

# Erasures of Defects

Juan Sebastian Valbuena Bermudez  
LMU-TUM Munich, Germany

Supervisor: Prof. Dr. Georgi Dvali

MPP PhD Recruiting Workshop  
3<sup>rd</sup> December 2019

# Outline

- ❖ GUTs and the Cosmological Monopole Problem
  - ❖ Sweeping Away the Monopole Problem
- ❖ Erasure of non-topological Defects: DLV Mechanism
- ❖ The  $\phi^6$ -Model
  - ❖ Domain Walls, and Coulomb Vacuum Layers
  - ❖ Vortex Lines
  - ❖ Erasure of Vortices by Coulomb Vacuum Layers
- ❖ Conclusions and Outlook

# Grand Unified Theories

- ꝝ In GUT, the three SM gauge interactions are merged into one-gauge interaction,  $G$ .
- ꝝ  $G$  is spontaneously broken at  $T_{GUT} \sim 10^{15}$ GeV, to the SM group when a Higgs-like field,  $\phi$ , acquires a VEV.
- ꝝ If  $T > T_{GUT}$ , in the very early universe there was a GUT-epoch

# The Cosmological Monopole Problem

- ❖ As the universe expands and cools down, it undergoes a phase transitions at  $T_{GUT}$ .
- ❖ During this phase transition different *topological defects* may be produced. For example:
  - ❖ Domain Walls
  - ❖ Cosmic Strings
  - ❖ Magnetic Monopoles

# The Cosmological Monopole Problem

- ❖ Monopole configurations are a common feature of GUTs.
- ❖ The estimated concentration of monopoles is unacceptably large in comparison to observations<sup>1</sup>.
- ❖ Different solutions for this tension were proposed:
  - ❖ Inflation  
A. H. Guth, Phys. Rev. D 23, 347 (1981)
  - ❖ Langacker & Pi mechanism  
P. Langacker and S.Y. Pi, Phys. Rev. Lett. 45, 1 (1980)
  - ❖ Non-restoration of the grand unified symmetry  
G. Dvali, A. Melfo, and G. Senjanovic, Phys. Rev. Lett. 75, 4559 (1995)

<sup>1</sup>Y. Zeldovich and M. Khlopov, Phys. Lett. B79, 239–241 (1978)  
J. Preskill, Phys. Rev. Lett. 43, 1365 (1979)

# Sweeping Away the Monopole Problem

❖ Solution: DLV Mechanism:<sup>3</sup>

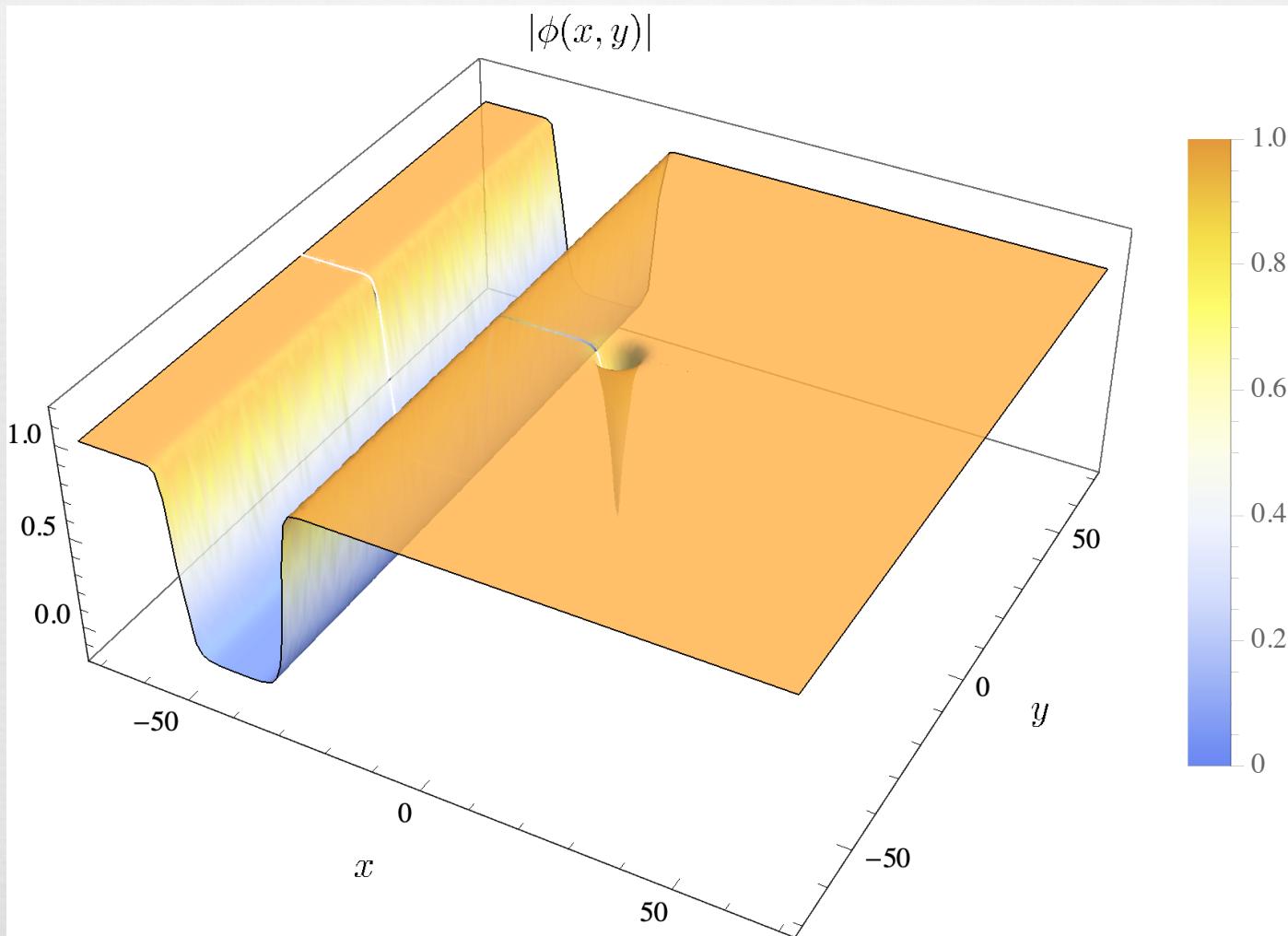
*Domain Walls + Magnetic Monopoles interaction  
→ Erasure of Monopoles*

❖ Basic idea:

- ❖ The same phase transition produces
  - ❖ **Magnetic Monopoles**
  - ❖ **Domain Walls**
- ❖ The domain walls move through space and sweep up the monopoles.
- ❖ When a monopole encounters a wall, it *unwinds* and *dissipates*.
- ❖ The walls are unstable at a lower energy scale, and hence collapse and go away.

<sup>3</sup>G. Dvali, H. Liu, and T. Vachaspati, Phys. Rev. Lett. 80, 2281–2284 (1998).

# Erasures of Defects: The DLV Mechanism



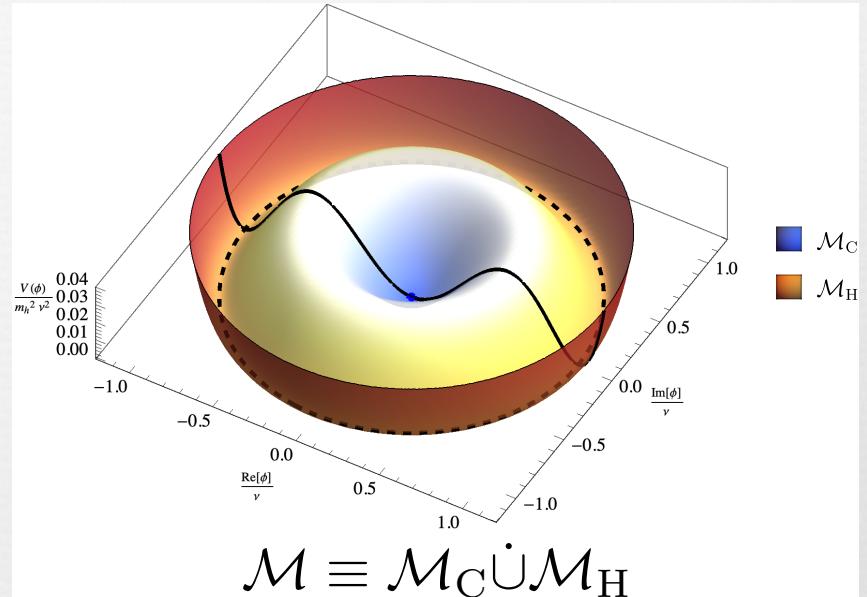
Erasure of a Vortex  
by  
a Coulomb Vacuum Layer

# The $\phi^6$ -Model

$$\mathcal{L}[\phi, A_\mu] = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_\mu\phi)^*D^\mu\phi - V(\phi),$$

$$G = U(1)$$

$$V(\phi) = \lambda^2 \phi \phi^* (\phi \phi^* - \nu^2)^2.$$



	<b>Coulomb Phase:</b> $\langle \phi \rangle = 0$	<b>Higgs Phase:</b> $\langle \phi \rangle = \nu$
Vacuum Manifold	$\mathcal{M}_C = G/H_0 = 1$	$\mathcal{M}_H = G/H_\nu = U(1)$
Spectrum	$m_\phi = \lambda \nu^2$	$m_h = 2 \lambda \nu^2$
	$m_A = 0$	$m_v = \sqrt{2} e \nu$

# Domain Walls in the $\phi^6$ -Model

# Domain Walls in the $\phi^6$ -Model

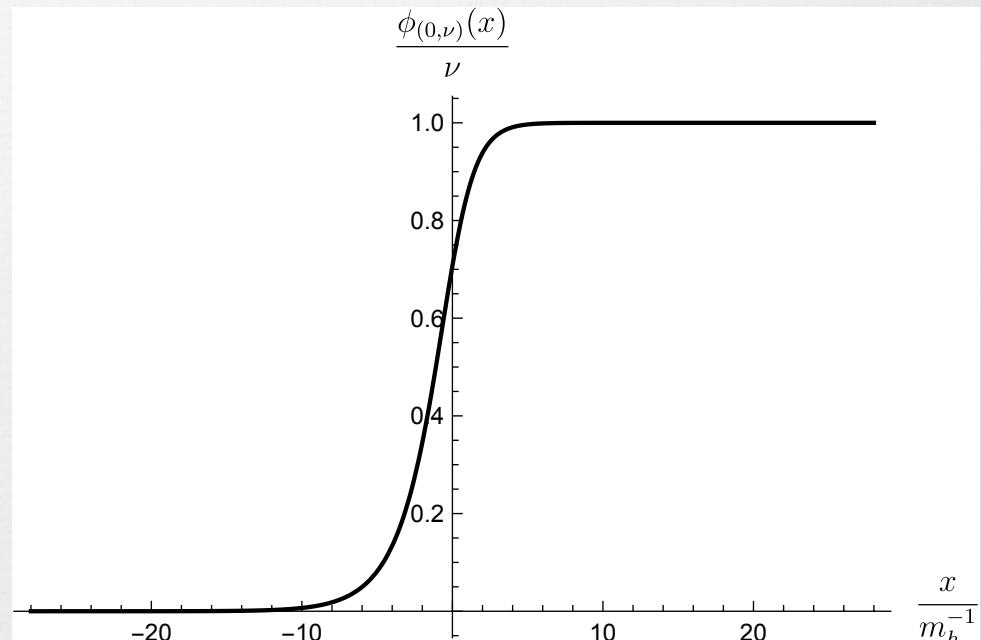
ꝝ  $\pi_0(\mathcal{M}) = \mathbb{Z}_2$

ꝝ (1+1)-dimensions

ꝝ Solitonic configurations  
interpolating between the  
Coulomb and the Higgs  
phases<sup>4</sup>

$$\phi_{(\pm\nu, 0)}(x) \equiv \pm\nu\sqrt{\frac{1}{1 + e^{m_h x}}}$$

$$\phi_{(0, \pm\nu)}(x) \equiv \pm\nu\sqrt{\frac{1}{1 + e^{-m_h x}}}$$



<sup>4</sup>V. A. Gani, A. E. Kudryavtsev, and M. A. Lizunov, Phys. Rev. D 89, 125009,(2014)

# Coulomb Vacuum Layer

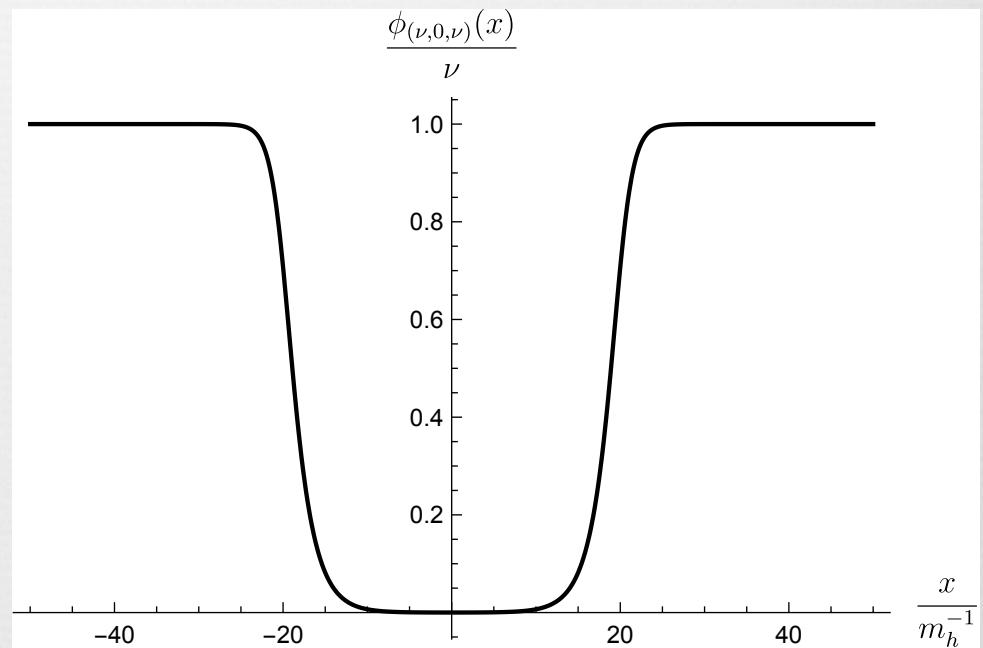
## $(\nu, 0, \nu)$ -Domain Wall

❖ Profile ansatz

$$\phi_{(\nu,0,\nu)}(x) = \phi_{(\nu,0)}\left(x + \frac{l}{2}\right) + \phi_{(0,\nu)}\left(x - \frac{l}{2}\right)$$

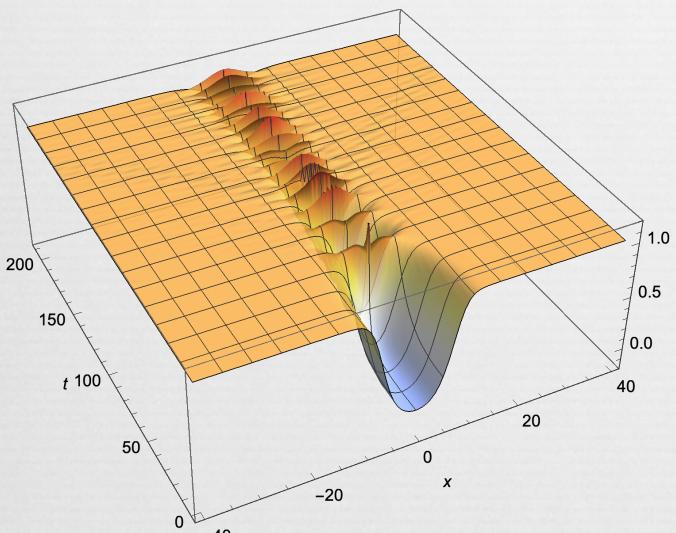
❖ Distance between the walls:

$$l > m_h^{-1}$$

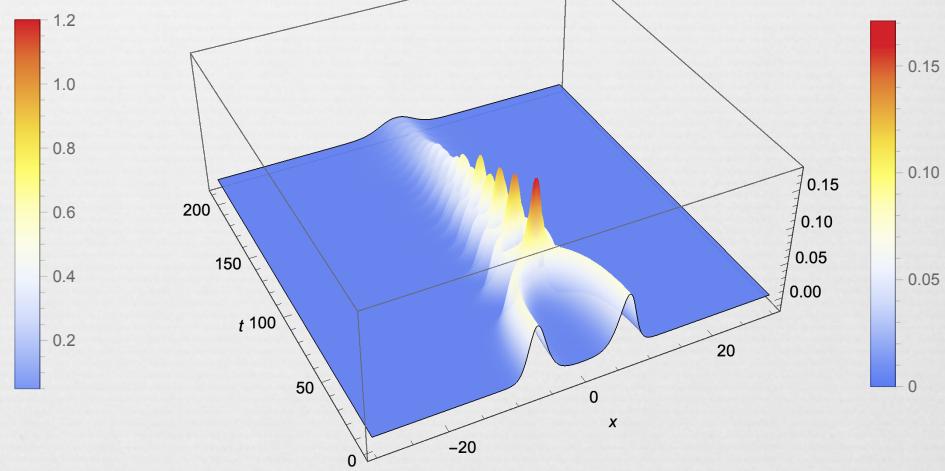


# Time Evolution: $(\nu, 0, \nu)$ -Domain Wall

$$l = 15m_h^{-1}$$



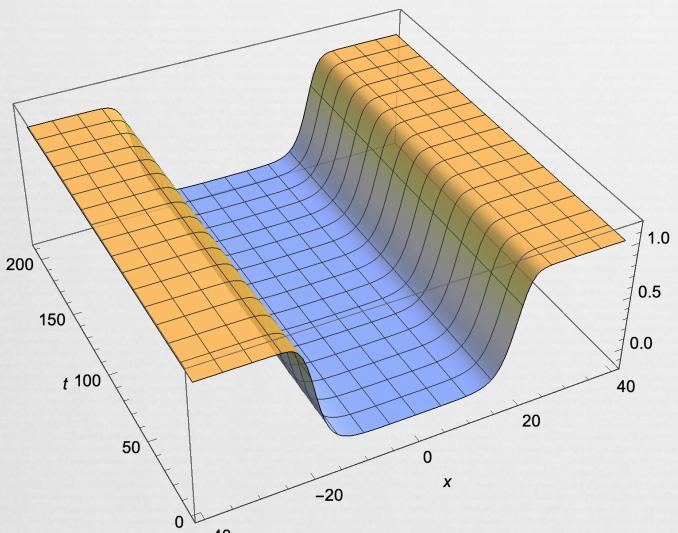
Field Profile  
 $\phi(x, t)$



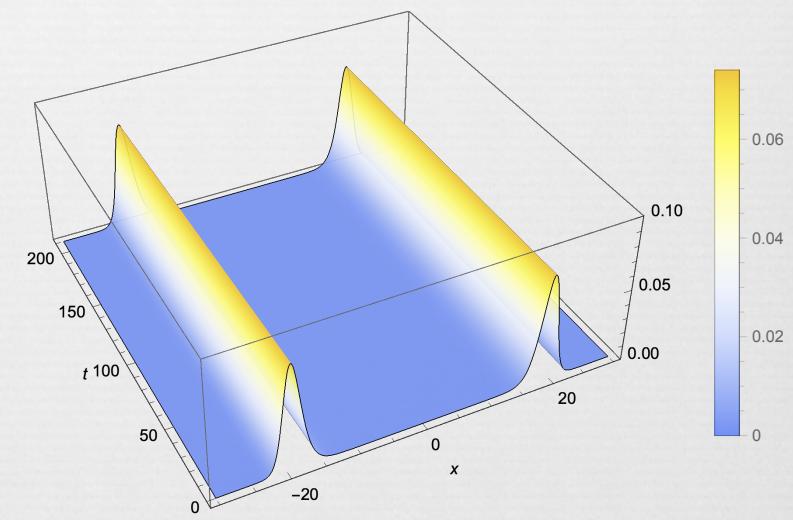
Energy density  
 $E(x, t)$

# Time Evolution: $(\nu, 0, \nu)$ -Domain Wall.

$$l = 40m_h^{-1}$$



Field Profile  
 $\phi(x, t)$



Energy density  
 $E(x, t)$

# Vortex in the $\phi^6$ -Model

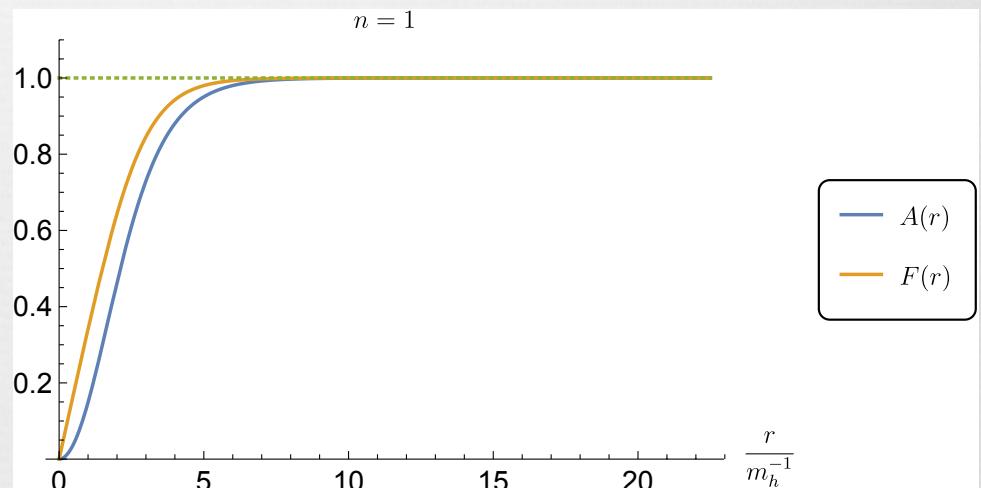
# Vortex Profile

- ꝝ  $\pi_1(\mathcal{M}_H) = \mathbb{Z}$
- ꝝ In (2+1) dimensions
- ꝝ Ansatz<sup>4</sup>

$$\phi(r, \theta) = \nu e^{in\theta} F(r)$$

$$A_i(r, \theta) = -\frac{n}{er} \epsilon_{ij} n_j A(r)$$

where  $n \in \mathbb{Z}$  is the  
*winding number*.



<sup>4</sup>H.B. Nielsen, P. Olesen, Nuc. Phys. B, Vol. 61, 45-61, (1973)

# Vortex: $\phi$

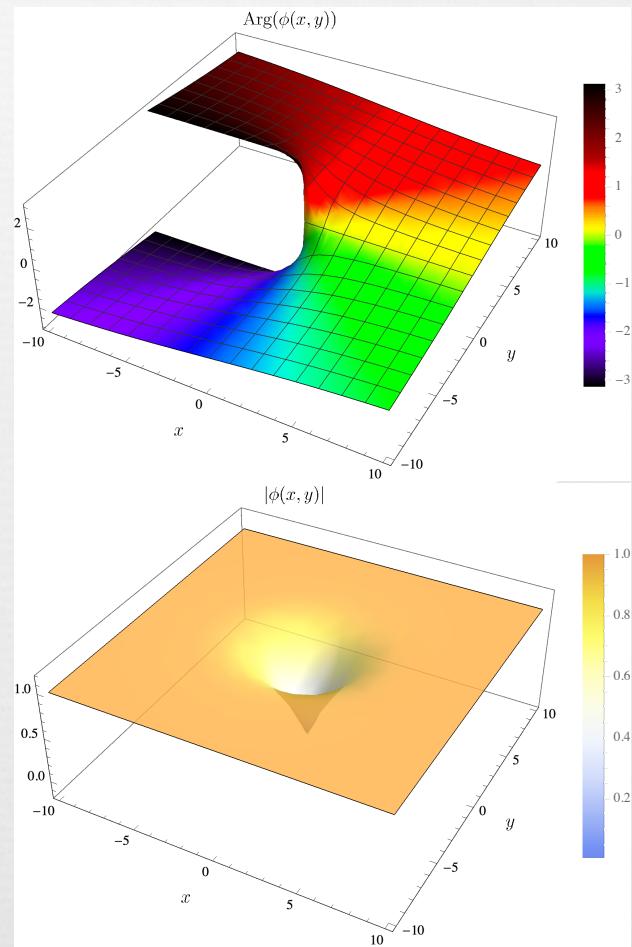
ꝝ  $m_h = 1$

ꝝ  $m_\nu = 1$

ꝝ  $n = 1$

ꝝ Scalar Field:

$$\phi(r, \theta) = \nu e^{in\theta} F(r)$$



# Vortex: $\phi$

$$\heartsuit m_h = 1$$

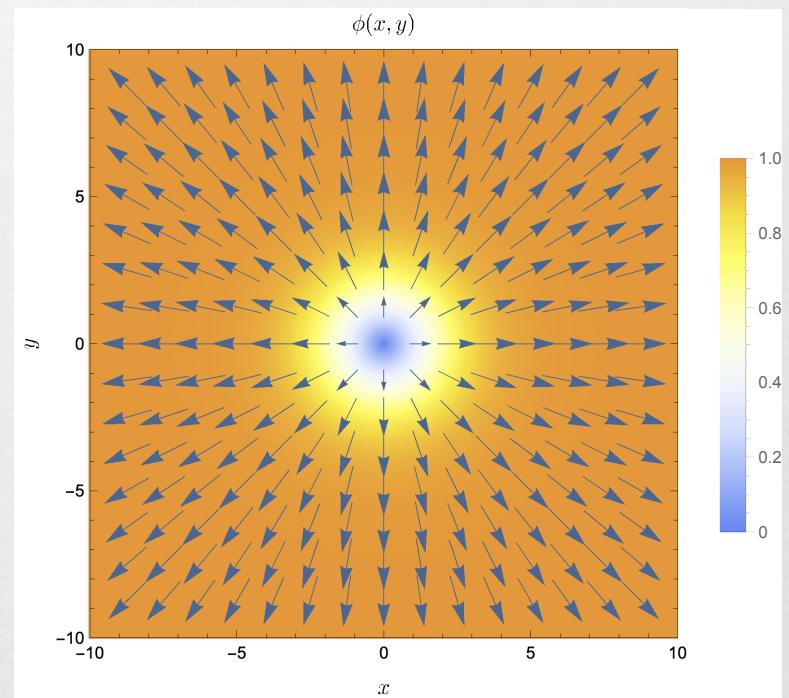
$$\heartsuit m_\nu = 1$$

$$\heartsuit n = 1$$

$\heartsuit$  Scalar Field:

$$\phi(r, \theta) = \nu e^{in\theta} F(r)$$

$$\heartsuit n = \frac{1}{2\pi i\nu^2} \oint dx^i \phi^* \partial_i \phi$$



# Vortex: $A_i$

ꝝ  $m_h = 1$

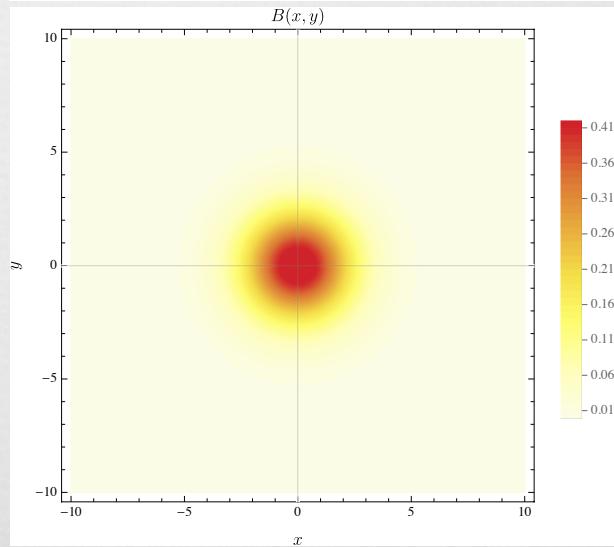
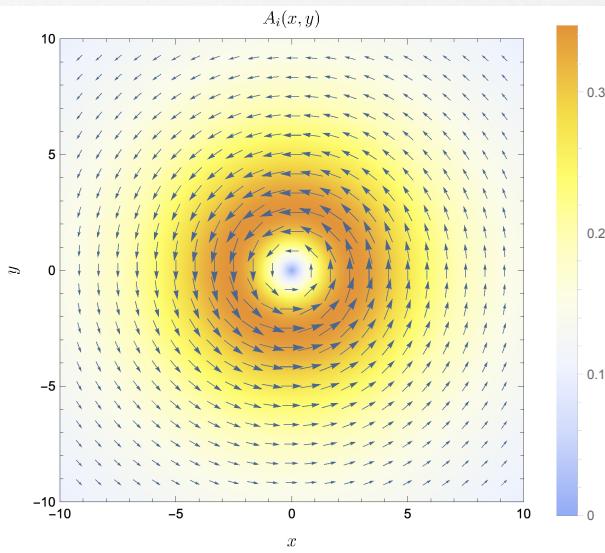
ꝝ  $m_v = 1$

ꝝ  $n = 1$

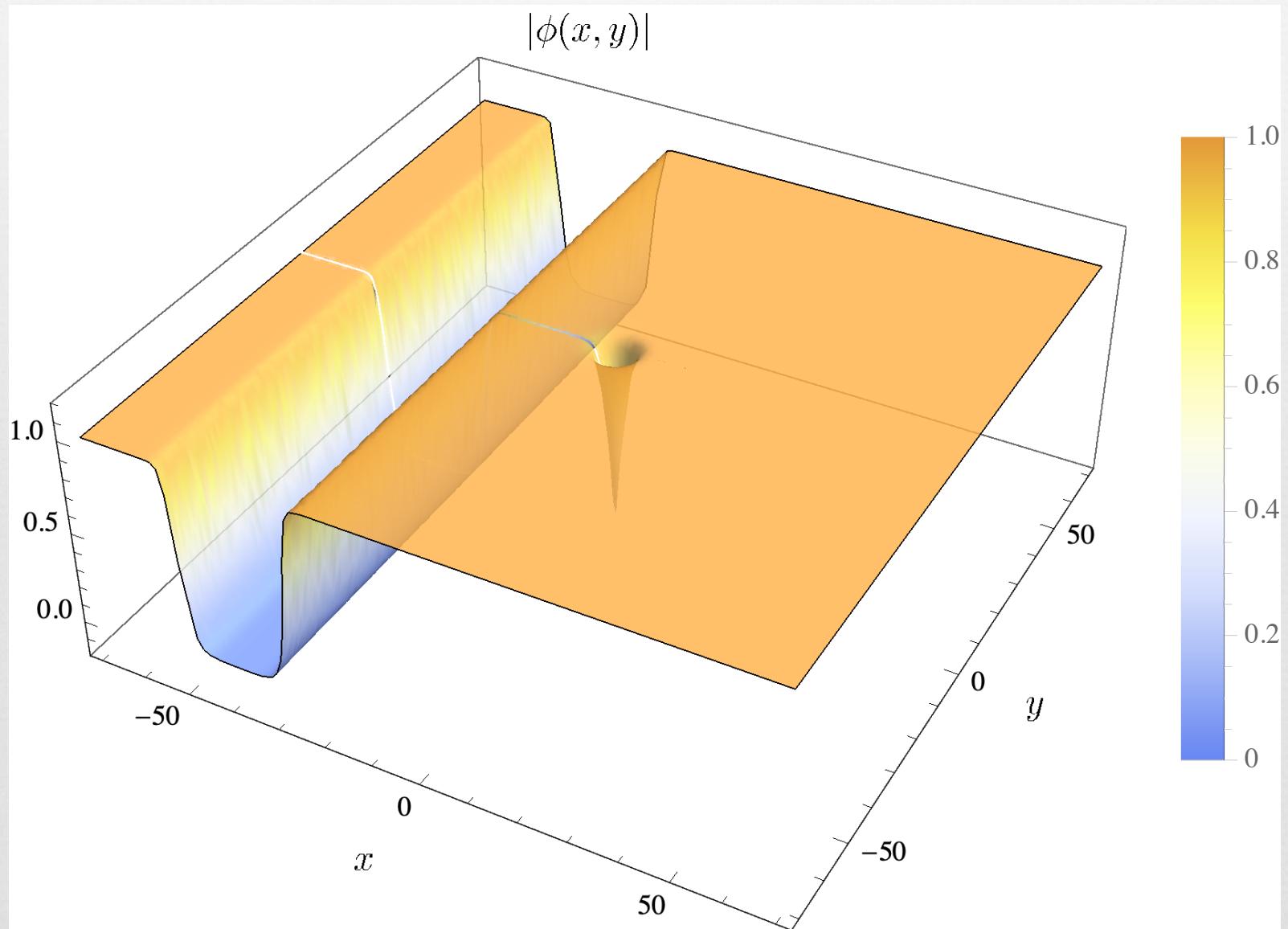
ꝝ Gauge Field:

$$A_i(r, \theta) = -\frac{n}{er} \epsilon_{ij} n_j A(r)$$

$$\text{ꝝ } n = \frac{e}{2\pi} \int B d^2x$$



# Erasure of a Vortex by a Coulomb Vacuum Layer



# Initial Conditions

## $\phi(x, y)$

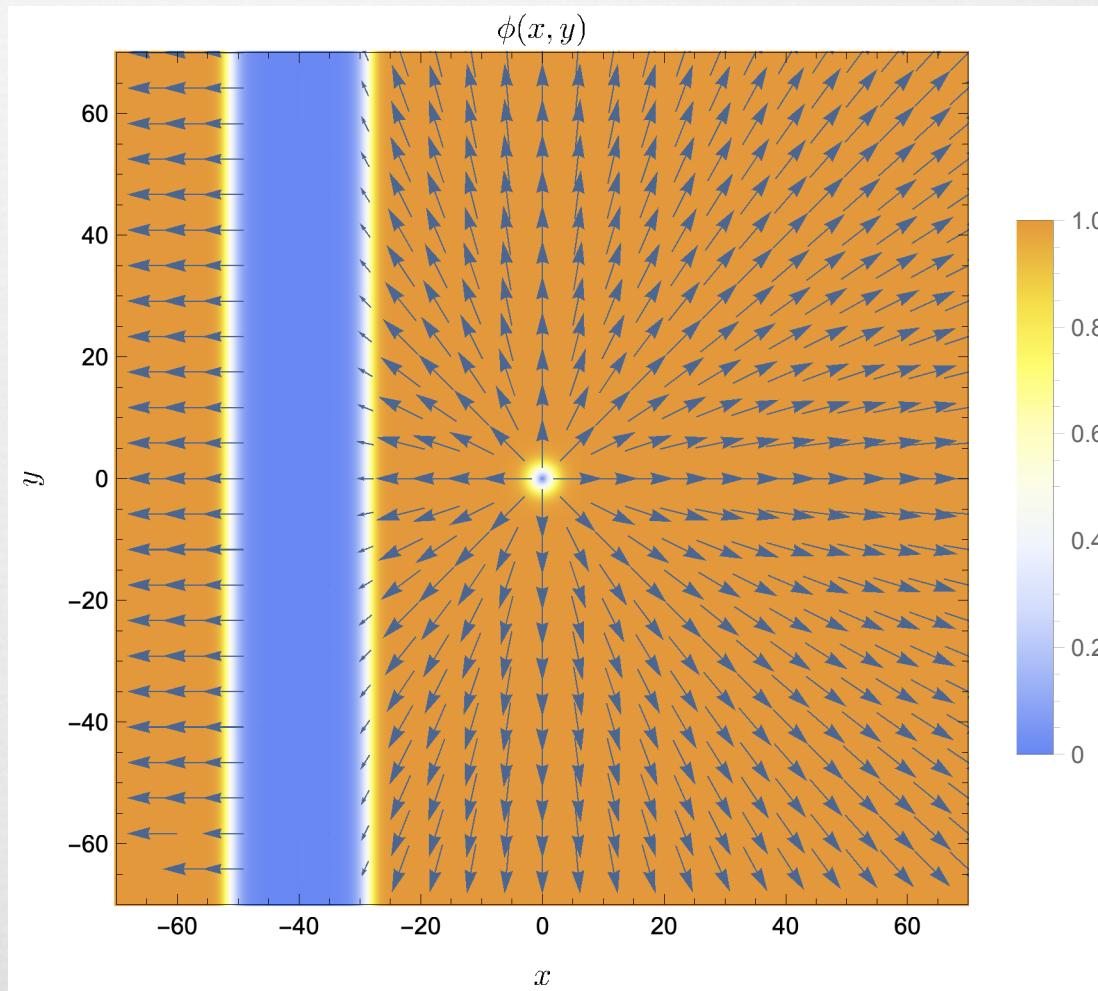
$$m_h = 1$$

$$m_v = 1$$

$$n = 1$$

$$l = 40 m_h^{-1}$$

$$v_0 = 1$$



# Initial Conditions

## $\text{Arg}(\phi(x, y))$

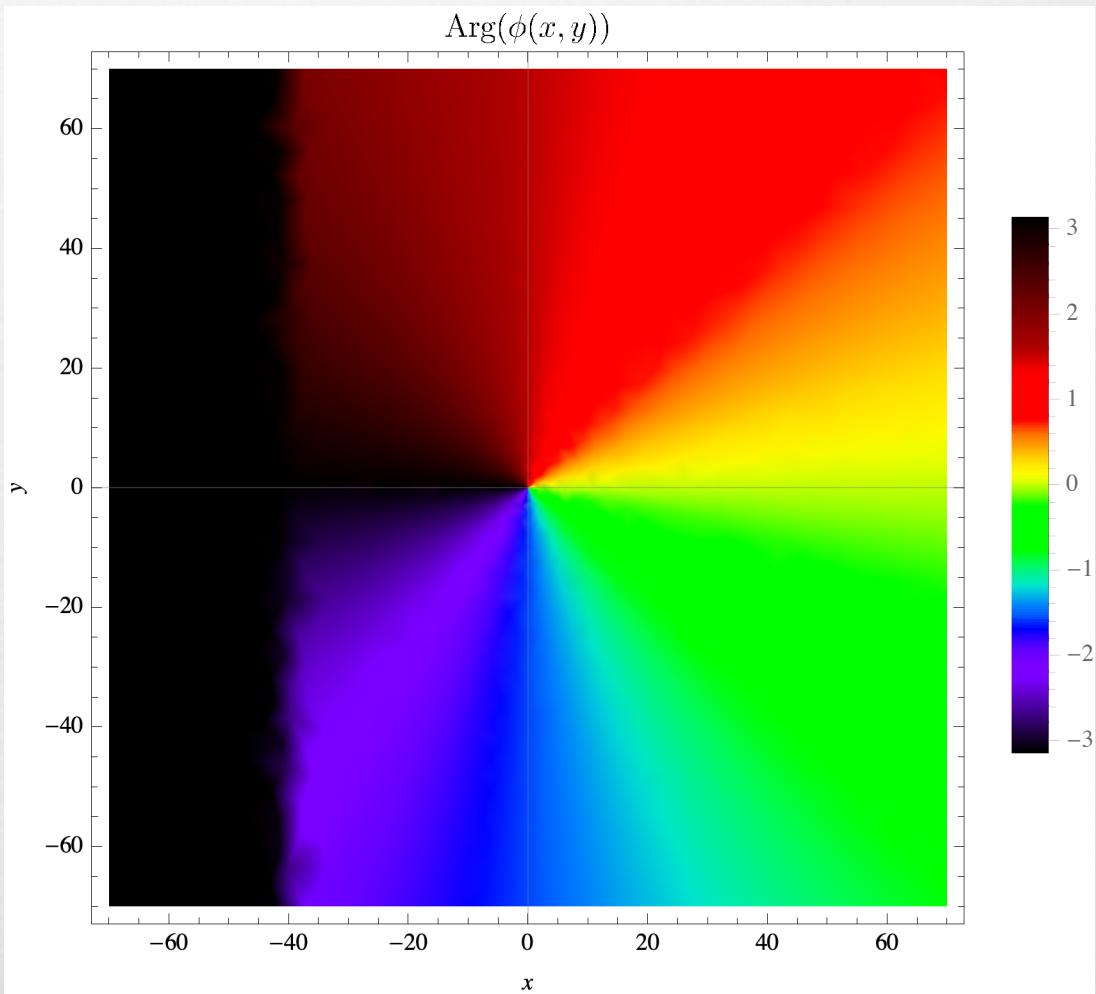
$$m_h = 1$$

$$m_v = 1$$

$$n = 1$$

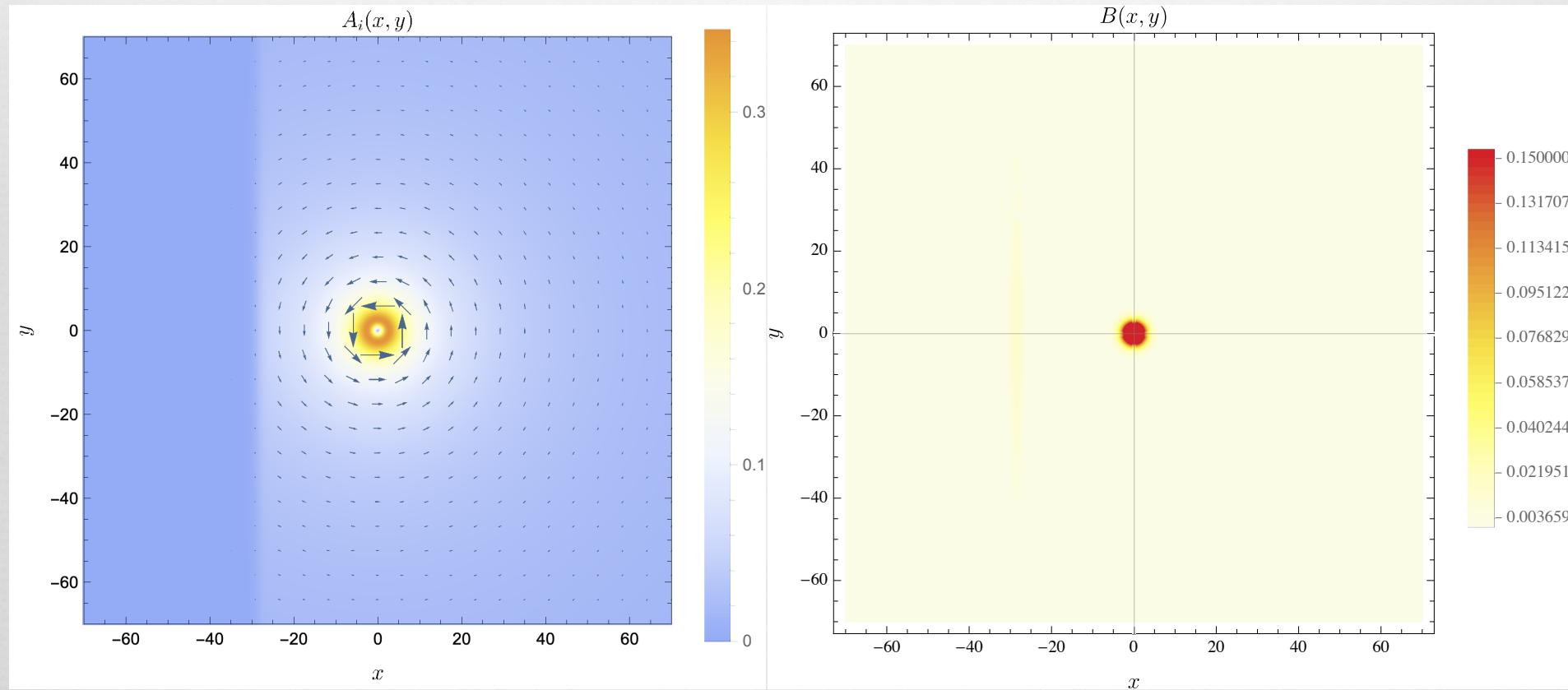
$$l = 40 m_h^{-1}$$

$$v_0 = 1$$



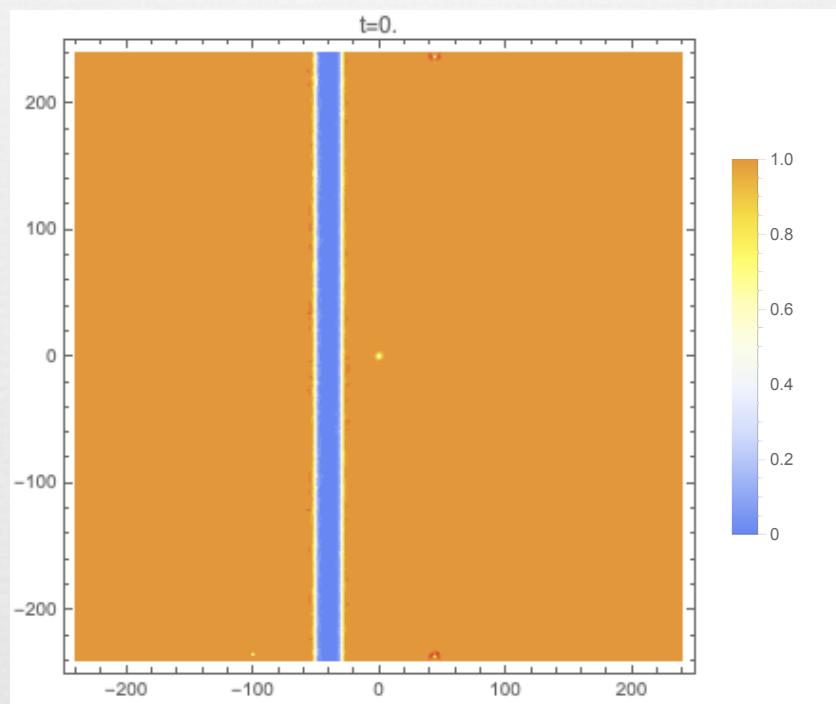
# Initial Conditions

## $A_i(x, y)$ & $B(x, y)$



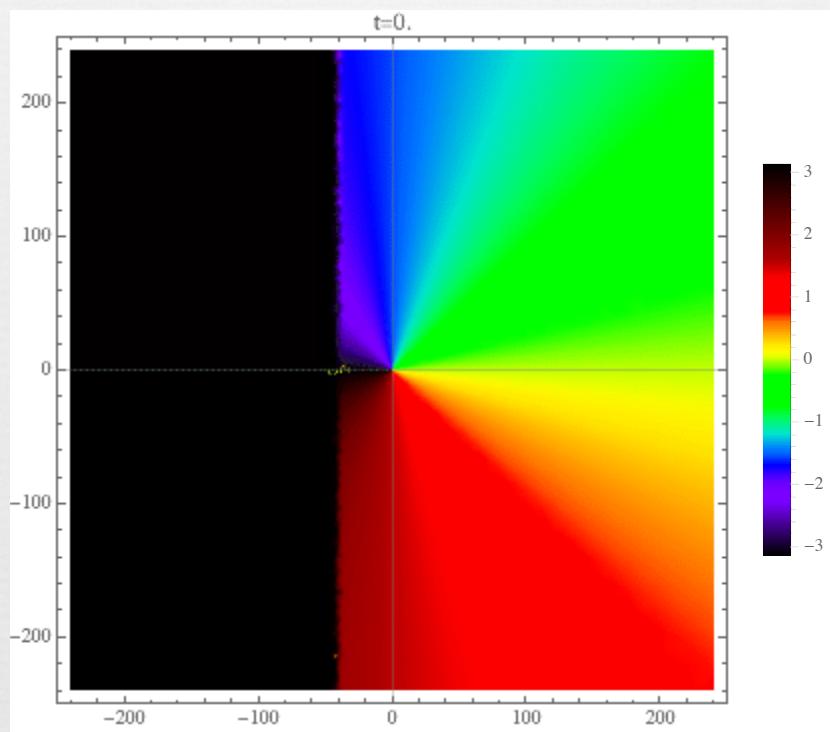
# Time Evolution

## $|\phi(x, y)|$



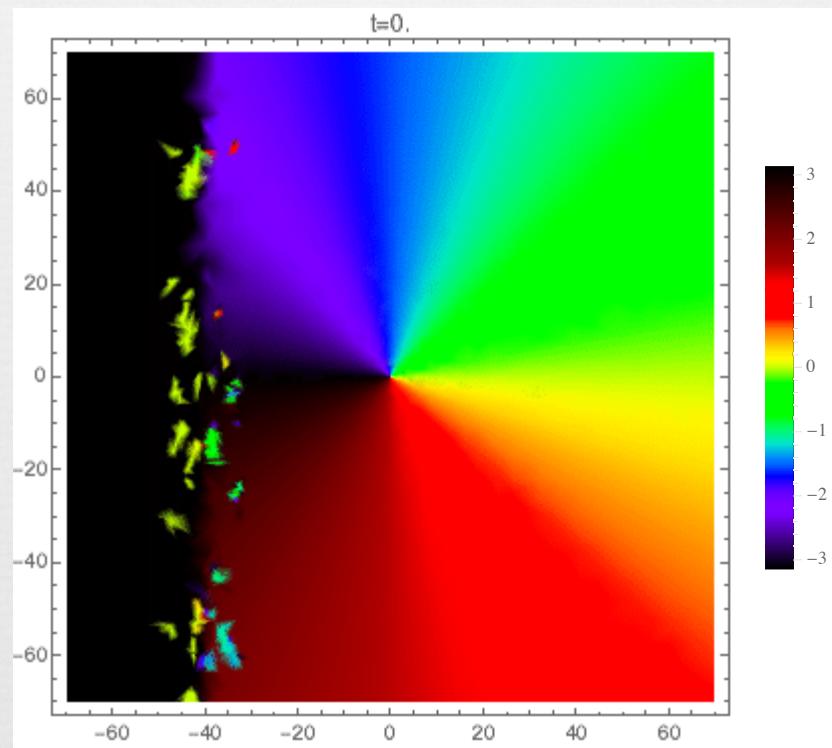
# Time Evolution

## $\text{Arg}(\phi(x, y))$



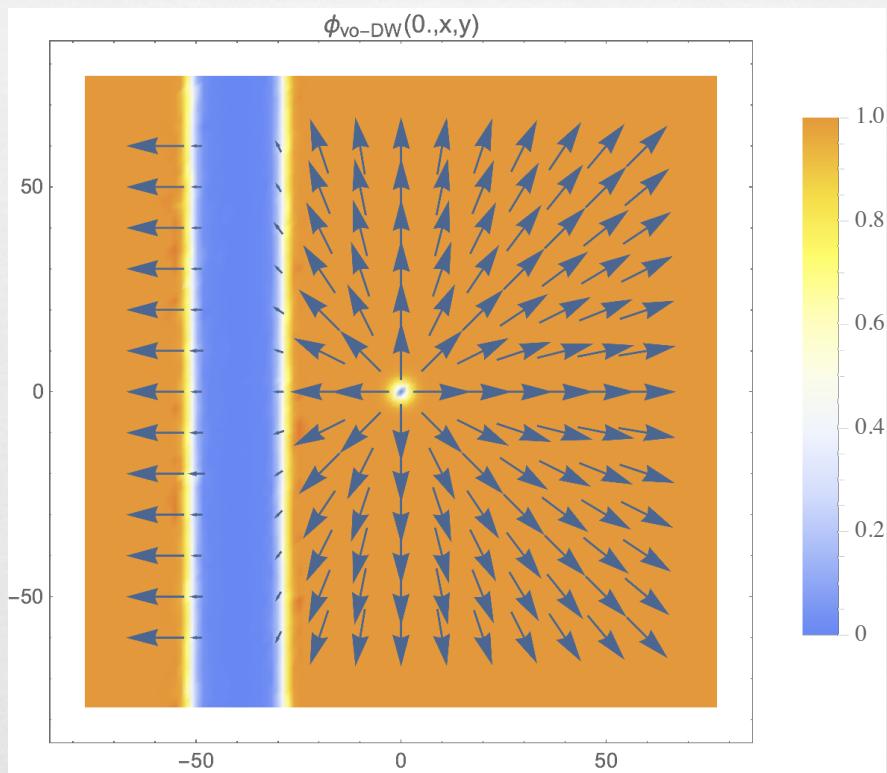
# Time Evolution

## $\text{Arg}(\phi(x, y))$



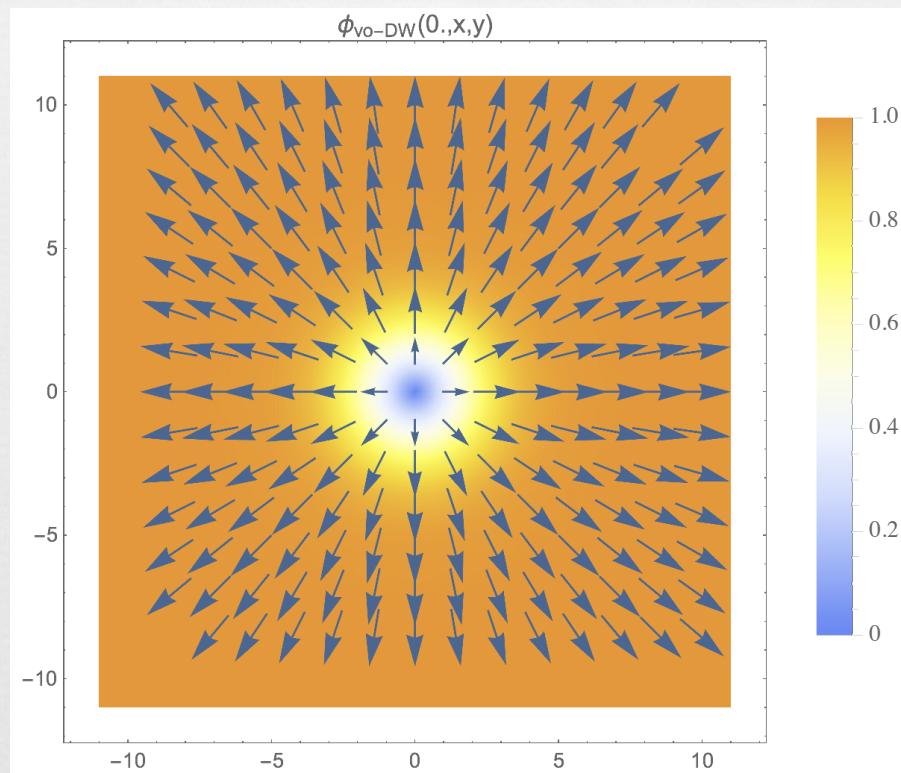
# Time Evolution

## $\phi(x, y)$



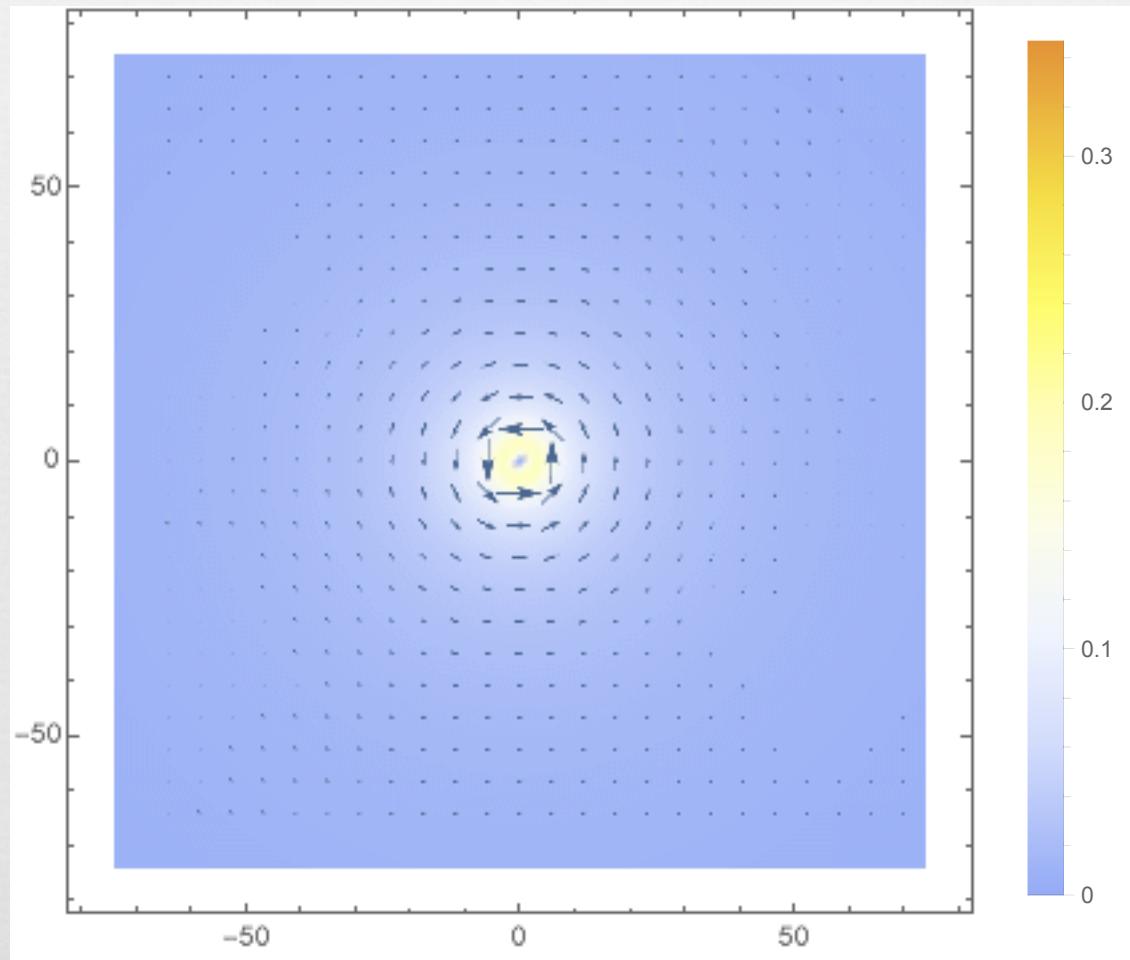
# Time Evolution

## $\phi(x, y)$



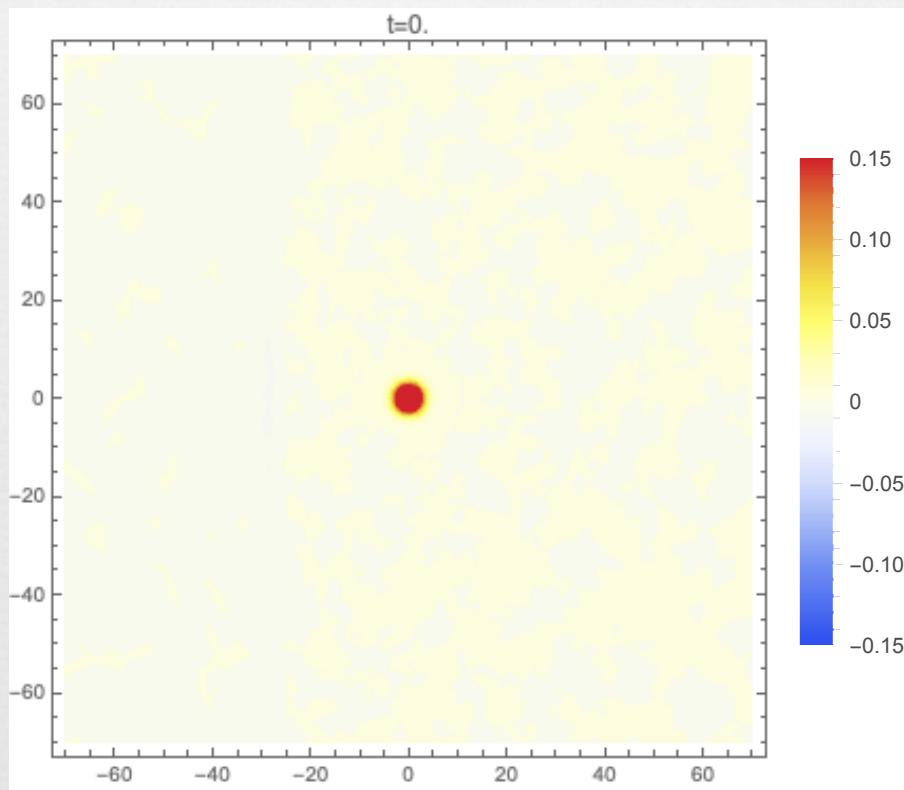
# Time Evolution

$$A_\mu(x, y)$$



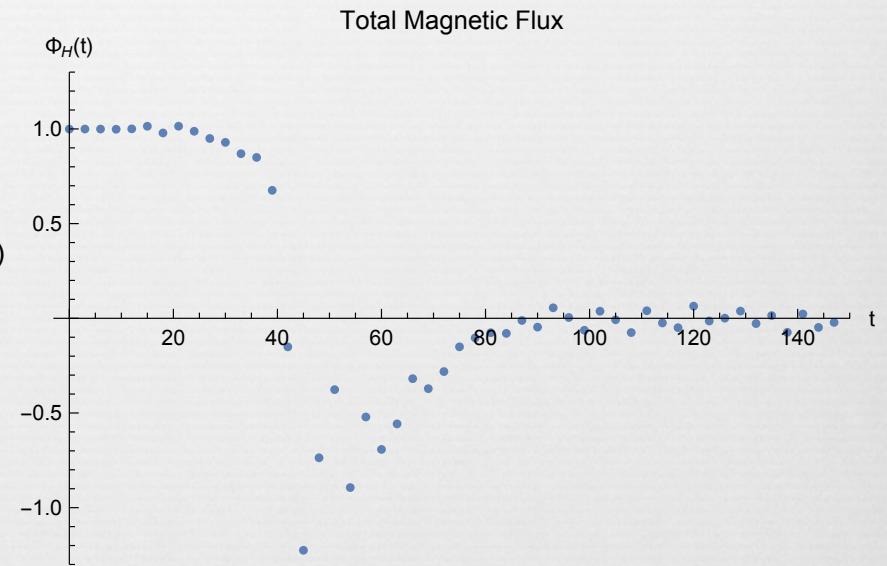
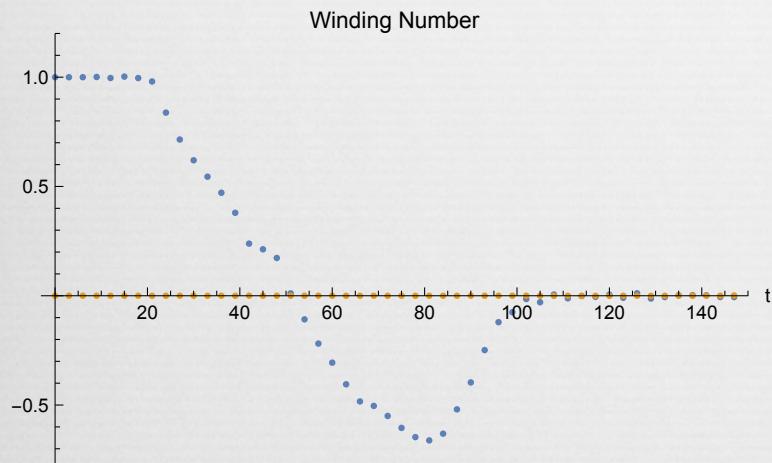
# Time Evolution

## $B(x, y)$



# Time Evolution

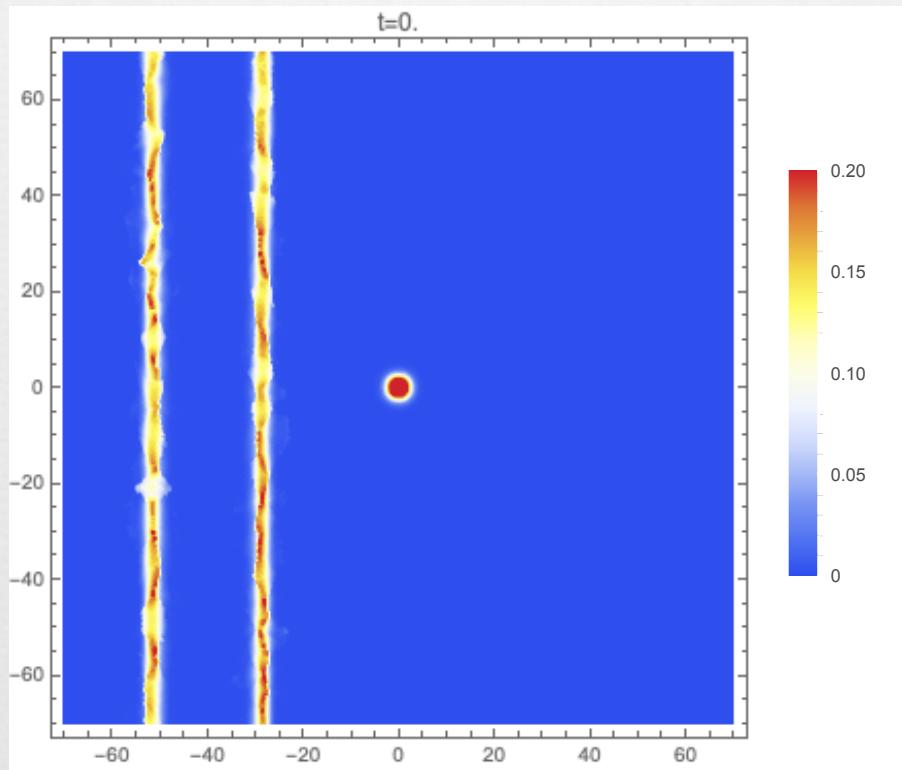
*n*



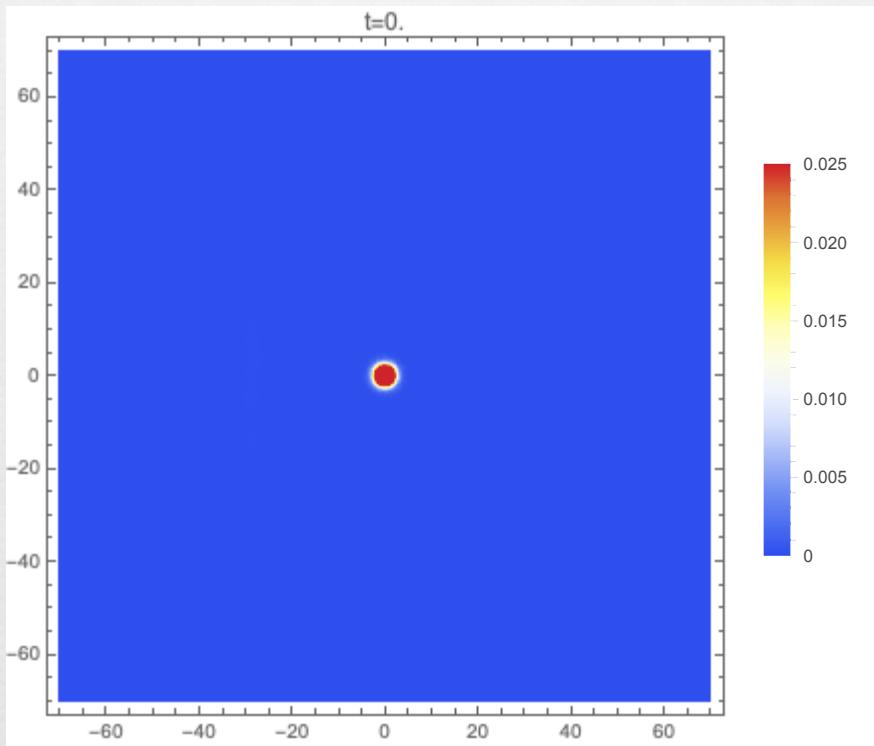
$$n = \frac{1}{2\pi i\nu^2} \oint dx^i \phi^* \partial_i \phi$$

$$n = \frac{e}{2\pi} \int B d^2x$$

# Time Evolution: Energy Density



# Time Evolution: EM Energy Density

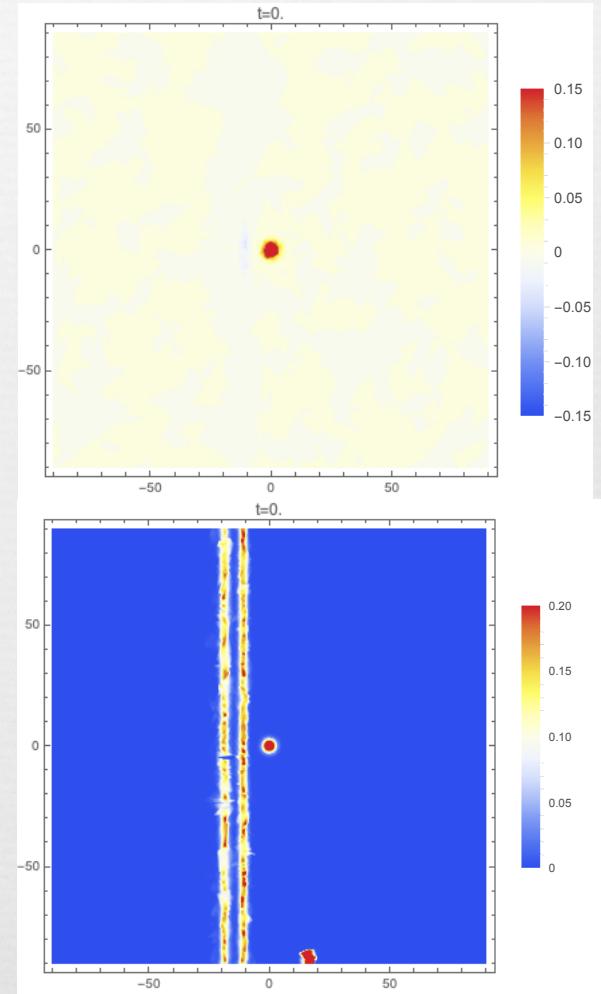
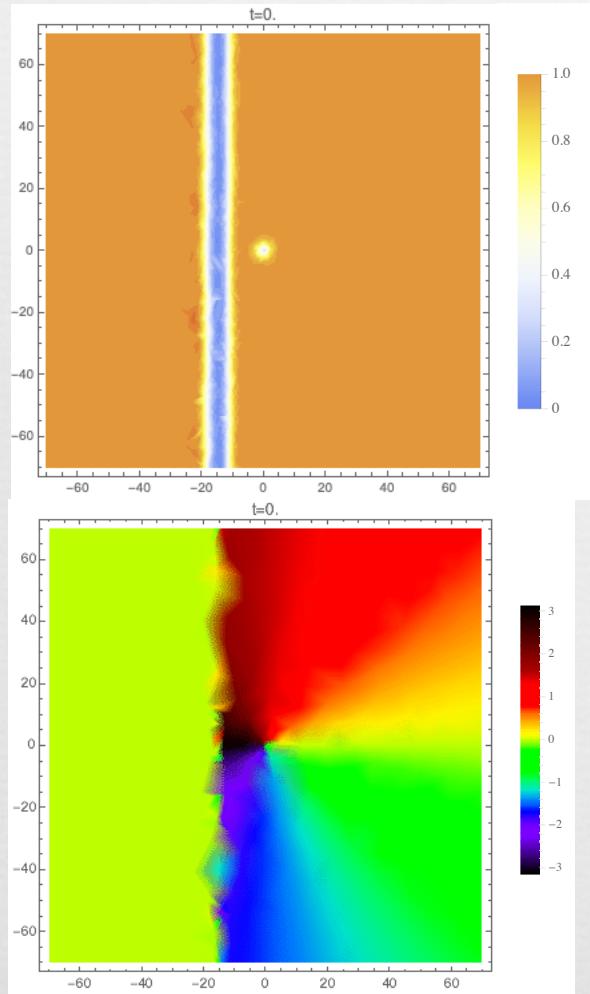


# Conclusions and Outlook

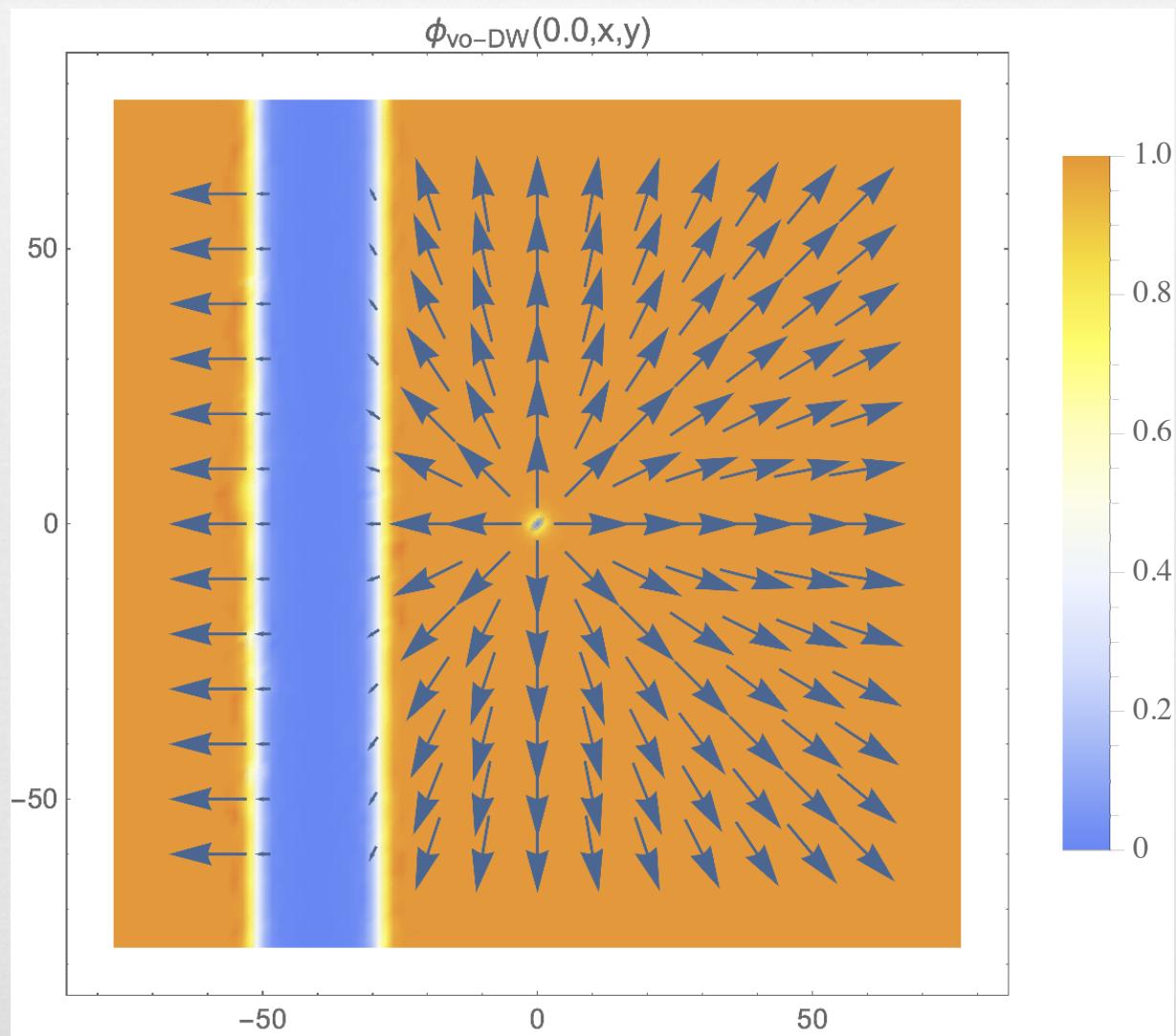
- ❖ Dvali-Liu-Vachaspati mechanism is borne out in the  $\phi^6$ -model.
- ❖ None of the studied regimes presented a case in which the vortex survive.
- ❖ This mechanism can be generalized to other models, and higher dimensions.
- ❖ Further evidence for this mechanism includes Interactions between
  - ❖ Monopoles and Domain Walls,
  - ❖ Skyrmions and Domain Walls,
  - ❖ Vortices and Domain Walls in  ${}^3He$
- ❖ The DLV mechanism suggests that interactions of topological defects lead to the erasure of defects.

Thank You

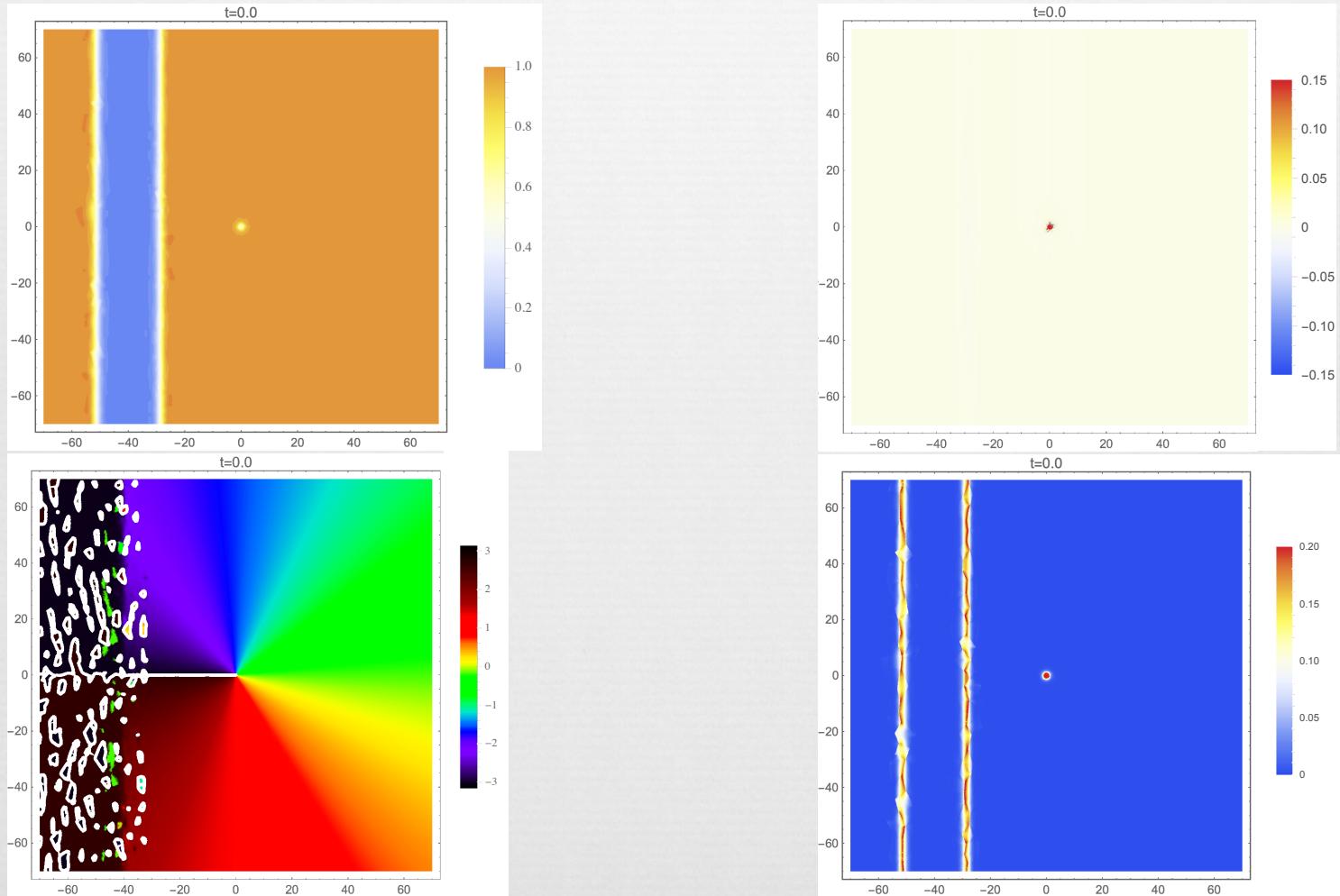
$$l = 15m_h^{-1}$$



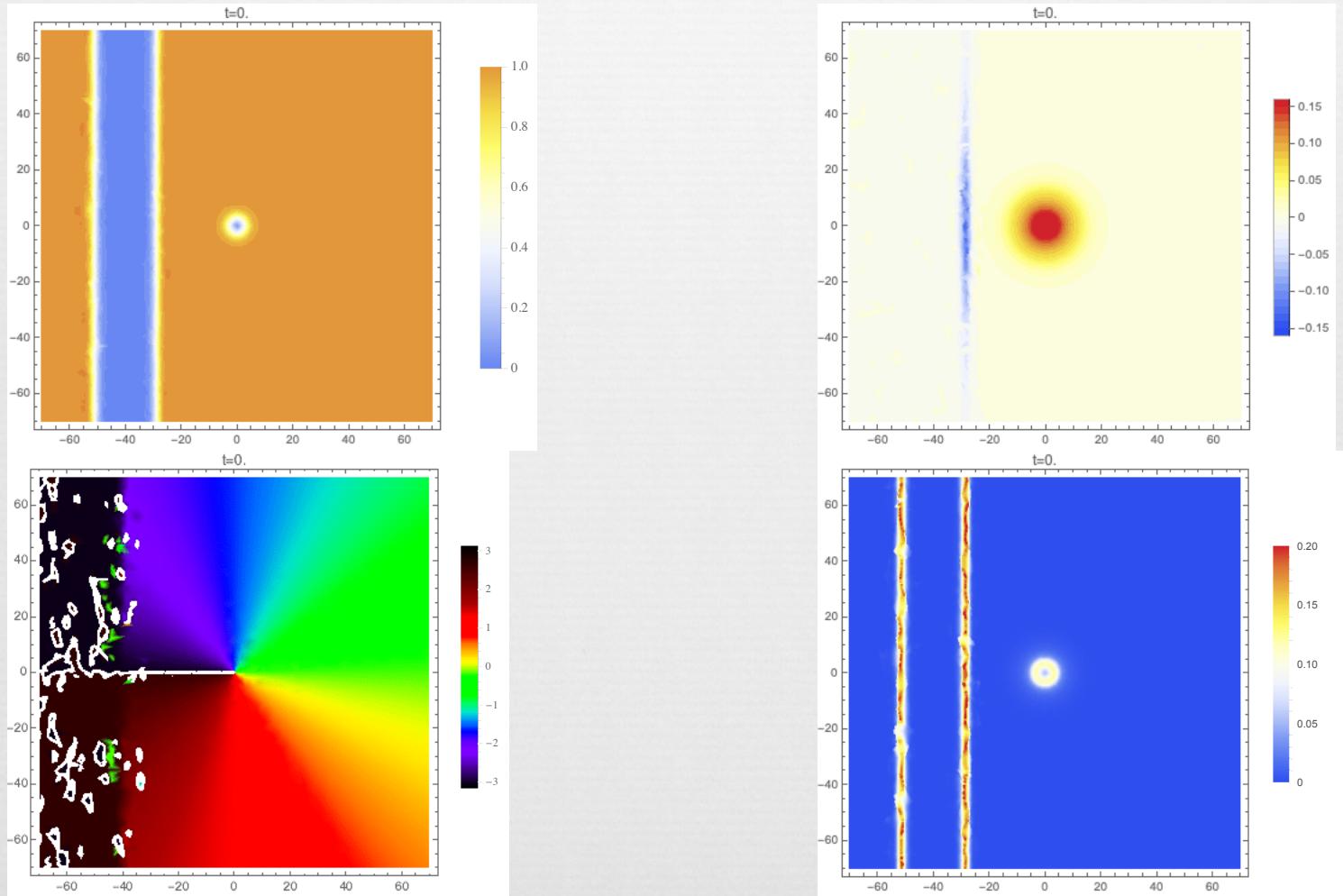
$$m_h < m_v$$



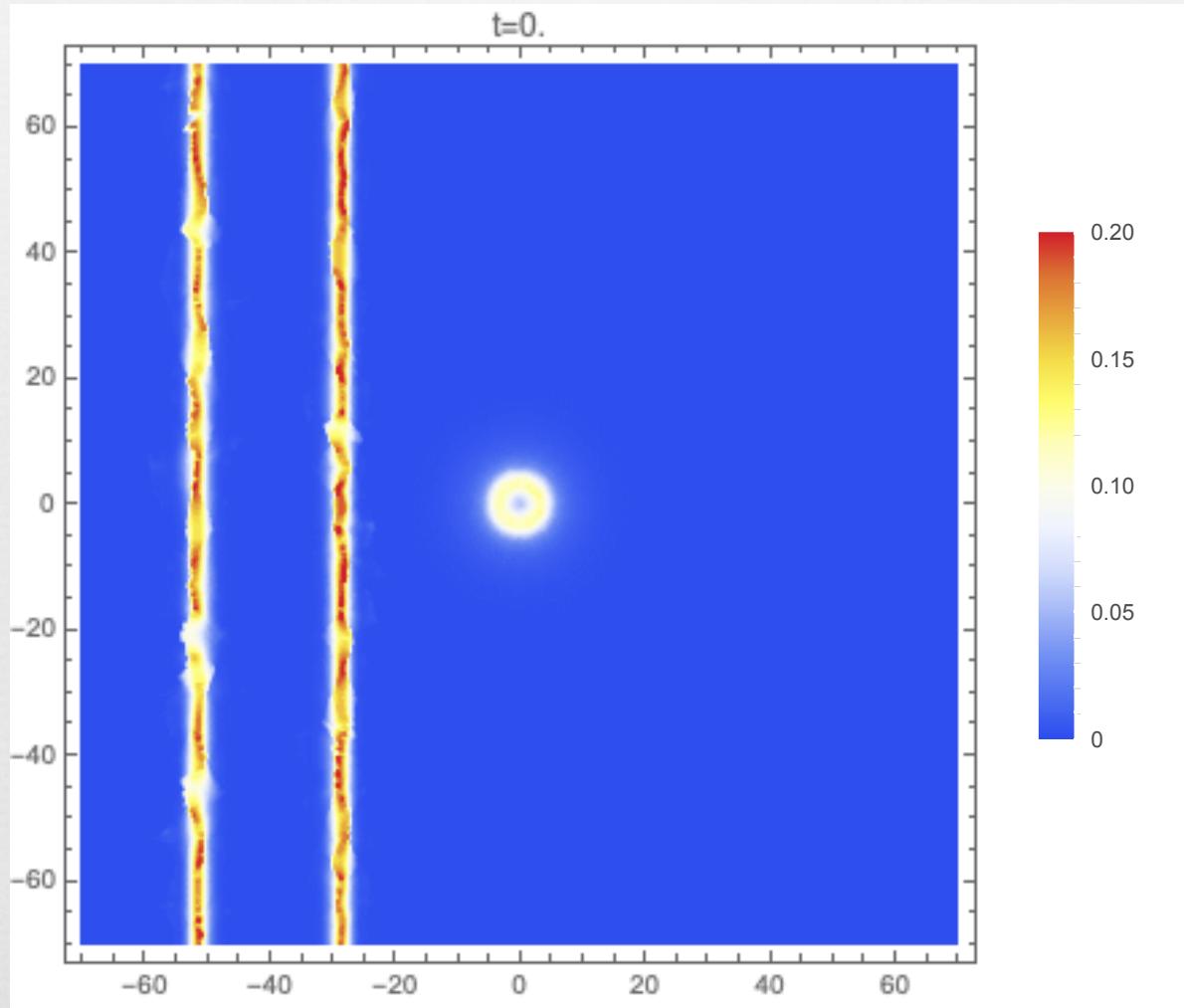
$$m_h < m_\nu$$



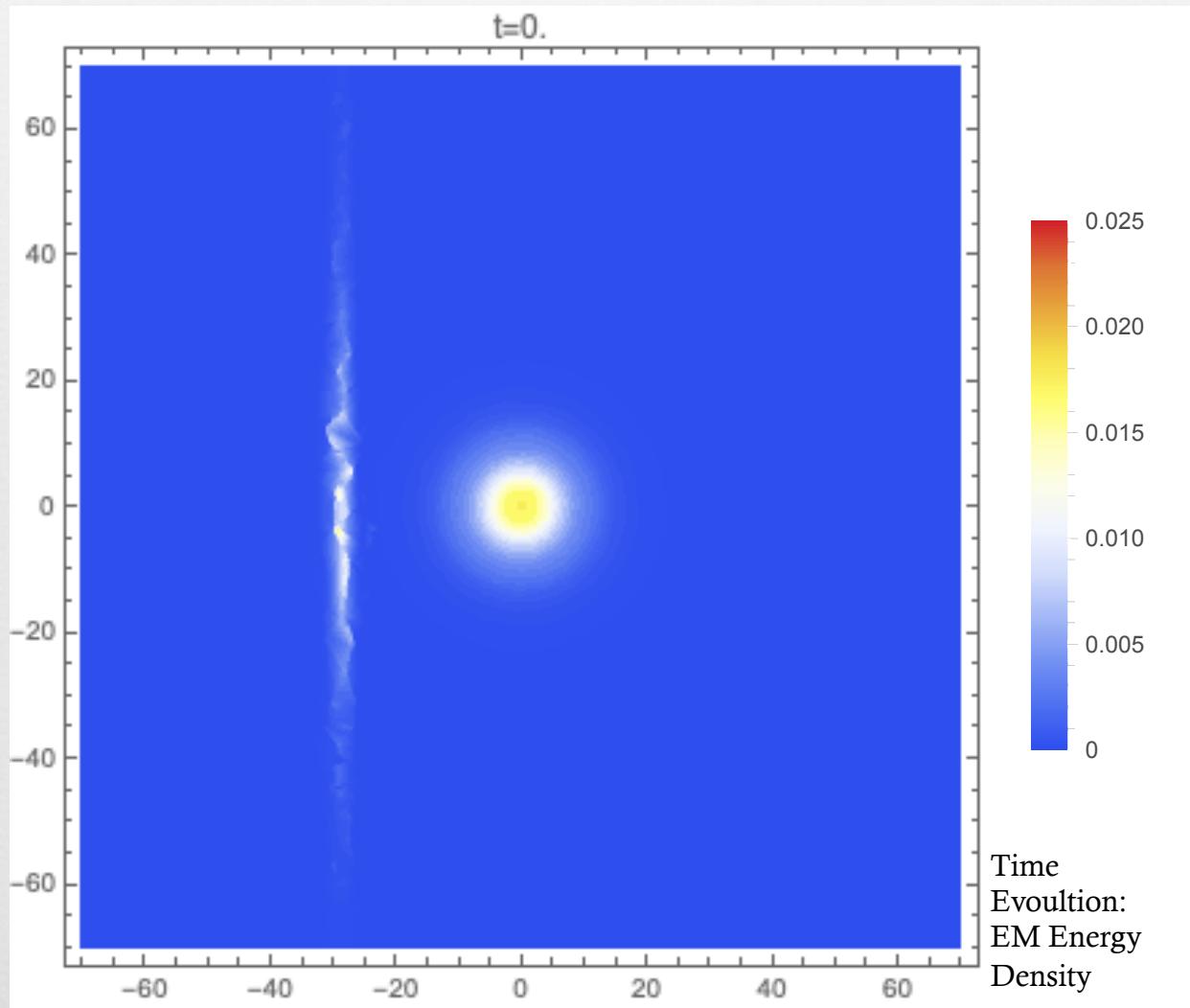
$$m_h > m_\nu$$



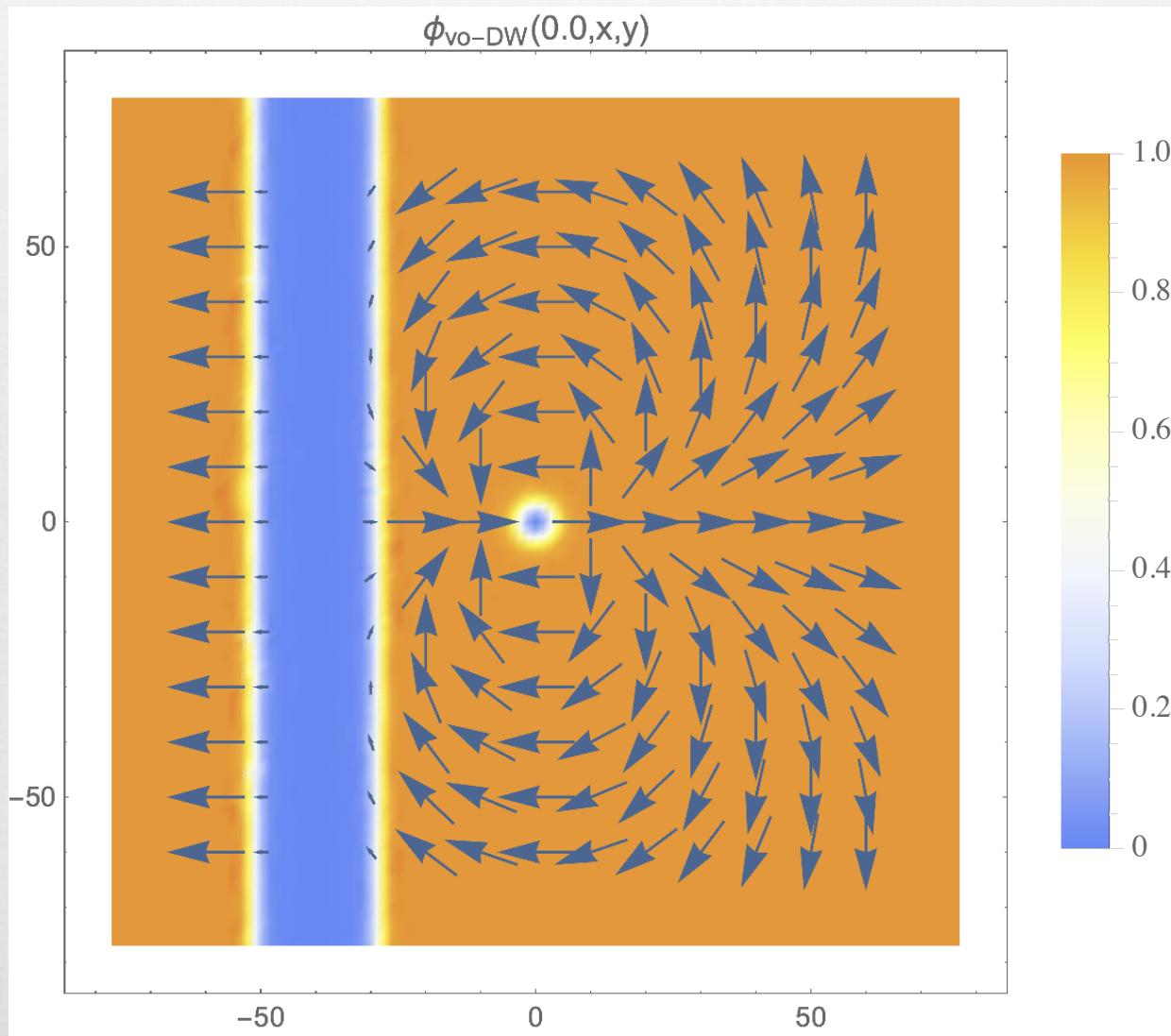
$$m_h > m_\nu$$



$$m_h > m_\nu$$



$n=2$



# Topological Defects in the $\phi^6$ -Model



- ❖ The topology of the vacuum manifold  $\mathcal{M}$  determines the spectrum of defects of the model.

Domain Walls:  $\pi_0(\mathcal{M}) = \mathbb{Z}_2$

$$\pi_1(\mathcal{M}_C) = 0$$

Vortex Lines:  $\pi_1(\mathcal{M}_H) = \mathbb{Z}_2$

$$\pi_n(\mathcal{M}_C) = 0, \quad n \geq 2$$

$$\pi_n(\mathcal{M}_H) = 0, \quad n \geq 2$$

# Domain Walls in the $\phi^6$ -Model

❖ (1+1)-dimensions

❖ Constraining the field  $\phi$  to be a real field:  $\phi = \frac{\chi}{\sqrt{2}}$

$$\mathcal{L}_{DW} = \frac{1}{2} \partial_\mu \chi \partial^\mu \chi - \lambda^2 \frac{\chi^2}{2} \left( \frac{\chi^2}{2} - \nu^2 \right)^2.$$

❖ Domain walls are solitonic configurations of  $\chi(x)$ , with the following asymptotic behavior

$$\lim_{x \rightarrow -\infty} \chi(x) = 0$$

$$\lim_{x \rightarrow \infty} \chi(x) = \pm \sqrt{2\nu}$$

# Vortex in the $\phi^6$ -Model

- ꝝ In (2+1) dimensions
- ꝝ Asymptotic Field Configuration as  $r \rightarrow \infty$

$$\phi \rightarrow \nu e^{ig(\theta)} = \nu e^{in\theta}$$

$$A_i \rightarrow \frac{1}{e} \partial_i g(\theta) = -\frac{n}{er} \epsilon_{ij} n_j$$

where  $n \in \mathbb{Z}$  is the *winding number*.

# Vortex in the $\phi^6$ -Model

⊗ Ansatz<sup>4</sup> for Vortex Profile

$$\phi(r, \theta) = \nu e^{in\theta} F(r)$$

$$A_i(r, \theta) = -\frac{n}{er} \epsilon_{ij} n_j A(r)$$

⊗ The field equations

$$\square \phi + \frac{\partial V(\phi)}{\partial \phi^*} = 0,$$

$$\partial_\mu F^{\mu\nu} = j^\nu,$$

<sup>4</sup>H.B. Nielsen, P. Olesen, Nuc. Phys. B, Vol. 61, 45-61, (1973)

# Vortex in the $\phi^6$ -Model

ꝝ The field equations become:

$$0 = -rF''(r) - F'(r) + \frac{(1 - A(r))^2}{r}n^2F(r)$$

$$+ \frac{m_h^2}{4}rF(r) \left(F(r)^2 - 1\right) \left(3F(r)^2 - 1\right)$$

$$0 = -\frac{A''(r)}{r} + \frac{A'(r)}{r^2} - \frac{m_v^2}{r}(1 - A(r))F(r)^2$$