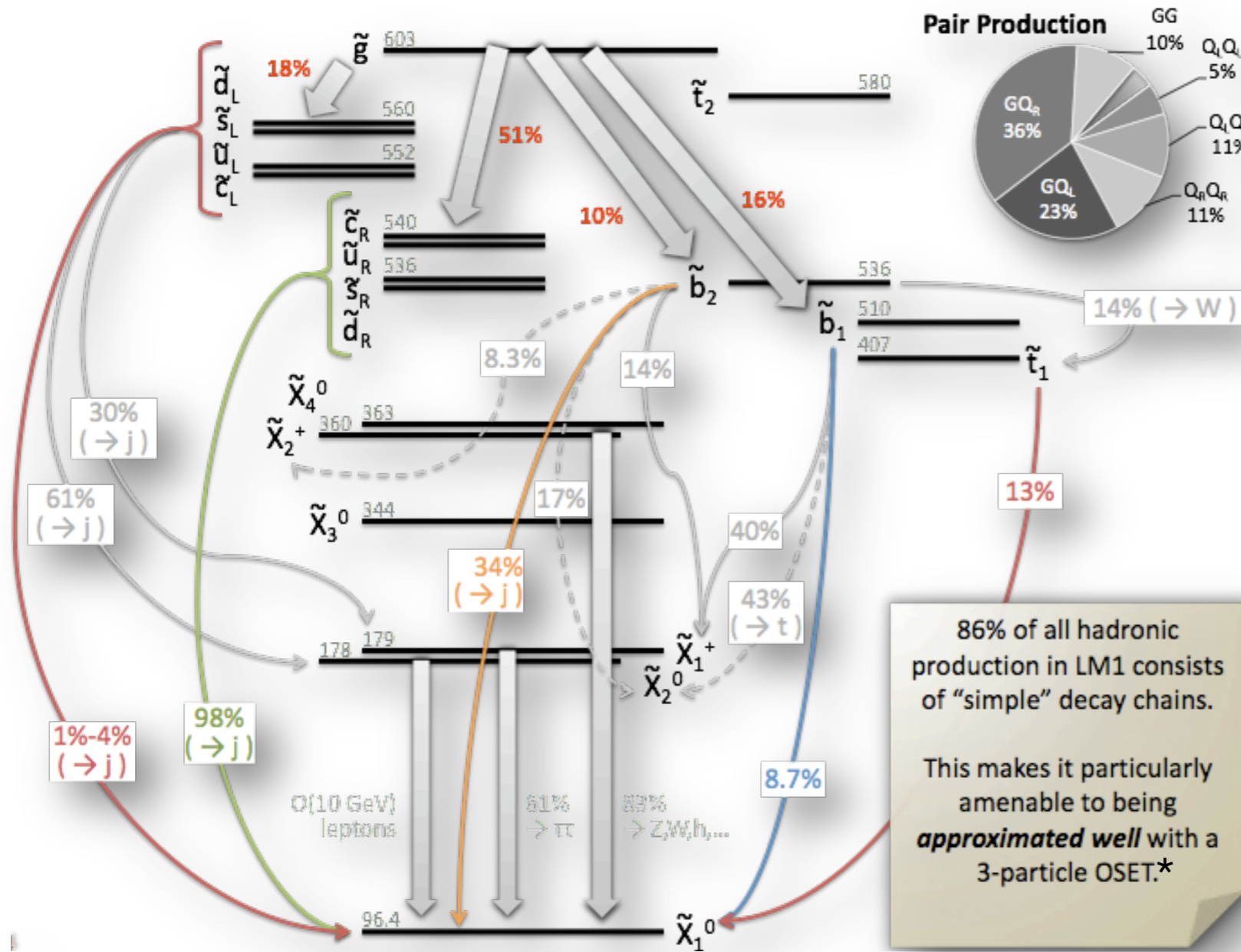


+ other combinations
 of $\tilde{q} + \chi^0 / \chi^{+-}$
 decay chains

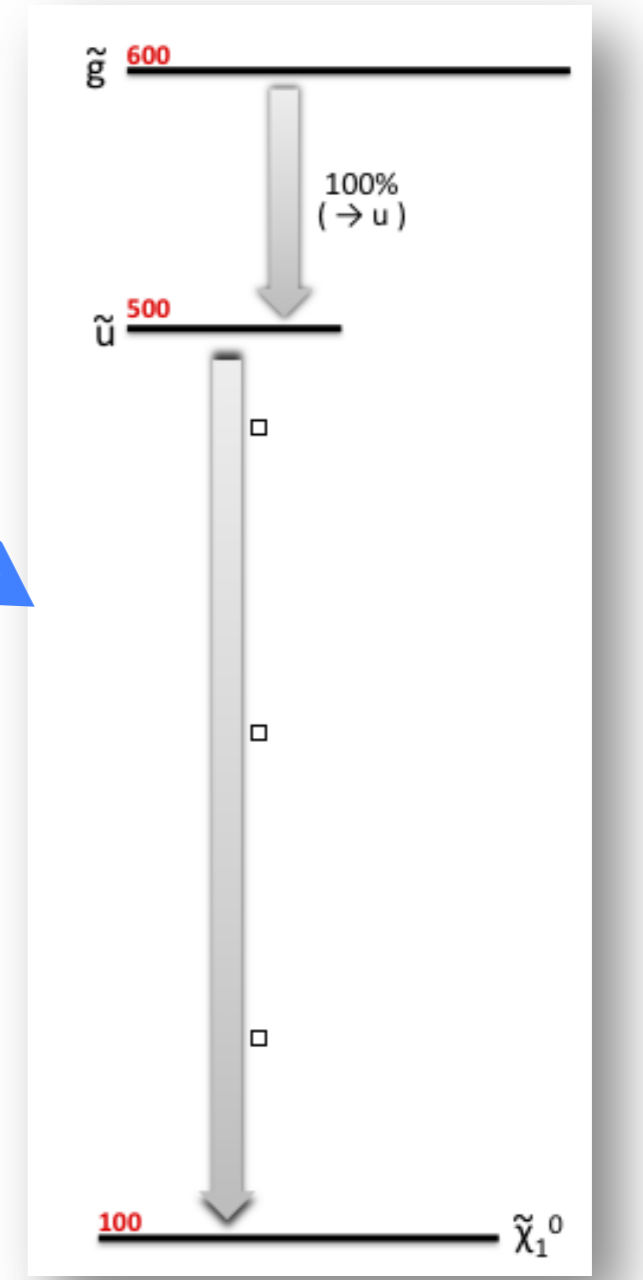
For R-parity conservation:
 Final particles always include χ_1^0 (LSP)

Interpretation of Results - Simplified Models

CMSSM



What the individual searches are sensitive to is much more simple...



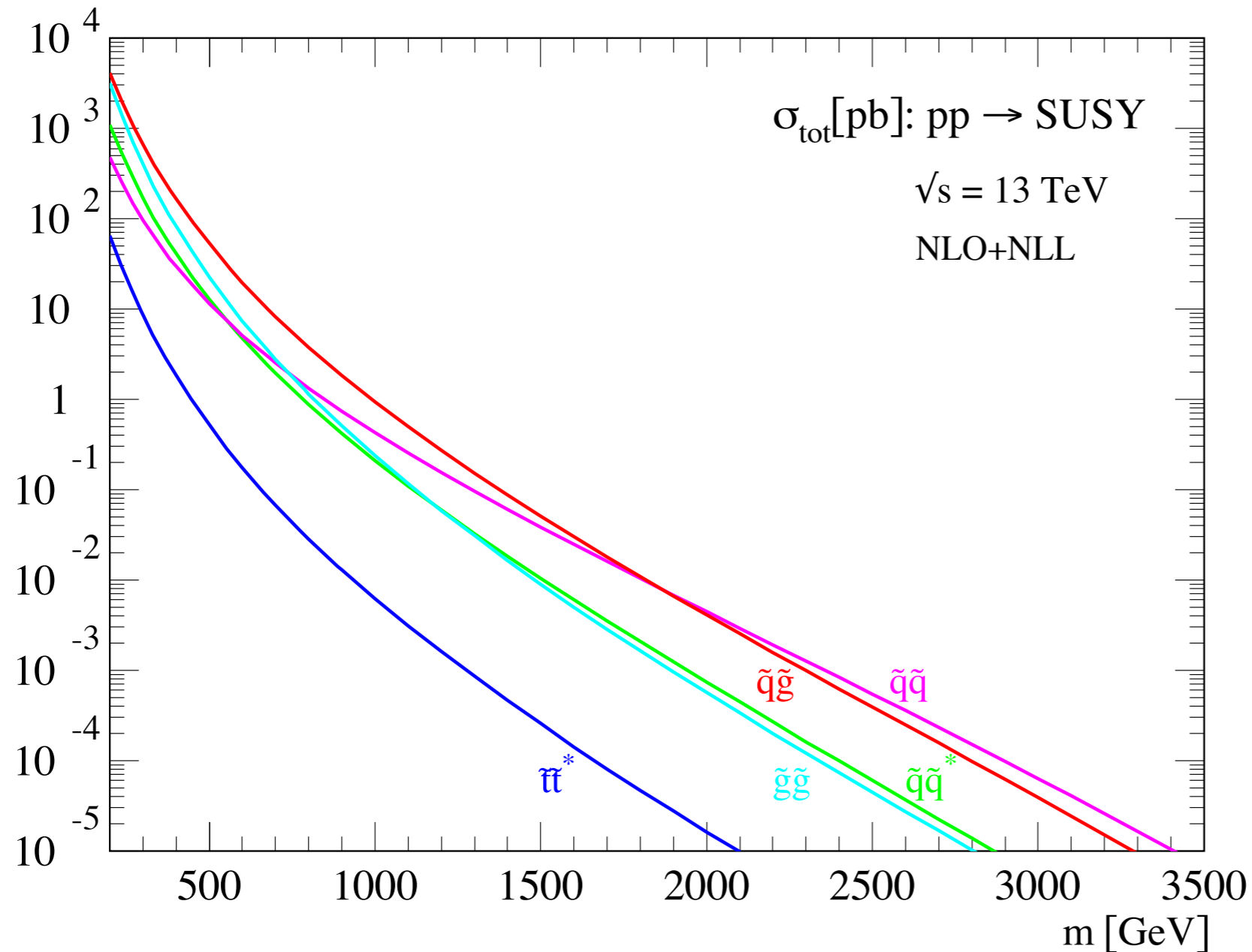
Simplified model spectrum (SMS)
with 3 particles, 2 decay modes

86% of all hadronic production in LM1 consists of "simple" decay chains. This makes it particularly amenable to being approximated well with a 3-particle OSET.*

* OSET = On-Shell Effective Theory

Production of SUSY Particles: General Features

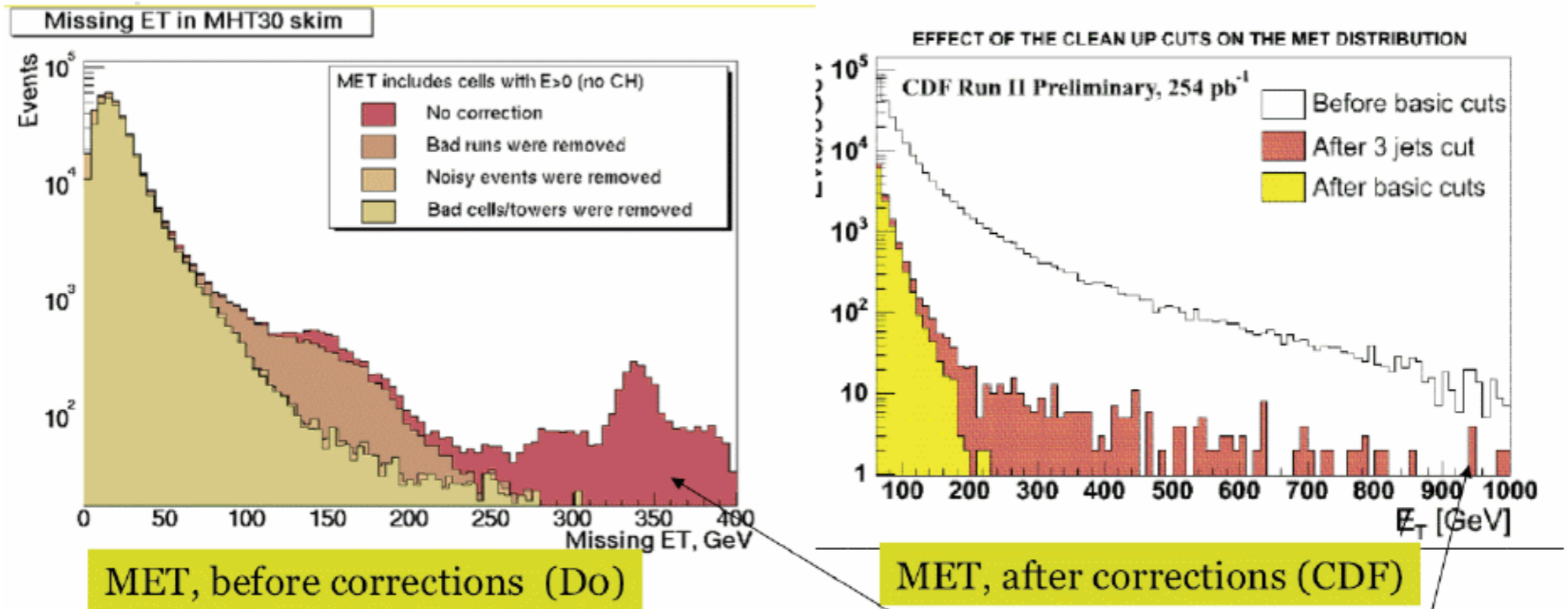
- Cross sections (and with that the mass reach at the LHC) depend strongly on the production mechanism: Highest reach via the strong interaction:
New particles carrying color charge - gluinos, squarks



EPJ C74, 3174 (2014)

Experimental Searches for SUSY - Requirements

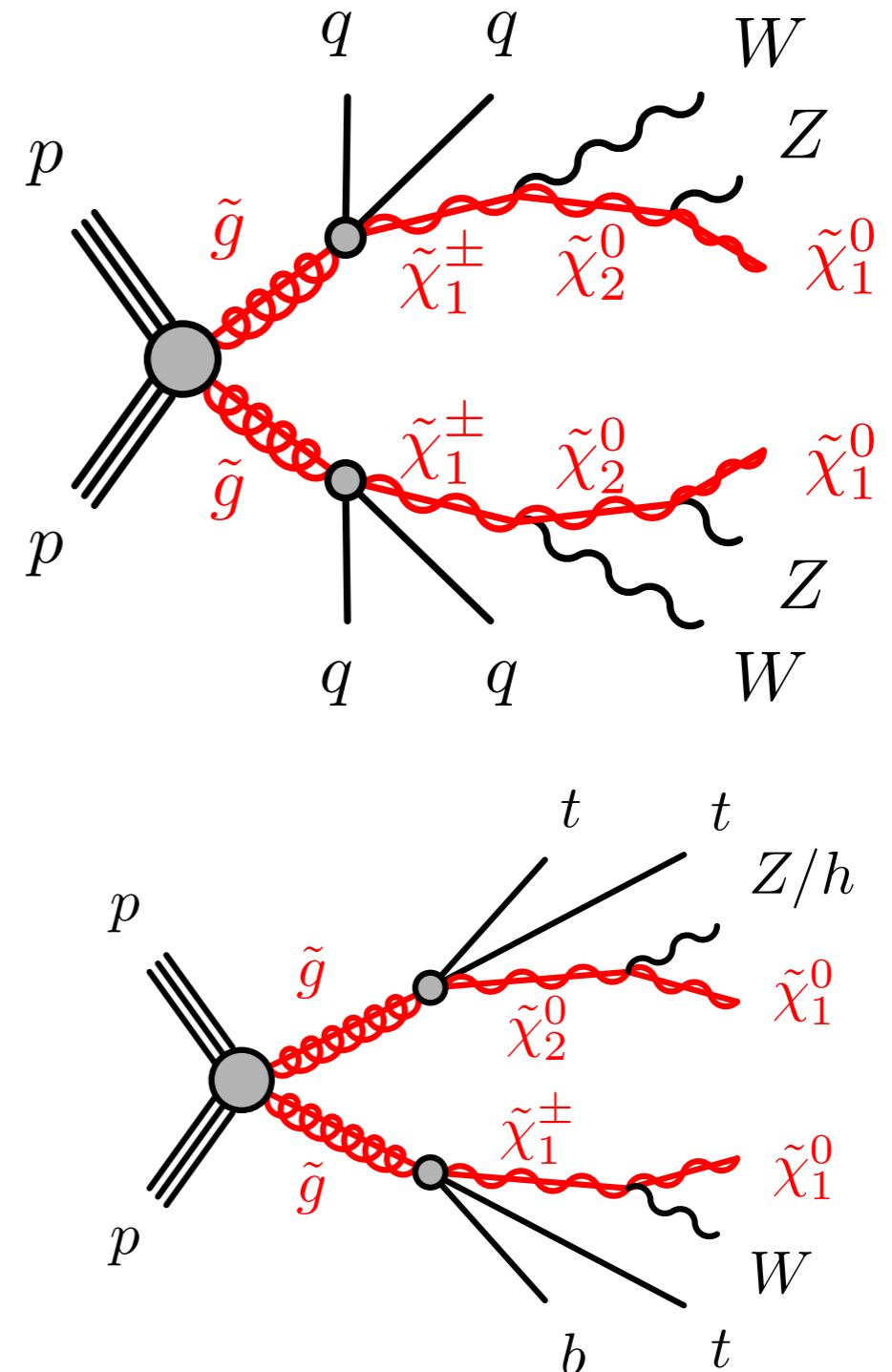
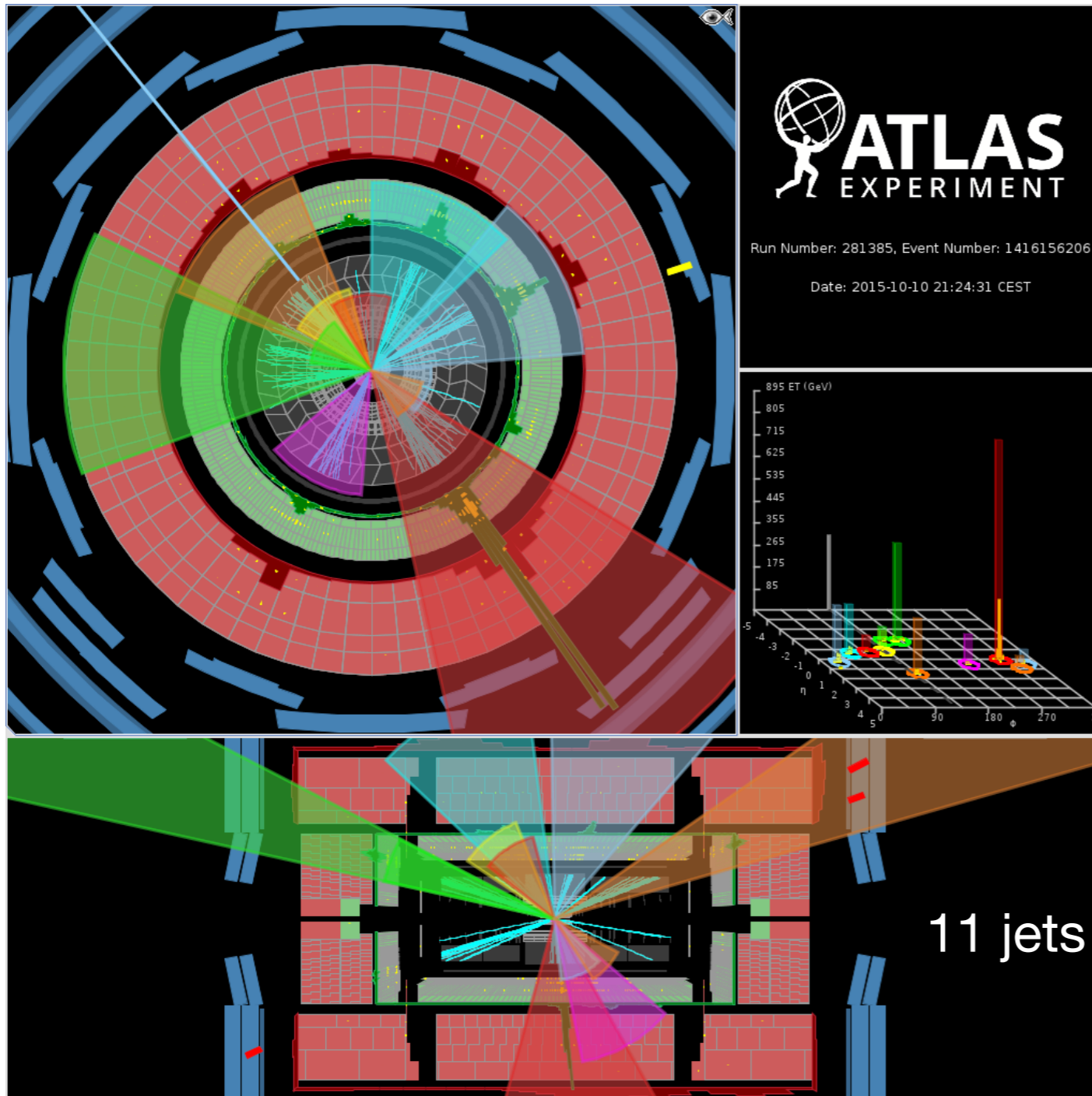
- Key experimental capabilities (and uncertainties) for new physics searches:
 - Jet and lepton reconstruction
 - Hermetic coverage of the events
 - Control of backgrounds and pile-up
 - reconstruction and resolution of missing (transverse) energy



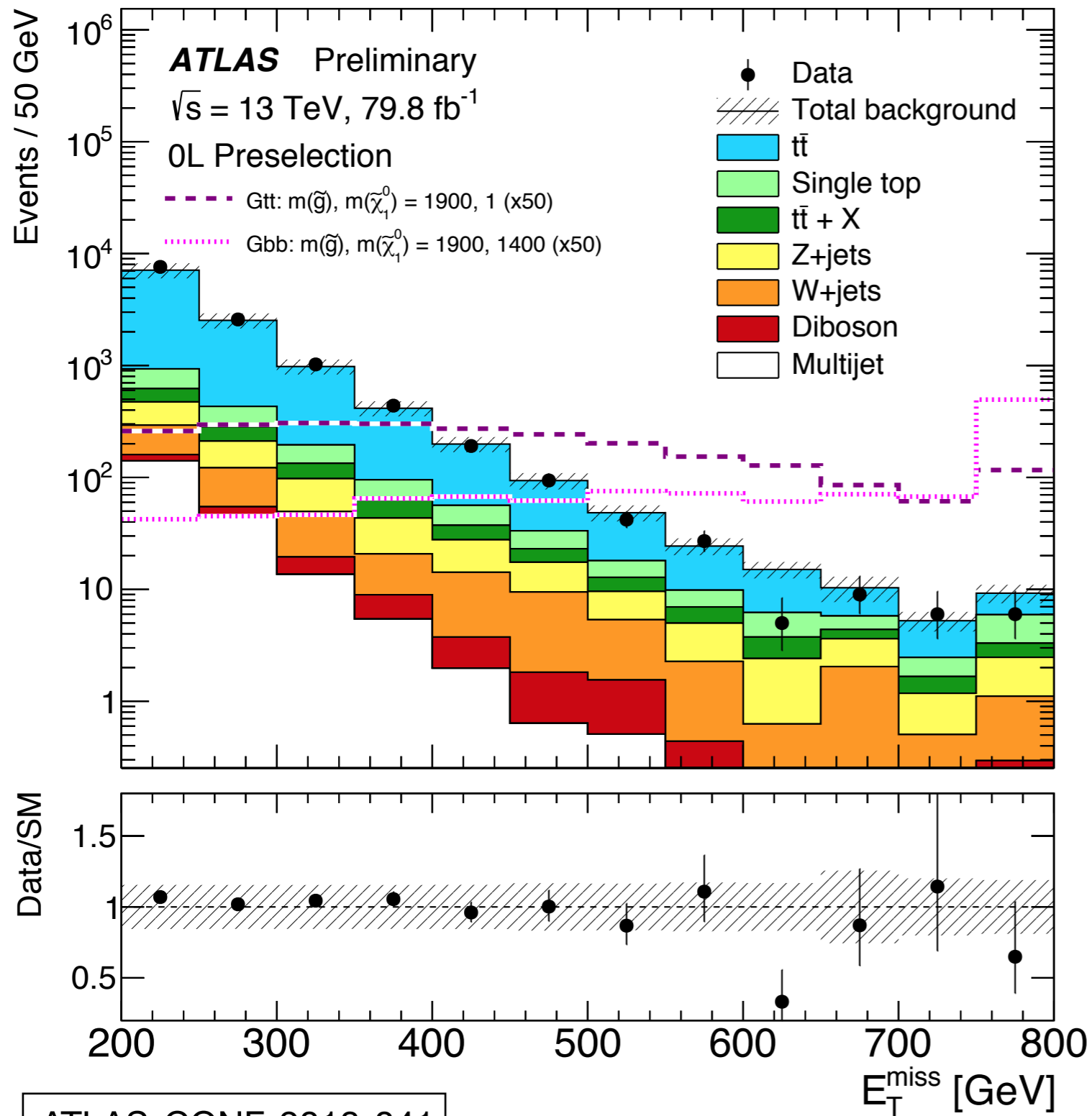
This is where new physics may sit

Experimental Searches for SUSY - Examples

- Search for strongly interacting SUSY: Characterized by multi-jet final states from cascade decays

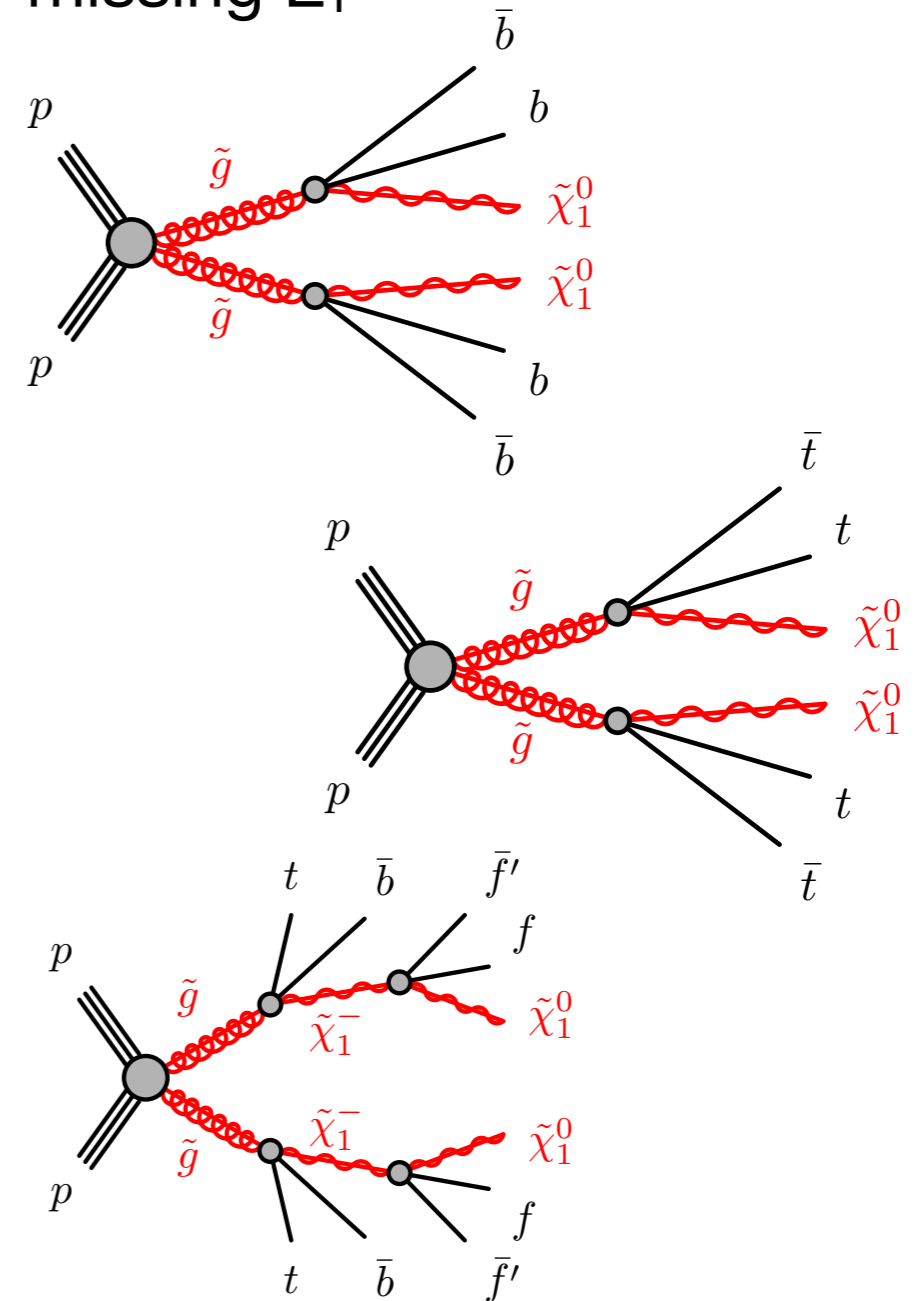


Experimental Searches for SUSY - Examples



ATLAS-CONF-2018-041

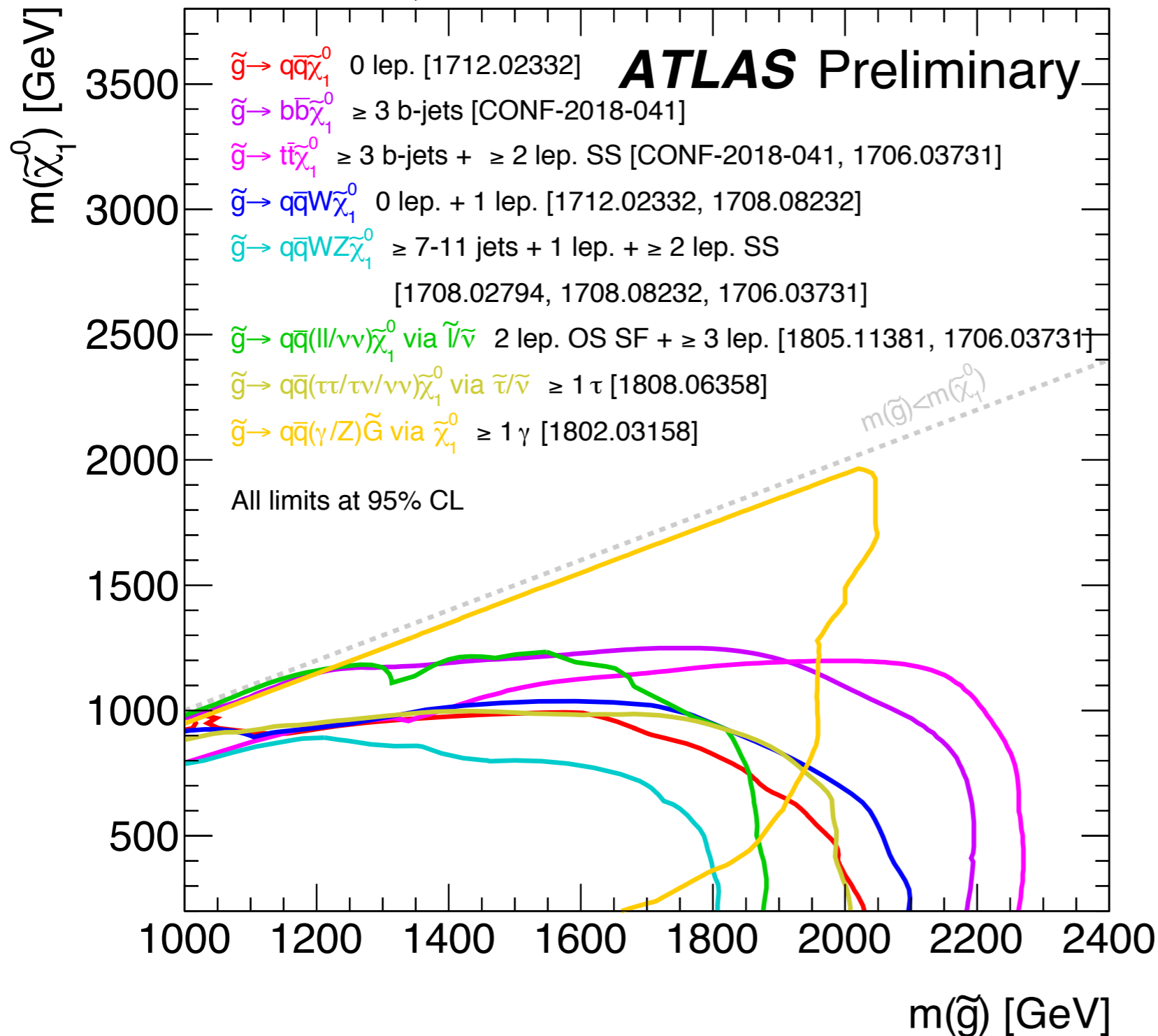
- Looking for a signal excess in different distributions - such as missing E_T



Experimental Searches for SUSY - Examples

$\sqrt{s}=13$ TeV, 36.1 - 79.8 fb⁻¹

September 2018

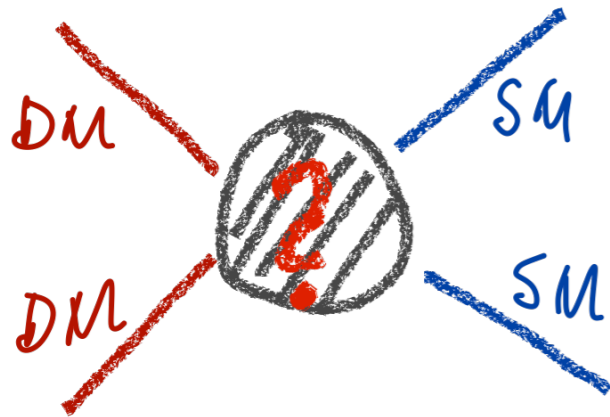


- Exclusion curves for a number of different analyses - as a function of gluino and neutralino (LSP) mass:

Limits beyond 2 TeV for some scenarios for gluinos

Searching for Dark Matter at Colliders - Principles

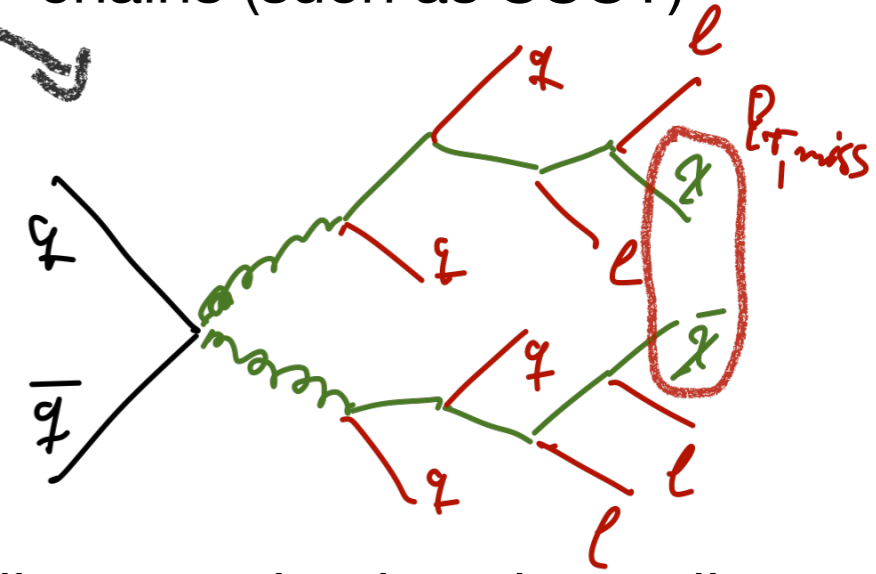
- Related to SUSY searches - but also more general



From cosmology: Expect some interaction between Dark Matter and SM particles - thermal equilibrium in early universe, followed by freeze-out

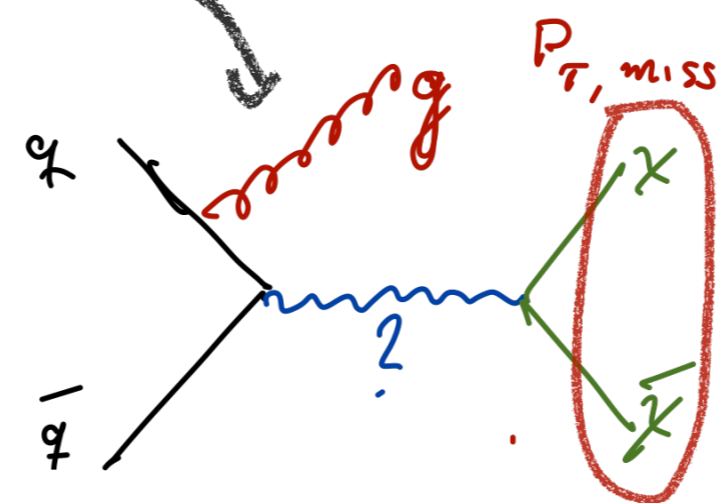
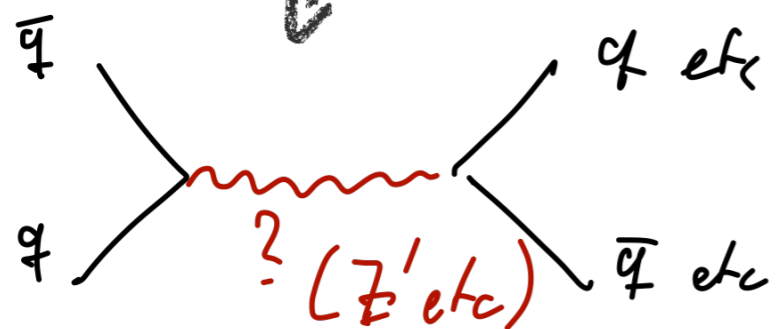
Could produce DM at colliders:
3 different search approaches

detection in complex decay chains (such as SUSY)

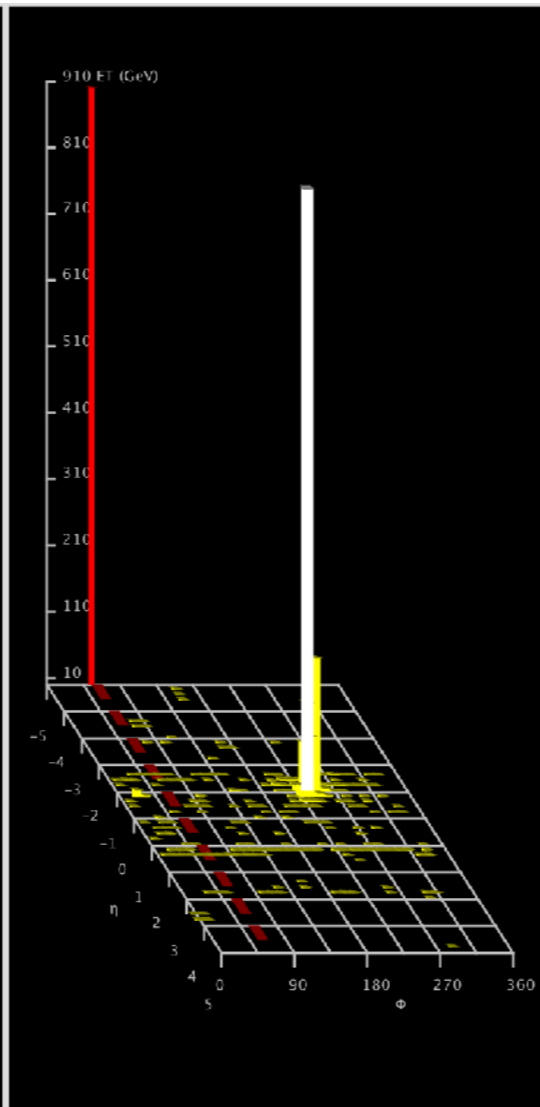
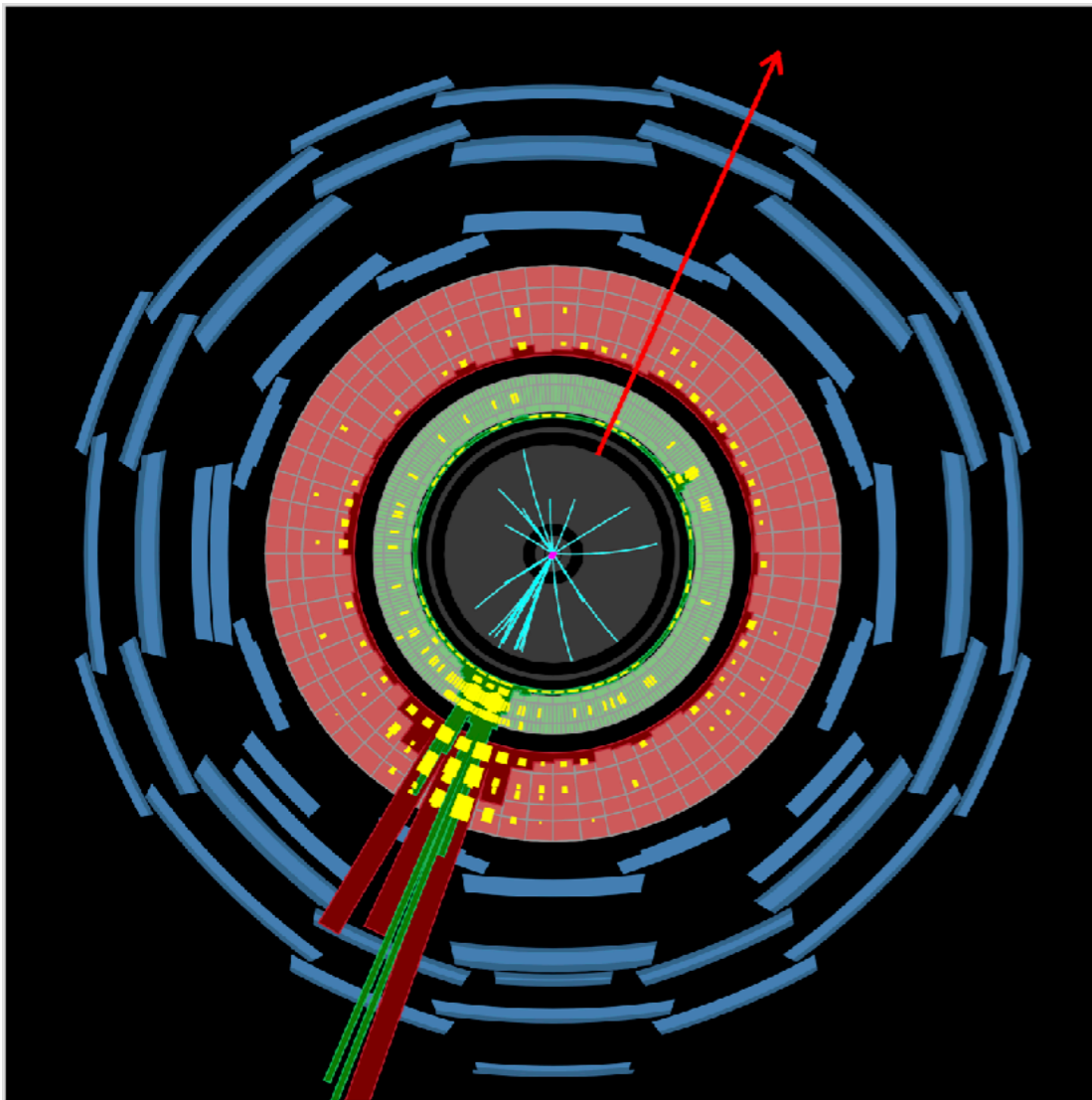


direct production via mediators, tagging of interaction

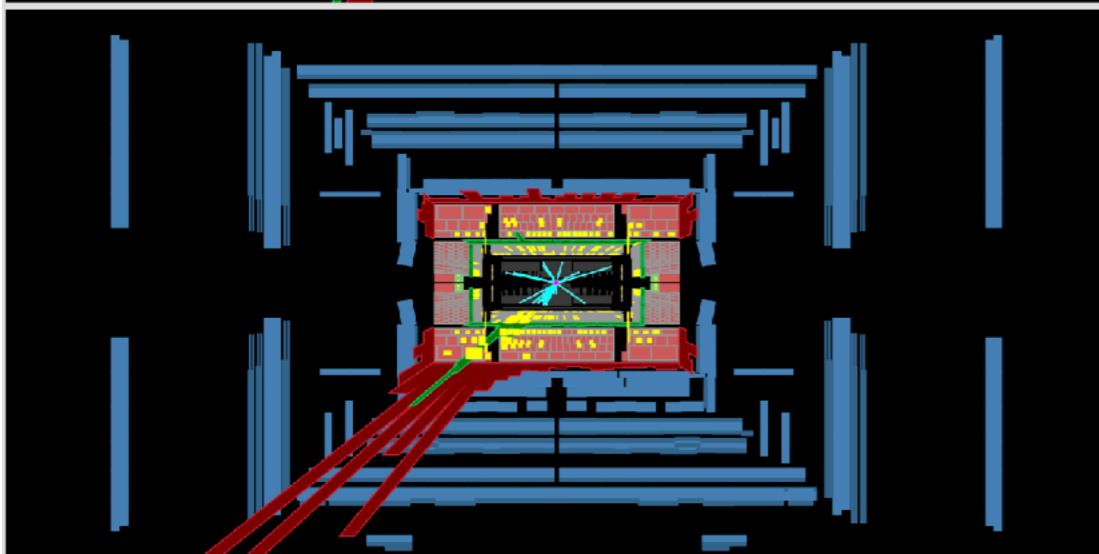
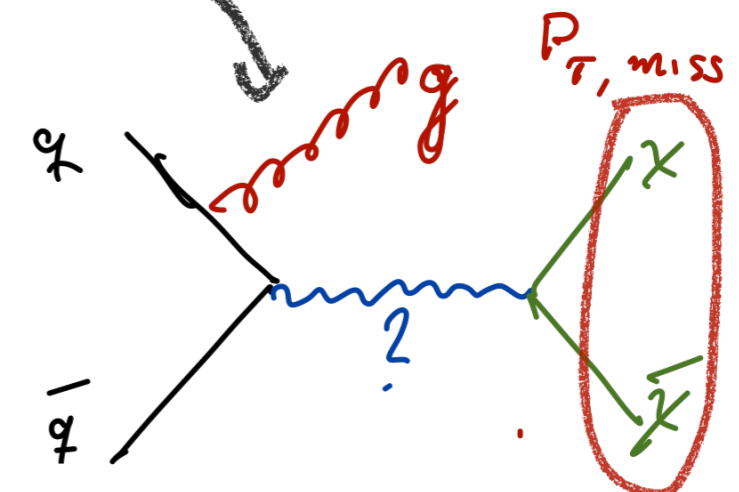
search for possible mediators



Dark Matter / SUSY Searches: Monojets



- Mono-jet (or mono-photon) signatures: Detecting the production of invisible final states through initial state radiation





ATLAS

 EXPERIMENT

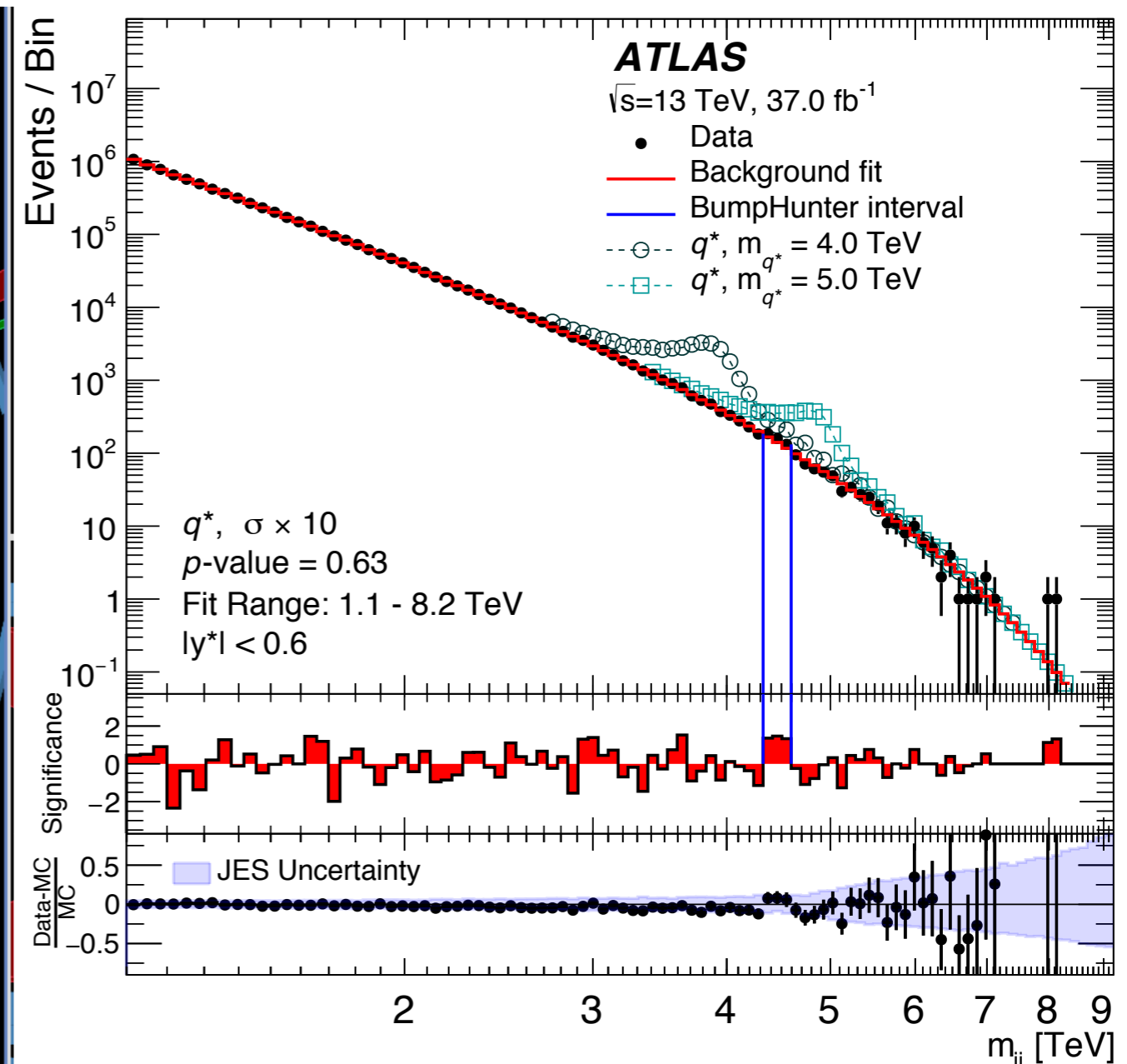
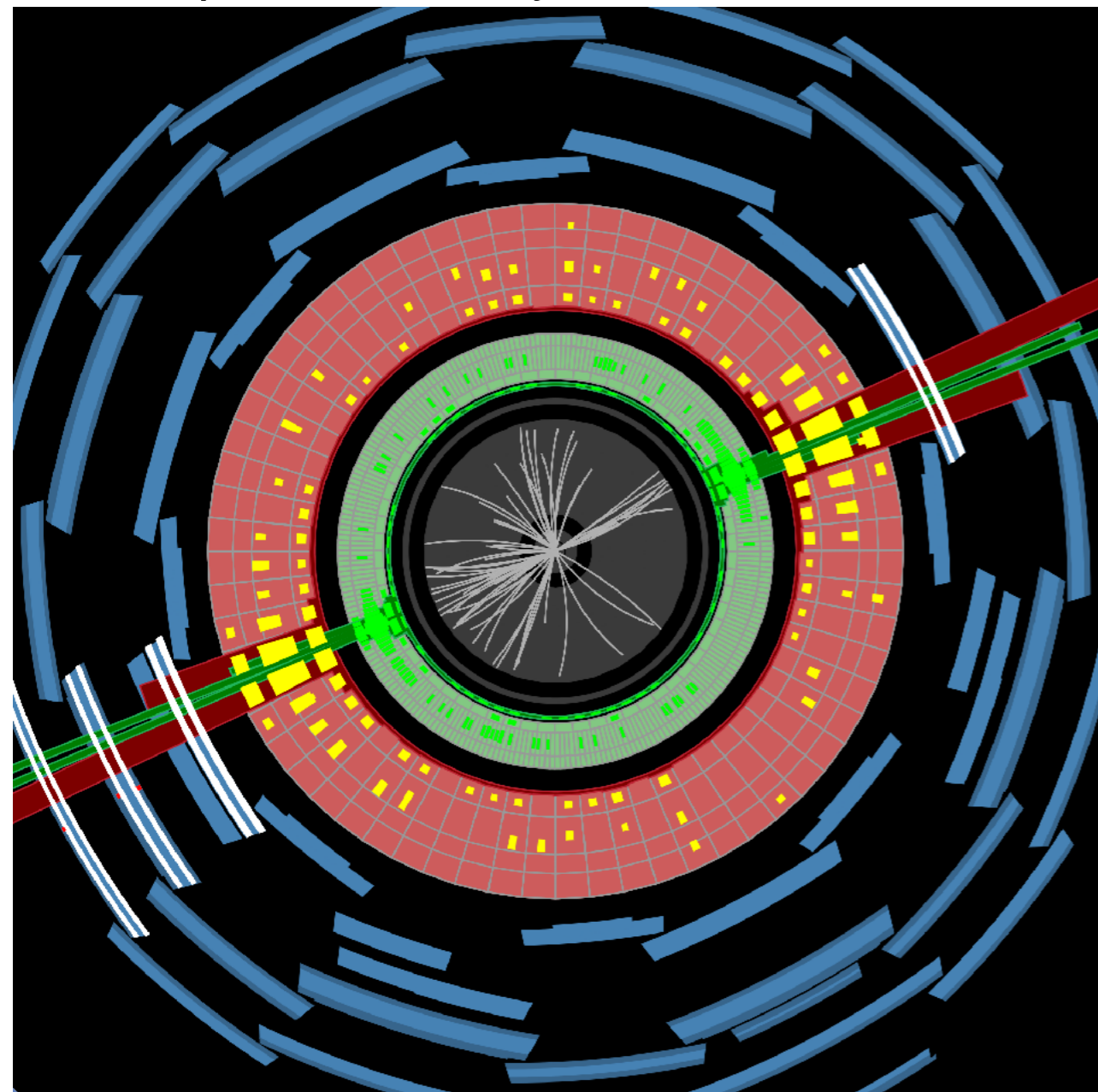
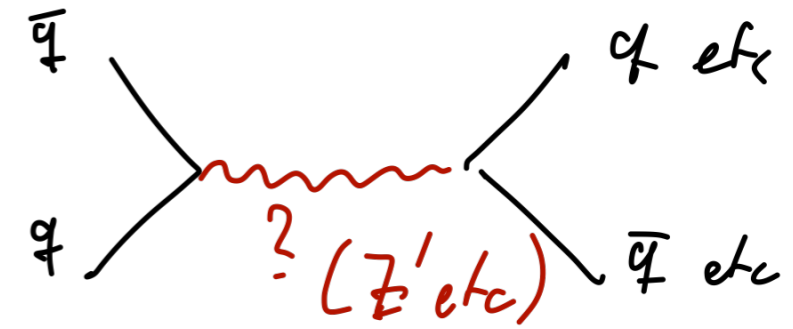
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Date: 2015-09-14 12:05:34 CEST

Searching for Dark Matter at Colliders - Mediators

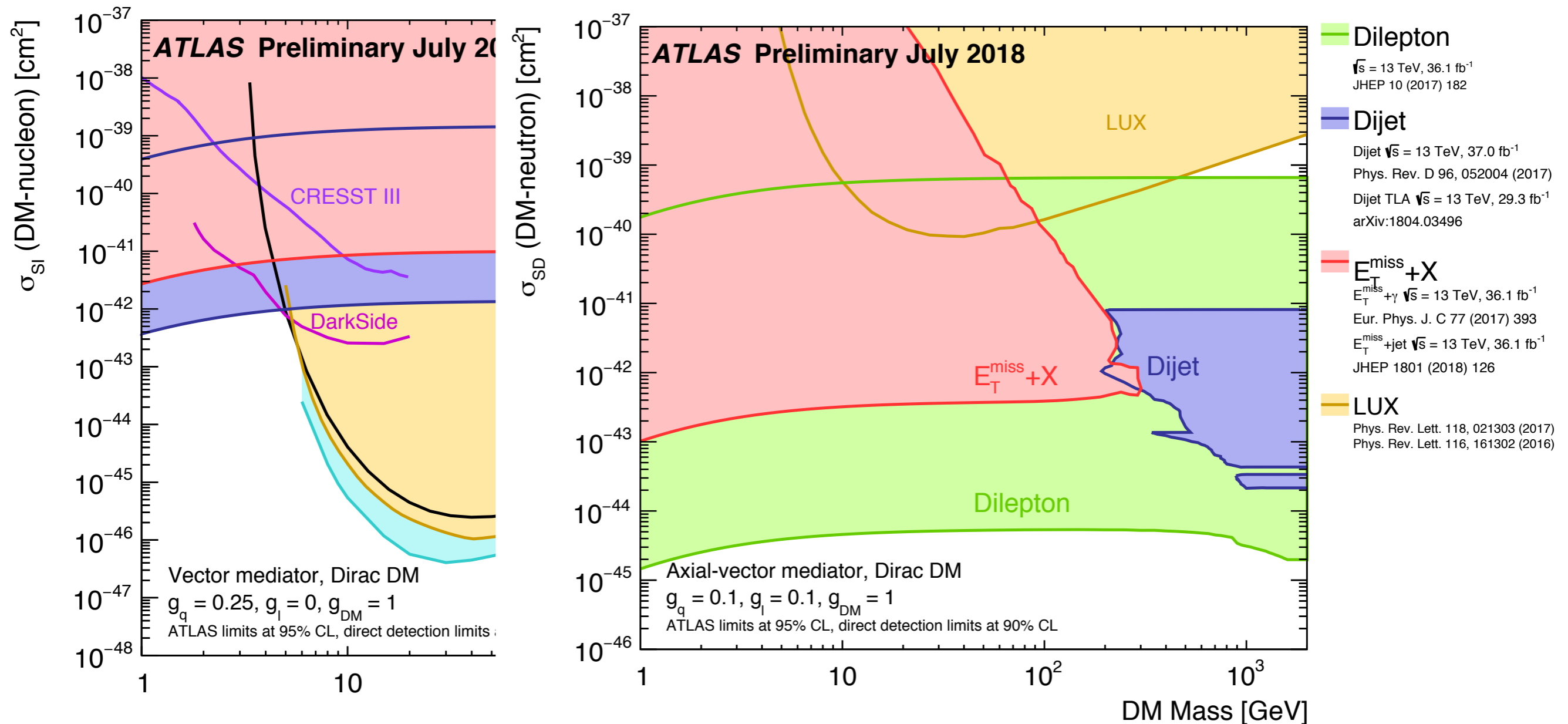
- A search for new force carriers - looking for high-mass resonances

Jet p_T 3.8 TeV, di-jet invariant mass 8.12 TeV



Dark Matter Searches - Limits

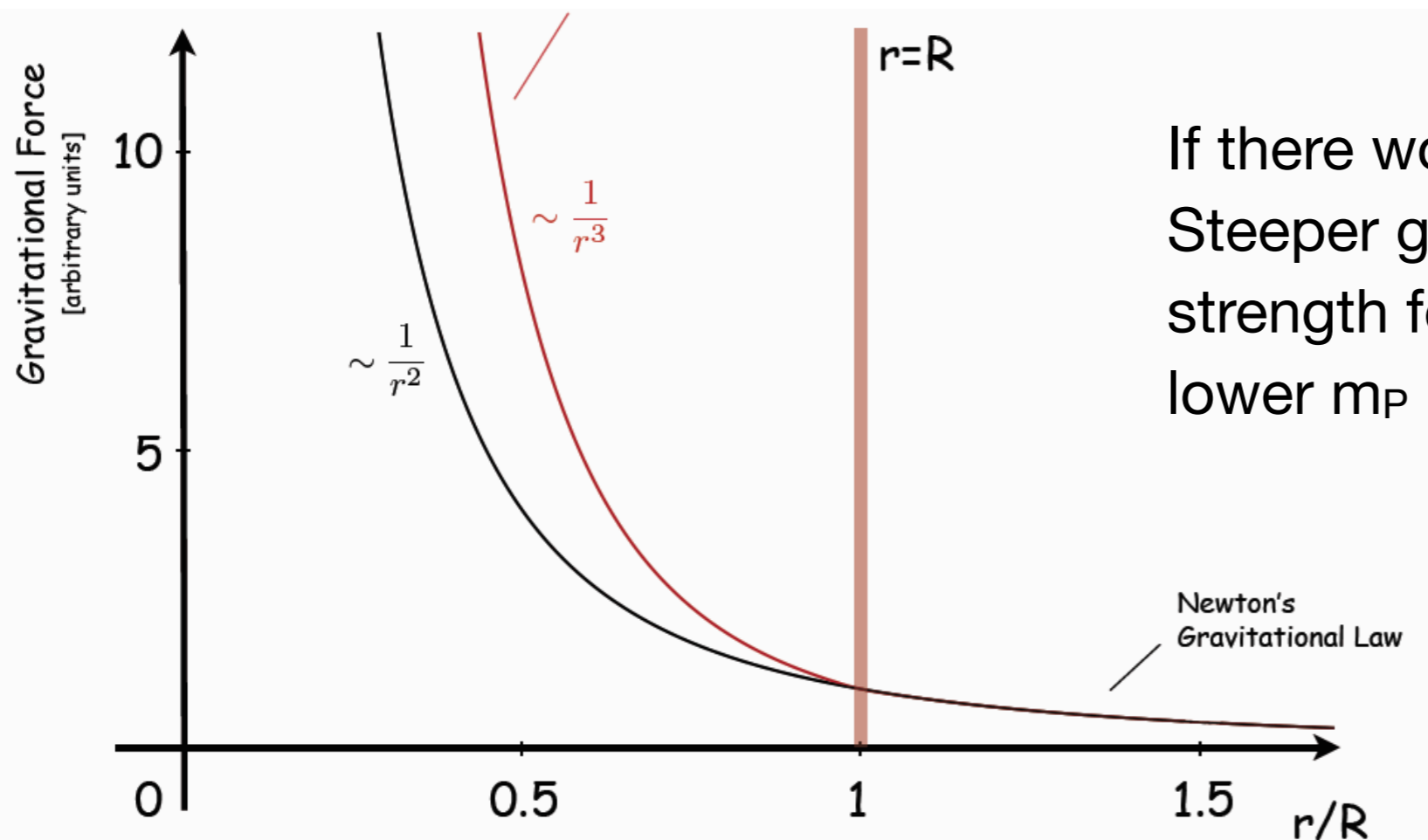
- Limits of collider searches are model-dependent: Assumptions on couplings between DM particles, mediators and SM particles, form of interaction, ...



Extra Space Dimensions

- An approach to solve the hierarchy problem from the side of gravity:
Lowering the Planck scale
 - Naively: The Planck scale is the mass that is required for an elementary particle such that its gravity is “strong” (comparable to other forces):

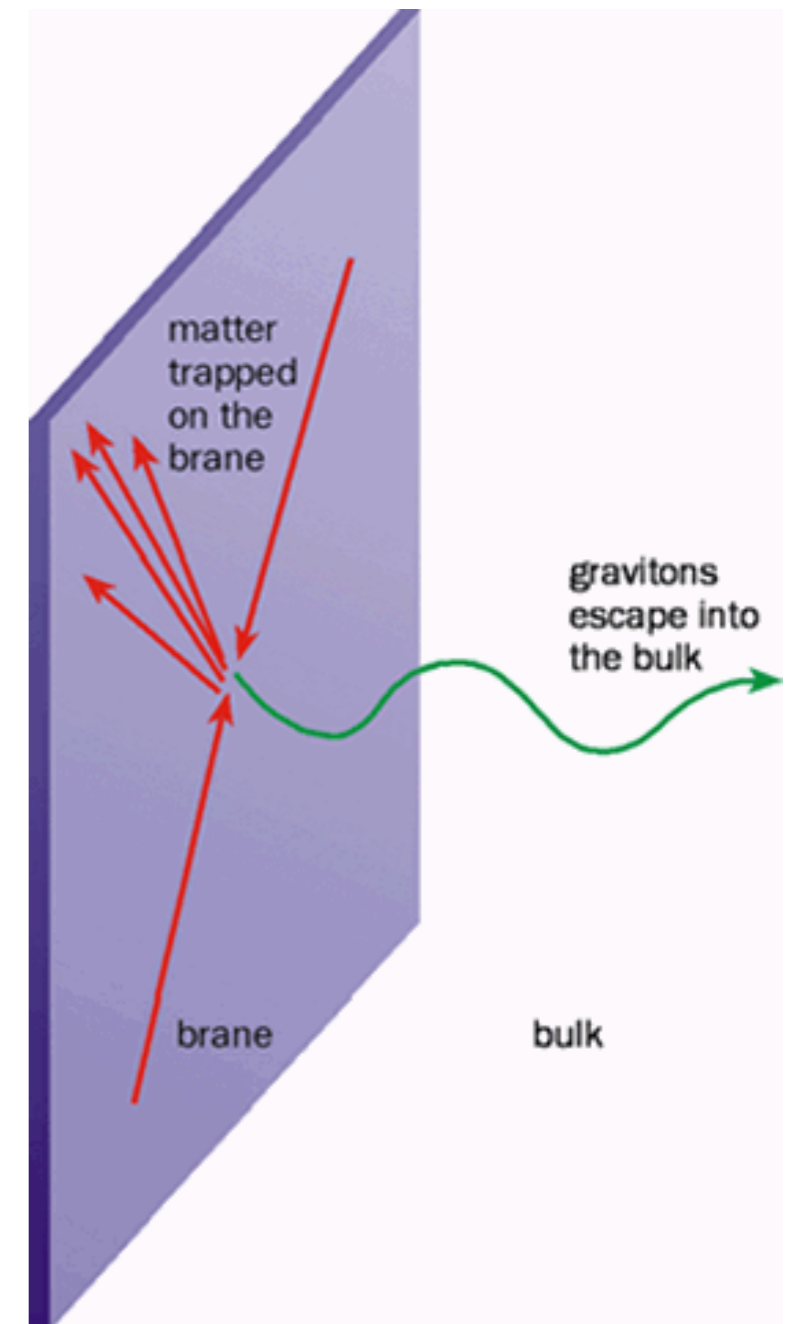
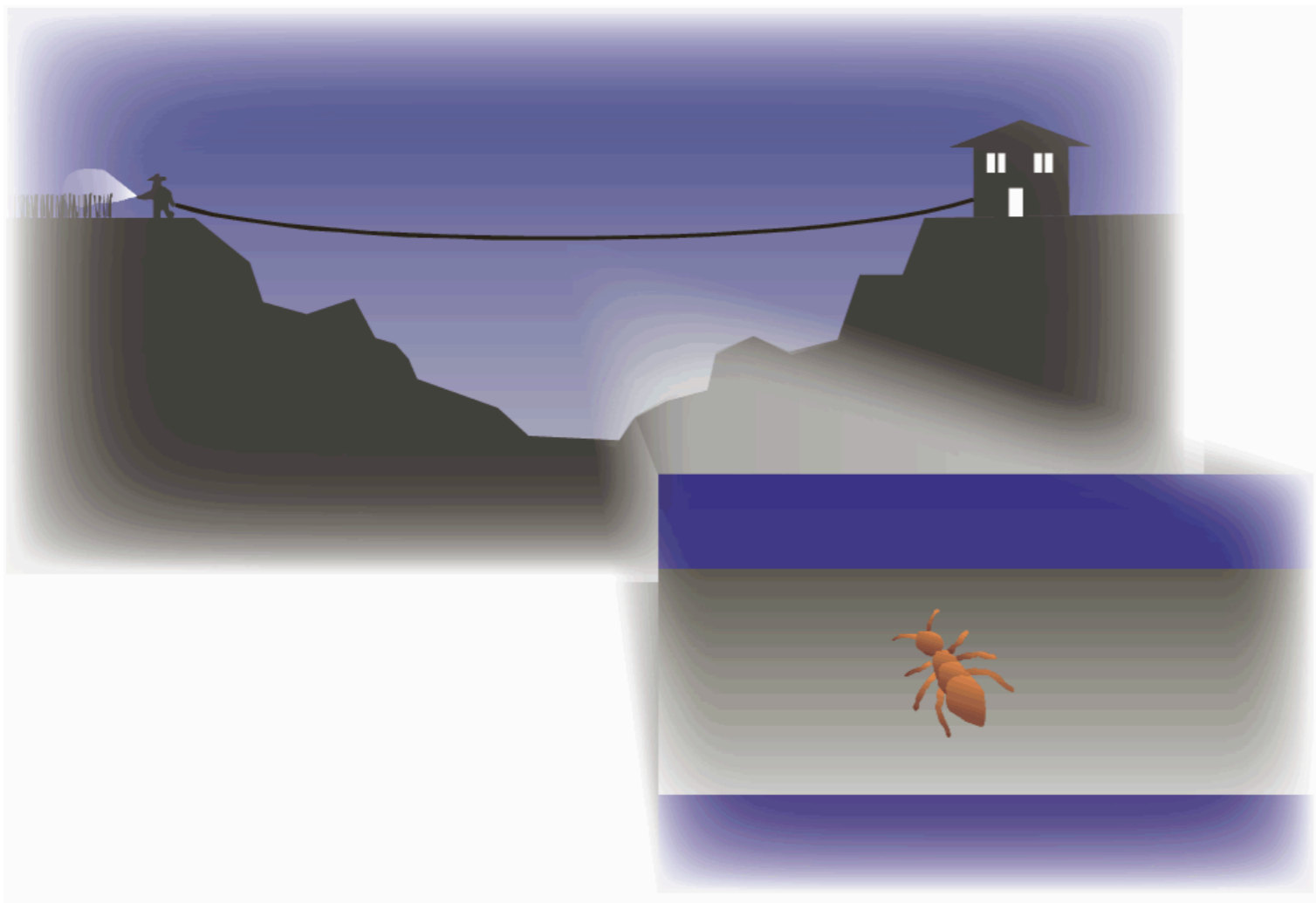
$$m_p = \sqrt{\frac{\hbar c}{G}} \sim 1 \times 10^{19} \text{ GeV} \quad \text{assuming 3 space dimensions}$$



If there would be more dimensions:
Steeper growth of gravitational
strength for smaller separations,
lower m_P

Extra Space Dimensions

- The idea: Extra dimensions are “compactified” - only relevant at small scales, and only visible to gravity



effective lowering of Planck mass by $R^{-n/2}$:

n : number of extra dimensions

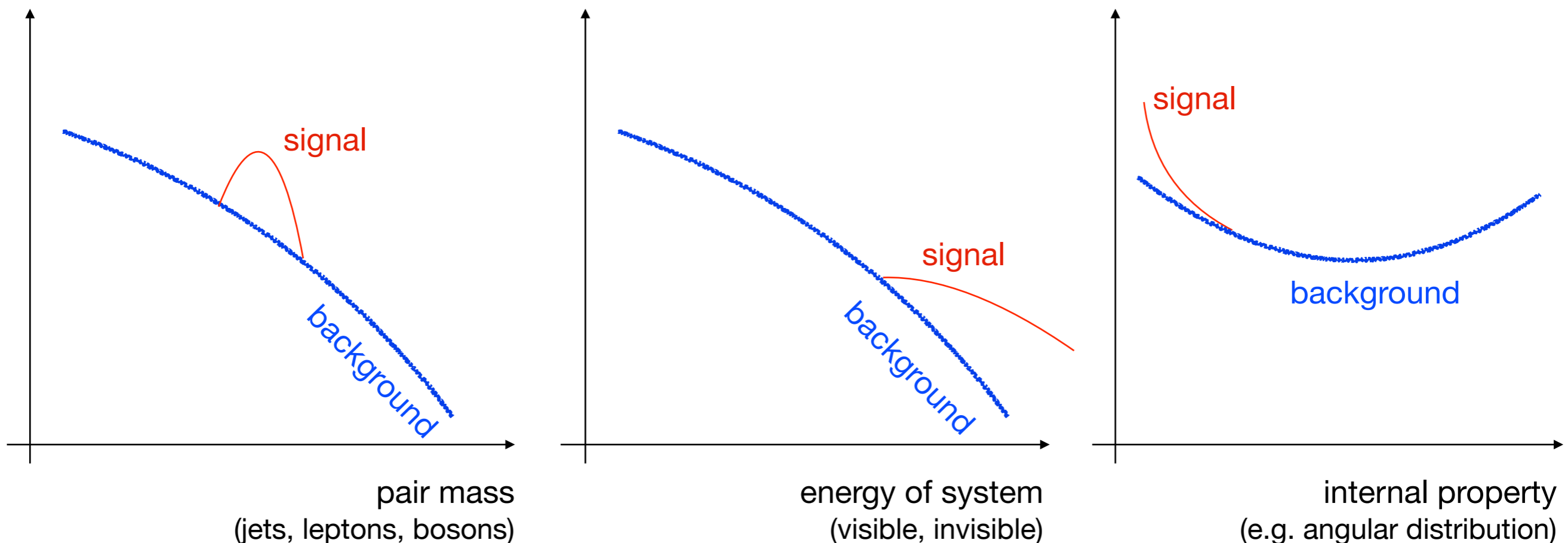
R : radius

For large n and large R smaller m_P , can reach TeV scale for nm - scale extra dimensions with $n > 3$

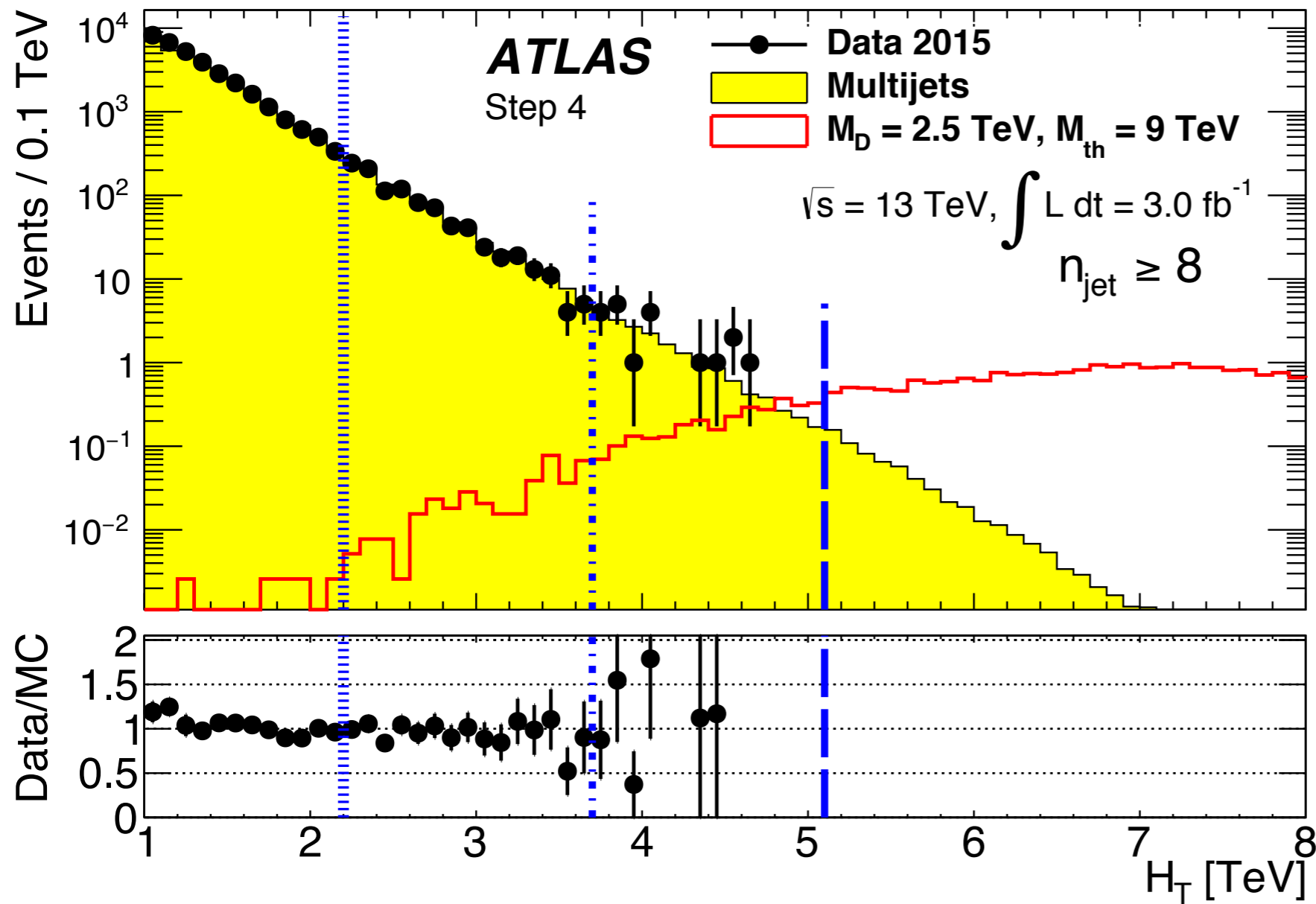
Extra Dimensions: Experimental Signatures

- Excitations in extra dimensions: High-mass resonances
- Creation and decay of micro - Black Holes: High-energy many-particle final states

Generic signatures for a number of “exotic” New Physics scenarios



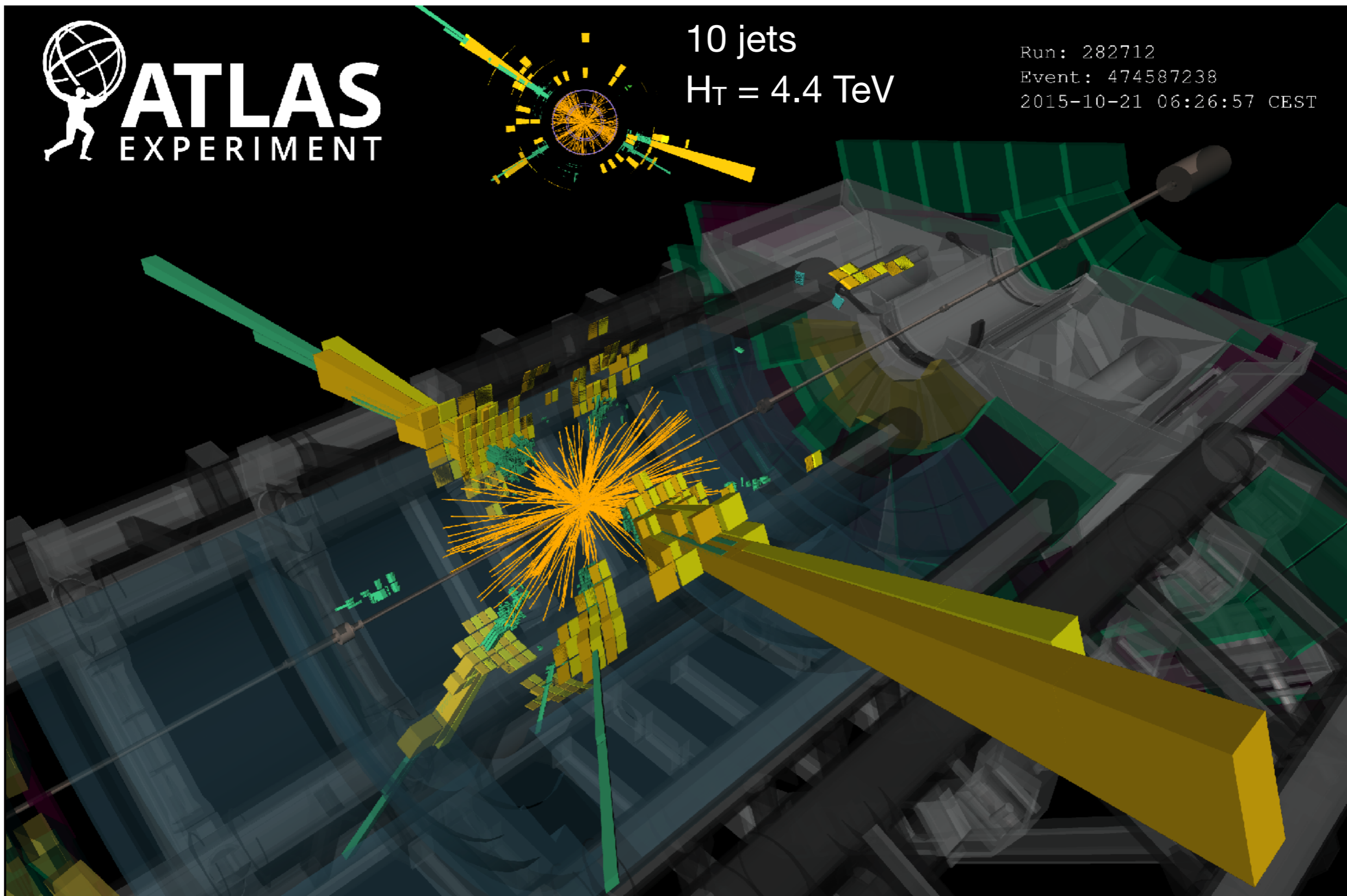
Black Hole Search: One Example



- Looking for an excess of signals in high jet multiplicities (here 8 or more jets)
- High total transverse Energy (H_T)

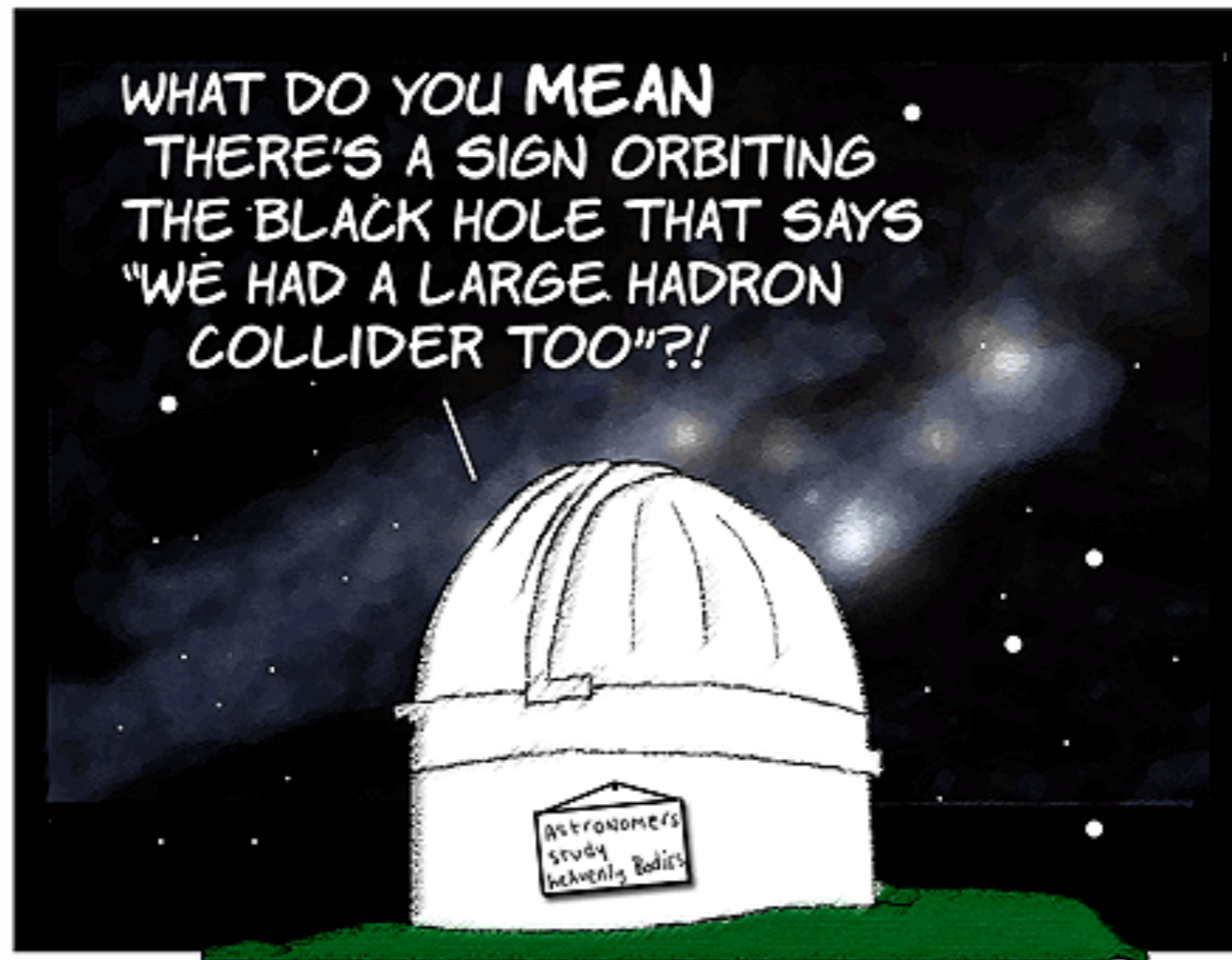
JHEP 03, 026 (2016)

Black Hole Search: One Example



Ideas that Capture the Imagination...

THE SMALLEST BLACK HOLE YET
DISCOVERED BY HUMANS
LOCATED AT BINARY XTE J1650-500.

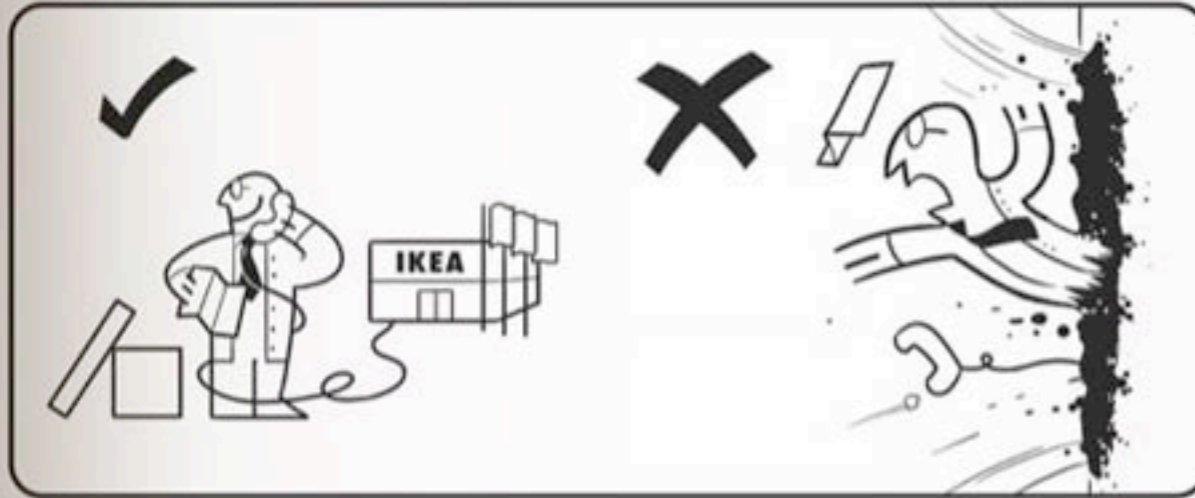
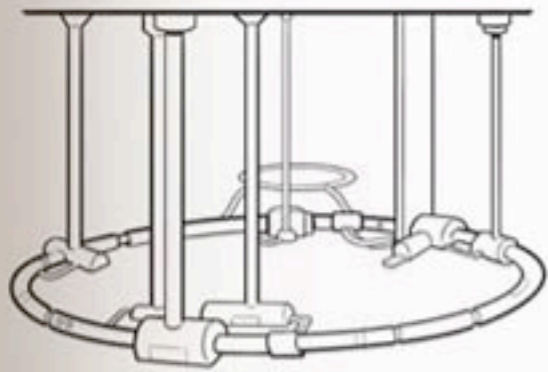


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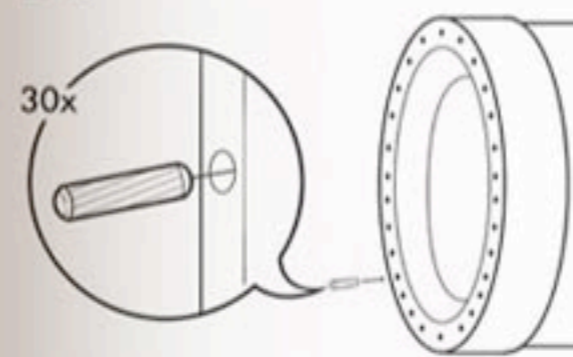
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Ideas that Capture the Imagination...

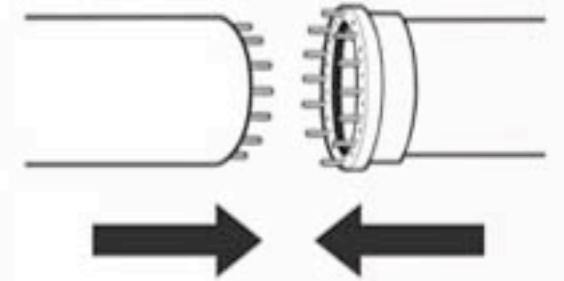
HÄDRÖNN CJÖLIDDER



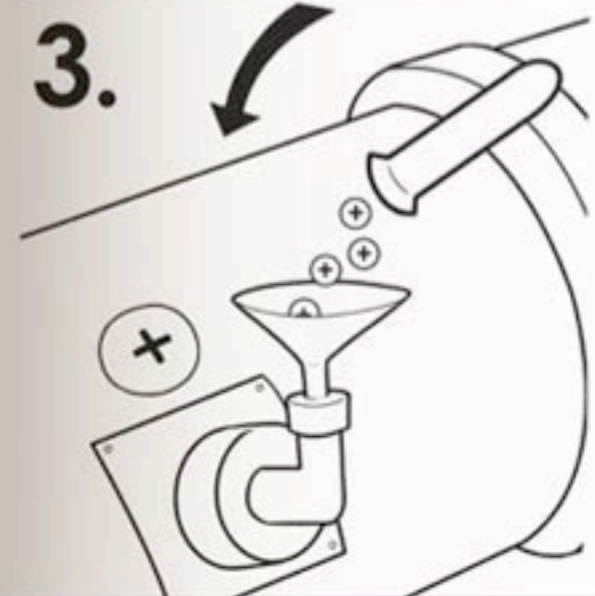
1.



2.



3.



4.

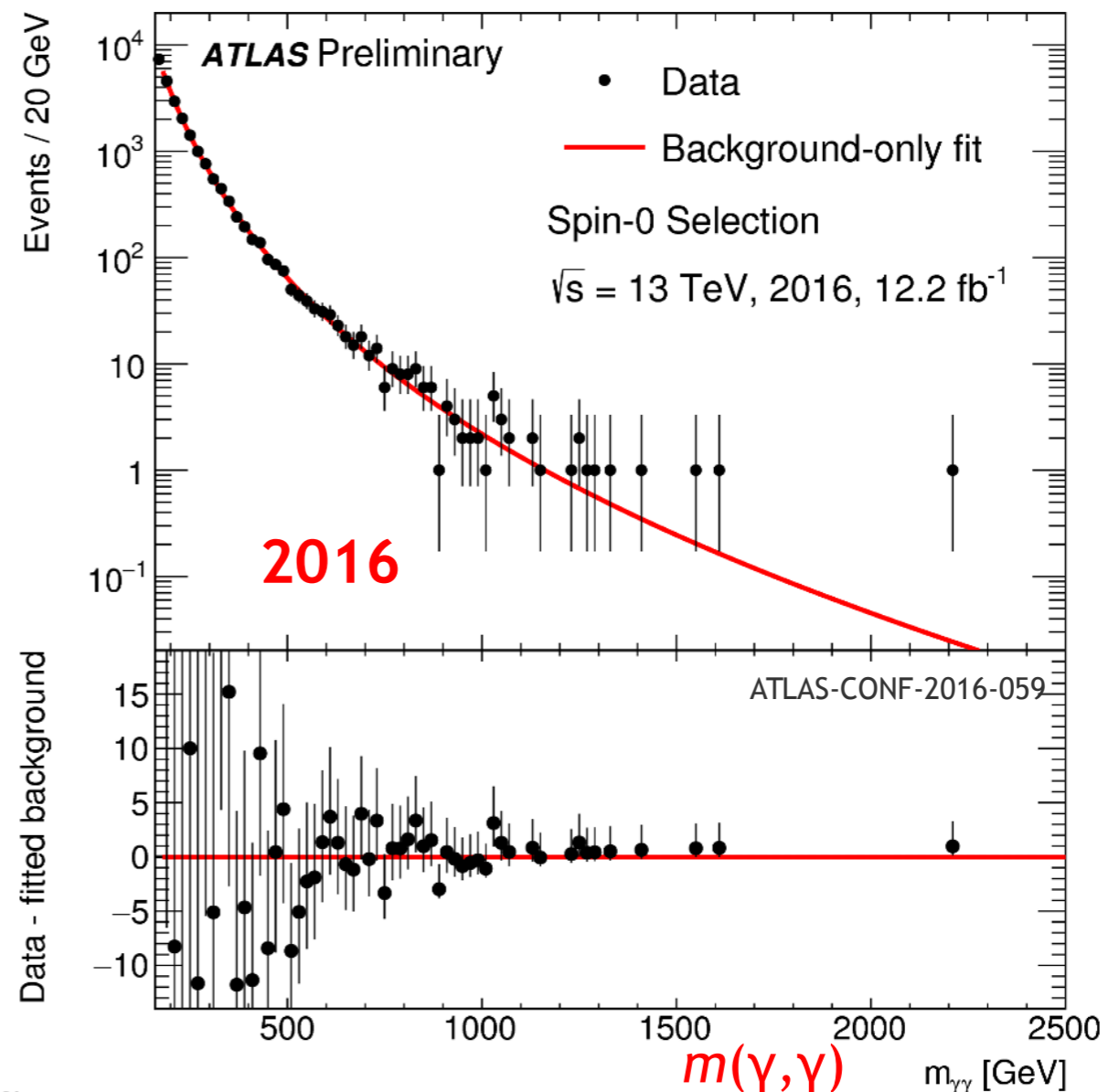
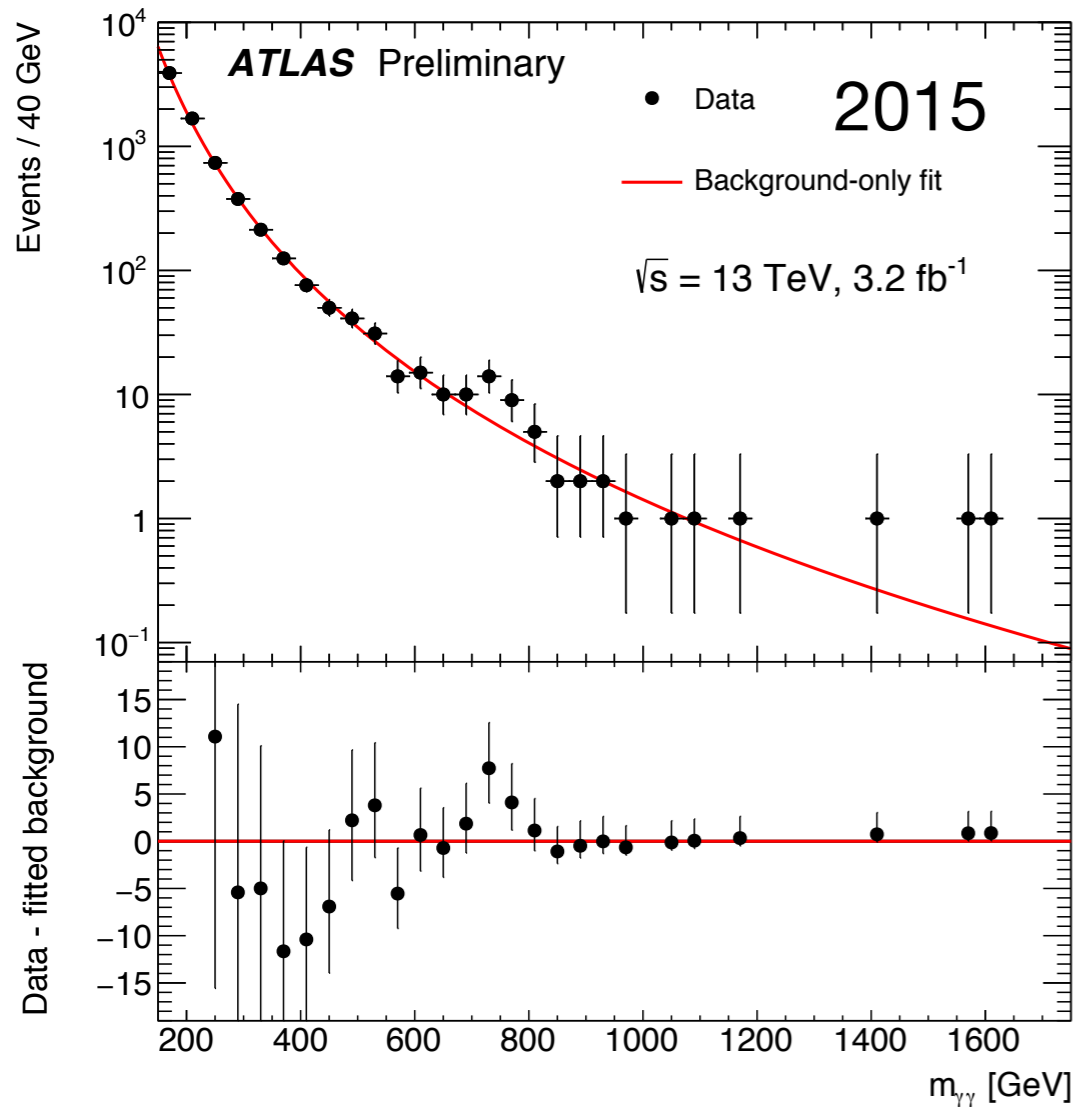


Ideas that Capture the Imagination...



Periods of Excitement

- Observed excess around 750 GeV in $\gamma\gamma$ invariant mass - a narrow resonance? also seen in similar (but not identical) mass range by CMS



... turned out to be a fluctuation.

A Broader View: Status of SUSY Searches at LHC

ATLAS SUSY Searches* - 95% CL Lower Limits

October 2019

ATLAS Preliminary

$\sqrt{s} = 13$ TeV

| Model | Signature | $\int \mathcal{L} dt$ [fb ⁻¹] | Mass limit | Reference | | | |
|---|--|--|----------------------------------|--|--|--|---|
| Inclusive Searches | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 e, μ mono-jet | 2-6 jets 1-3 jets | E_T^{miss} 139 E_T^{miss} 36.1 | \tilde{q} [10x Degen.] 1.9 \tilde{q} [1x, 8x Degen.] 0.43 0.71 | $m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 5$ GeV | ATLAS-CONF-2019-040 1711.03301 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$ | 0 e, μ | 2-6 jets | E_T^{miss} 139 | \tilde{g} 2.35 Forbidden 1.15-1.95 | $m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV | ATLAS-CONF-2019-040 ATLAS-CONF-2019-040 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$ | 3 e, μ $ee, \mu\mu$ | 4 jets 2 jets | E_T^{miss} 36.1 E_T^{miss} 36.1 | \tilde{g} 1.85 \tilde{g} 1.2 | $m(\tilde{\chi}_1^0) < 800$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 50$ GeV | 1706.03731 1805.11381 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ | 0 e, μ SS e, μ | 7-11 jets 6 jets | E_T^{miss} 36.1 E_T^{miss} 139 | \tilde{g} 1.8 \tilde{g} 1.15 | $m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV | 1708.02794 1909.08457 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0$ | 0-1 e, μ SS e, μ | 3 b 6 jets | E_T^{miss} 79.8 E_T^{miss} 139 | \tilde{g} 2.25 \tilde{g} 1.25 | $m(\tilde{\chi}_1^0) < 200$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV | ATLAS-CONF-2018-041 ATLAS-CONF-2019-015 |
| | 3 rd gen. squarks direct production | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0/\tilde{\chi}_1^\pm$ | Multiple Multiple Multiple | Multiple Multiple Multiple | E_T^{miss} 36.1 E_T^{miss} 36.1 E_T^{miss} 139 | Forbidden 0.9 Forbidden 0.58-0.82 Forbidden 0.74 | $m(\tilde{\chi}_1^0) = 300$ GeV, BR($b\tilde{\chi}_1^0$) = 1 $m(\tilde{\chi}_1^0) = 300$ GeV, BR($b\tilde{\chi}_1^\pm$) = BR($\tilde{\chi}_1^\pm$) = 0.5 $m(\tilde{\chi}_1^0) = 200$ GeV, $m(\tilde{\chi}_1^\pm) = 300$ GeV, BR($\tilde{\chi}_1^\pm$) = 1 |
| $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$ | | 0 e, μ | 6 b | E_T^{miss} 139 | Forbidden 0.23-0.48 \tilde{b}_1 0.23-1.35 | $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV | 1908.03122 1908.03122 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$ | | 0-2 e, μ | 0-2 jets/1-2 b | E_T^{miss} 36.1 | \tilde{t}_1 1.0 | $m(\tilde{\chi}_1^0) = 1$ GeV | 1506.08616, 1709.04183, 1711.11520 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ | | 1 e, μ | 3 jets/1 b | E_T^{miss} 139 | \tilde{t}_1 0.44-0.59 | $m(\tilde{\chi}_1^0) = 400$ GeV | ATLAS-CONF-2019-017 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$ | | 1 τ + 1 e, μ, τ | 2 jets/1 b | E_T^{miss} 36.1 | \tilde{t}_1 1.16 | $m(\tilde{\tau}_1) = 800$ GeV | 1803.10178 |
| $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$ | | 0 e, μ | 2 c | E_T^{miss} 36.1 | \tilde{t}_1 0.85 \tilde{t}_1 0.46 \tilde{t}_1 0.43 | $m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 50$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV | 1805.01649 1805.01649 1711.03301 |
| $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + h$ | | 1-2 e, μ | 4 b | E_T^{miss} 36.1 | \tilde{t}_2 0.32-0.88 | $m(\tilde{\chi}_1^0) = 0$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 180$ GeV | 1706.03986 |
| $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | | 3 e, μ | 1 b | E_T^{miss} 139 | Forbidden 0.86 | $m(\tilde{\chi}_1^0) = 360$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 40$ GeV | ATLAS-CONF-2019-016 |
| EW direct | $\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via WZ | 2-3 e, μ $ee, \mu\mu$ | ≥ 1 | E_T^{miss} 36.1 E_T^{miss} 139 | $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.6 $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ 0.205 | $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\chi}_1^\pm) - m(\tilde{\chi}_1^0) = 5$ GeV | 1403.5294, 1806.02293 ATLAS-CONF-2019-014 |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via WW | 2 e, μ | | E_T^{miss} 139 | $\tilde{\chi}_1^\pm$ 0.42 | $m(\tilde{\chi}_1^0) = 0$ | 1908.08215 |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_2^0$ via Wh | 0-1 e, μ | 2 $b/2 \gamma$ | E_T^{miss} 139 | Forbidden 0.74 | $m(\tilde{\chi}_1^0) = 70$ GeV | ATLAS-CONF-2019-019, 1909.09226 |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ via $\tilde{\ell}_L/\tilde{\nu}$ | 2 e, μ | | E_T^{miss} 139 | $\tilde{\chi}_1^\pm$ 1.0 | $m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^\pm) + m(\tilde{\chi}_1^0))$ | ATLAS-CONF-2019-008 |
| | $\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$ | 2 τ | | E_T^{miss} 139 | $\tilde{\tau}$ [L, R] 0.16-0.3 0.12-0.39 | $m(\tilde{\chi}_1^0) = 0$ | ATLAS-CONF-2019-018 |
| | $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$ | 2 e, μ 2 e, μ | 0 jets ≥ 1 | E_T^{miss} 139 E_T^{miss} 139 | $\tilde{\ell}$ 0.7 $\tilde{\ell}$ 0.256 | $m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 10$ GeV | ATLAS-CONF-2019-008 ATLAS-CONF-2019-014 |
| | $\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$ | 0 e, μ 4 e, μ | ≥ 3 b 0 jets | E_T^{miss} 36.1 E_T^{miss} 36.1 | \tilde{H} 0.13-0.23 0.29-0.88 \tilde{H} 0.3 | BR($\tilde{\chi}_1^0 \rightarrow h\tilde{G}$) = 1 BR($\tilde{\chi}_1^0 \rightarrow Z\tilde{G}$) = 1 | 1806.04030 1804.03602 |
| Long-lived particles | Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet | E_T^{miss} 36.1 | $\tilde{\chi}_1^\pm$ 0.46 $\tilde{\chi}_1^\pm$ 0.15 | Pure Wino Pure Higgsino | 1712.02118 ATL-PHYS-PUB-2017-019 |
| | Stable \tilde{g} R-hadron | Multiple | Multiple | E_T^{miss} 36.1 | \tilde{g} 2.0 | | 1902.01636, 1808.04095 |
| | Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$ | Multiple | Multiple | E_T^{miss} 36.1 | \tilde{g} [$\tau(\tilde{g}) = 10$ ns, 0.2 ns] 2.05 2.4 | $m(\tilde{\chi}_1^0) = 100$ GeV | 1710.04901, 1808.04095 |
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e\mu/\tau\mu$ | $e\mu, e\tau, \mu\tau$ | | 3.2 | $\tilde{\nu}_\tau$ 1.9 | $\lambda'_{511} = 0.11, \lambda'_{132/133/233} = 0.07$ | 1607.08079 |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\nu\nu$ | 4 e, μ | 0 jets | E_T^{miss} 36.1 | $\tilde{\chi}_1^\pm/\tilde{\chi}_2^0$ [$\lambda'_{333} \neq 0, \lambda'_{12k} \neq 0$] 0.82 1.33 | $m(\tilde{\chi}_1^0) = 100$ GeV | 1804.03602 |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qq\tilde{q}$ | 4-5 large- R jets Multiple | | 36.1 36.1 | \tilde{g} [$m(\tilde{\chi}_1^0) = 200$ GeV, 1100 GeV] 1.3 1.9 \tilde{g} [$\lambda'_{112} = 2e-4, 2e-5$] 1.05 2.0 | Large λ'_{112} $m(\tilde{\chi}_1^0) = 200$ GeV, bino-like | 1804.03568 ATLAS-CONF-2018-003 |
| | $\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$ | Multiple | | 36.1 | \tilde{g} [$\lambda'_{323} = 2e-4, 1e-2$] 0.55 1.05 | $m(\tilde{\chi}_1^0) = 200$ GeV, bino-like | ATLAS-CONF-2018-003 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$ | 2 jets + 2 b | | 36.7 | \tilde{t}_1 [qq, bs] 0.42 0.61 | | 1710.07171 |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$ | 2 e, μ 1 μ | 2 b DV | 36.1 136 | \tilde{t}_1 1.0 0.4-1.45 \tilde{t}_1 [$1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9$] 1.0 1.6 | BR($\tilde{t}_1 \rightarrow b\ell$) > 20% BR($\tilde{t}_1 \rightarrow q\mu$) = 100%, $\cos\theta = 1$ | 1710.05544 ATLAS-CONF-2019-006 |

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]



A Broader View: Status of BSM Searches at LHC

ATLAS Exotics Searches* - 95% CL Upper Exclusion Limits

Status: May 2019

ATLAS Preliminary

$$\int \mathcal{L} dt = (3.2 - 139) \text{ fb}^{-1}$$

$$\sqrt{s} = 8, 13 \text{ TeV}$$

| Model | ℓ, γ | Jets [†] | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Limit | Reference | | |
|------------------|--|--|------------------------|--|-------------------------|--|---|---------------------|
| Extra dimensions | ADD $G_{KK} + g/q$ | 0 e, μ | 1-4 j | Yes | 36.1 | M_D 7.7 TeV | $n = 2$ | 1711.03301 |
| | ADD non-resonant $\gamma\gamma$ | 2 γ | - | - | 36.7 | M_S 8.6 TeV | $n = 3$ HLZ NLO | 1707.04147 |
| | ADD QBH | - | 2 j | - | 37.0 | M_{th} 8.9 TeV | $n = 6$ | 1703.09127 |
| | ADD BH high $\sum p_T$ | $\geq 1 e, \mu$ | $\geq 2 j$ | - | 3.2 | M_{th} 8.2 TeV | $n = 6, M_D = 3 \text{ TeV}$, rot BH | 1606.02265 |
| | ADD BH multijet | - | $\geq 3 j$ | - | 3.6 | M_{th} 9.55 TeV | $n = 6, M_D = 3 \text{ TeV}$, rot BH | 1512.02586 |
| | RS1 $G_{KK} \rightarrow \gamma\gamma$ | 2 γ | - | - | 36.7 | G_{KK} mass 4.1 TeV | $k/\overline{M}_{Pl} = 0.1$ | 1707.04147 |
| | Bulk RS $G_{KK} \rightarrow WW/ZZ$ | multi-channel | - | - | 36.1 | G_{KK} mass 2.3 TeV | $k/\overline{M}_{Pl} = 1.0$ | 1808.02380 |
| | Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\bar{q}\bar{q}$ | 0 e, μ | 2 J | - | 139 | G_{KK} mass 1.6 TeV | $k/\overline{M}_{Pl} = 1.0$ | ATLAS-CONF-2019-003 |
| | Bulk RS $g_{KK} \rightarrow t\bar{t}$ | 1 e, μ | $\geq 1 b, \geq 1J/2j$ | Yes | 36.1 | g_{KK} mass 3.8 TeV | $\Gamma/m = 15\%$ | 1804.10823 |
| | 2UED / RPP | 1 e, μ | $\geq 2 b, \geq 3 j$ | Yes | 36.1 | KK mass 1.8 TeV | Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow t\bar{t}) = 1$ | 1803.09678 |
| Gauge bosons | SSM $Z' \rightarrow \ell\ell$ | 2 e, μ | - | - | 139 | Z' mass 5.1 TeV | | 1903.06248 |
| | SSM $Z' \rightarrow \tau\tau$ | 2 τ | - | - | 36.1 | Z' mass 2.42 TeV | | 1709.07242 |
| | Leptophobic $Z' \rightarrow b\bar{b}$ | - | 2 b | - | 36.1 | Z' mass 2.1 TeV | | 1805.09299 |
| | Leptophobic $Z' \rightarrow t\bar{t}$ | 1 e, μ | $\geq 1 b, \geq 1J/2j$ | Yes | 36.1 | Z' mass 3.0 TeV | $\Gamma/m = 1\%$ | 1804.10823 |
| | SSM $W' \rightarrow \ell\nu$ | 1 e, μ | - | Yes | 139 | W' mass 6.0 TeV | | CERN-EP-2019-100 |
| | SSM $W' \rightarrow \tau\nu$ | 1 τ | - | Yes | 36.1 | W' mass 3.7 TeV | | 1801.06992 |
| | HVT $V' \rightarrow WZ \rightarrow qq\bar{q}\bar{q}$ model B | 0 e, μ | 2 J | - | 139 | V' mass 3.6 TeV | $g_V = 3$ | ATLAS-CONF-2019-003 |
| | HVT $V' \rightarrow WH/ZH$ model B | multi-channel | - | - | 36.1 | V' mass 2.93 TeV | $g_V = 3$ | 1712.06518 |
| | LRSM $W_R \rightarrow t\bar{b}$ | multi-channel | - | - | 36.1 | W_R mass 3.25 TeV | | 1807.10473 |
| | LRSM $W_R \rightarrow \mu N_R$ | 2 μ | 1 J | - | 80 | W_R mass 5.0 TeV | $m(N_R) = 0.5 \text{ TeV}, g_L = g_R$ | 1904.12679 |
| CI | CI $qq\bar{q}\bar{q}$ | - | 2 j | - | 37.0 | Λ 21.8 TeV η_{LL} | | 1703.09127 |
| | CI $\ell\ell q\bar{q}$ | 2 e, μ | - | - | 36.1 | Λ 40.0 TeV η_{LL} | | 1707.02424 |
| | CI $t\bar{t}t\bar{t}$ | $\geq 1 e, \mu$ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Λ 2.57 TeV | $ C_{4t} = 4\pi$ | 1811.02305 |
| DM | Axial-vector mediator (Dirac DM) | 0 e, μ | 1-4 j | Yes | 36.1 | m_{med} 1.55 TeV | $g_q = 0.25, g_\chi = 1.0, m(\chi) = 1 \text{ GeV}$ | 1711.03301 |
| | Colored scalar mediator (Dirac DM) | 0 e, μ | 1-4 j | Yes | 36.1 | m_{med} 1.67 TeV | $g = 1.0, m(\chi) = 1 \text{ GeV}$ | 1711.03301 |
| | $VV\chi\chi$ EFT (Dirac DM) | 0 e, μ | 1 J, $\leq 1 j$ | Yes | 3.2 | M_* 700 GeV | $m(\chi) < 150 \text{ GeV}$ | 1608.02372 |
| | Scalar reson. $\phi \rightarrow t\bar{t}$ (Dirac DM) | 0-1 e, μ | 1 b, 0-1 J | Yes | 36.1 | m_ϕ 3.4 TeV | $y = 0.4, \lambda = 0.2, m(\chi) = 10 \text{ GeV}$ | 1812.09743 |
| LQ | Scalar LQ 1 st gen | 1,2 e | $\geq 2 j$ | Yes | 36.1 | LQ mass 1.4 TeV | $\beta = 1$ | 1902.00377 |
| | Scalar LQ 2 nd gen | 1,2 μ | $\geq 2 j$ | Yes | 36.1 | LQ mass 1.56 TeV | $\beta = 1$ | 1902.00377 |
| | Scalar LQ 3 rd gen | 2 τ | 2 b | - | 36.1 | LQ_3^u mass 1.03 TeV | $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$ | 1902.08103 |
| | Scalar LQ 3 rd gen | 0-1 e, μ | 2 b | Yes | 36.1 | LQ_3^d mass 970 GeV | $\mathcal{B}(LQ_3^d \rightarrow t\tau) = 0$ | 1902.08103 |
| Heavy quarks | VLQ $TT \rightarrow Ht/Zt/Wb + X$ | multi-channel | - | - | 36.1 | T mass 1.37 TeV | SU(2) doublet | 1808.02343 |
| | VLQ $BB \rightarrow Wt/Zb + X$ | multi-channel | - | - | 36.1 | B mass 1.34 TeV | SU(2) doublet | 1808.02343 |
| | VLQ $T_{5/3} T_{5/3} \rightarrow Wt + X$ | $2(SS) \geq 3 e, \mu \geq 1 b, \geq 1 j$ | Yes | 36.1 | $T_{5/3}$ mass 1.64 TeV | $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$ | 1807.11883 | |
| | VLQ $Y \rightarrow Wb + X$ | 1 e, μ | $\geq 1 b, \geq 1 j$ | Yes | 36.1 | Y mass 1.85 TeV | $\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$ | 1812.07343 |
| | VLQ $B \rightarrow Hb + X$ | 0 $e, \mu, 2 \gamma$ | $\geq 1 b, \geq 1 j$ | Yes | 79.8 | B mass 1.21 TeV | $\kappa_B = 0.5$ | ATLAS-CONF-2018-024 |
| | VLQ $QQ \rightarrow WqWq$ | 1 e, μ | $\geq 4 j$ | Yes | 20.3 | Q mass 690 GeV | | 1509.04261 |
| Excited fermions | Excited quark $q^* \rightarrow qg$ | - | 2 j | - | 139 | q^* mass 6.7 TeV | only u^* and d^* , $\Lambda = m(q^*)$ | ATLAS-CONF-2019-007 |
| | Excited quark $q^* \rightarrow q\gamma$ | 1 γ | 1 j | - | 36.7 | q^* mass 5.3 TeV | only u^* and d^* , $\Lambda = m(q^*)$ | 1709.10440 |
| | Excited quark $b^* \rightarrow bg$ | - | 1 b, 1 j | - | 36.1 | b^* mass 2.6 TeV | | 1805.09299 |
| | Excited lepton ℓ^* | 3 e, μ | - | - | 20.3 | ℓ^* mass 3.0 TeV | $\Lambda = 3.0 \text{ TeV}$ | 1411.2921 |
| | Excited lepton ν^* | 3 e, μ, τ | - | - | 20.3 | ν^* mass 1.6 TeV | $\Lambda = 1.6 \text{ TeV}$ | 1411.2921 |
| Other | Type III Seesaw | 1 e, μ | $\geq 2 j$ | Yes | 79.8 | N^0 mass 560 GeV | | ATLAS-CONF-2018-020 |
| | LRSM Majorana ν | 2 μ | 2 j | - | 36.1 | N_R mass 3.2 TeV | $m(W_R) = 4.1 \text{ TeV}, g_L = g_R$ | 1809.11105 |
| | Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ | 2,3,4 e, μ (SS) | - | - | 36.1 | $H^{\pm\pm}$ mass 870 GeV | DY production | 1710.09748 |
| | Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$ | 3 e, μ, τ | - | - | 20.3 | $H^{\pm\pm}$ mass 400 GeV | DY production, $\mathcal{B}(H^{\pm\pm} \rightarrow \ell\tau) = 1$ | 1411.2921 |
| | Multi-charged particles | - | - | - | 36.1 | multi-charged particle mass 1.22 TeV | DY production, $ q = 5e$ | 1812.03673 |
| | Magnetic monopoles | - | - | - | 34.4 | monopole mass 2.37 TeV | DY production, $ g = 1g_D$, spin 1/2 | 1905.10130 |

$\sqrt{s} = 8 \text{ TeV}$

$\sqrt{s} = 13 \text{ TeV}$
partial data

$\sqrt{s} = 13 \text{ TeV}$
full data

10⁻¹

1

10

Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown.

†Small-radius (large-radius) jets are denoted by the letter j (J).



Finally...

Absence of evidence is not evidence of absence

meaning:

no sign of physics BSM from Run-I / Run-II data,
but unexplored phase space still large!

... and there are corners that cannot be fully explored with LHC,
even within its energy / mass reach.

Summary

- The Standard Model is incomplete - conceptual problems and failures to describe astrophysical observations. But so far experiments have not yet revealed concrete discoveries of New Physics
- A wide range of theoretical models:
 - Grand Unified Theories as an overarching theory at very high energies
 - Supersymmetry, Large Extra Dimensions, ... on the electroweak scale
- A rich array of experimental searches at the LHC, looking for:
 - New particles
 - Dark Matter candidates
 - New forces and unexpected phenomena
 - ...

Up to now: Nothing found!

Next (and final!) Lecture:

Physics beyond the Standard Model in the Early Universe - B. Majorovits, 03.02.2020

Lecture Overview

| | | |
|--------|---|----------------------|
| 14.10. | Introduction, Particle Physics Refresher | <i>F. Simon</i> |
| 21.10. | Introduction to Cosmology I | <i>B. Majorovits</i> |
| 28.10. | Introduction to Cosmology II | <i>B. Majorovits</i> |
| 04.11. | Particle Collisions at High Energy | <i>F. Simon</i> |
| 11.11. | The Higgs Boson | <i>F. Simon</i> |
| 18.11. | The Early Universe: Thermal Freeze-out of Particles | <i>B. Majorovits</i> |
| 25.11. | The Universe as a High Energy Laboratory: BBN | <i>B. Majorovits</i> |
| 02.12. | Particle Colliders | <i>F. Simon</i> |
| 09.12. | The Universe as a High Energy Laboratory: CMB | <i>B. Majorovits</i> |
| 16.12. | Cosmic Rays: Acceleration Mechanisms and Possible Sources | <i>B. Majorovits</i> |
| | Christmas Break | |
| 13.01. | Supernovae Accelerators for Charged Particles and Neutrinos | <i>B. Majorovits</i> |
| 20.01. | Detectors for Particle Colliders | <i>F. Simon</i> |
| 27.01. | Searching for New Physics at the Energy Frontier | <i>F. Simon</i> |
| 03.02. | Physics beyond the Standard Model in the Early Universe | <i>B. Majorovits</i> |