

ALMA MATER STUDIORUM · UNIVERSITY OF BOLOGNA
School of Science
Department of Physics and Astronomy
Master Degree in Physics

Backreaction from magnetogenesis during inflation

Sofia Corbà

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study of the inflationary evolution in the presence of backreaction effects due to the production of primordial magnetic fields

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Coupling:

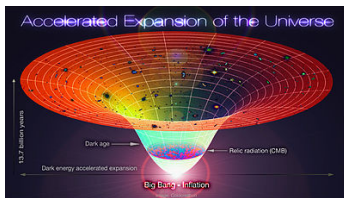
$$\mathcal{L} = - \int d^4x \frac{1}{4} \sqrt{-g} f^2(\phi) F_{\mu\nu} F^{\mu\nu}$$

Inflation

Inflation

Initial period of accelerated expansion governed by a scalar field

$$S = -\frac{M_{pl}^2}{2} \int d^4x \sqrt{-g} R + \int d^4x \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right)$$



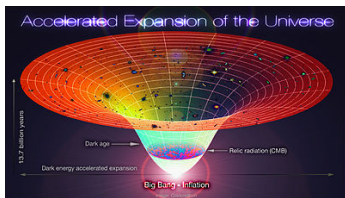
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$$H^2 = \frac{1}{3M_{pl}^2} \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right)$$

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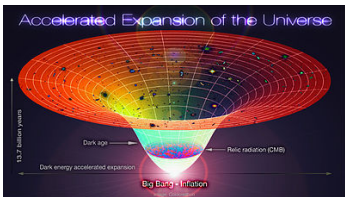
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- Horizon Problem
- Flatness Problem
- CMB Anisotropies
- Large-scale structure

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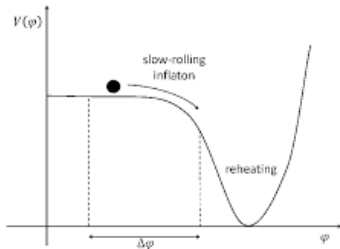
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Reheating → Standard Model degrees of freedom

Primordial Magnetogenesis

The action of the system is: $S = -\frac{M_{pl}^2}{2} \int d^4x \sqrt{-g} \mathcal{R} + S_M$

$$S_M = \int d^4x \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) - \frac{\alpha}{4\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} - \frac{f^2(\phi)}{4} F_{\mu\nu} F^{\mu\nu} \right)$$

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$$\nabla \cdot \mathbf{E} = 0 \quad \partial_0(a^2 \mathbf{E}) + \frac{1}{f^2(\phi)} \frac{d(f^2(\phi))}{d\phi} \dot{\phi} a^2 \mathbf{E} - a \nabla \wedge \mathbf{B} = 0$$

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$$\frac{d^2 u_{\mathbf{k}}^\lambda(\eta)}{d\eta^2} + \left(k^2 - \frac{1}{f(\phi)} \frac{d^2 f(\phi)}{d\eta^2} \right) u_{\mathbf{k}}^\lambda(\eta) = 0$$

Primordial Magnetogenesis

$$\rho_E = \frac{H^2 f^2(\phi)}{4\pi^2 a^2} \int_0^\infty dk k^2 \left| \frac{1}{f(\phi)} \frac{du_{\mathbf{k}}(N)}{dN} - \frac{1}{f^2(\phi)} \frac{df(\phi)}{d\phi} \frac{d\phi}{dN} u_{\mathbf{k}}(N) \right|^2$$

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$$\rho_B = \frac{f^2(\phi)}{4\pi^2 a^4} \int_0^\infty dk k^4 \left| \frac{u_{\mathbf{k}}(N)}{f(\phi)} \right|^2$$

$$N(t) = \ln \frac{a(t_{end})}{a(t)}$$

Backreaction from magnetogenesis in string inflation

Inflation:
ultraviolet sensitive phenomenon



String Theory

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ultraviolet sensitive phenomenon → String Theory

Cosmological applications:

- compactification of the 6 extra dimensions on a Calabi-Yau;

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$N = 1, D = 4$ SUPERGRAVITY

By expanding the
10D fields
in harmonic forms

- Complex Structure Moduli U_α
- Axio-dilaton S
- Kähler Moduli $T_i = \tau_i + i\theta_i$

Moduli Stabilization

Complex Structure Moduli
Axio-dilaton



Stabilized at tree level

Moduli Stabilization

Complex Structure Moduli
Axio-dilaton \rightarrow Stabilized at tree level

Kähler Moduli \rightarrow Perturbative corrections
to the Kähler Potential
Non perturbative corrections
to the Superpotential

$$W = W_0 + \sum_i A_i e^{-a_i T_i}$$

$$K = -2 \ln \left(\mathcal{V} + \frac{\xi}{2} \right)$$

Moduli Stabilization

Complex Structure Moduli
Axio-dilaton

→

Stabilized at tree level

Kähler Moduli

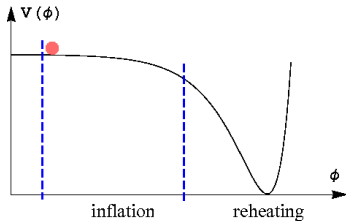
→

Perturbative corrections
to the Kähler Potential
Non perturbative corrections
to the Superpotential

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$$K = -2 \ln \left(\mathcal{V} + \frac{\xi}{2} \right)$$

$$V_F = e^K \left((K^{-1})^{\bar{i}j} D_{\bar{i}} W D_j \bar{W} - 3|W|^2 \right)$$

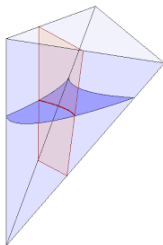


Why are we interested in Backreaction effects?

[Cicoli, Ciupke, Mayrhofer, Shukla]

$$\mathcal{V}(\tau_i) = \langle \mathcal{V} \rangle$$

$$\tau_s = \langle \tau_s \rangle$$



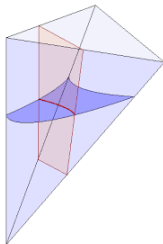
$$\frac{\Delta\phi}{M_{pl}} \leq c \ln(\mathcal{V}) \Rightarrow \text{restriction on } N_e$$

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We need a mechanism which slows down the inflaton



COUPLING OF THE INFLATON WITH ELECTROMAGNETIC
FIELDS

String inflationary models

- Kähler Inflation

$$\mathcal{V} = \alpha \left(\tau_1^{3/2} - \sum_{i=2}^n \lambda_i \tau_i^{3/2} \right) \quad f_K(\phi) = \sqrt{1 + \gamma \left(\frac{\phi^{4/3}}{\langle \phi \rangle^{4/3}} - 1 \right)}$$

$$V(\tau_n) = \frac{W_0^2 \beta}{\mathcal{V}^3} + \frac{8aA\sqrt{\tau_n} e^{-2a\tau_n}}{3\mathcal{V}\lambda\alpha} - \frac{4aAW_0\tau_n e^{-a\tau_n}}{\mathcal{V}^2}$$

- Fibre Inflation

$$\mathcal{V} = \alpha \left(\sqrt{\tau_1} \tau_2 - \gamma \tau_3^{3/2} \right) \quad f_F(\phi) = \sqrt{1 + \gamma \left(\frac{e^{k\phi}}{e^{k\langle \phi \rangle}} - 1 \right)}$$

$$V(\hat{\phi}) = 3 - 4e^{-\frac{\hat{\phi}}{\sqrt{3}}} + e^{-\frac{4\hat{\phi}}{\sqrt{3}}}$$

Equations

- $$\frac{d^2\phi}{dN^2} + \left(3 - \frac{1}{2}\left(\frac{d\phi}{dN}\right)^2\right) \left(\frac{d\phi}{dN} + \frac{V'(\phi)}{V(\phi)}\right) = 0$$

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Backreaction

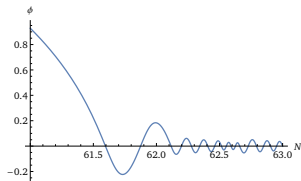
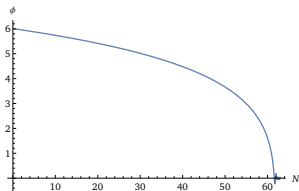
- $$\frac{d^2\phi}{dN^2} + \left(3 - \frac{1}{2} \left(\frac{d\phi}{dN}\right)^2\right) \left(\frac{1}{6} \frac{d\rho_E}{dN} \frac{d\phi}{dN} + \frac{d\phi}{dN} (V(\phi) + \rho_E) + \frac{1}{3} \left(\frac{d\phi}{dN}\right)^2 \rho_E \frac{f'(\phi)}{f(\phi)} + V'(\phi) - 2 \frac{f'(\phi)}{f(\phi)} \rho_E\right) \frac{1}{V(\phi) + \rho_E} = 0$$

- $$\frac{d\rho_E}{dN} + 4\rho_E + 2 \frac{df(\phi)}{dN} \frac{1}{f(\phi)} \rho_E =$$

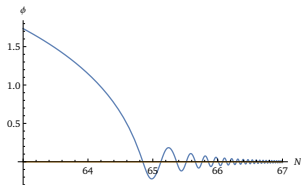
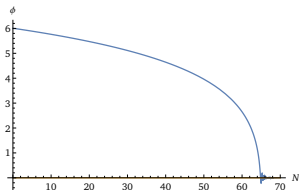
$$\left(V(\phi) + \rho_E\right)^2 \frac{1}{4\pi^2} \left(1 + \left(\frac{df(\phi)}{dN} \frac{1}{f(\phi)}\right)^2\right) \left(3 - \frac{1}{2} \left(\frac{d\phi}{dN}\right)^2\right)^{-2}$$

Numerical results: Fibre Inflation

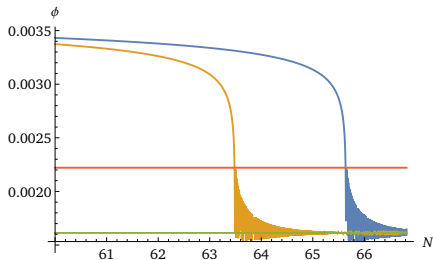
In absence of Backreaction



In presence of Backreaction



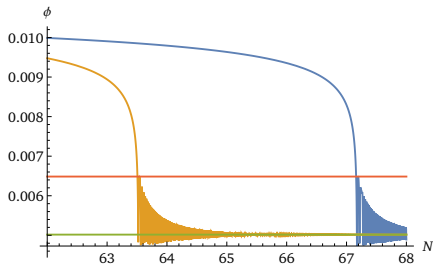
Numerical results: Kähler Inflation



$$W_0 = 400 \quad \nu = 10^4$$

$$\alpha = 1/100 \quad \lambda = 1$$

$$A = 1 \quad a = 2\pi$$



$$W_0 = 500 \quad \nu = 10^4$$

$$\alpha = 1/100 \quad \lambda = 10$$

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Conclusions

Backreaction due to the coupling of the inflaton with the electromagnetic field



slowdown of the inflationary evolution

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Backreaction due to the coupling of the inflaton with the electromagnetic field



slowdown of the inflationary evolution

- compatibility of string inflationary models with phenomenology;
- production of primordial magnetic fields.

Thanks for your attention

The Standard Cosmological Model

Friedmann-Robertson-Walker metric:

$$dS^2 = dt^2 - a^2(t) \left(\frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right)$$

Einstein Equations:

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi GT_{\mu\nu}$$

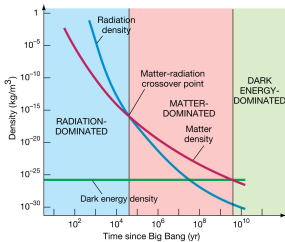
$$T_{\mu\nu} = (\rho + p)u_\mu u_\nu - pg_{\mu\nu}$$

$$p = w\rho$$

Friedmann Equations:

$$\dot{a}^2 + k = \frac{8\pi G}{3}\rho a^2$$

$$\ddot{a} = -\frac{4\pi G}{3}(\rho + 3p)a$$



$$\left\{ \begin{array}{ll} w = 0 & \text{matter} \\ w = 1/3 & \text{radiation} \\ w = -1 & \text{cosmological constant} \end{array} \right.$$

$$\rho(t) = \rho_0 \left(\frac{a_0}{a(t)} \right)^{3(1+w)}$$

Problems in the Standard Cosmological Model

- **Horizon Problem**

The high degree of homogeneity
of the CMB



thermal equilibrium at decoupling

- **Flatness Problem**

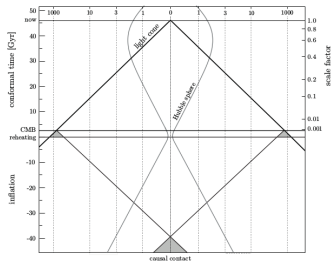
$$\Omega = \frac{\rho}{\rho_c} \quad \frac{d\Omega}{d \log a} = \Omega(\Omega - 1)(1 + 3w)$$

$\Omega = 1$: unnatural fine tuning of the initial conditions

Solution → Inflation

Inflation: initial period of accelerated expansion governed by a scalar field that violates the strong energy condition $w > -1/3$

- Large scale structure
- CMB anisotropies



Inflation

Scalar field with negative pressure: $T_{\mu\nu} = \text{diag}(\rho, -p, -p, -p)$

$$S = -\frac{M_{pl}^2}{2} \int d^4x \sqrt{-g} R + \underbrace{\int d^4x \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right)}_{S_M}$$

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{\delta S_M}{\delta g^{\mu\nu}} \rightarrow \begin{cases} \rho = \frac{1}{2} \dot{\phi}^2 + V(\phi) \\ p = \frac{1}{2} \dot{\phi}^2 - V(\phi) \end{cases}$$

$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0 \quad H^2 = \frac{1}{3M_{pl}^2} \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right)$$

SLOW ROLL: $3H\dot{\phi} \simeq -V'(\phi) \quad H^2 \simeq \frac{1}{3M_{pl}^2} V(\phi)$

Inflation

Slow Roll parameters:

$$\epsilon = \frac{M_{pl}^2}{2} \left(\frac{V'(\phi)}{V(\phi)} \right)^2 \quad \eta = M_{pl}^2 \frac{V''(\phi)}{V(\phi)} \quad \epsilon, |\eta| \ll 1$$

Number of e-foldings:

$$N_e = \ln \frac{a(t_{end})}{a(t)} \simeq \frac{1}{M_{pl}^2} \int_{\phi_{end}}^{\phi} \frac{V}{V'} d\phi$$

End of inflation:

$$\epsilon \sim 1$$

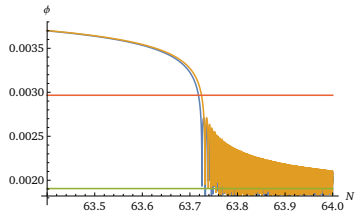
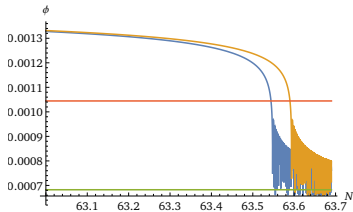
$$|\eta| \sim 1$$

$$N_e \simeq 50 - 60$$

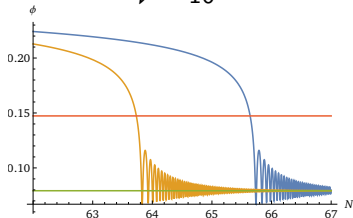
Reheating \rightarrow Standard Model degrees of freedom

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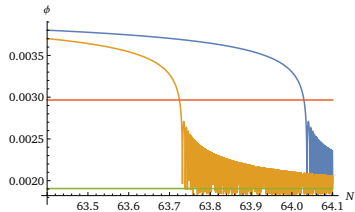
$$W_0 = 57 \quad \alpha = 1 \quad \lambda = 1 \quad A = 1.87 \quad a = 2\pi$$



$$\nu = 10^7$$



$$\nu = 10^6$$



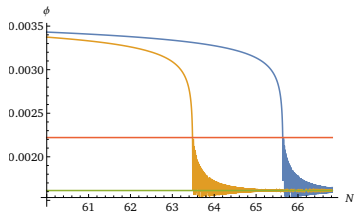
$$\nu = 10^2$$

$$\nu = 10^6, \quad \rho(0) = 10^{-22}$$

$$W_0 = 130 \quad \mathcal{V} = 10^4$$

$$\alpha = 1/100 \quad \lambda = 1$$

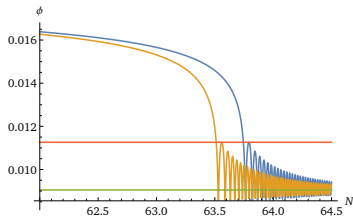
$$A = 10 \quad a = 2\pi/10$$



$$W_0 = 400 \quad \mathcal{V} = 10^4$$

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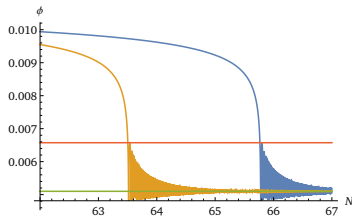
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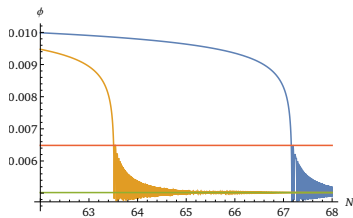
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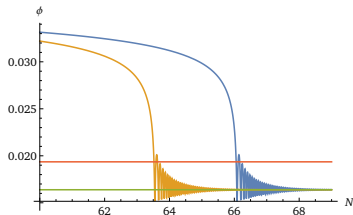
$$A = 10 \quad a = 2\pi/30$$



$$W_0 = 500 \quad \nu = 10^5$$

$$\alpha = 1/100 \quad \lambda = 10$$

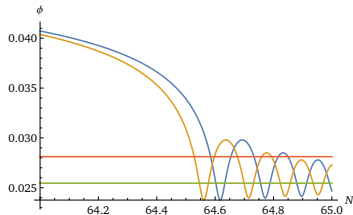
$$A = 50 \quad a = 2\pi/50$$

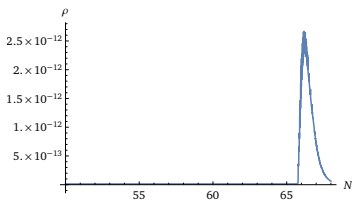


$$W_0 = 500 \quad \nu = 10^4$$

$$\alpha = 1/100 \quad \lambda = 10$$

$$A = 10 \quad a = 2\pi$$





Electric energy density

$$\frac{d^2\phi}{dN^2} = - \left(3 - \frac{1}{2} \left(\frac{d\phi}{dN} \right)^2 \right) \left(\frac{1}{6} \frac{d\rho_E}{dN} \frac{d\phi}{dN} + \frac{d\phi}{dN} (V(\phi) + \rho_E) + \frac{1}{3} \left(\frac{d\phi}{dN} \right)^2 \rho_E \frac{f'(\phi)}{f(\phi)} + V'(\phi) - 2 \frac{f'(\phi)}{f(\phi)} \rho_E \right) \frac{1}{V(\phi) + \rho_E}$$

$$\frac{d\rho_E}{dN} + 4\rho_E + 2 \frac{df(\phi)}{dN} \frac{1}{f(\phi)} \rho_E =$$

$$\left(V(\phi) + \rho_E \right)^2 \frac{1}{4\pi^2} \left(1 + \left(\frac{df(\phi)}{dN} \frac{1}{f(\phi)} \right)^2 \right) \left(3 - \frac{1}{2} \left(\frac{d\phi}{dN} \right)^2 \right)^{-2}$$