

LHC powered by factorization: some resummation and some fixed-order QCD

Tobias Neumann, BNL

April 22nd, 2021

Factorization

Loops
& Legs

RGE

PDF's,

α_s

EFT

Collider pheno

....

N3LL qT resummation for
color singlets

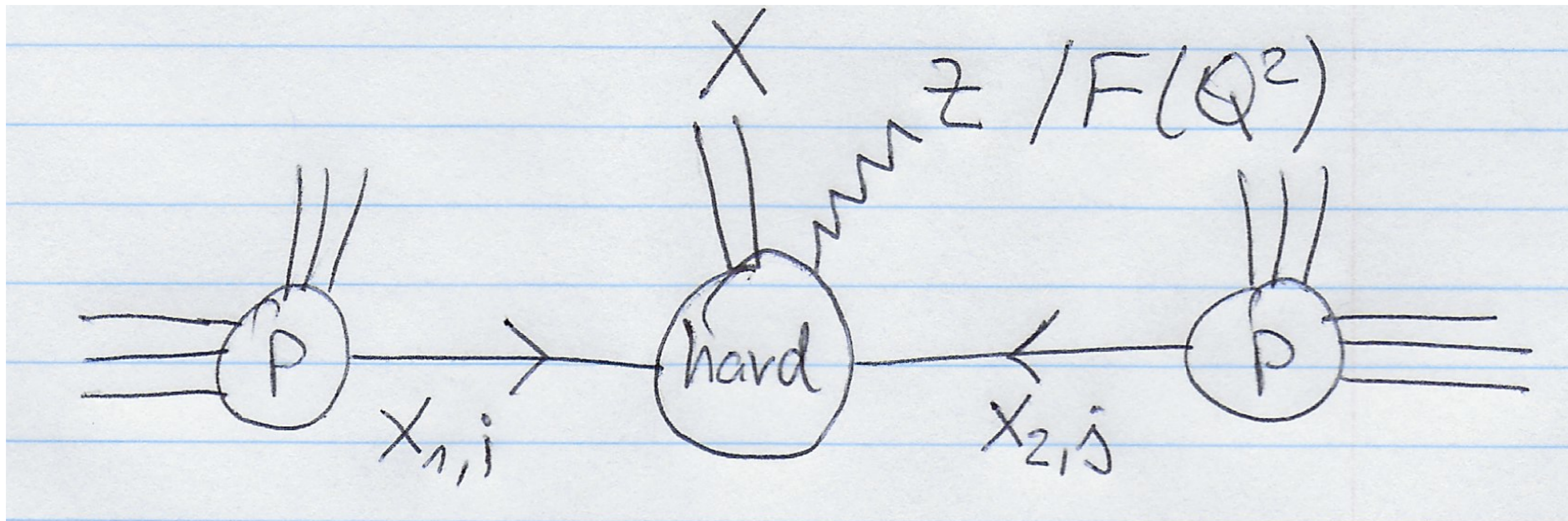
....

NNLO t-channel
single top-quark

....

Factorization

$$\sigma_{i,j} = f_i(x_i, \mu_F) \otimes f_j(x_j, \mu_F) \otimes \sigma_{i,j}^H(\mu_R, \mu_F, \vec{p}) + \mathcal{O}(\Lambda_{\text{QCD}}$$



$$\sigma_{i,j} = \int d\mathbf{x}_1 d\mathbf{x}_2 f_a(\mathbf{x}_i, \mu_F) f_b(\mathbf{x}_j, \mu_F) \cdot \sigma_{i,j}^H(\mu_R, \mu_F, \vec{p})$$

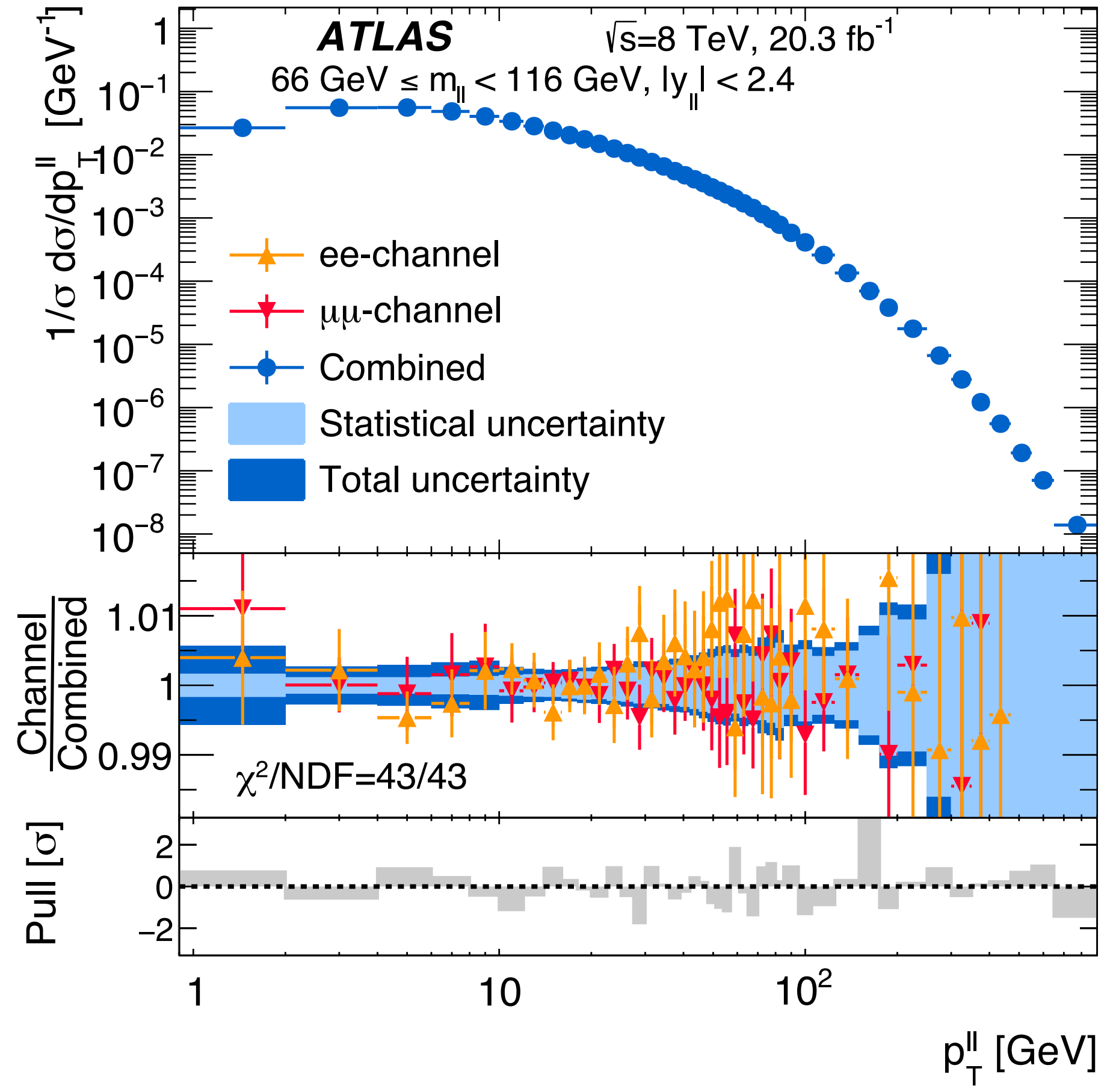
transverse dependence

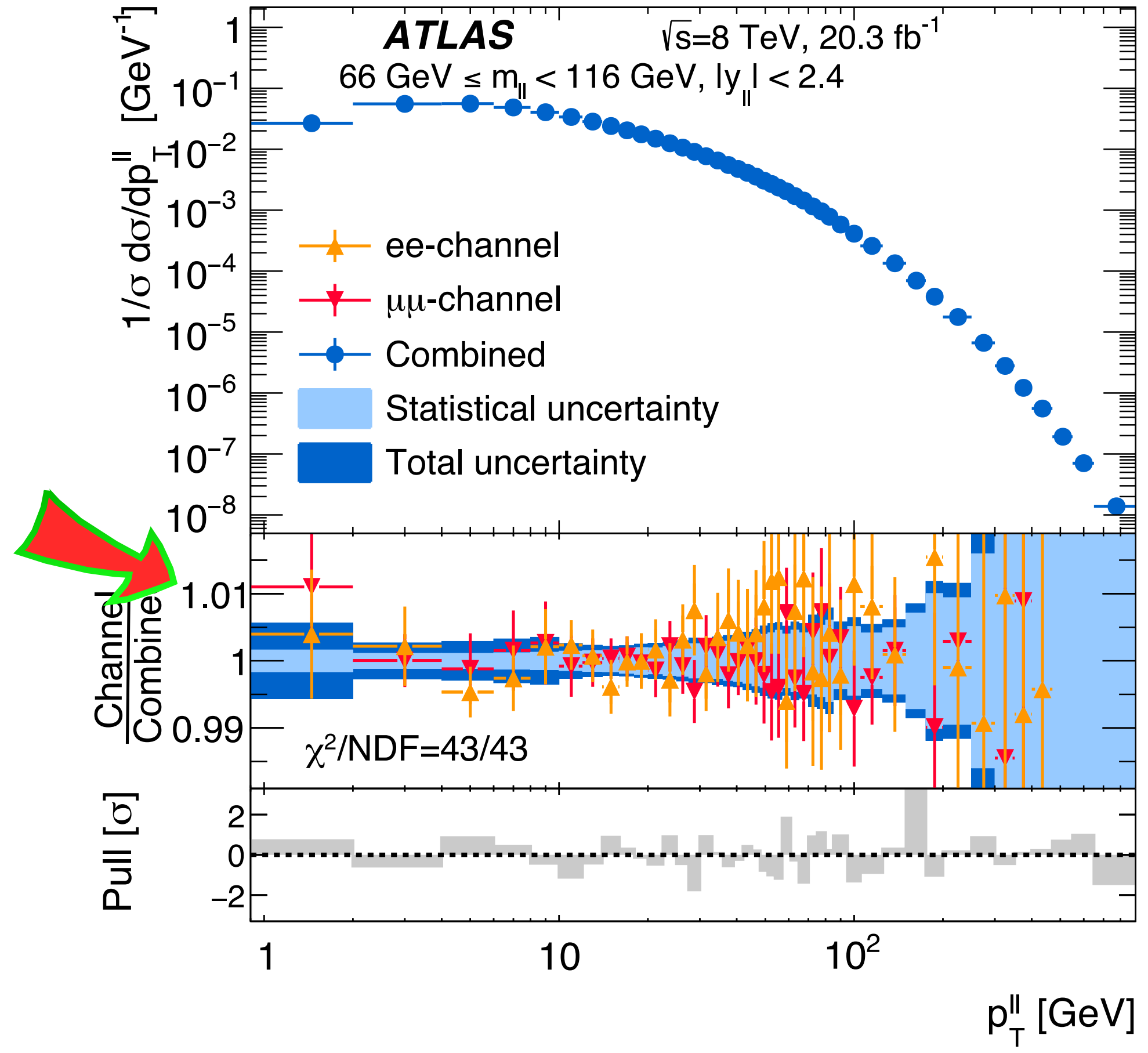
$$\sigma_{i,j} = \int dx_1 dx_2 f_i(x_i, \mu_F) f_j(x_j, \mu_F) \cdot \sigma_{i,j}^H(\mu_R, \mu_F, \vec{p})$$

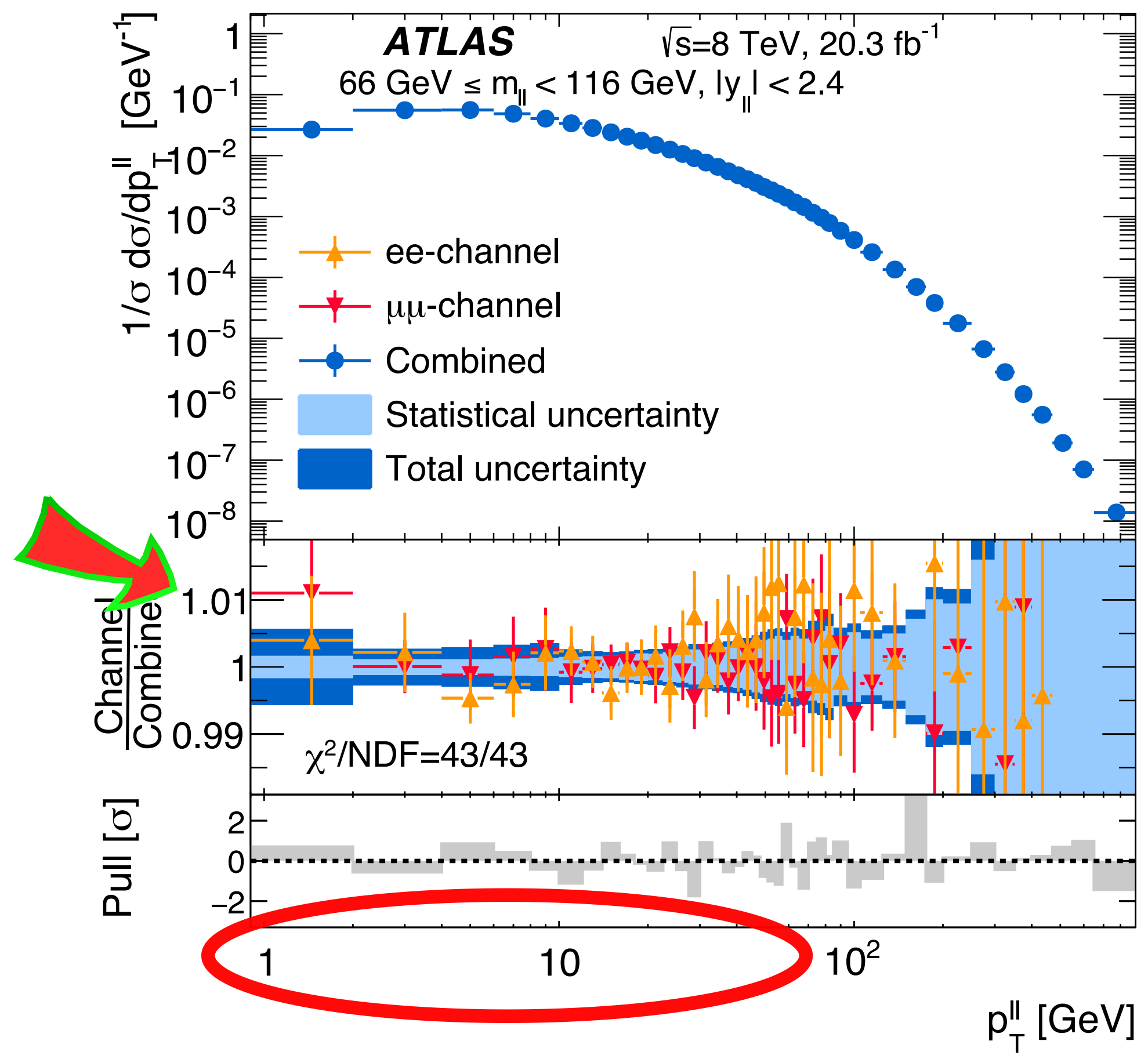
transverse dependence

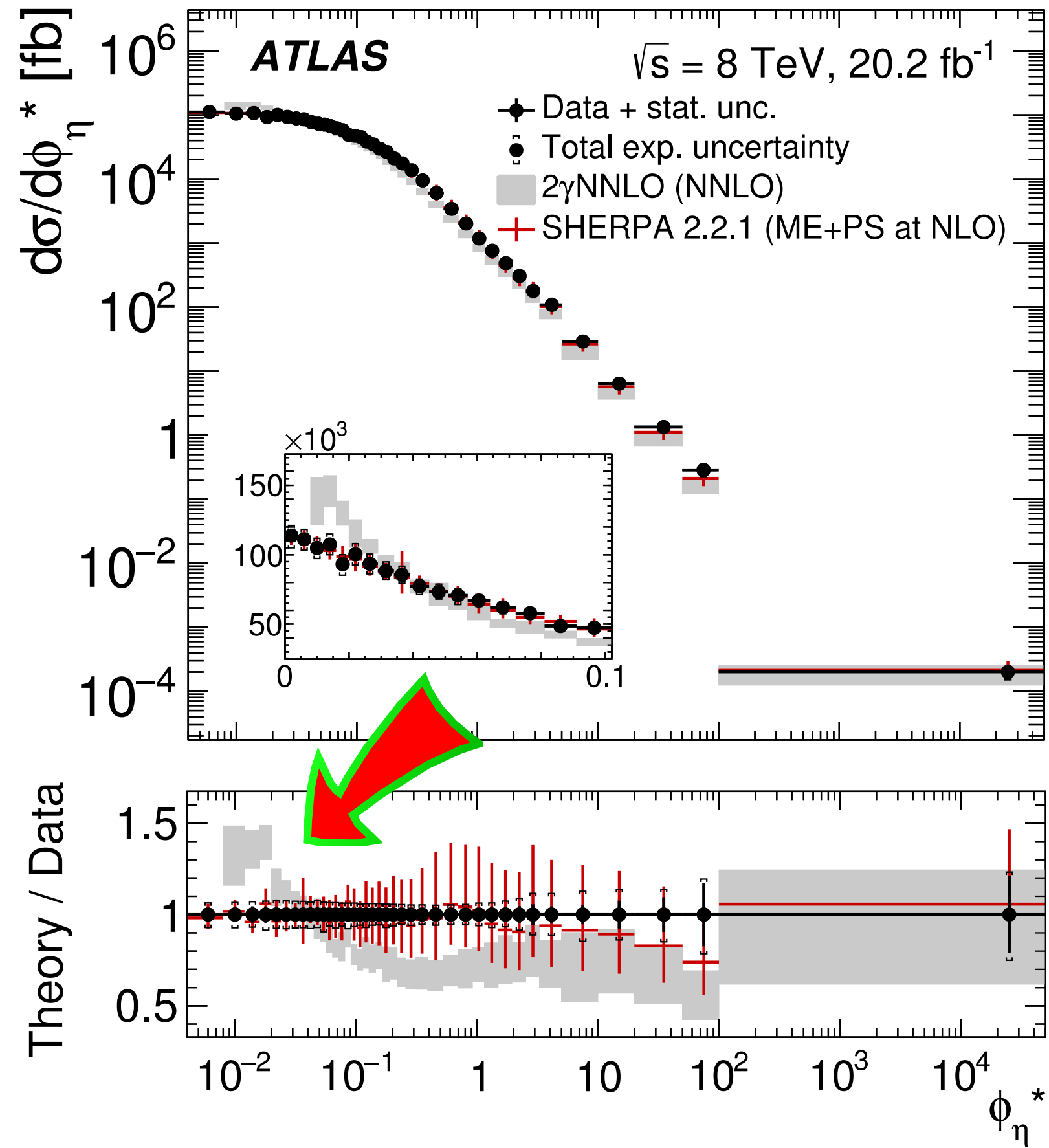
spin dependence

Why q_T resummation?

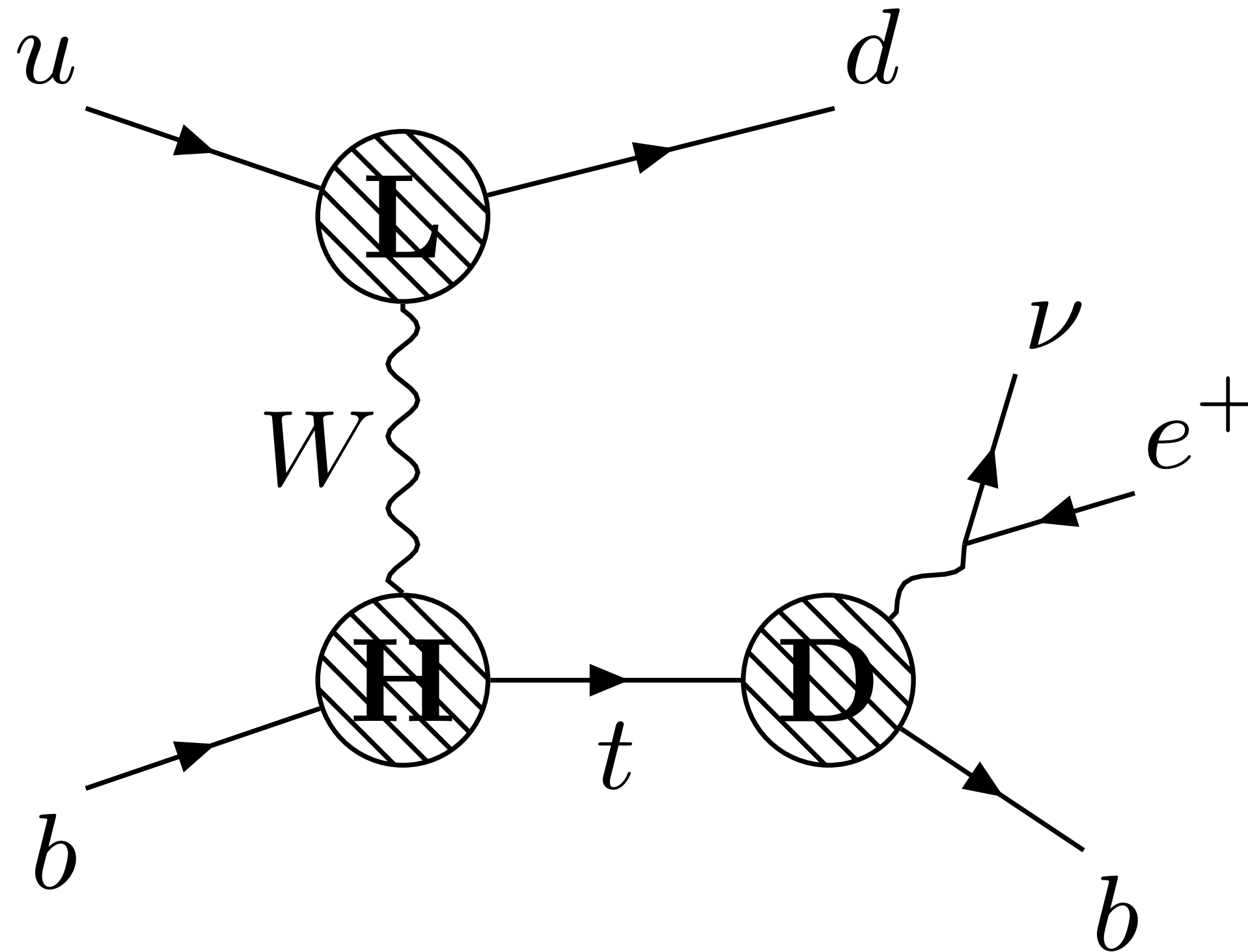








Why t-channel single-top-quark?



// We found a difference of $\sim 1\%$ on the NNLO cross sections. [...] It has not been possible to further pin down the differences //

— Berger, Gao, Zhu '07

$$\sigma_{i,j} = \int dx_1 dx_2 f_i(\mathbf{x}_i, \mu_F) f_j(\mathbf{x}_j, \mu_F) \cdot \sigma_{i,j}^H(\mu_R, \mu_F, \vec{p})$$

$$\mathbf{f}_i(\mathbf{x}_i, \mu_F) \mathbf{f}_j(\mathbf{x}_j, \mu_F)$$


Part 1



***q_T* resummation of color singlets at N³LL+NNLO**

q_T factorization from SCET

$$\begin{aligned} d\sigma_{ij}(p_1, p_2, \{\underline{q}\}) &= \int_0^1 d\xi_1 \int_0^1 d\xi_2 d\sigma_{ij}^0(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}) \mathcal{H}_{ij}(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}, \mu) \cdot \\ &\quad \frac{1}{4\pi} \int d^2 x_\perp e^{-i q_\perp x_\perp} \left(\frac{x_T^2 Q^2}{b_0^2} \right)^{-F_{ij}(x_\perp, \mu)} B_i(\xi_1, x_\perp, \mu) \cdot B_j(\xi_2, x_\perp, \mu) \end{aligned}$$

(Becher, Neubert '10; Becher, Neubert, Wilhelm '11 '12), (Mantry, Petriello '09)

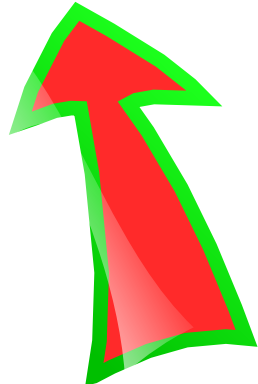
$$\begin{aligned}
d\sigma_{ij}(p_1, p_2, \{\underline{q}\}) &= \int_0^1 d\xi_1 \int_0^1 d\xi_2 d\sigma_{ij}^0(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}) \mathcal{H}_{ij}(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}, \mu) \cdot \\
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\end{aligned}$$

Collinear anomaly

Becher, Neubert '10, 11

$$\begin{aligned} d\sigma_{ij}(p_1, p_2, \{\underline{q}\}) &= \int_0^1 d\xi_1 \int_0^1 d\xi_2 d\sigma_{ij}^0(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}) \mathcal{H}_{ij}(\xi_1 p_1, \xi_2 p_2, \{\underline{q}\}, \mu) \cdot \\ &\frac{1}{4\pi} \int d^2 x_{\perp} e^{-iq_{\perp} x_{\perp}} \left(\frac{x_T^2 Q^2}{b_0^2} \right)^{-F_{ij}(x_{\perp}, \mu)} B_i(\xi_1, x_{\perp}, \mu) \cdot B_j(\xi_2, x_{\perp}, \mu) \end{aligned}$$


Rapidity anomalous dimension

$$\sigma \sim H \otimes B \otimes B + \mathcal{O}(q_T/Q)$$

$$\frac{d}{d \ln \mu} C_V(-q^2, \mu) = \left[\Gamma_{\text{cusp}}^F(a_s) \ln \frac{-q^2}{\mu^2} + 2\gamma^q(a_s) \right] C_V(-q^2, \mu), \quad (5)$$

$$\frac{d}{d \ln \mu} F_{q\bar{q}}(L_\perp, a_s) = 2\Gamma_{\text{cusp}}^F(a_s) \quad (6)$$

$$\frac{d}{d \ln \mu} I_{q\leftarrow i}(z, L_\perp, a_s) = \left[\Gamma_{\text{cusp}}^F(a_s) L_\perp - 2\gamma^q(a_s) \right] I_{q\leftarrow i}(z, L_\perp, a_s) - \sum_j \int_z^1 \frac{du}{u} I_{q\leftarrow j}(u, L_\perp, a_s) \mathcal{P}_{j\leftarrow i}(z/u, a_s), \quad (7)$$

(Becher, Neubert '06 '10 '11)

Three and four loop ingredients have been published in the past 2-3 years.

Thank you MPI (Henn, Korchemsky, Mistlberger '19; Manteuffel, Panzer, Schabinger '20)
(also for NNLO beamfunctions)

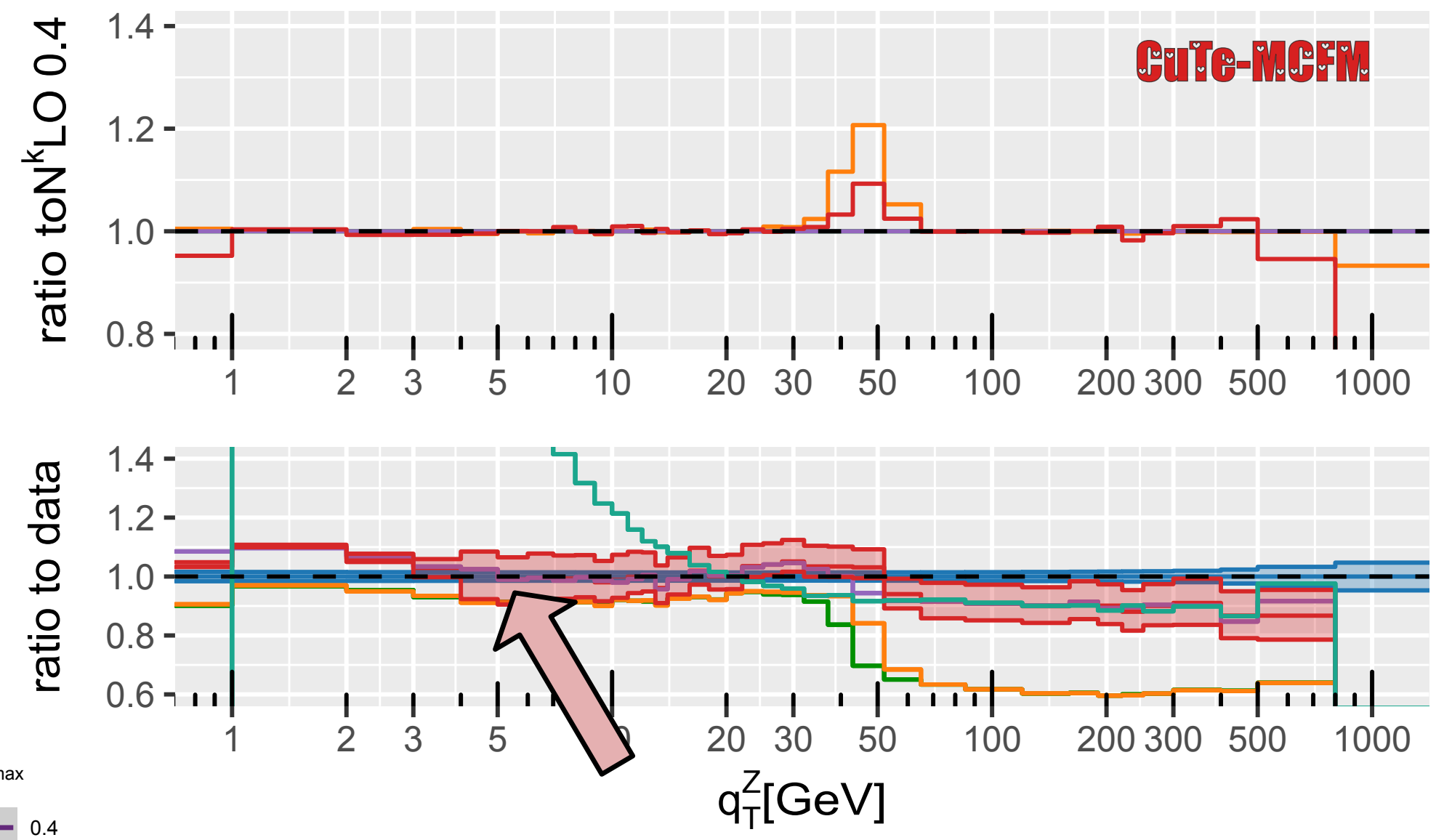
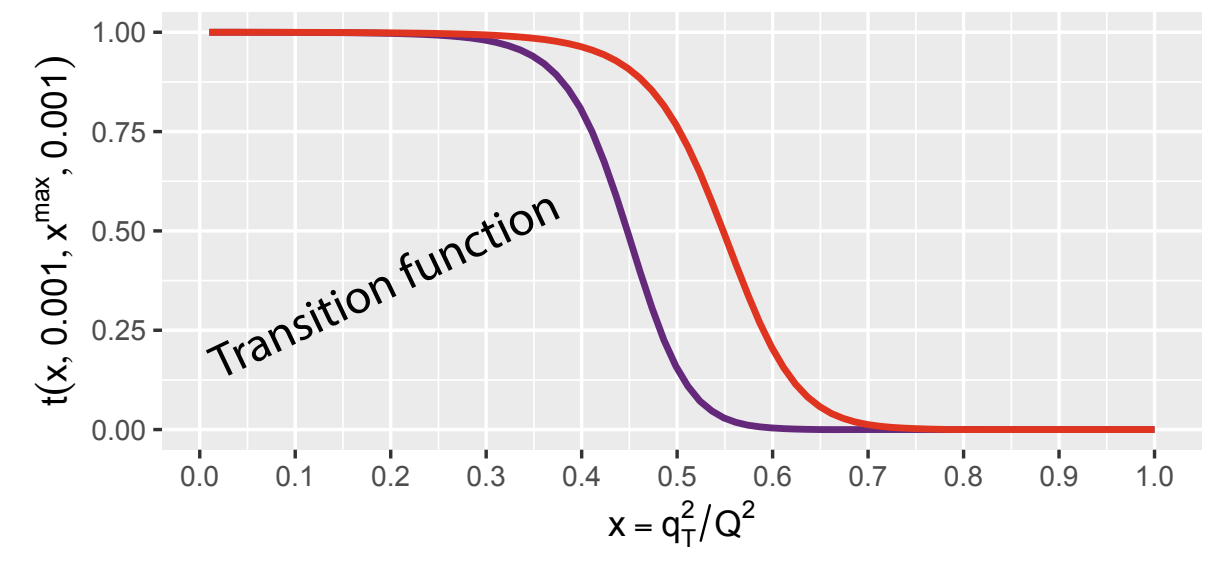
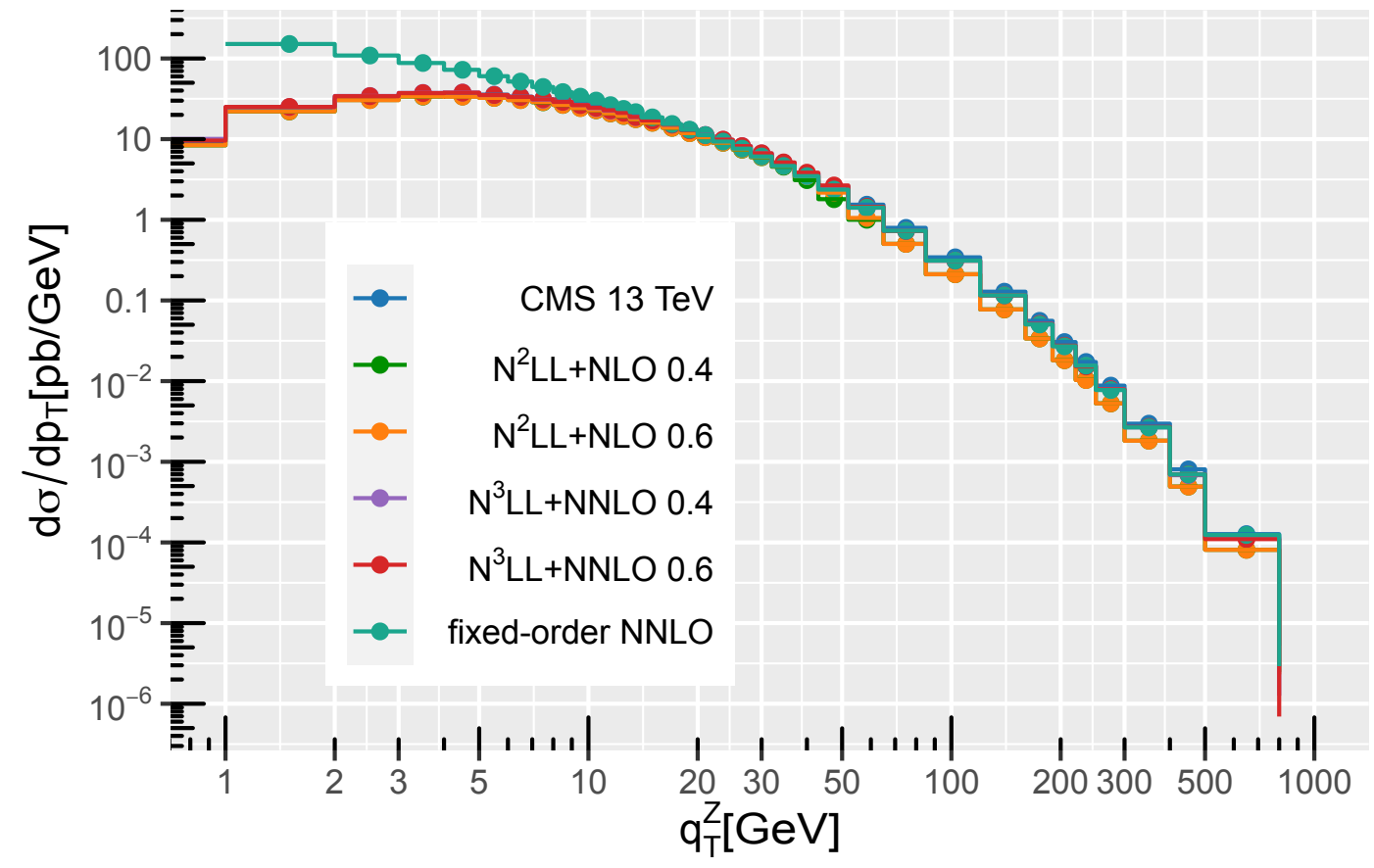
CUTE-MCFM

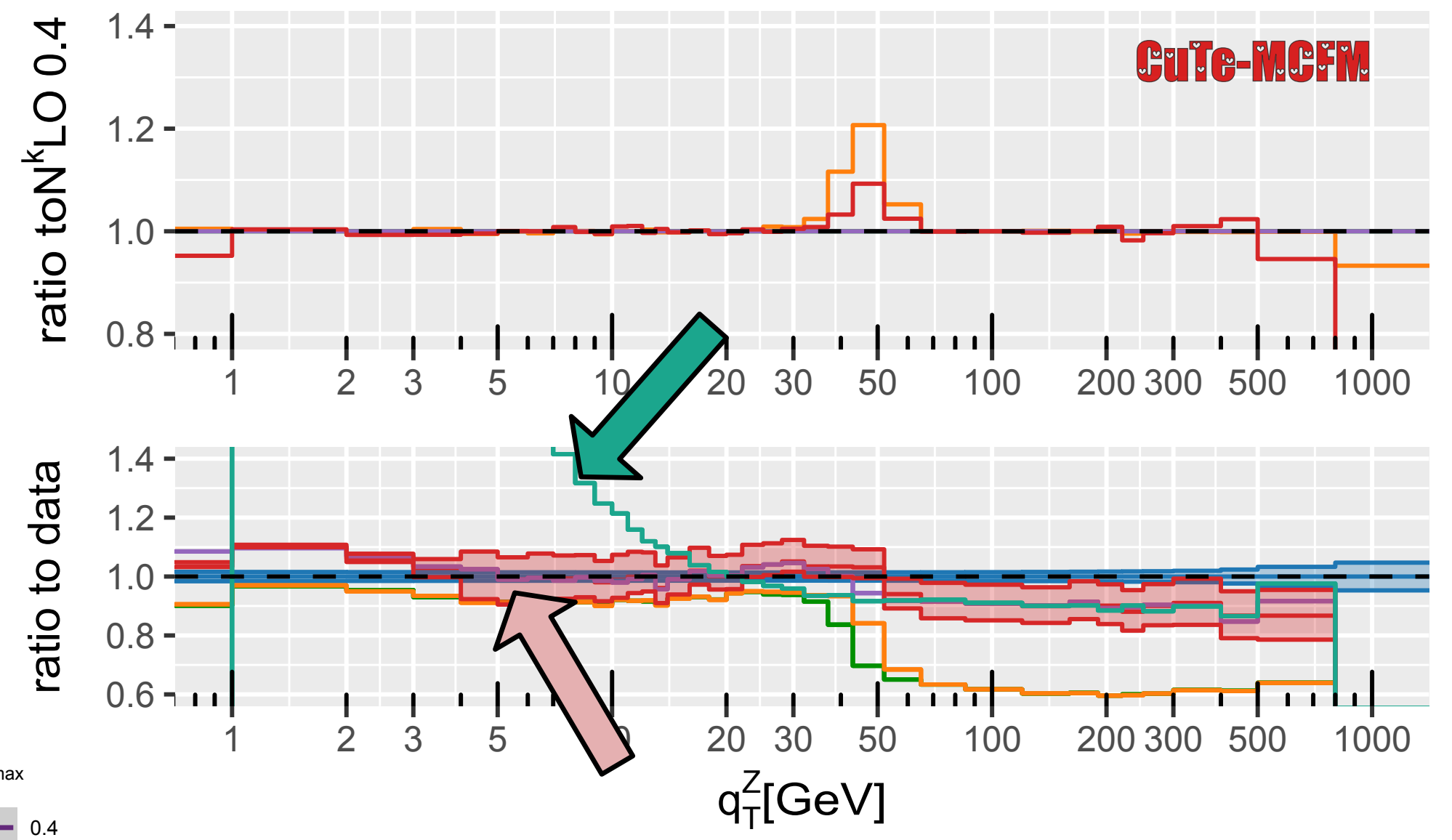
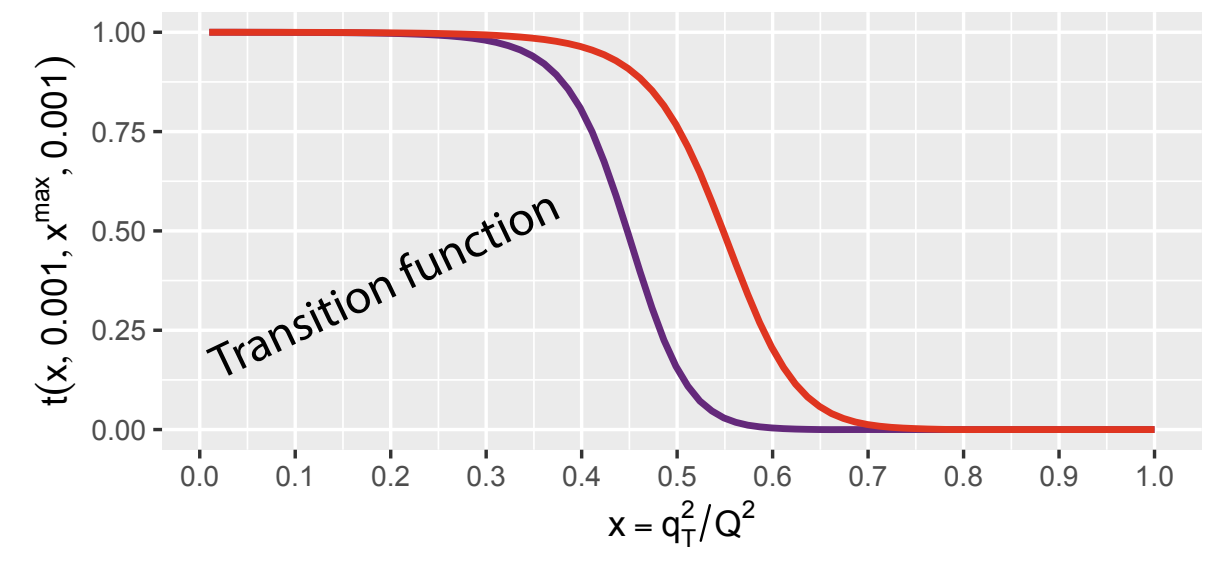
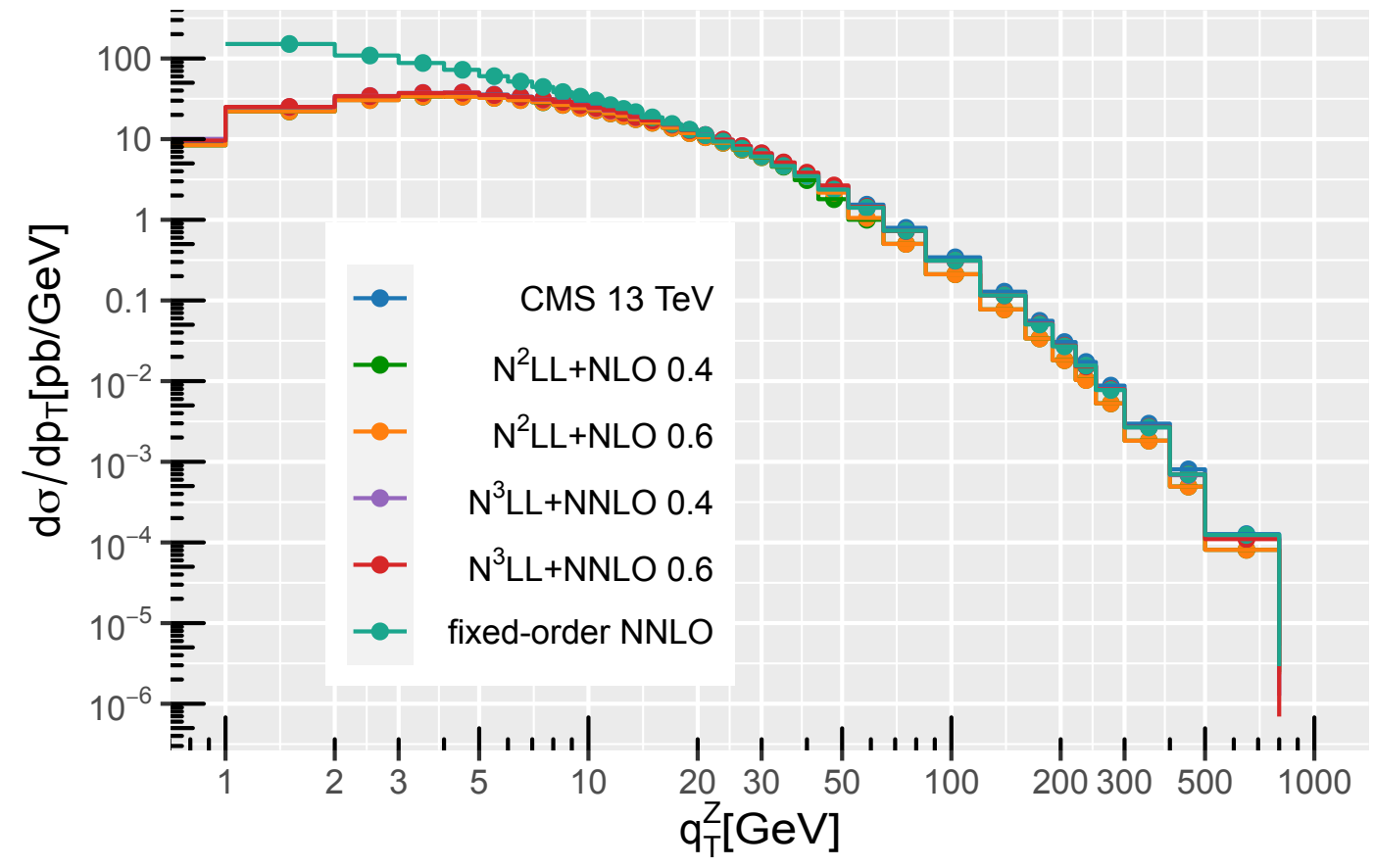
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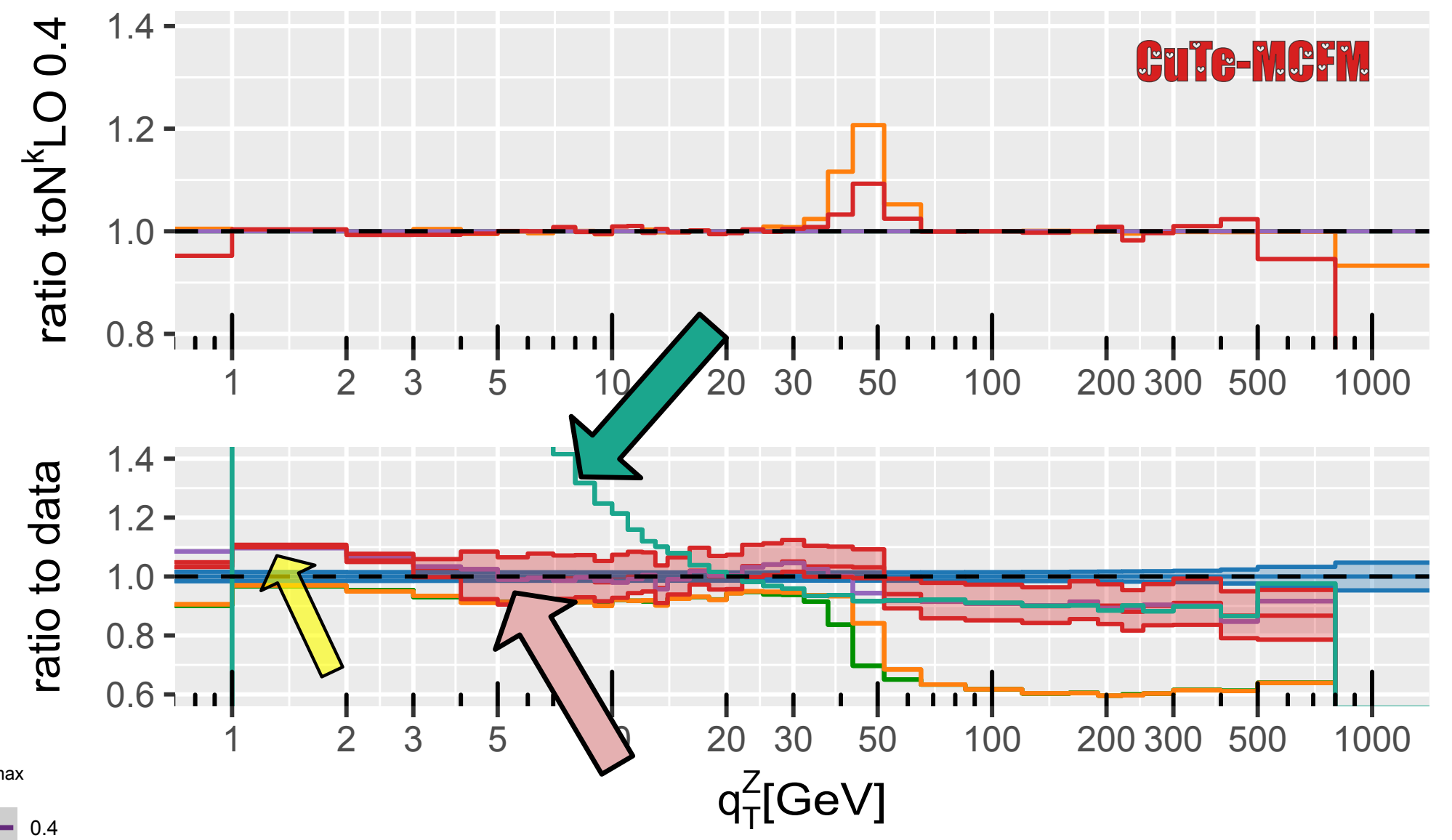
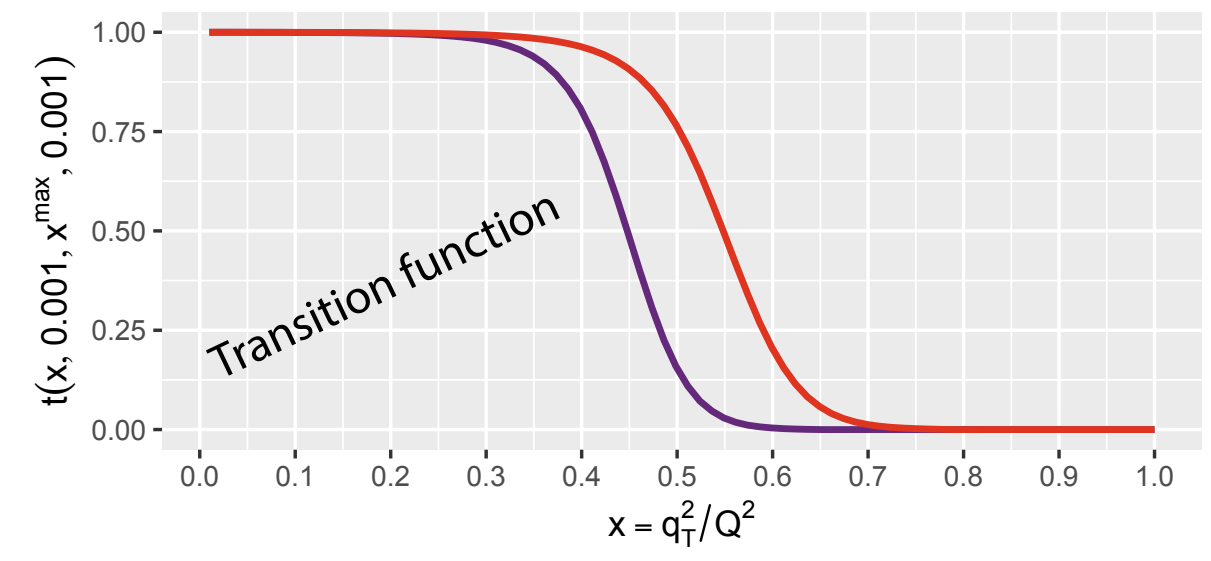
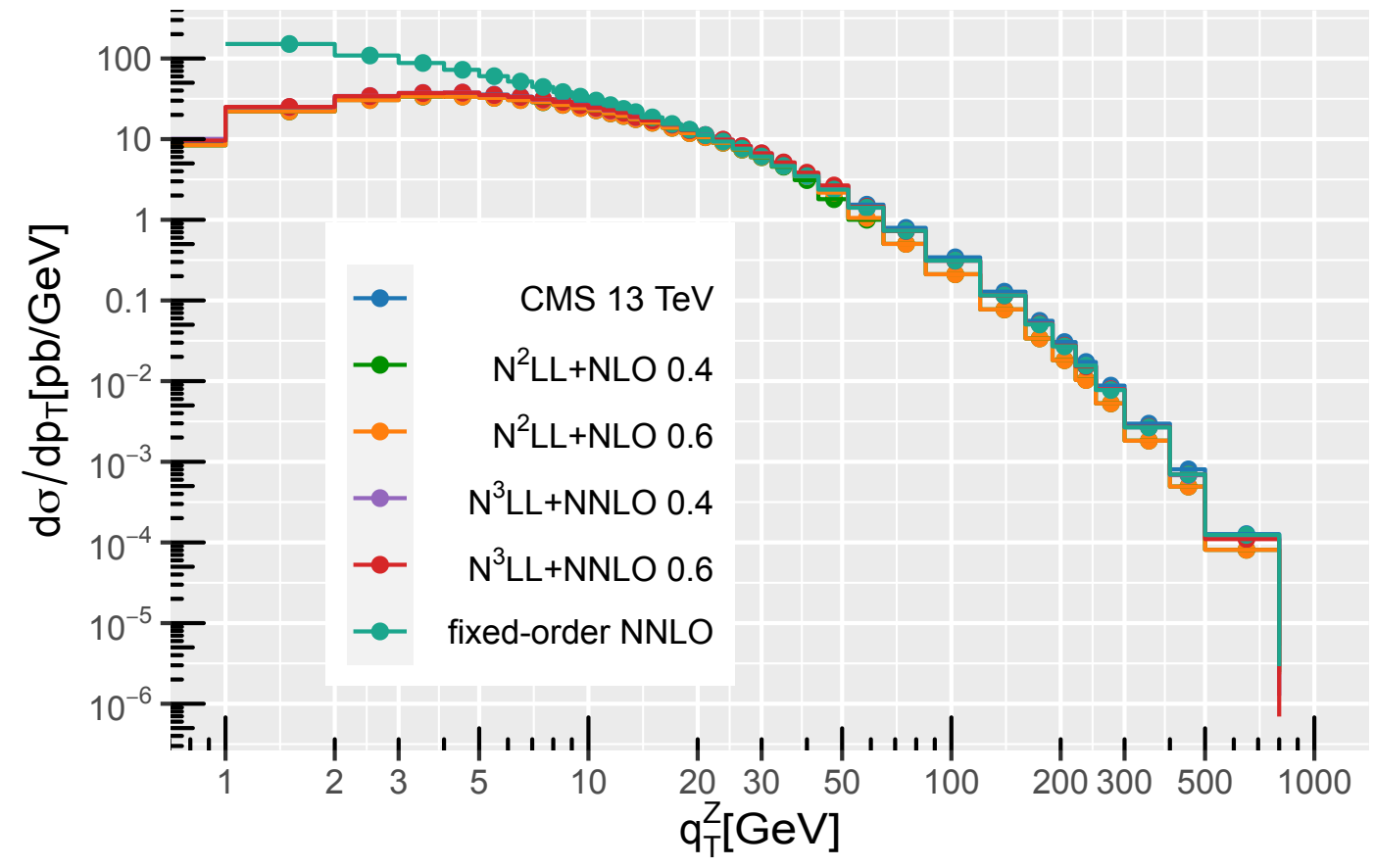
fixed-order MCFM + SCET-based N^3 LL resummation framework

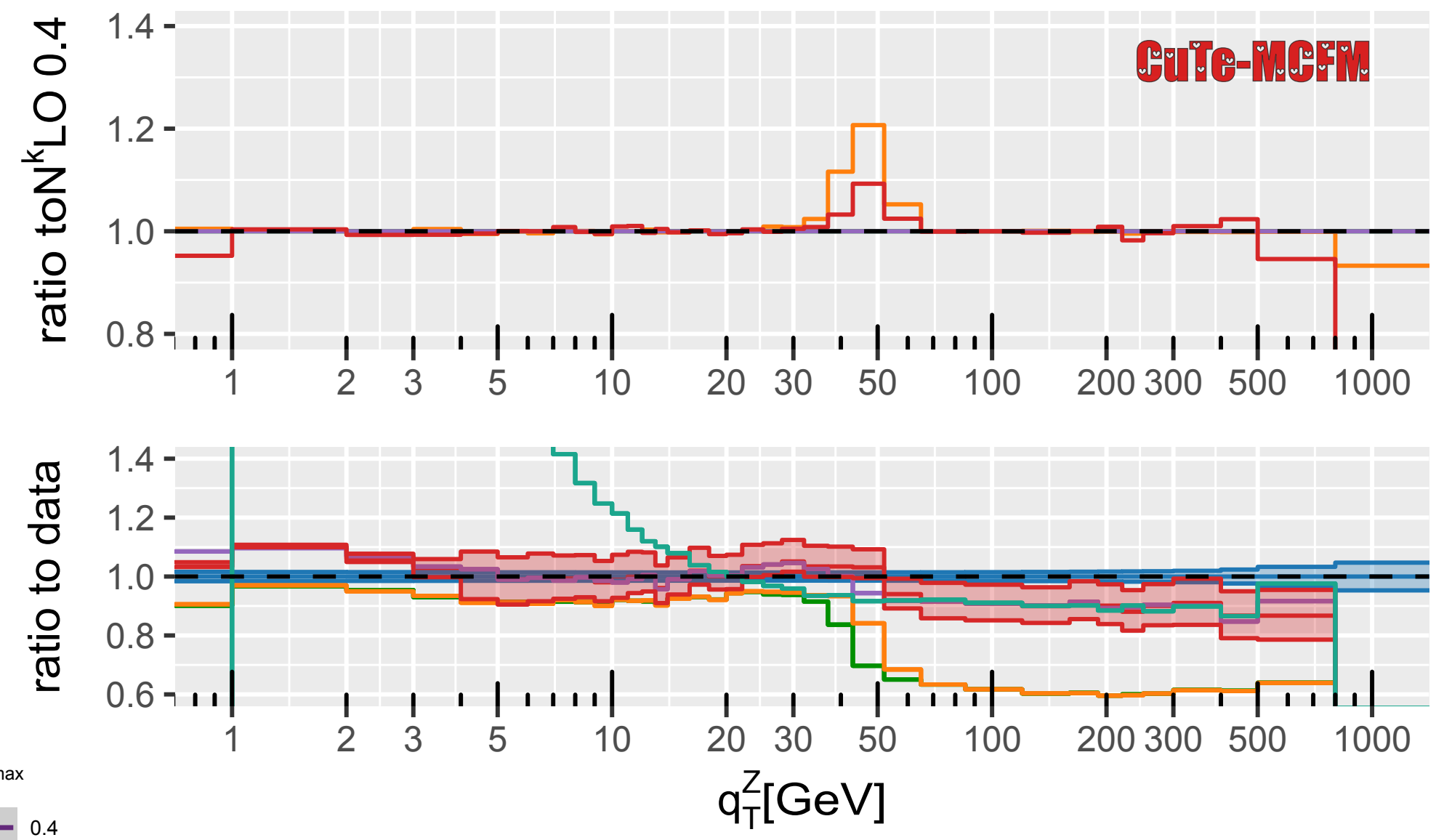
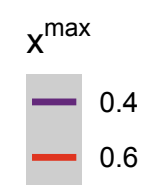
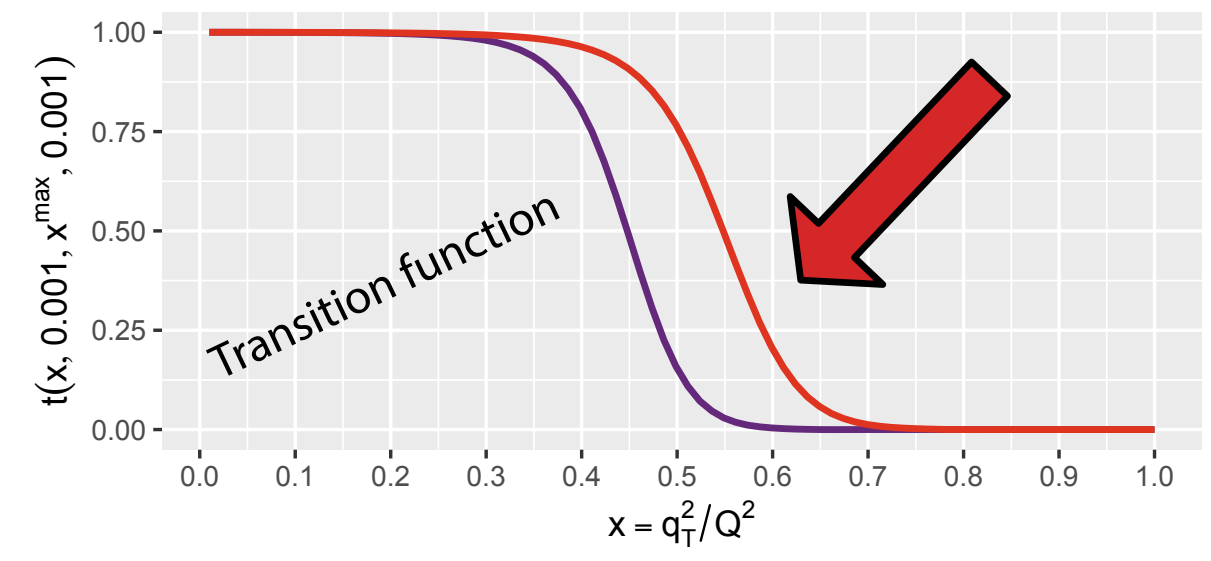
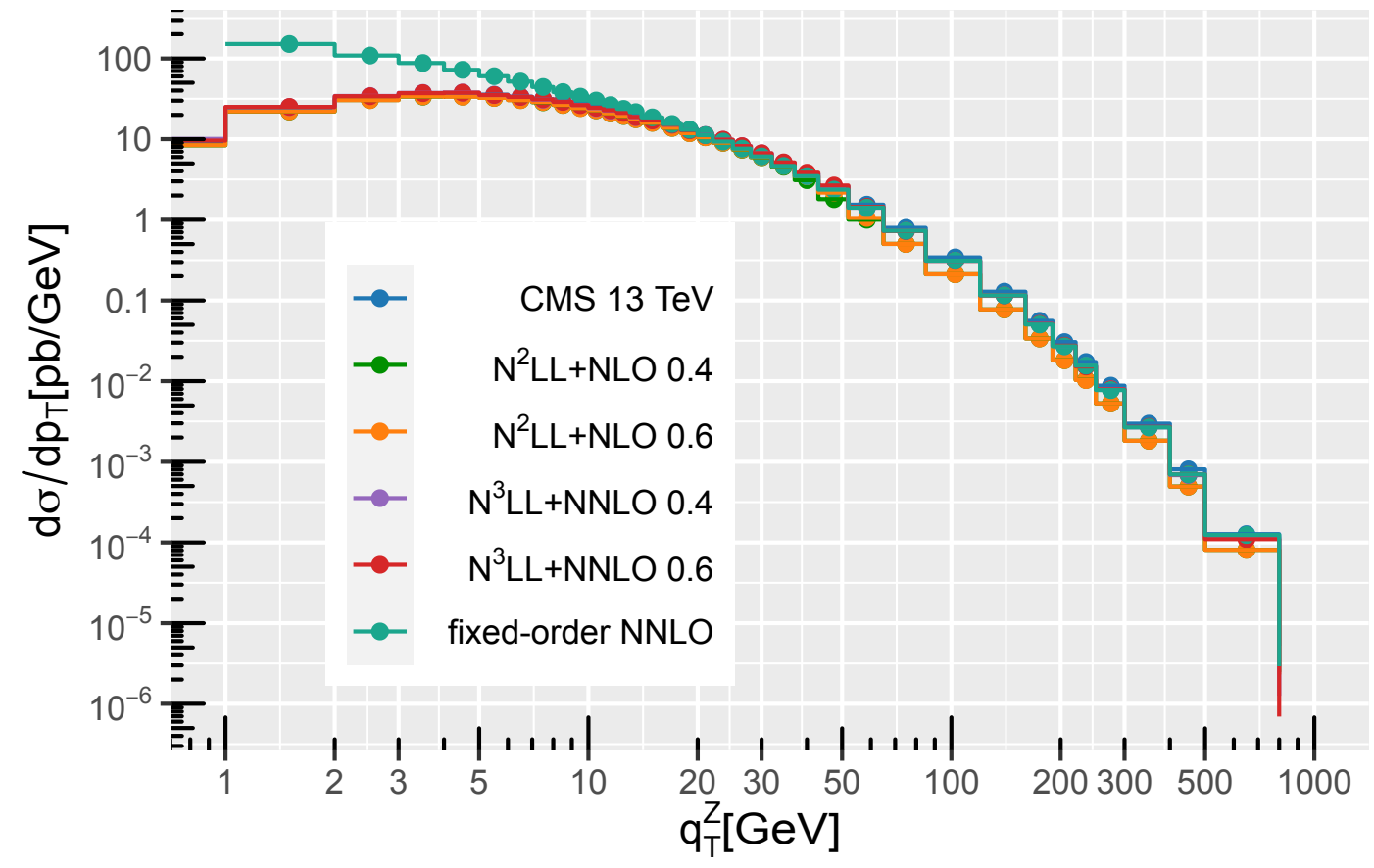
$W^\pm, Z, H, \gamma\gamma, Z\gamma, ZH, W^\pm H$
@ N^3 LL+NNLO

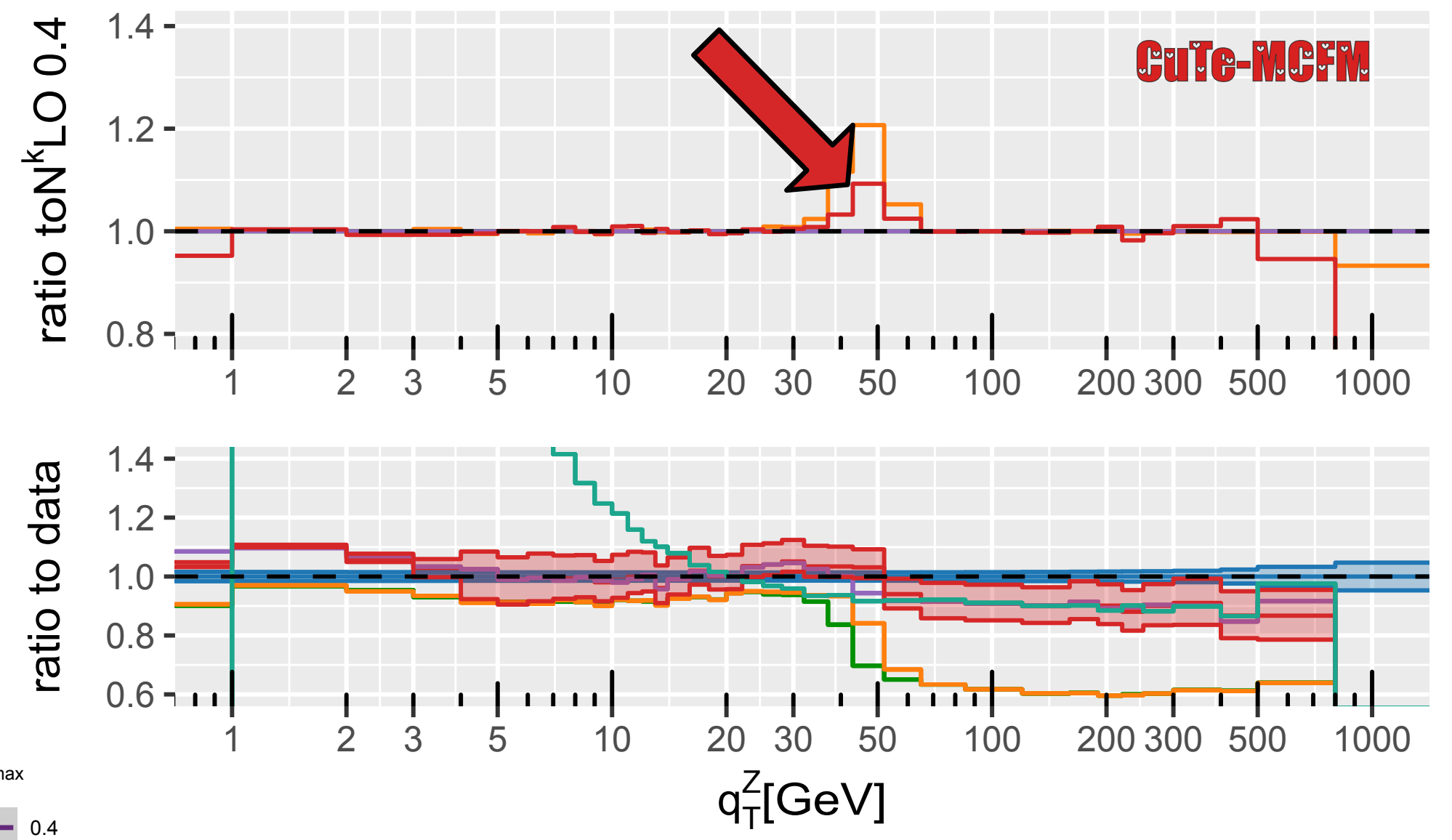
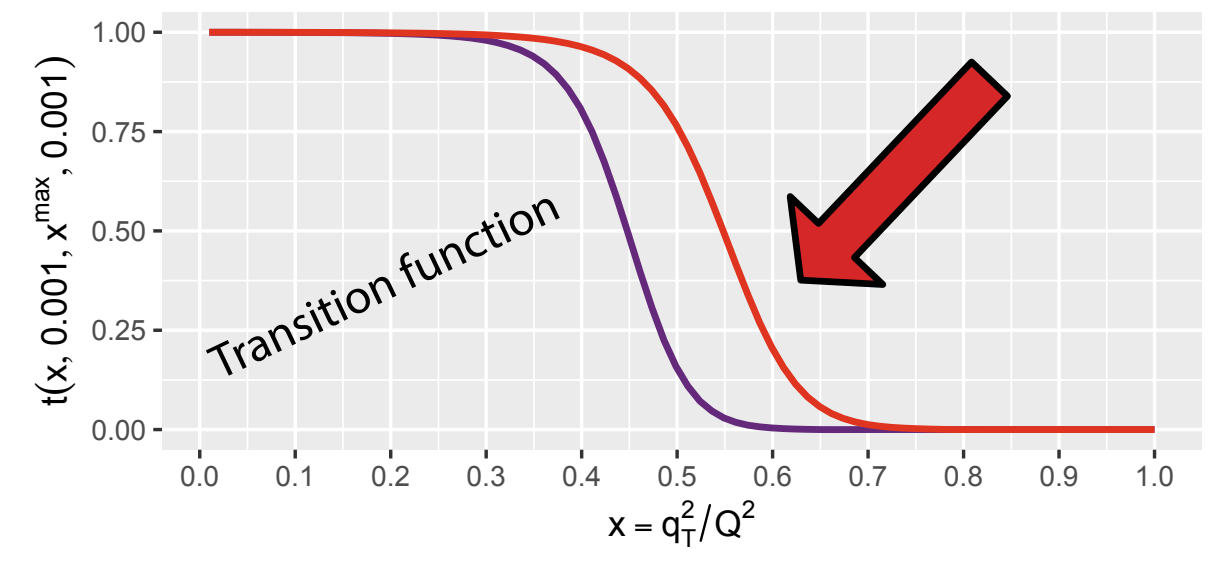
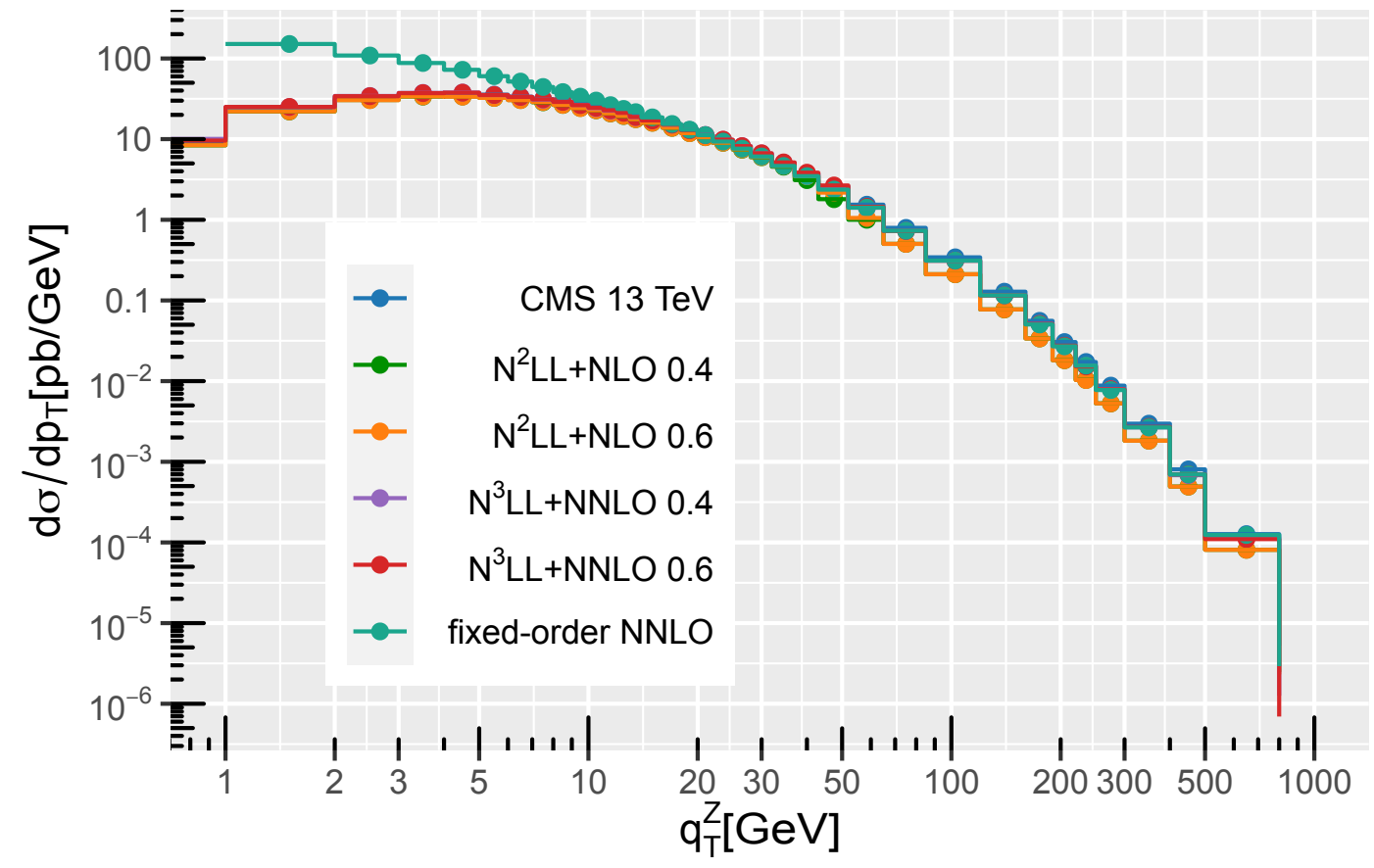
with Thomas Becher

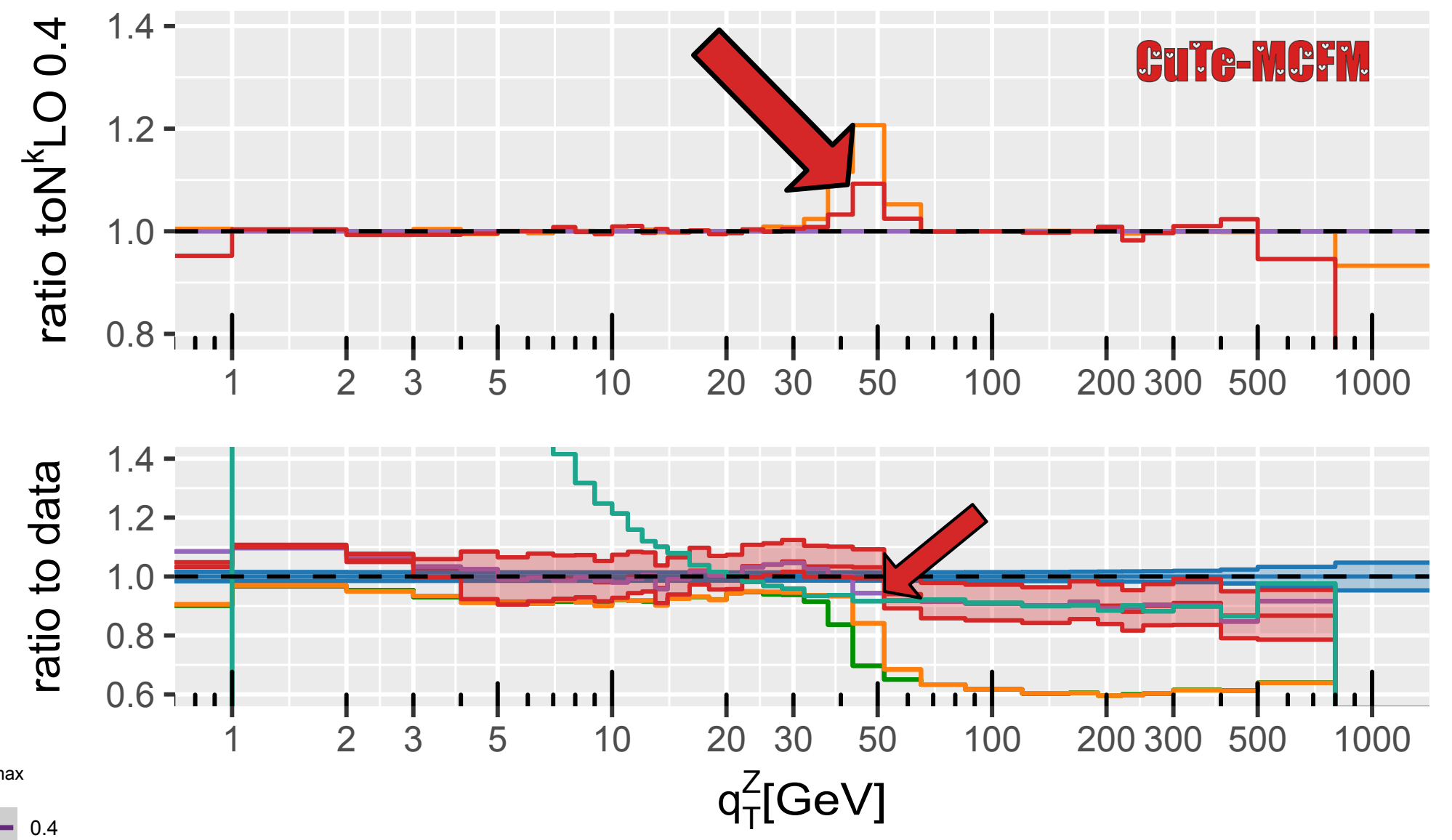
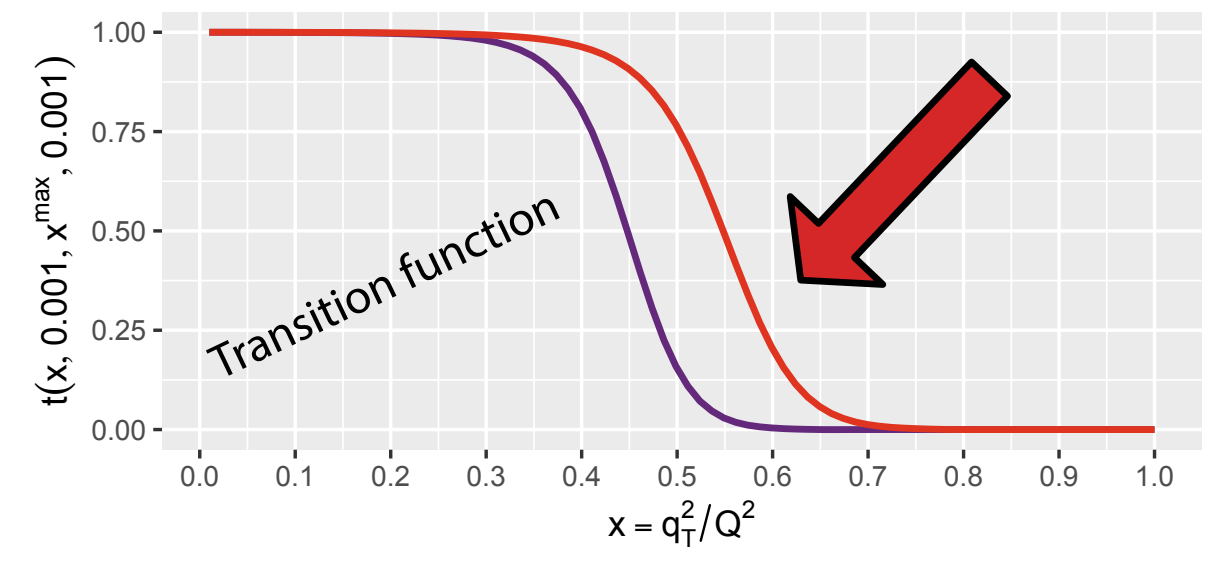
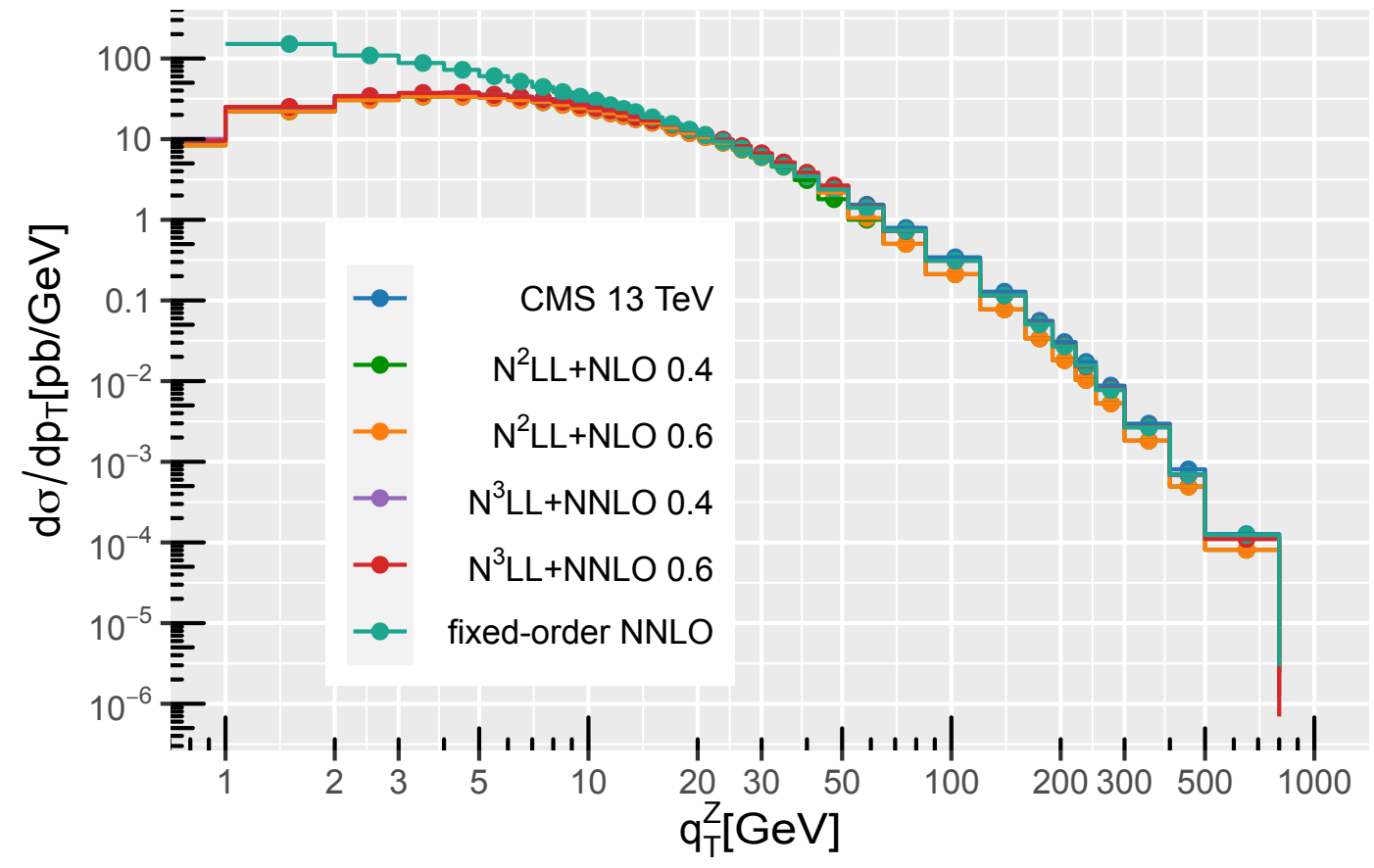


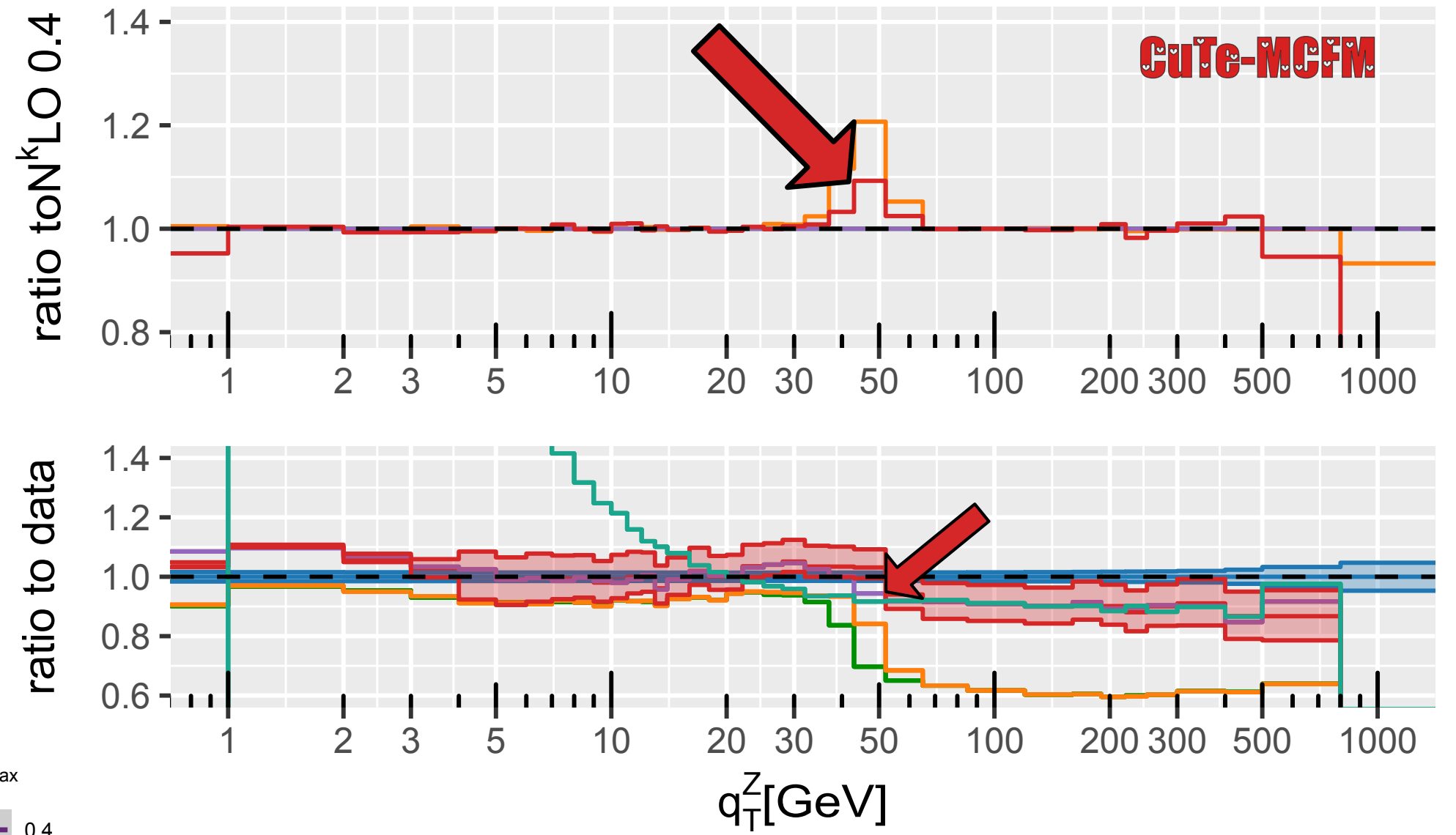
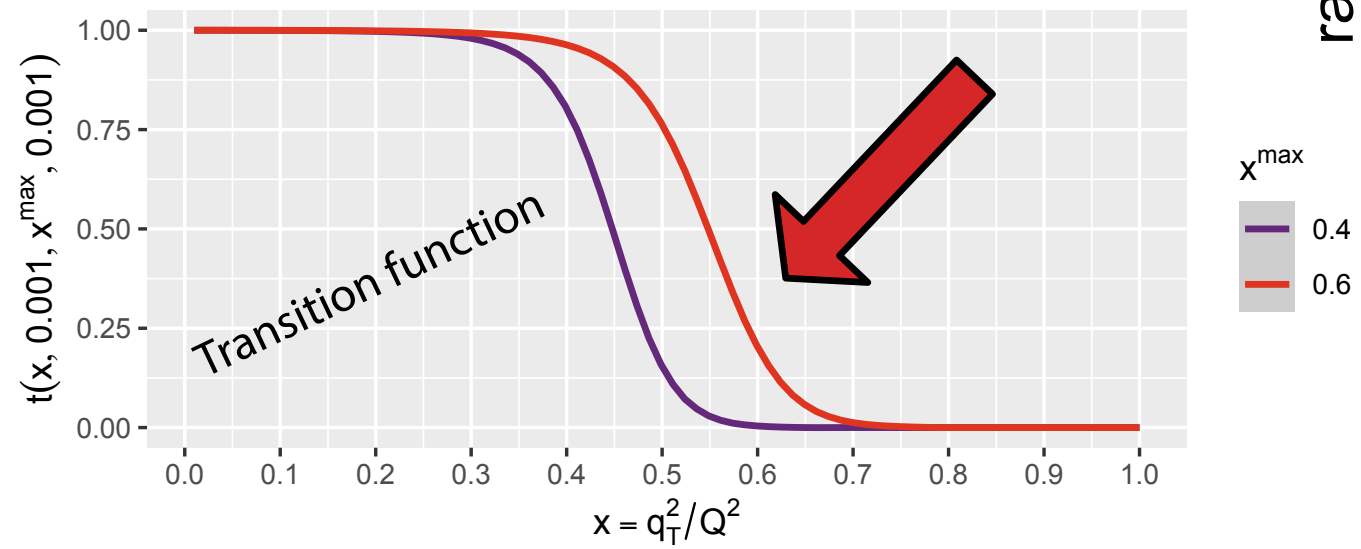
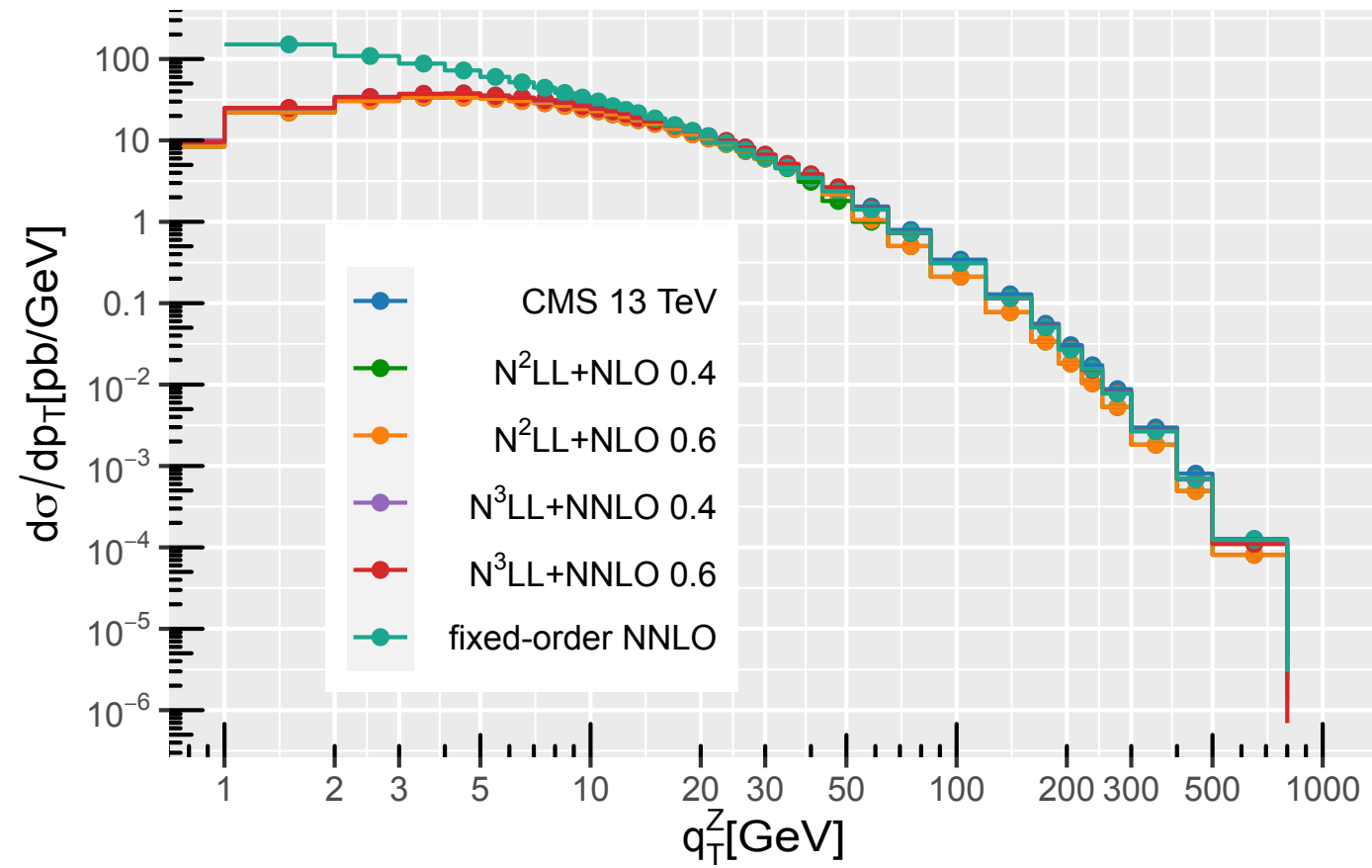






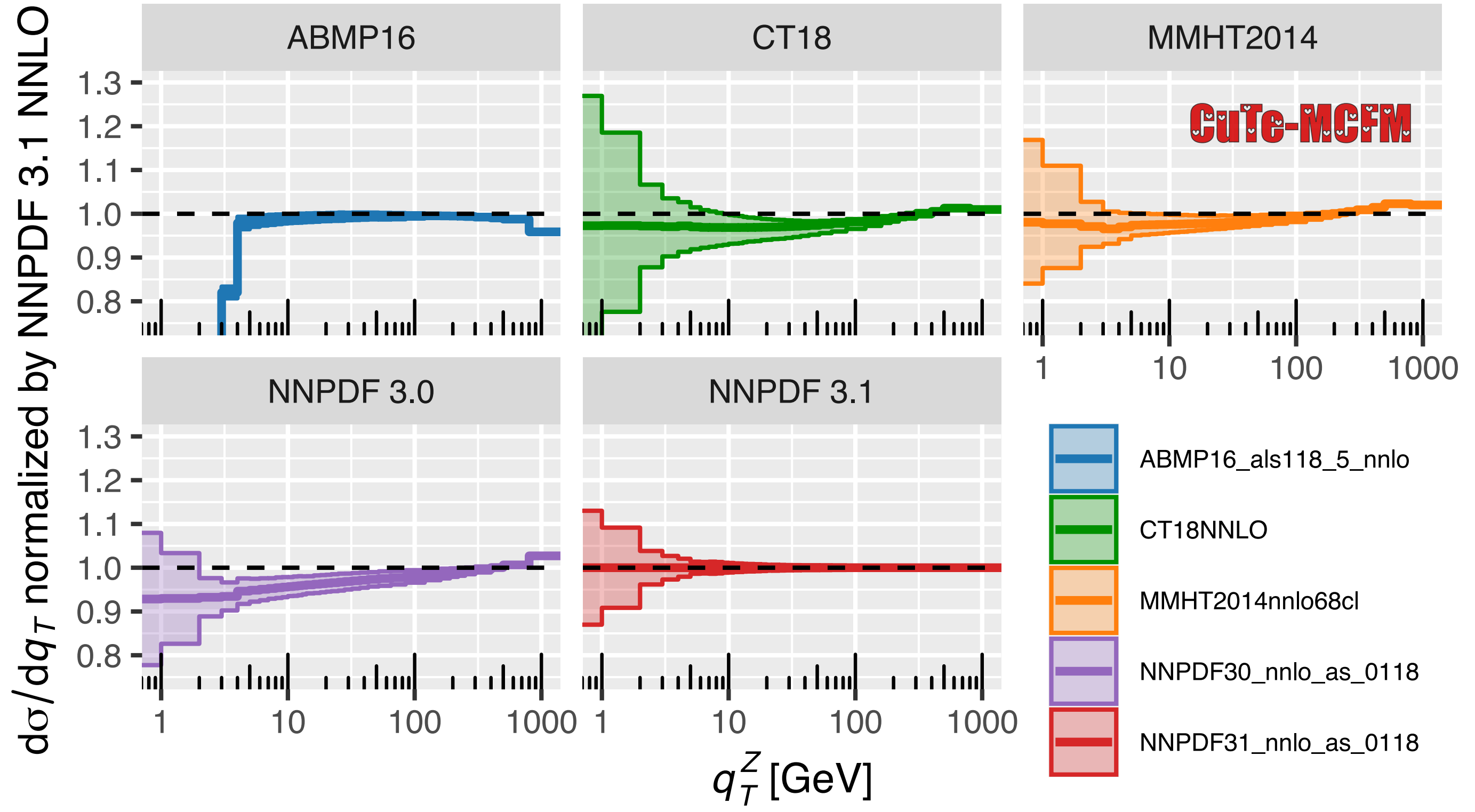






Recoil prescription takes into account $\mathcal{O}(q_T/Q)$ power corrections

(Ebert, Michel, Stewart, Tackmann '20)

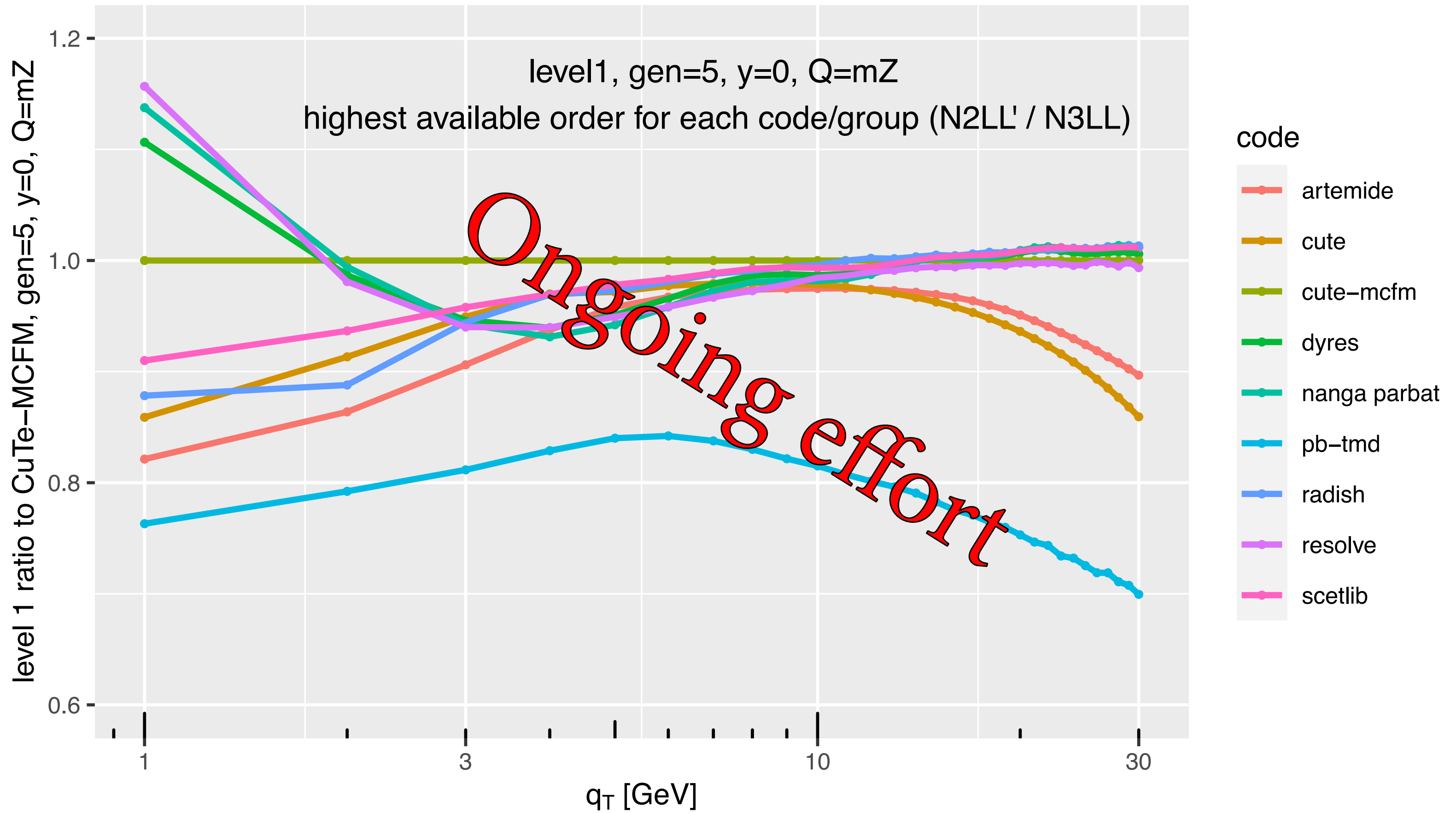


CuTe-MCFM is not alone

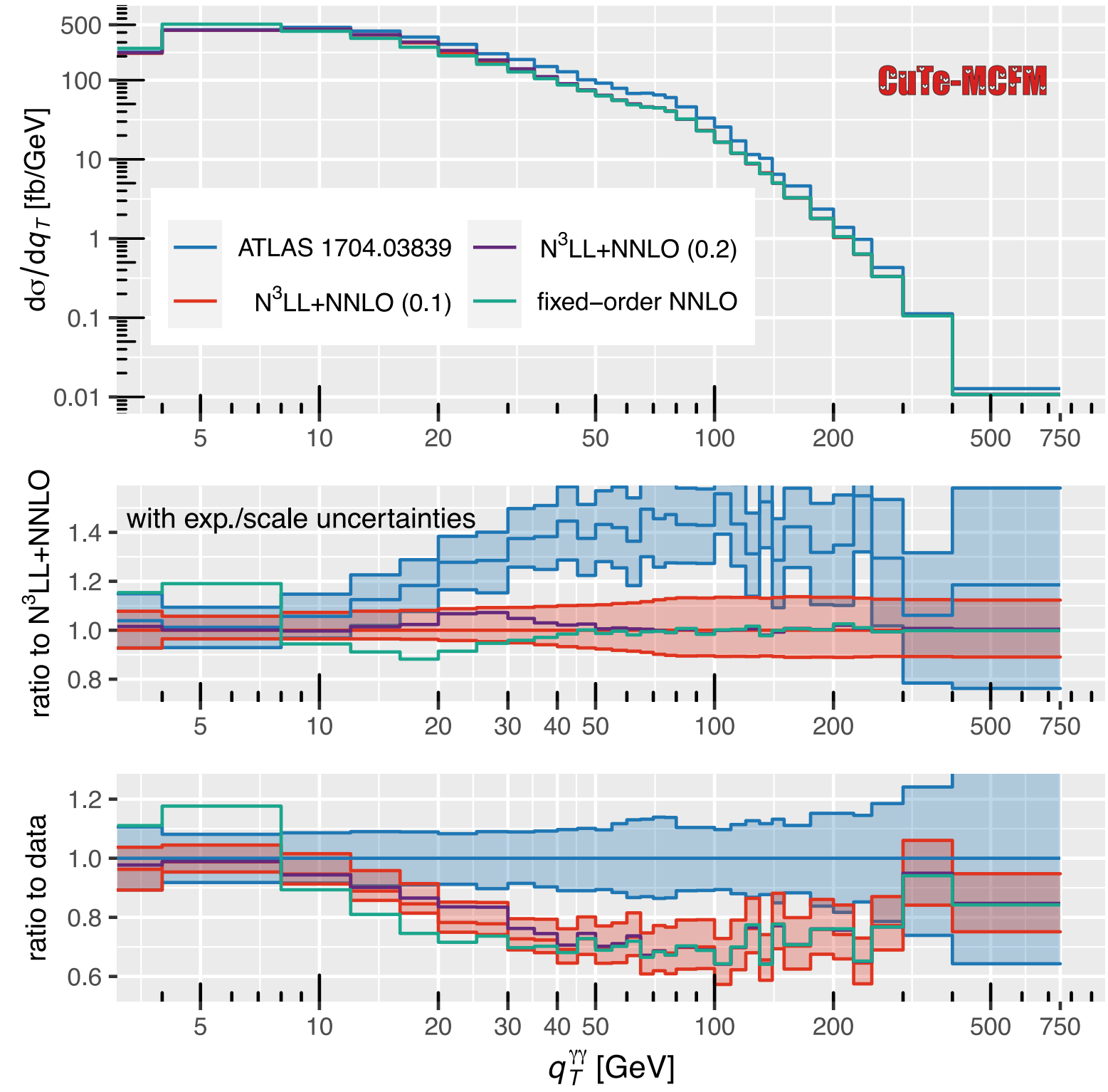
Lots of public and private codes and implementations...

different orders, formalisms, matching prescriptions, processes, limitations ...

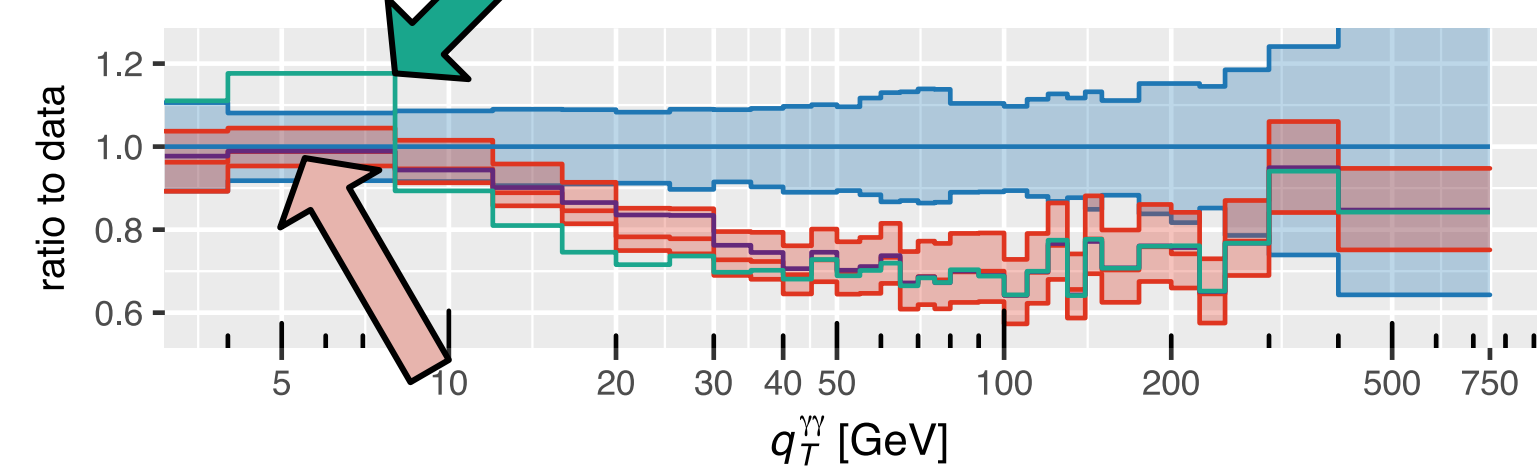
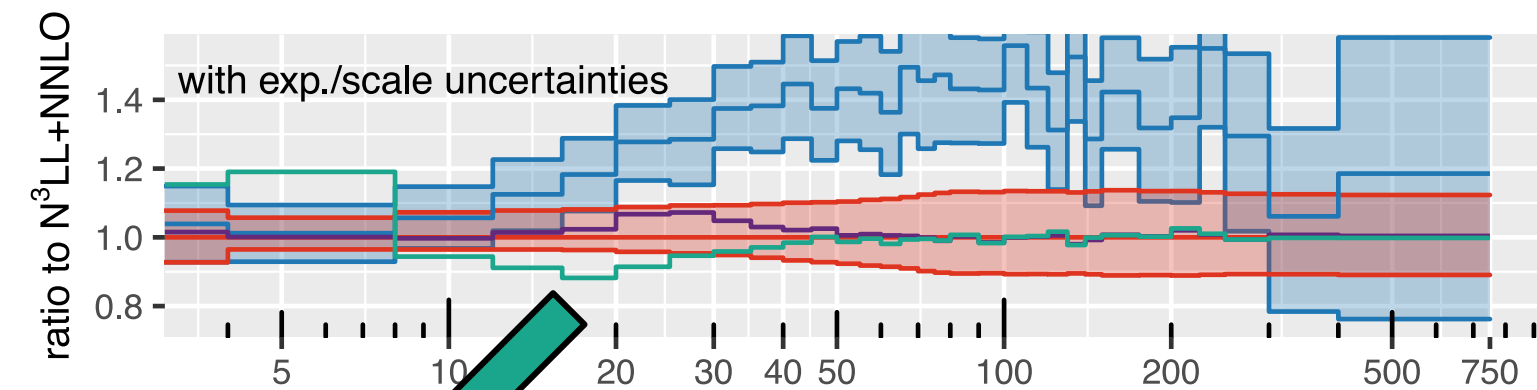
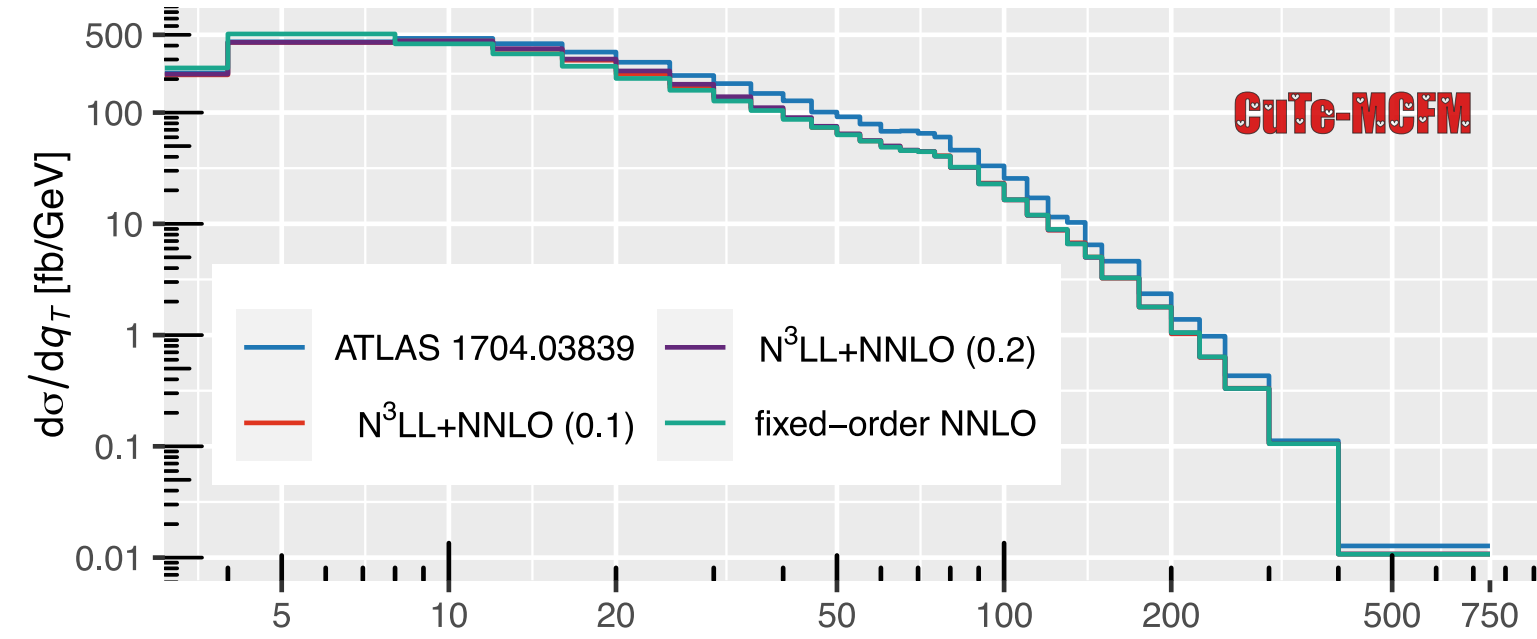
CuTe, DYRes, HRes, Resbos, ReSolve, **RadISH+MATRIX**, RadISH+NNLOjet, SCETlib, TMD focus:
NangaParbat, Artemide,...



$$\gamma\gamma, q_T^{\gamma,1} > 40, q_T^{\gamma,2} > 30$$

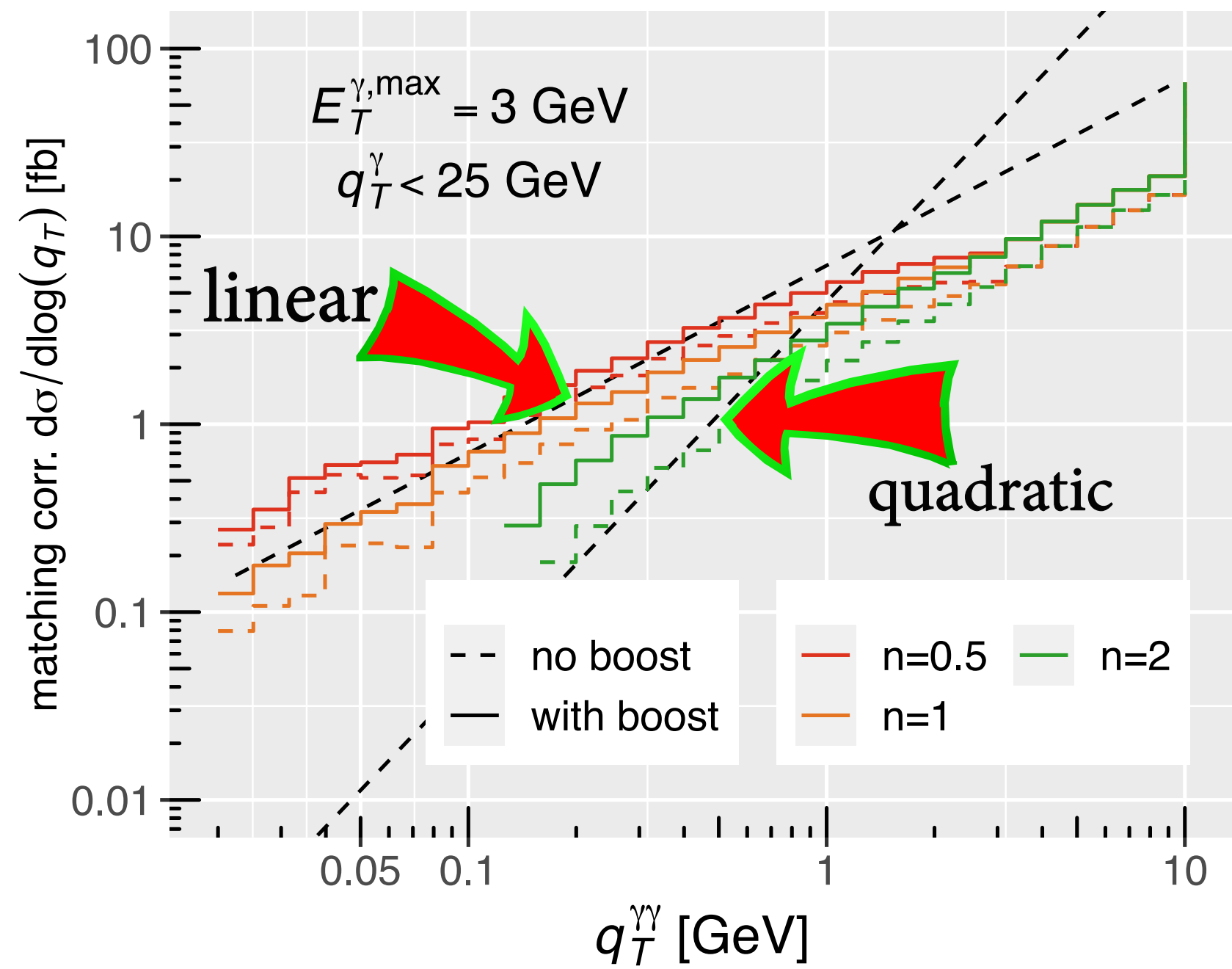


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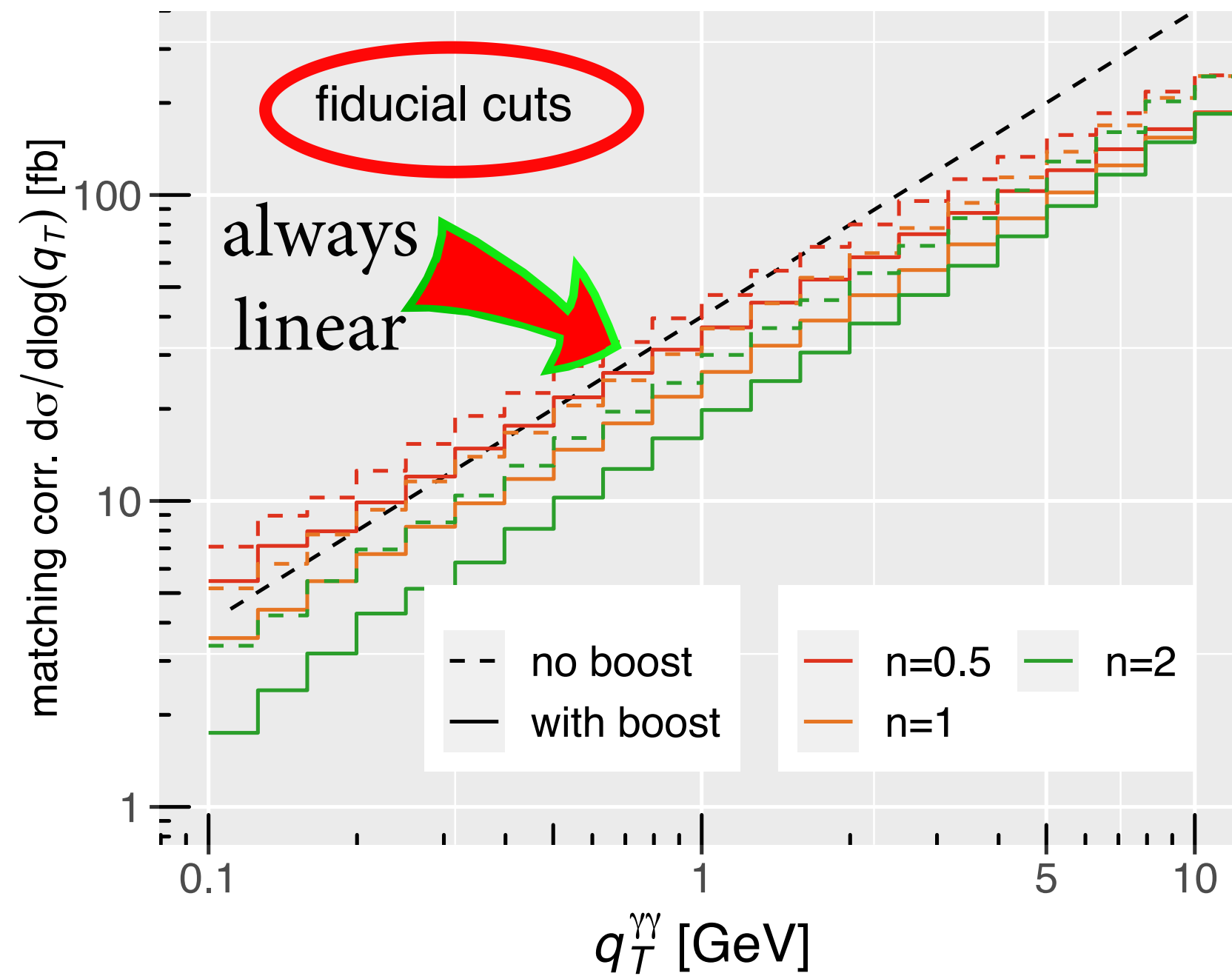
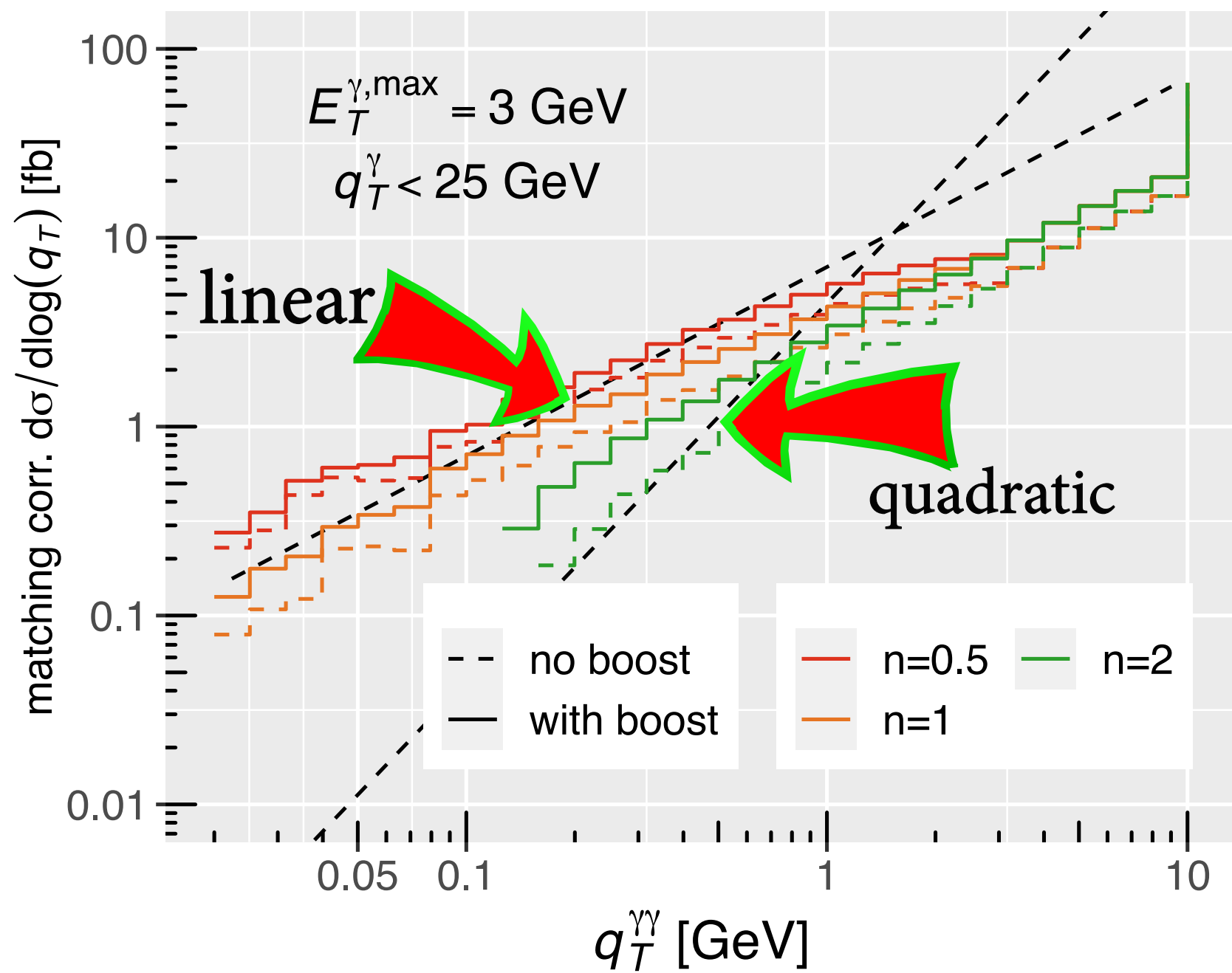


Power corrections in the factorization theorem

$$d\sigma = \dots + \mathcal{O}(q_T/Q) + \mathcal{O}(q_T^2/Q^2) + \dots$$



(see also Ebert; Tackmann '19)



CuTe-MCFM

https://mcfm.fnal.gov/cute-mcfm.html

CuTe-MCFM

MCFM **CuTe-MCFM**

N^3LL q_T resummation for color-singlet processes in MCFM

Authors: [Thomas Becher](#), [Tobias Neumann](#)

CuTe-MCFM is a program and framework for q_T resummation at N^3LL accuracy for color-singlet processes based on a factorization theorem in SCET. It is fully differential in the Born kinematics and matches to large- q_T fixed-order predictions in MCFM at relative order α_s^2 . It provides an efficient way to estimate uncertainties from fixed-order truncation, resummation, and parton distribution functions. In addition to W^\pm , Z and H production, also the diboson processes $\gamma\gamma$, $Z\gamma$, ZH and $W^\pm H$ are available, including decays.

The program is based on the publication

- *Fiducial q_T resummation of color-singlet processes at $N^3LL+NNLO$*
Thomas Becher, Tobias Neumann, JHEP 03 (2021) 199, [arXiv:2009.11437](#)

We release CuTe-MCFM and MCFM as a combined code: Please refer to it as CuTe-MCFM if you use the resummation and simply as MCFM otherwise.

Download

A quickstart manual for CuTe-MCFM is included in the distribution in Doc/cute-mcfm.pdf. The manual for MCFM itself is included in Doc/manual.pdf. Example input files for all color-singlet processes are included in the Bin/ directory and can be used to reproduce the results in our publication.

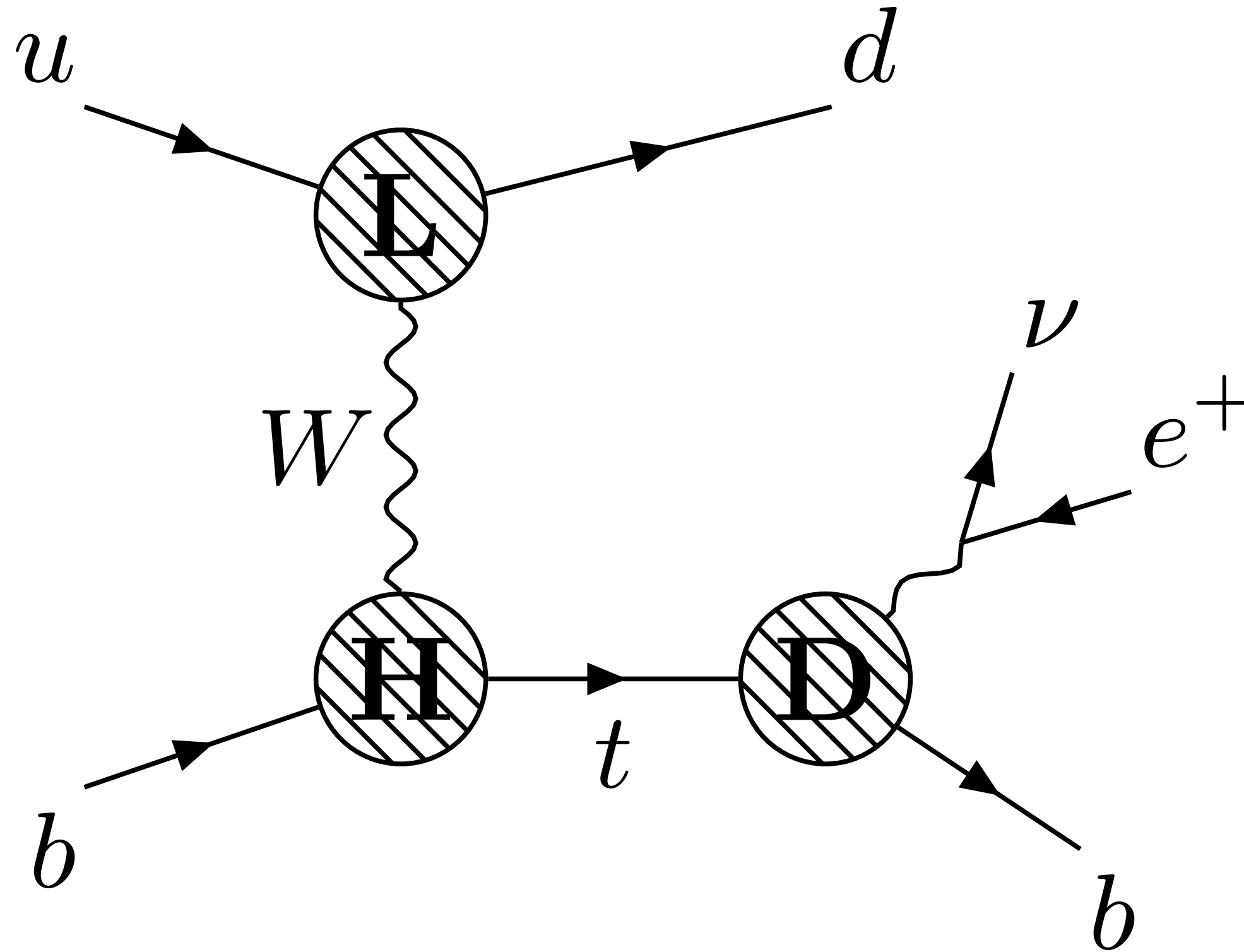
- Initial release [MCFM 10.0 / CuTe-MCFM 1.0](#) (March 2021)

Please feel free to contact the authors for any questions.

CuTe-MCFM now released: <https://mcfm.fnal.gov/cute-mcfm.html>

resummation

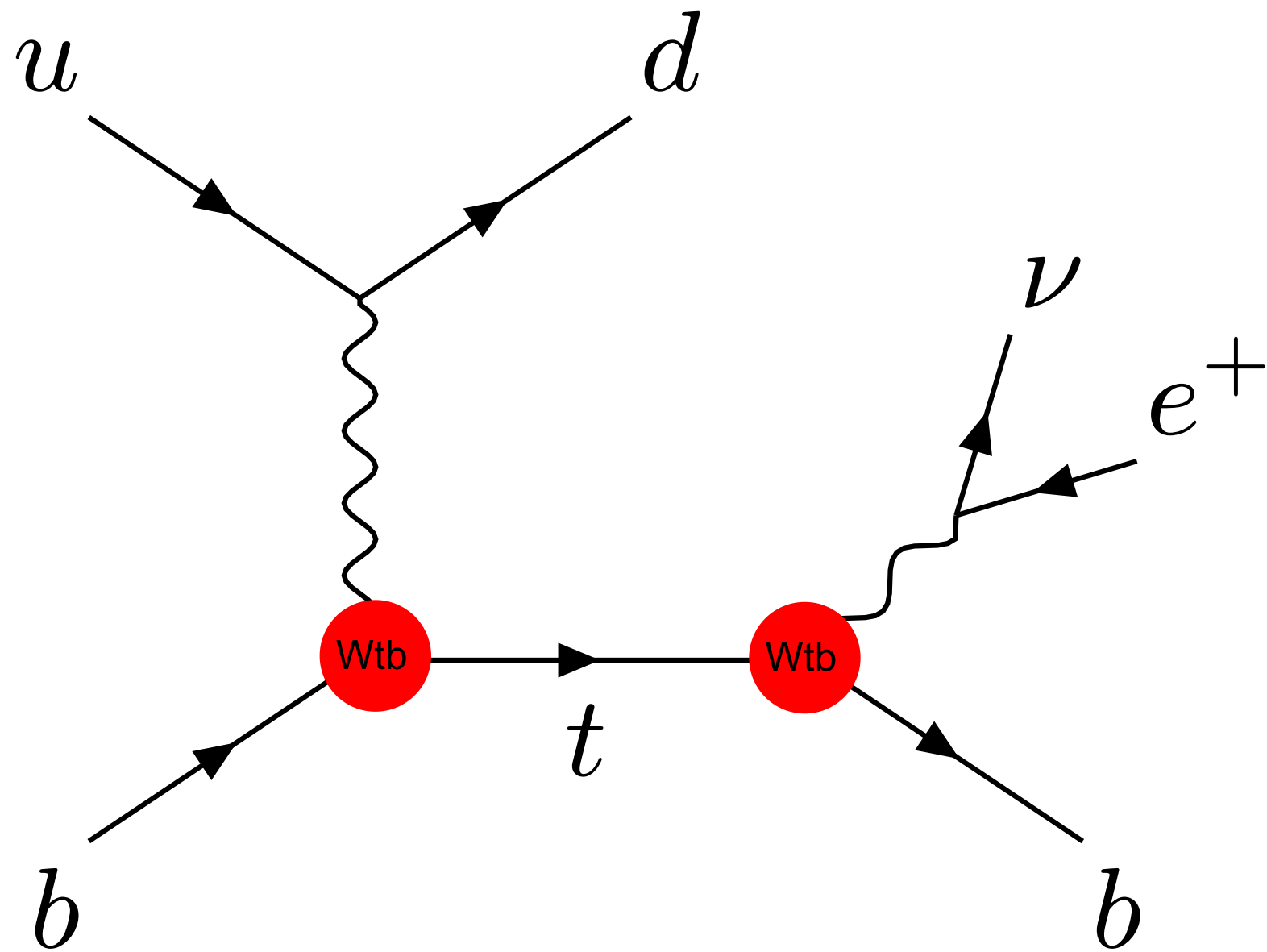
fixed-order



@ NNLO; with John Campbell and Zack Sullivan

Why is this process relevant?

... is it because ... ?



Access to V_{tb} : $\propto |V_{tb}|^2$

Top-quark mass: m_{bl} lineshape

As background with signature $W, b + \text{light jets}$

Prime process to test V-A structure $\gamma^\mu P_L$

1. non-decaying top, needs decay

(Brucherseifer, Caola, Melnikov '14)

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(NNLO corrections are ~1-2%)

$$Z = \int dx f_q^{LO}(x) \text{ --- } q \text{ --- } (LO) \text{ --- } q' = \boxed{\text{physics}}$$

$$Z = \int dx f_q^{NLO} \text{ --- } q \text{ --- } (NLO) \text{ --- } q' = \boxed{\text{physics}}$$

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$$\delta = \int dx f_q^{LO}(x) \text{ --- } q \text{ --- } (LO) \text{ --- } q' = \boxed{\text{physics}}$$

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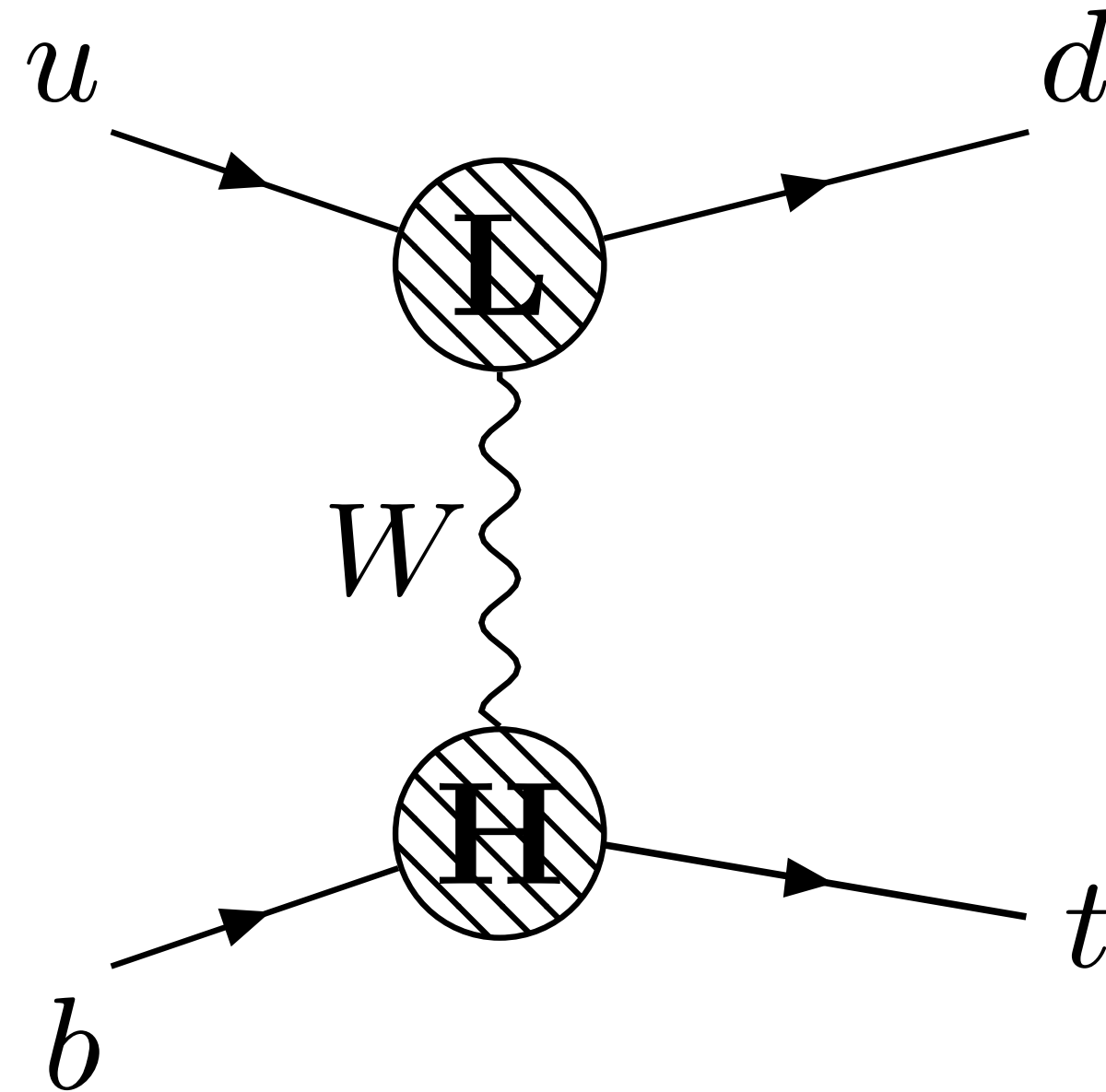
||

$$\delta = \int dx f_q^{NLO}(x) \text{ --- } \textcircled{NLO} \text{ --- } q' = \boxed{\text{physics}}$$

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$$\delta = \int dx f_q^{NNLO}(x) \text{ --- } \textcircled{NNLO} \text{ --- } q' = \boxed{\text{physics}}$$

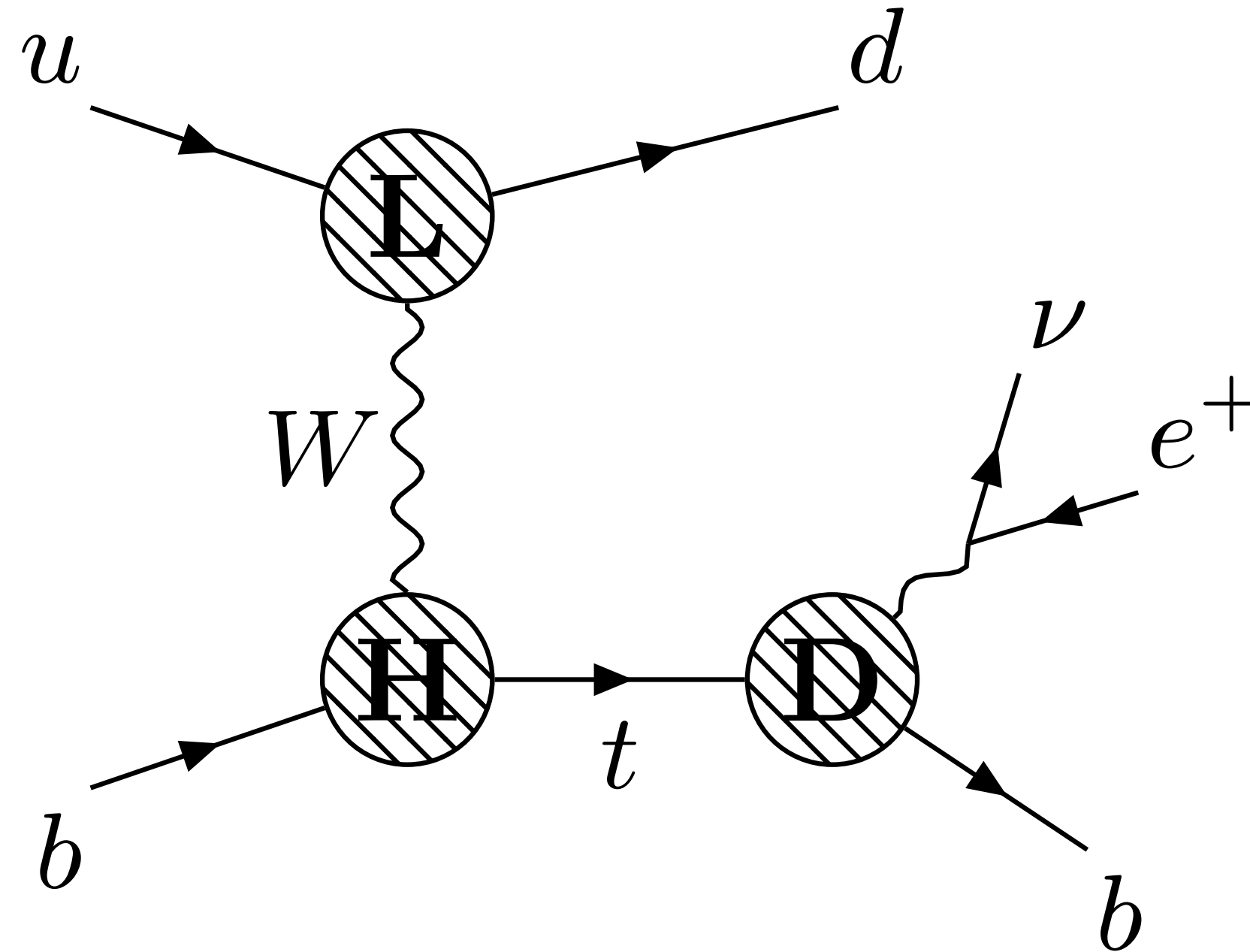
Double DIS = DDIS



DDIS constraint at NLO: "Are PDFs still consistent with Tevatron data?"

(Sullivan '17)

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DDIS constraint at NLO: "Are PDFs still consistent with Tevatron data?"

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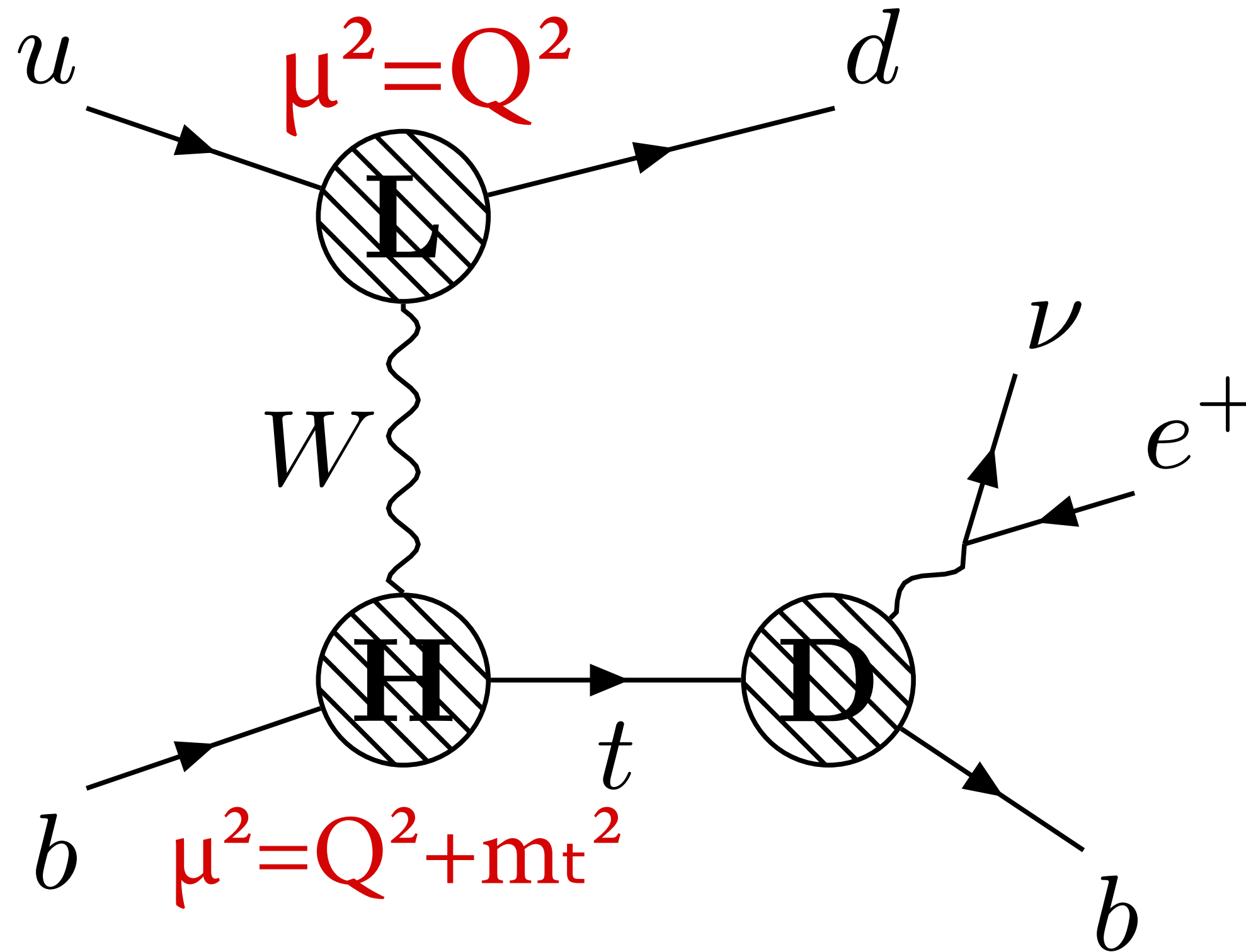
"found discrepancy of ~1%"

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(NNLO corrections are ~1-2%)

Why is the same scale everywhere a problem?

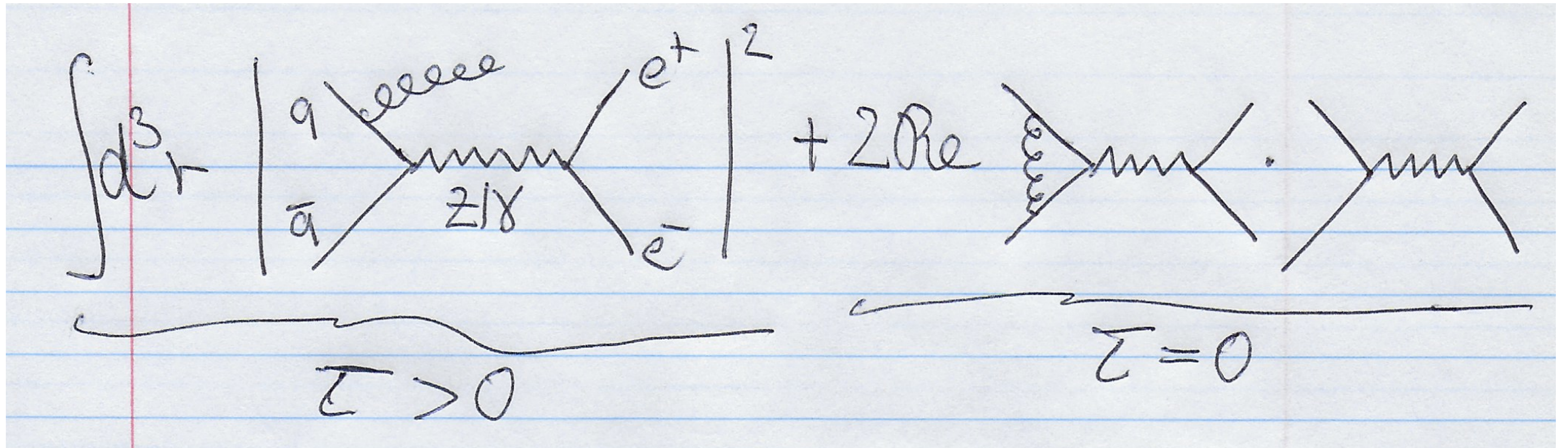
Double DIS



Mini interlude: higher-order pQCD

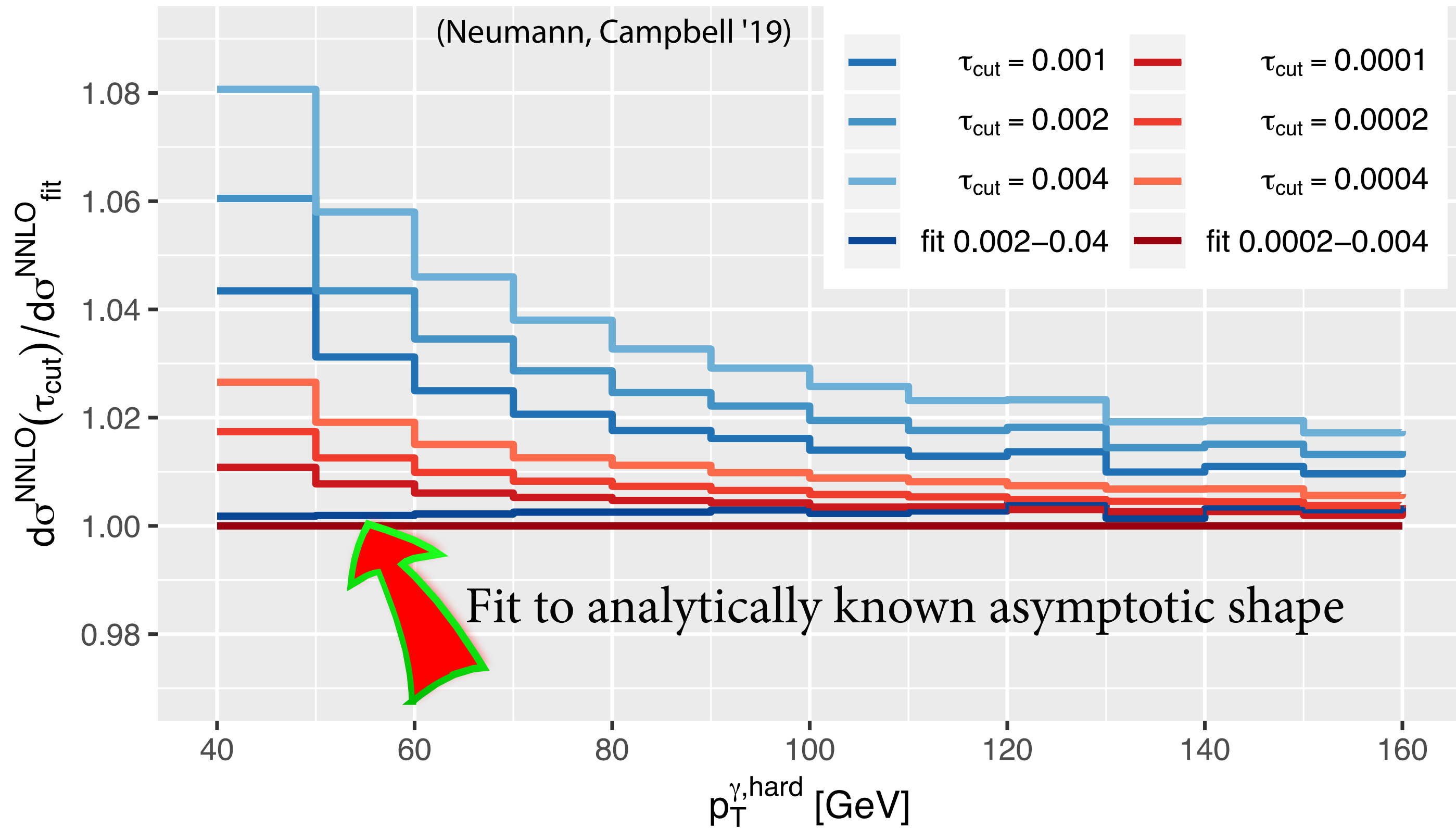
$$\int d^3r \left| \begin{array}{c} q \\ \text{gluon} \\ q \end{array} \right| \left| \begin{array}{c} e^+ \\ \text{gluon} \\ e^- \end{array} \right|^2 + 2 \text{Re} \left[\begin{array}{c} \text{gluon} \\ \text{gluon} \end{array} \right] \cdot \left[\begin{array}{c} \text{gluon} \\ \text{gluon} \end{array} \right]$$

The image shows a handwritten mathematical expression on lined paper. On the left, there is an integral $\int d^3r$ followed by a vertical bar. Inside the bar, a quark line (labeled 'q') and an antiquark line (labeled 'q') meet at a vertex, with a gluon line (represented by a wavy line) connecting this vertex to another vertex. From the second vertex, an electron line (labeled 'e+') and a positron line (labeled 'e-') emerge. The entire diagram is enclosed in a vertical bar with a superscript '2'. To the right of this is a plus sign followed by '2 Re' and a dot product of two diagrams. Each diagram in the dot product consists of a gluon line (wavy) connecting two vertices, with a quark line (solid) and an antiquark line (solid) meeting at each vertex.



$$\sigma^{\text{NNLO}} = \lim_{\tau_c \rightarrow 0} \left((H^{\text{NNLO loop}} \otimes S \otimes B \otimes B \otimes J)(\tau_c) + \sigma_{\text{X+jet}}^{\text{NLO}}(\tau_c) + \mathcal{O}(\tau_c) \right) + \mathcal{O}(\Lambda_{\text{QCD}}/Q)$$

- $H \sim$ two loop amplitude
- B collinear singularities for radiation close to beam
- J collinear singularities associated with final-state jets
- S all remaining soft singularities
- $\sigma_{\text{X+jet}}^{\text{NLO}}(\tau_c)$ radiation pieces, finite for $\tau > 0$



Systematically improvable: subleading power corrections

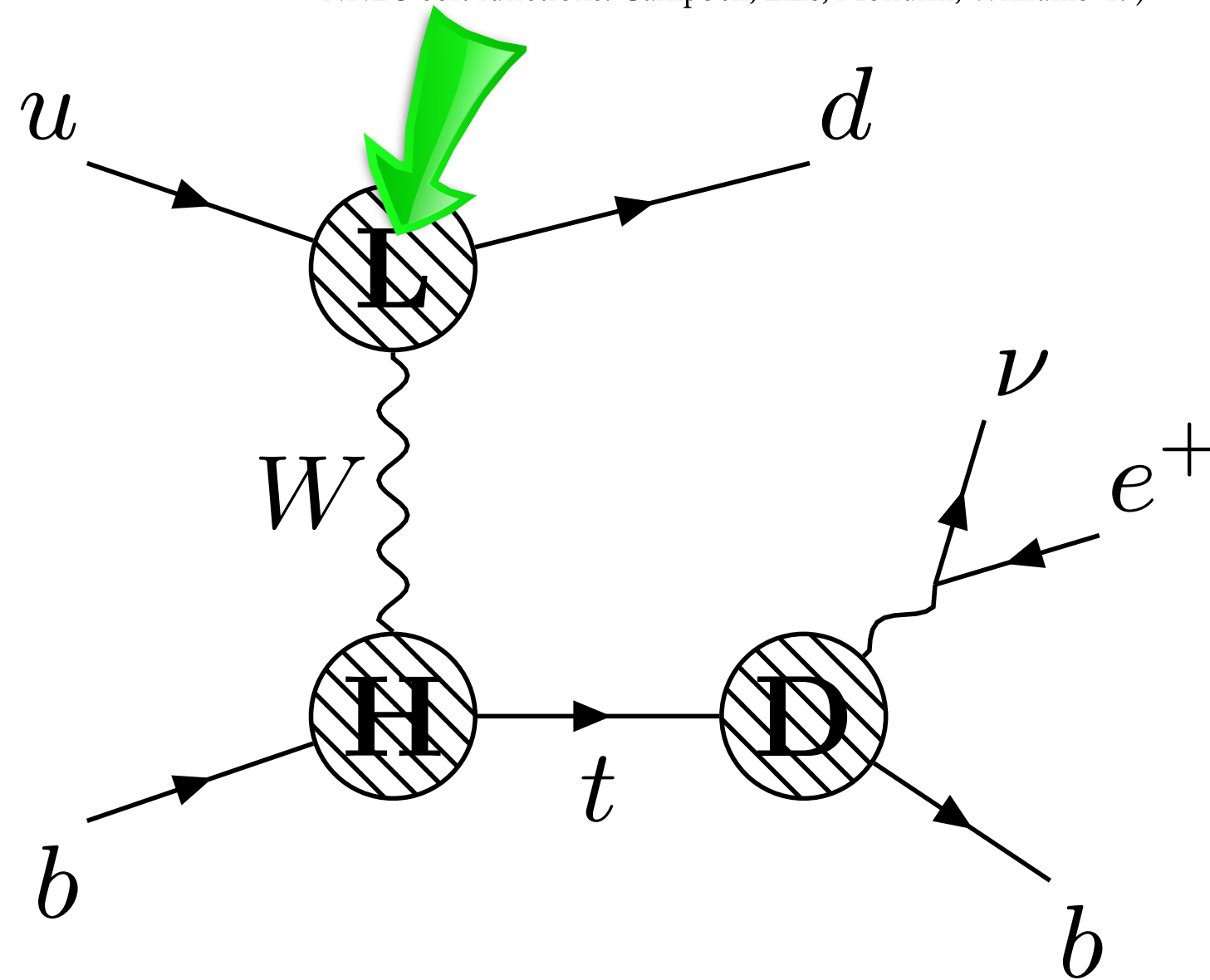
(Moult, Rothen, Stewart, Tackmann '16; Ebert, Moult, Stewart, Tackmann, Vita, Zhu '18)

Back to t-channel single-top-quark

Three NNLO calculations

massless 1-jettiness

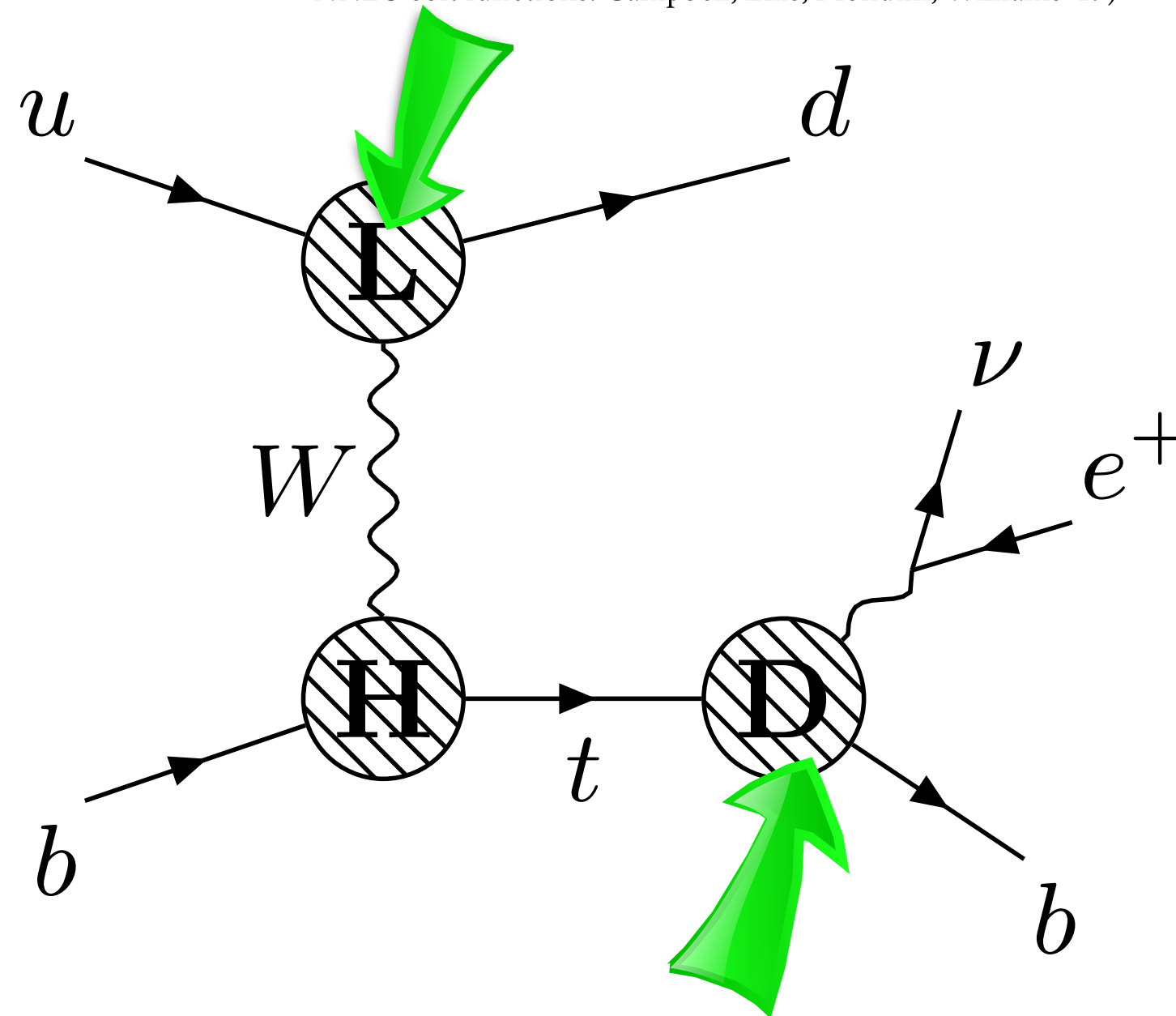
(Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15;
NNLO soft functions: Campbell, Ellis, Mondini, Williams '17)



Three NNLO calculations

massless 1-jettiness

(Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15;
NNLO soft functions: Campbell, Ellis, Mondini, Williams '17)



"massive" jettiness

(Liu '10; Bosch, Lange; Neubert, Paz '04; Bauer, Manohar '04)

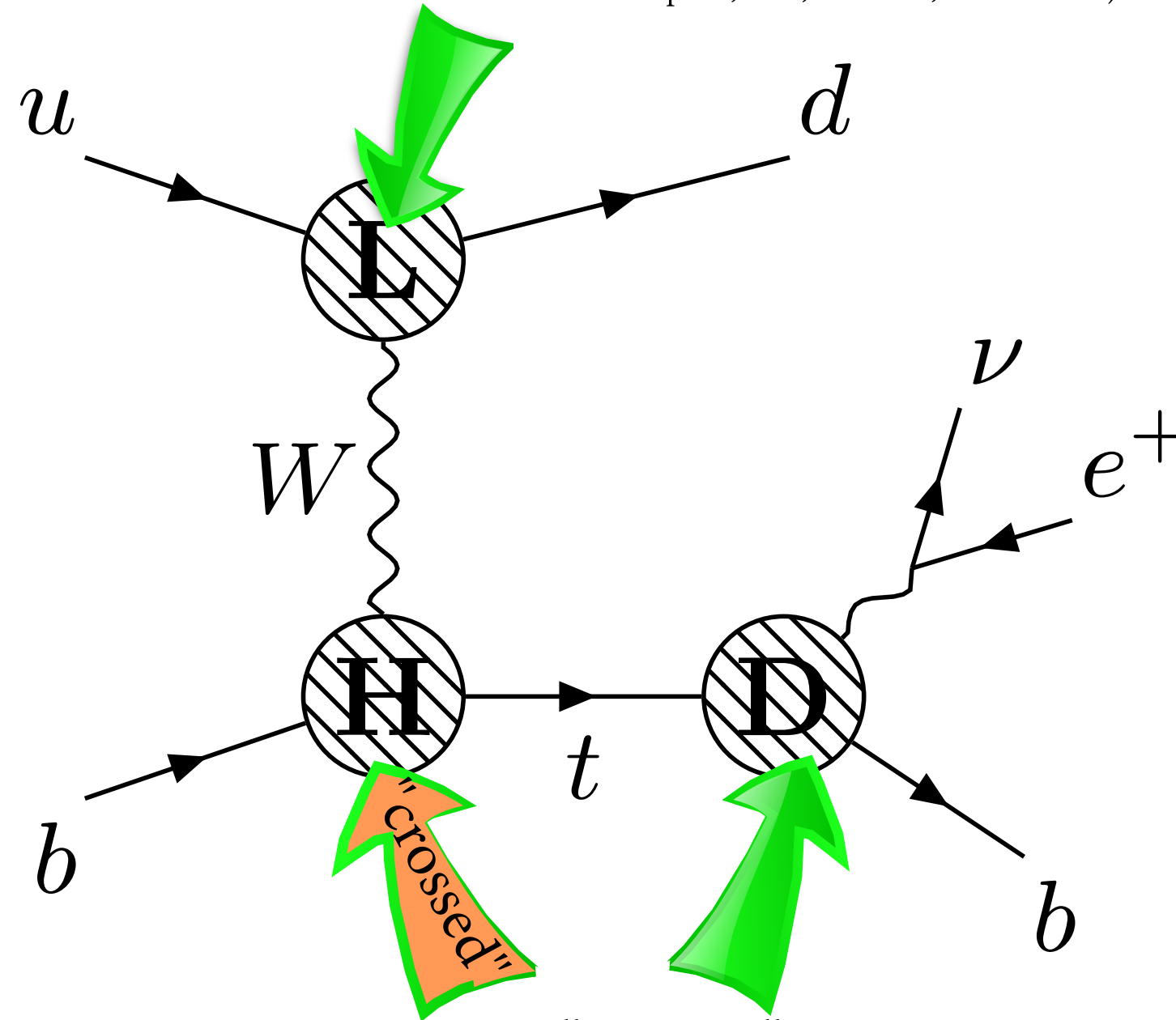
(soft: Becher, Neubert '05; jet: Becher, Neubert '06)

(hard: Bonciani, Ferroglia '08; Asatryan, Greub, Pecjak '08; Beneke, Huber, Li '08; Bell '08)

Three NNLO calculations

massless 1-jettiness

(Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15;
NNLO soft functions: Campbell, Ellis, Mondini, Williams '17)



"massive" jettiness

(Liu '10; Bosch, Lange; Neubert, Paz '04; Bauer, Manohar '04)

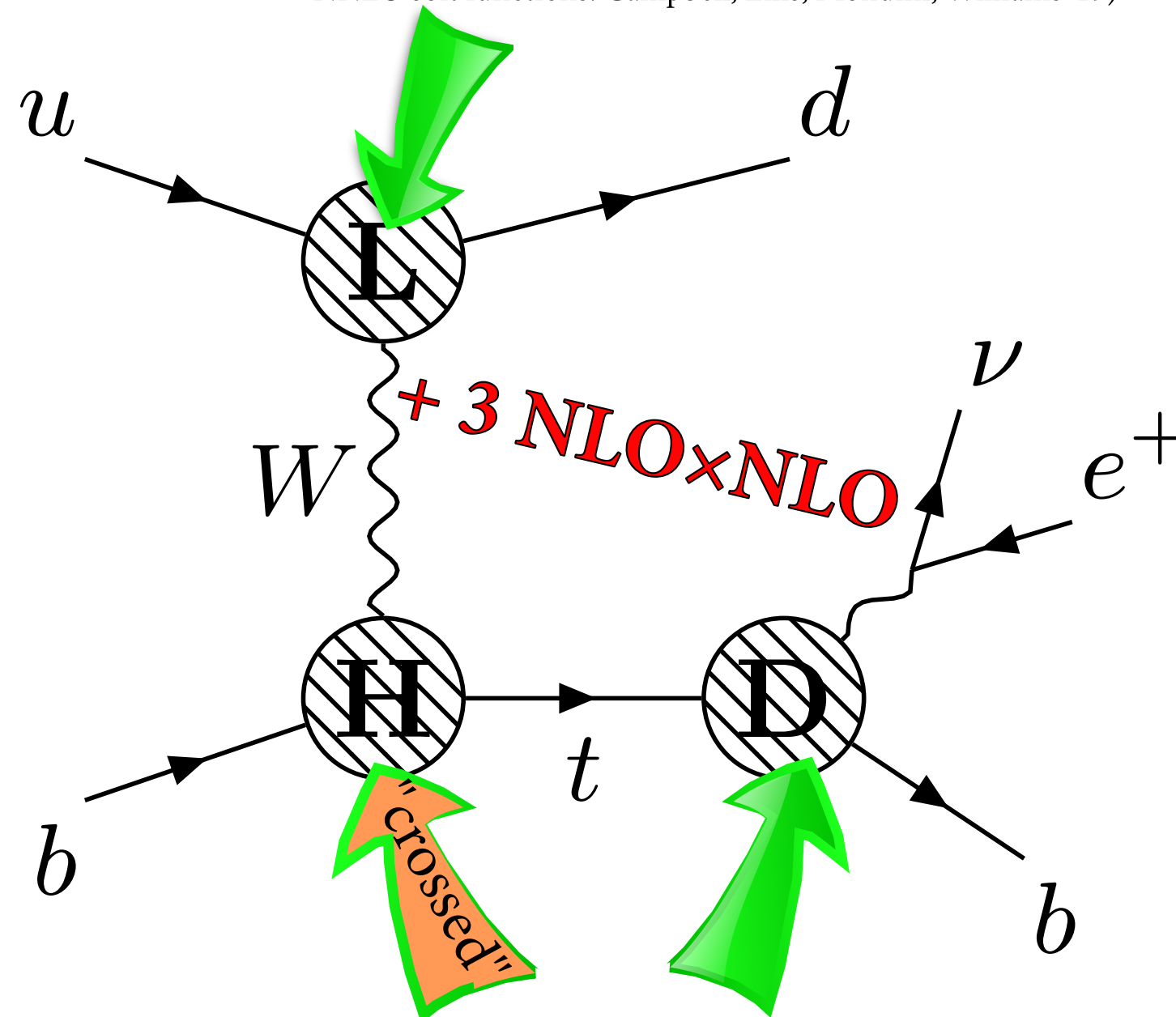
(soft: Becher, Neubert '05; jet: Becher, Neubert '06)

(hard: Bonciani, Ferroglia '08; Asatryan, Greub, Pecjak '08; Beneke, Huber, Li '08; Bell '08)

Three NNLO calculations + three NLOxNLO calculations

massless 1-jettiness

(Boughezal, Focke, Liu, Petriello '15; Gaunt, Stahlhofen, Tackmann, Walsh '15;
NNLO soft functions: Campbell, Ellis, Mondini, Williams '17)

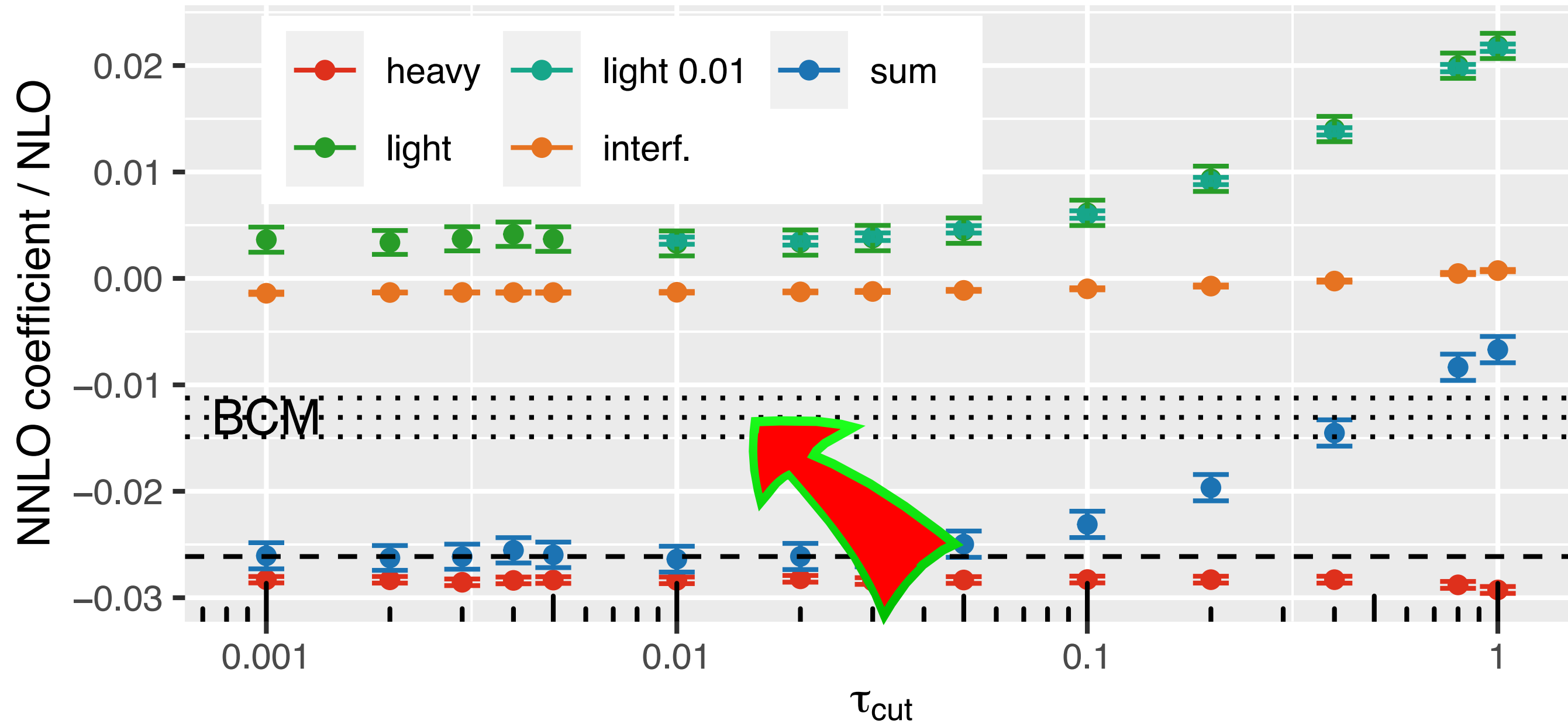


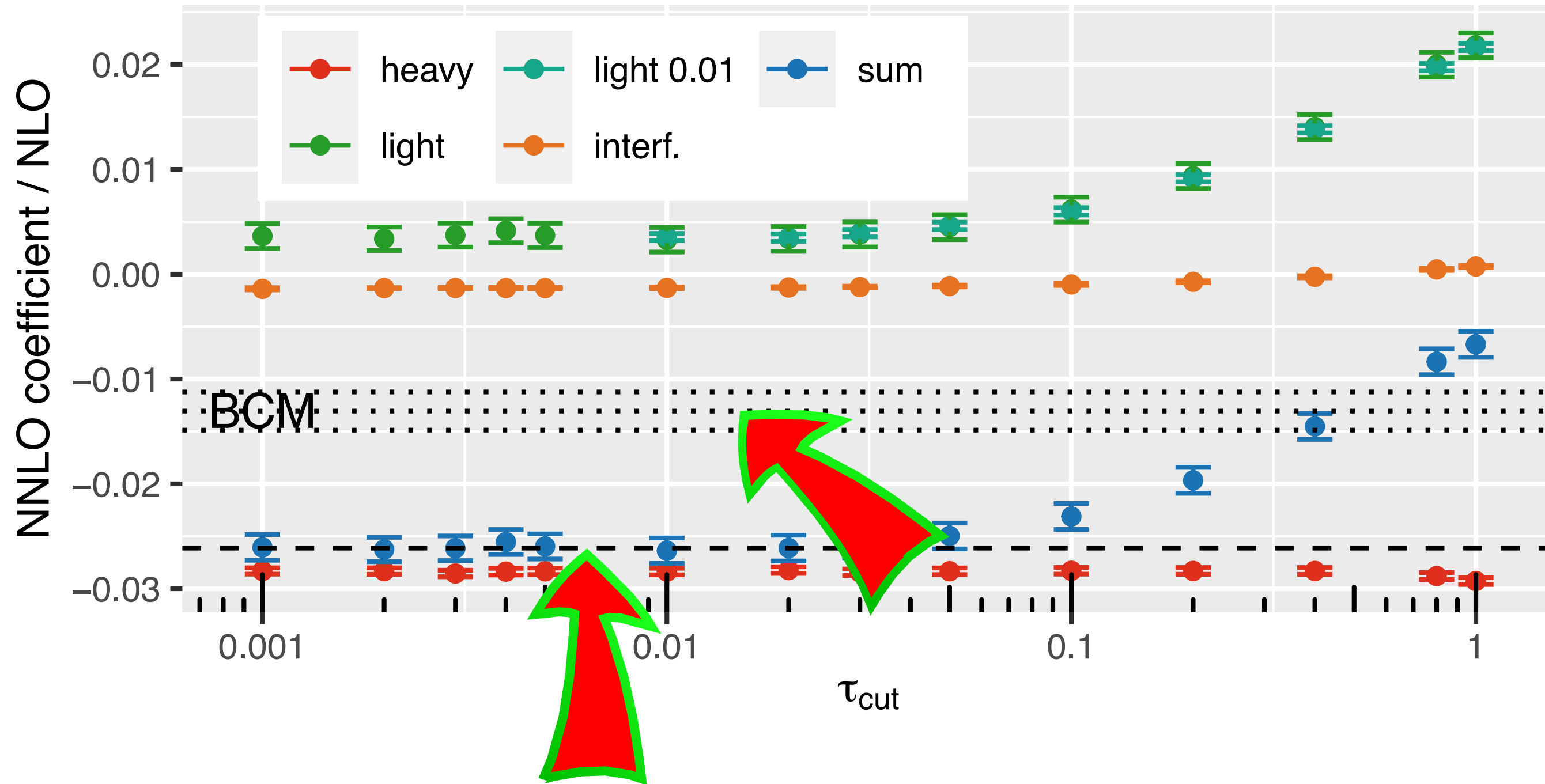
"massive" jettiness

(Liu '10; Bosch, Lange; Neubert, Paz '04; Bauer, Manohar '04)

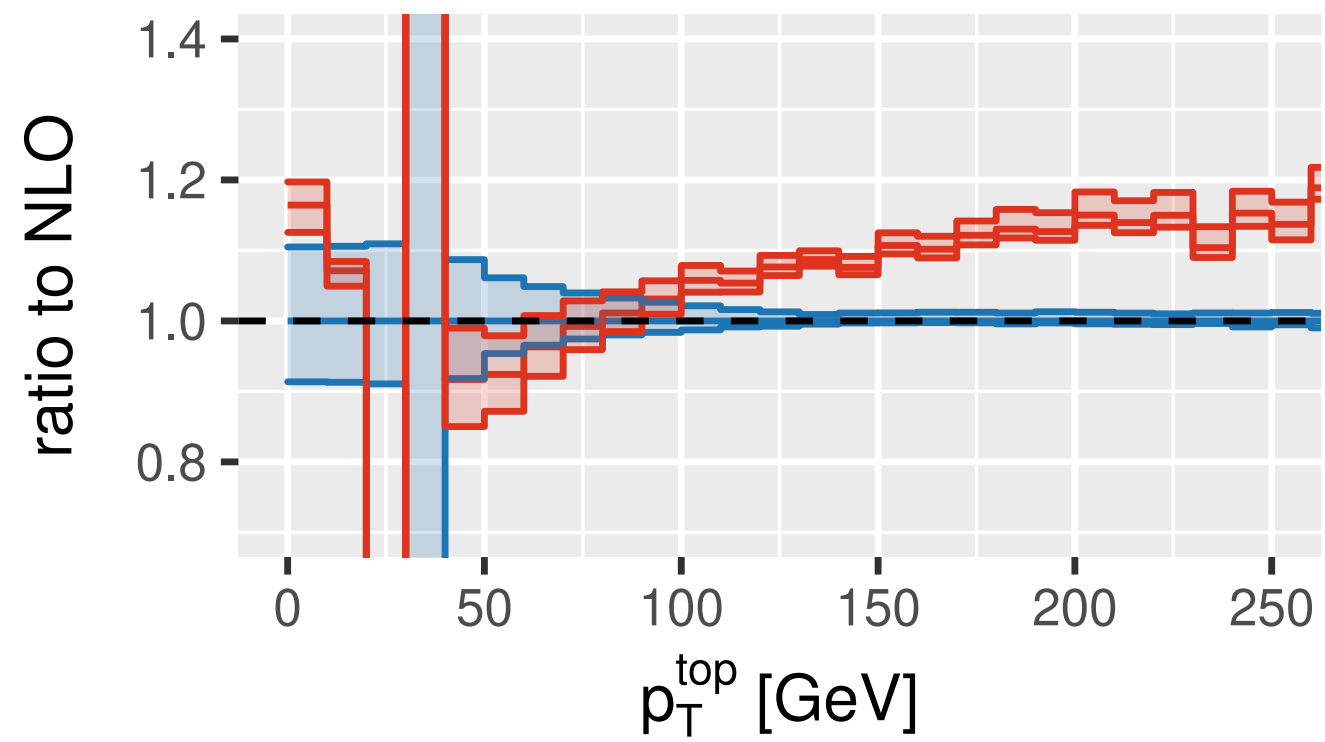
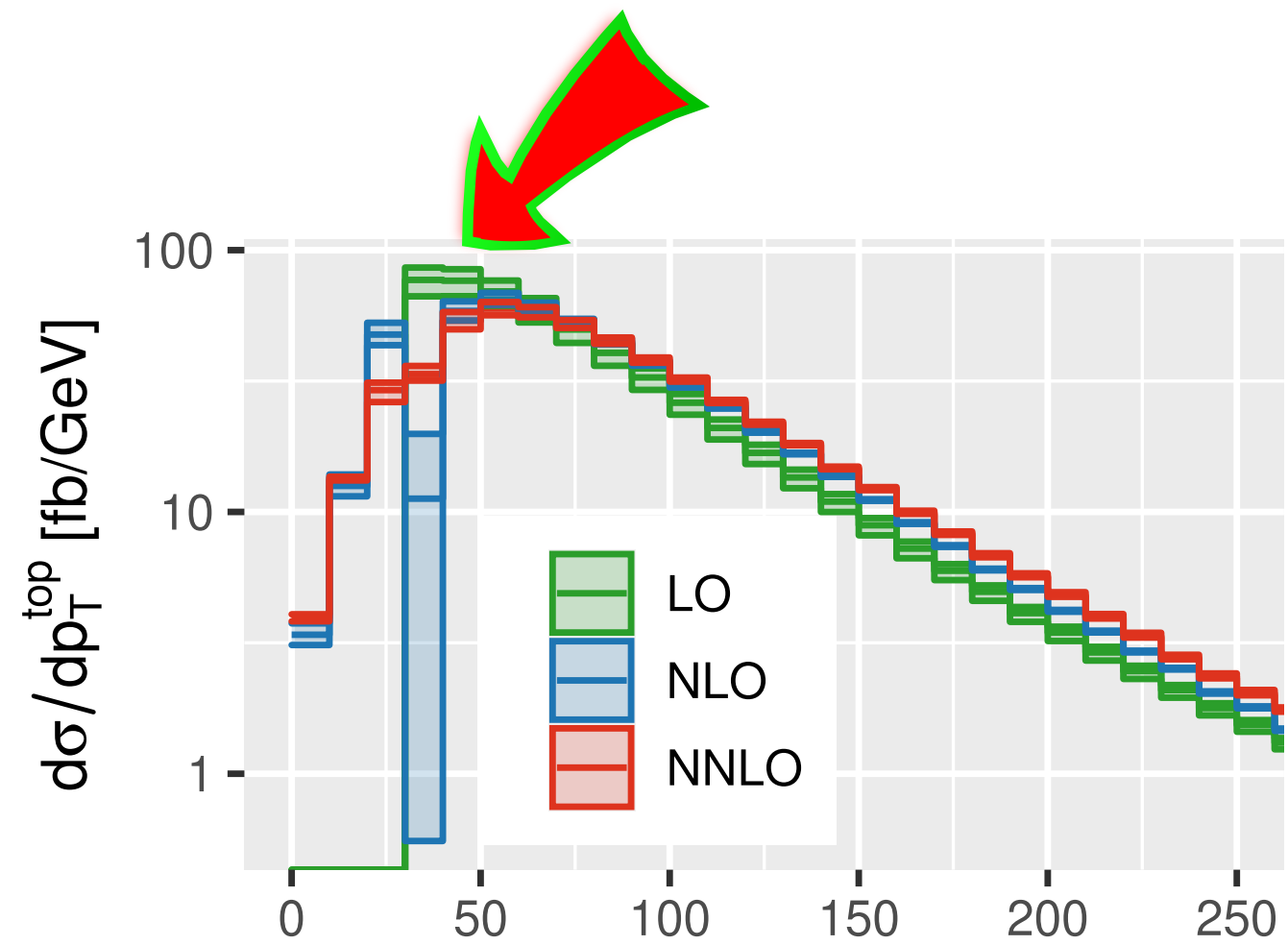
(soft: Becher, Neubert '05; jet: Becher, Neubert '06)

(hard: Bonciani, Ferroglia '08; Asatryan, Greub, Pecjak '08; Beneke, Huber, Li '08; Bell '08)

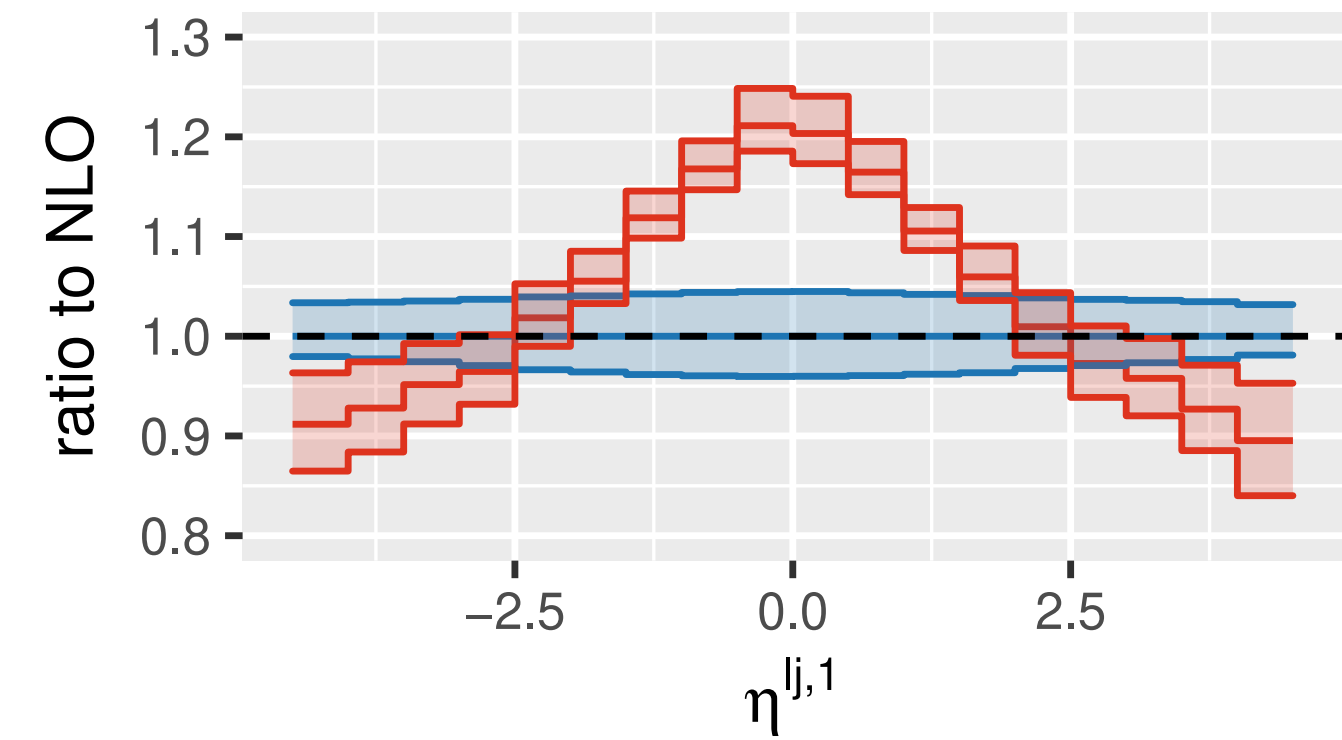
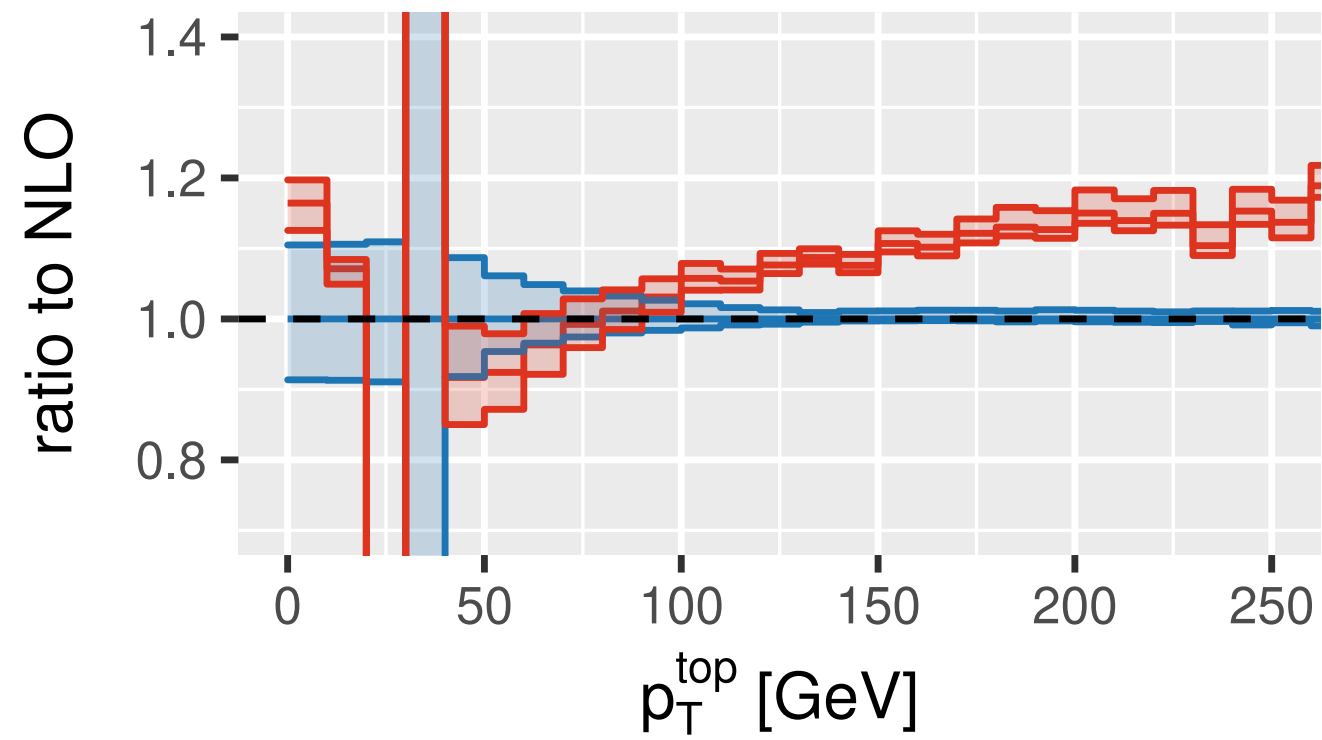
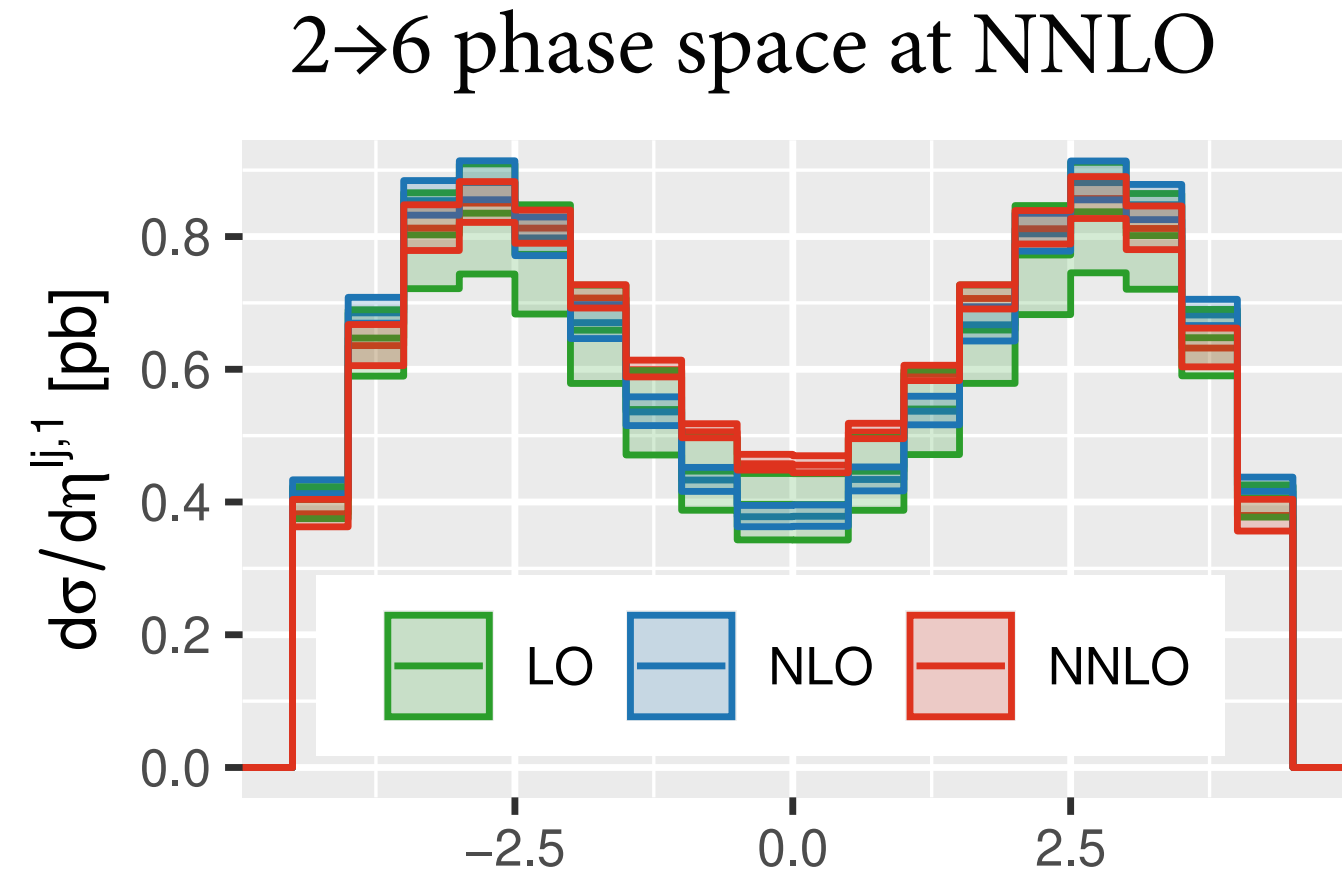
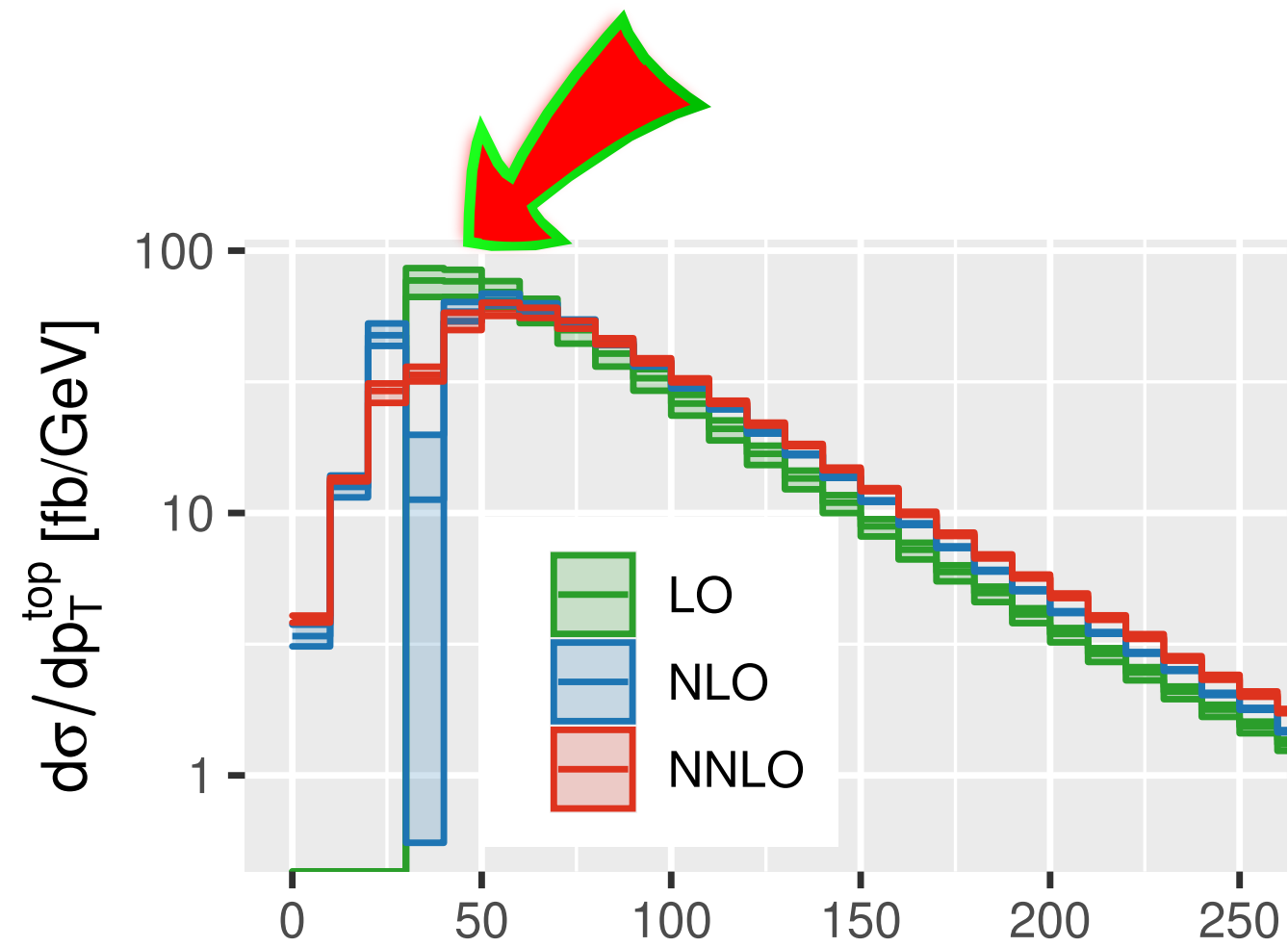




On the other hand: We find full agreement with Berger, Gao, Zhu in extensive comparisons.



(LHC, typical fiducial cuts, DDIS scales)



(LHC, typical fiducial cuts, DDIS scales)

Study of PDFs discrepancies with DDIS in upcoming study!

$$\delta = \int dx f_q^{LO}(x) \text{ --- } q \text{ --- } \textcircled{LO} \text{ --- } q' \text{ --- } = \boxed{\text{physics}}$$

||

$$\delta = \int dx f_q^{NLO}(x) \text{ --- } q \text{ --- } \textcircled{NLO} \text{ --- } q' \text{ --- } = \boxed{\text{physics}}$$

||

$$\delta = \int dx f_q^{NNLO}(x) \text{ --- } q \text{ --- } \textcircled{NNLO} \text{ --- } q' \text{ --- } = \boxed{\text{physics}}$$

Contributions

CuTe-MCFM: N^3 LL+NNLO SCET-based resummation framework

Prime tool to contribute to W/Z precision physics (unique SCET formulation)

Improvements for diboson processes at small q_T ($\gamma\gamma$, $Z\gamma$, ZH , WH)

Demonstrated power corrections using recoil-scheme
in presence of fiducial cuts and photon isolation

General framework for future processes and other codes

Calculation of t-channel single-top-quark production at NNLO

Resolved discrepancies between previous two calculations

Large fiducial corrections make NNLO important

Double-DIS scales allow for novel PDF constraint

Backup/Details

CuTe-MCFM

q_T -resummation framework

- All-order formula for resummation of q_T -logs first by Collins, Soper, Sterman '85

- Semi-analytical / PS-like approach: Radish

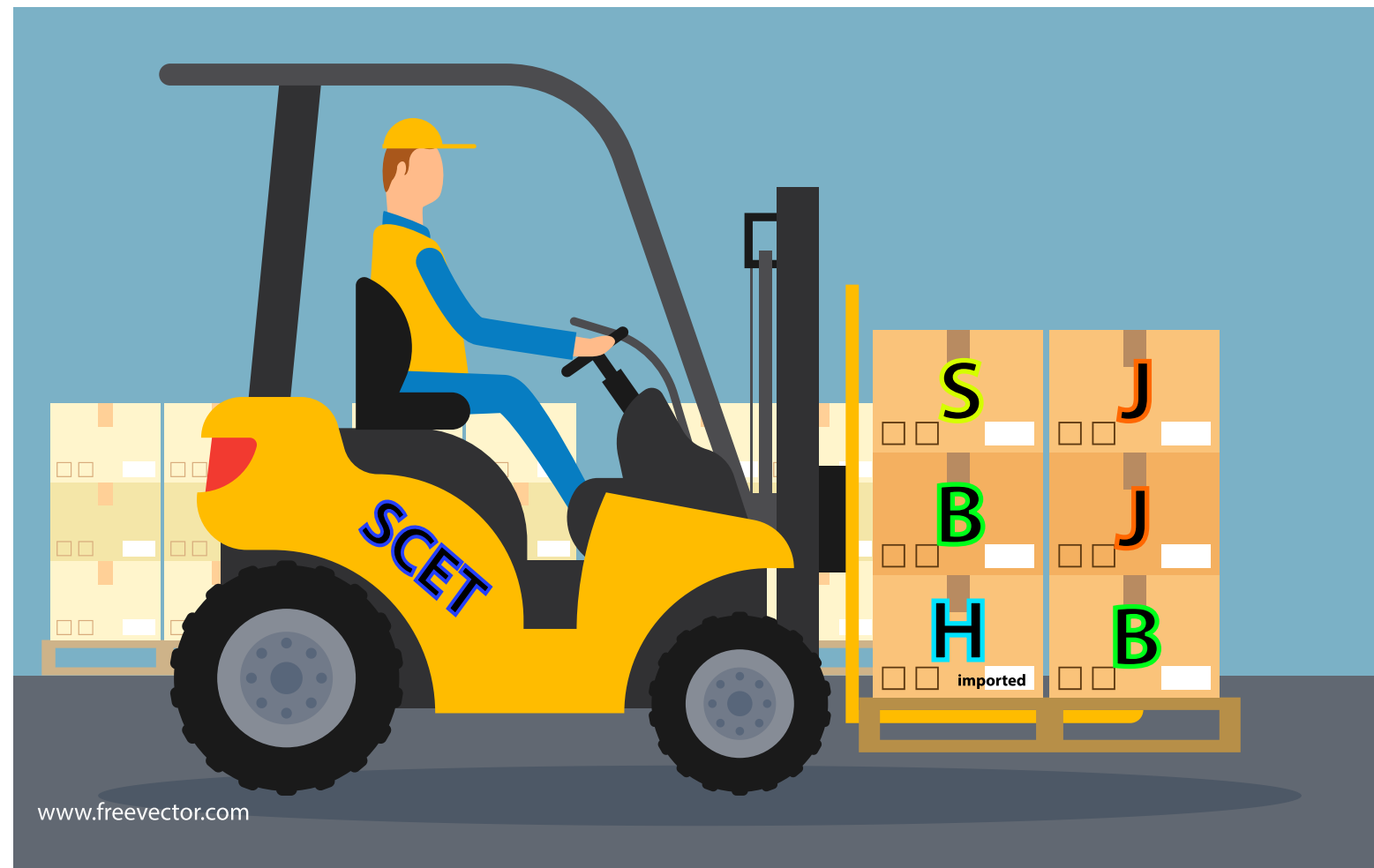
Monni, Re, Torrielli '16; Bizon, Monni, Re, Rottoli, Torrielli '17; Banfi, Salam, Zanderighi '04; ...

- Modern development in SCET: (re-)derivation of factorization theorem

Becher, Neubert '10; Chiu, Jain, Neill, Rothstein '12; see also Gao, Li, Liu '05, Idilbi, Ji, Yuan '05, Mantry, Petriello '09; rapidity RGE in q_T space: Ebert, Tackmann '16

$$\begin{aligned} d\sigma = & \sum_{i,j} \iint d\xi_1 d\xi_2 d\sigma_{ij}^0(p) H_{ij}(p, \mu) \times \\ & \frac{1}{4\pi} \int_{-\infty}^{\infty} d^2 x_{\perp} e^{-iq_{\perp} x_{\perp}} \left(\frac{x_T^2 Q^2}{2e^{-\gamma_E}} \right)^{-F_{ij}(x_{\perp}, \mu)} B_i(\xi_1, x_{\perp}, \mu) B_j(\xi_2, x_{\perp}, \mu) \end{aligned}$$

N^3 LL color-singlet q_T resummation



N^3 LL+NNLO: overall α_s^2 accuracy for
 $\log(q_T^2/Q^2) \sim 1/\alpha_s$

A pot-pourri of cutting-edge ingredients:

- NNLO / Hard function: process dependent
qq/gg-formfactors at 2-loop; diboson Gehrman, Manteuffel, Tancredi '15
- Hard function RGE evolution
4-loop cusp: Henn, Korchemsky, Mistlberger '19; Manteuffel, Panzer, Schabinger '20
other ingredients: see Becher, Neubert '09 for overview
- Collinear anomaly/rapidity anom. dim.
 α_s^3 : *Li, Zhu '16; Vladimirov '16*
- NNLO beam functions
Gehrman, Lubbert, Lin Yang '14; Echevarria, Scimemi, Vladimirov '16;
Luebbert, Oredsson, Stahlhofen '16,

Meanwhile also N^3 LO beam functions

Ebert, Mistlberger, Vita '20; Luo, Yang, H.X. Zhu, Y.J. Zhu '20

Our implementation: CuTe-MCFM

Uses broad availability of fixed/higher-order processes in MCFM

- MCFM v9 as basis

Boughezal, Campbell, Ellis, Focke, Giele, Liu, Petriello, Williams '16
Neumann, Campbell '19;

- Set of standard NNLO processes: W^\pm , Z, H, $\gamma\gamma$, $Z\gamma$, $W^\pm H$, ZH + many NLO processes
- q_T resummation: Same formalism as CuTe (and cross checks against it), but completely independent implementation and different subleading choices
- Beam-function grid generation
- Scale variation, PDF uncertainties
- Precise control over all numerical uncertainties, OMP+MPI



MCFM-RE N^2 LL jet-veto based on JetVHeto

Arpino, Banfi, Jäger, Kauer '19; Banfi, Monni, Salam, Zanderighi '12

(Subleading) Choices

$$d\sigma = \sum_{i,j} \iint d\xi_1 d\xi_2 d\sigma_{ij}^0(p) H_{ij}(p, \mu) \times$$

$$\frac{1}{4\pi} \int_{-\infty}^{\infty} d^2 x_{\perp} e^{-iq_{\perp} x_{\perp}} \left(\frac{x_T^2 Q^2}{2e^{-\gamma_E}} \right)^{-F_{ij}(x_{\perp}, \mu)} B_i(\xi_1, x_{\perp}, \mu) B_j(\xi_2, x_{\perp}, \mu)$$

- Each ingredient separately expanded in RG improved perturbation theory (compare CuTe)
- Collinear anomaly formalism (equivalent to rapidity RGE with fixed ν)
- Reorganized resummation for $q_t \rightarrow 0$ (" ϵ "-expansion)
Becher, Neubert, Wilhelm '11, follows Parisi, Petronzio '79
- No non-perturbative (Landau-pole) treatment necessary, resummation scale always within perturbative regime
- Tree-level kinematics boosted to have $q_{\perp}^{\mu} = (0, q_T \cos \Phi, q_T \sin \Phi, 0)$

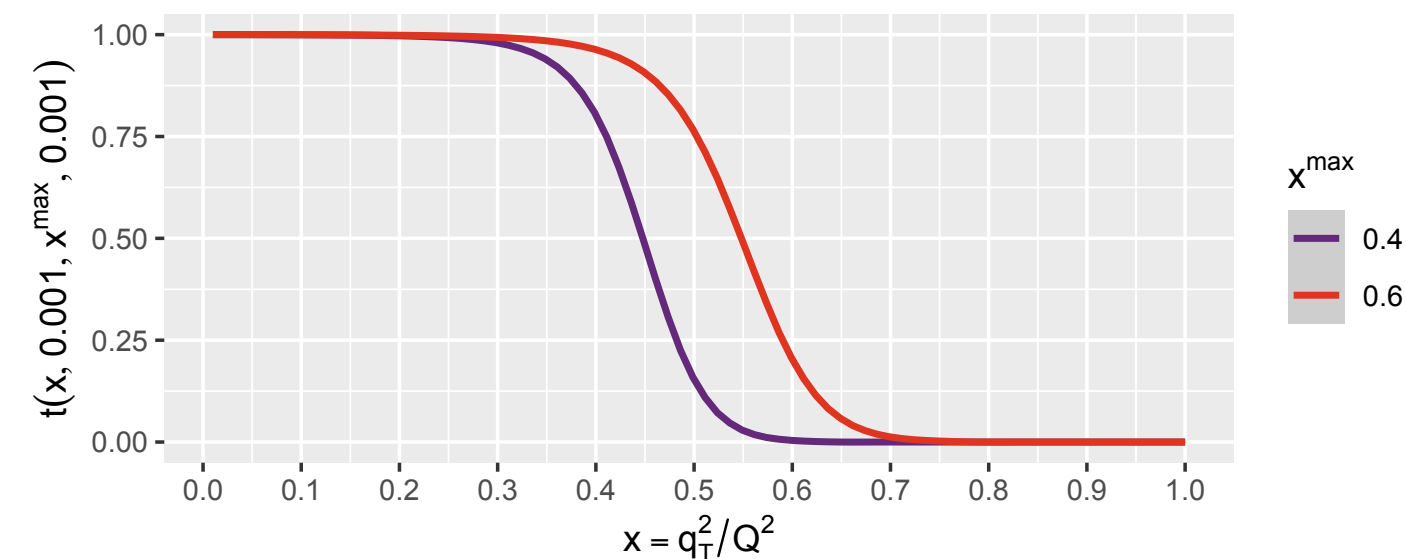
Matching to fixed-order (NNLO)

Naive matching

$$\frac{d\sigma}{dq_T}^{\text{N}^3\text{LL, naive}} = \frac{d\sigma}{dq_T}^{\text{N}^3\text{LL}} + \left(\frac{d\sigma}{dq_T}^{\text{f.o.}} - \frac{d\sigma}{dq_T}^{\text{N}^3\text{LL, exp.}} \right)$$

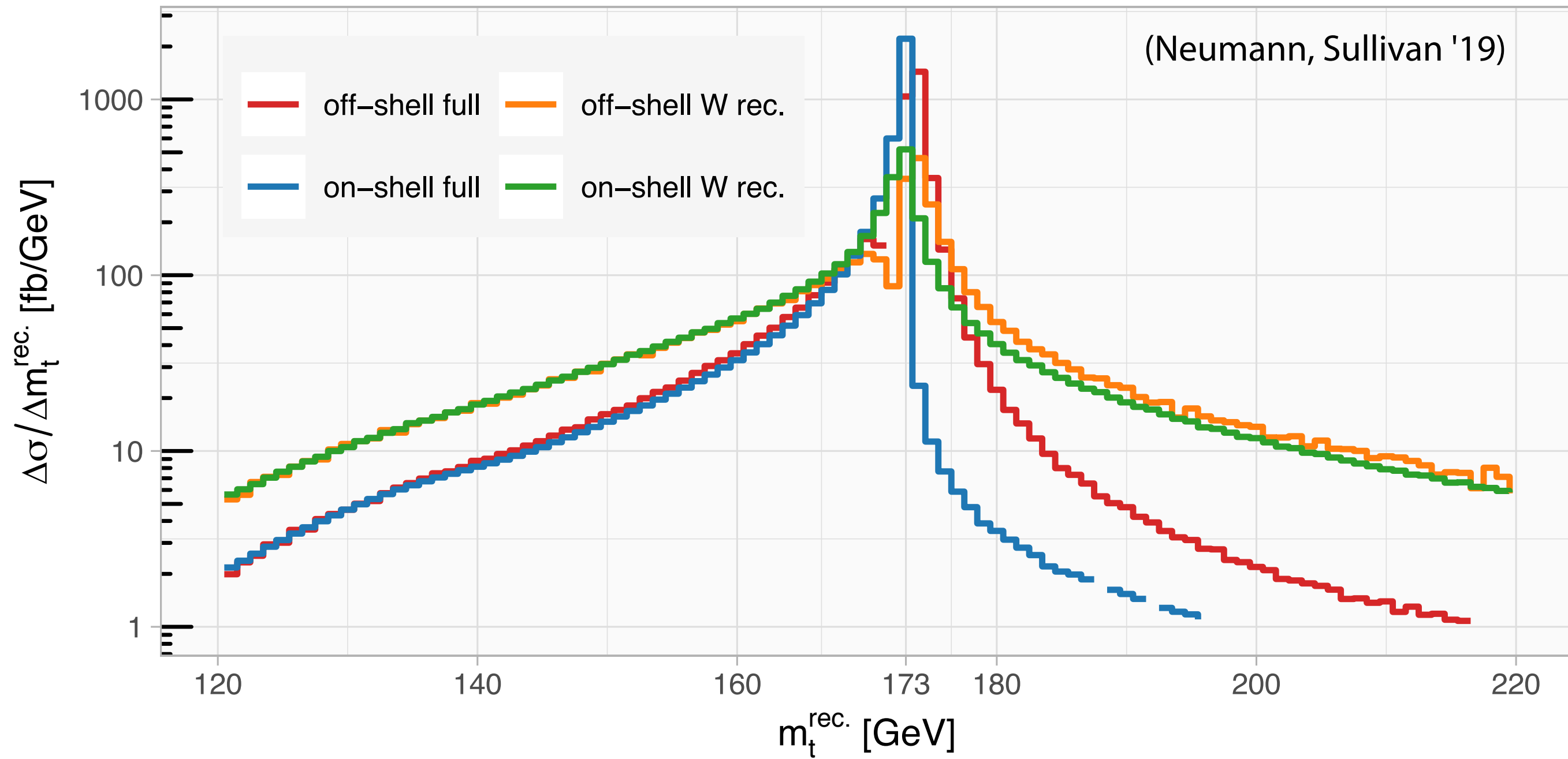
Resummed result only valid for $q_t^2 / Q^2 \ll 1$: Transition-function

$$\frac{d\sigma}{dq_T}^{\text{N}^3\text{LL, full}} = t\left(\frac{q_T^2}{Q^2}\right) \frac{d\sigma}{dq_T}^{\text{N}^3\text{LL, naive}} + \left(1 - t\left(\frac{q_T^2}{Q^2}\right)\right) \frac{d\sigma}{dq_T}^{\text{f.o.}}$$



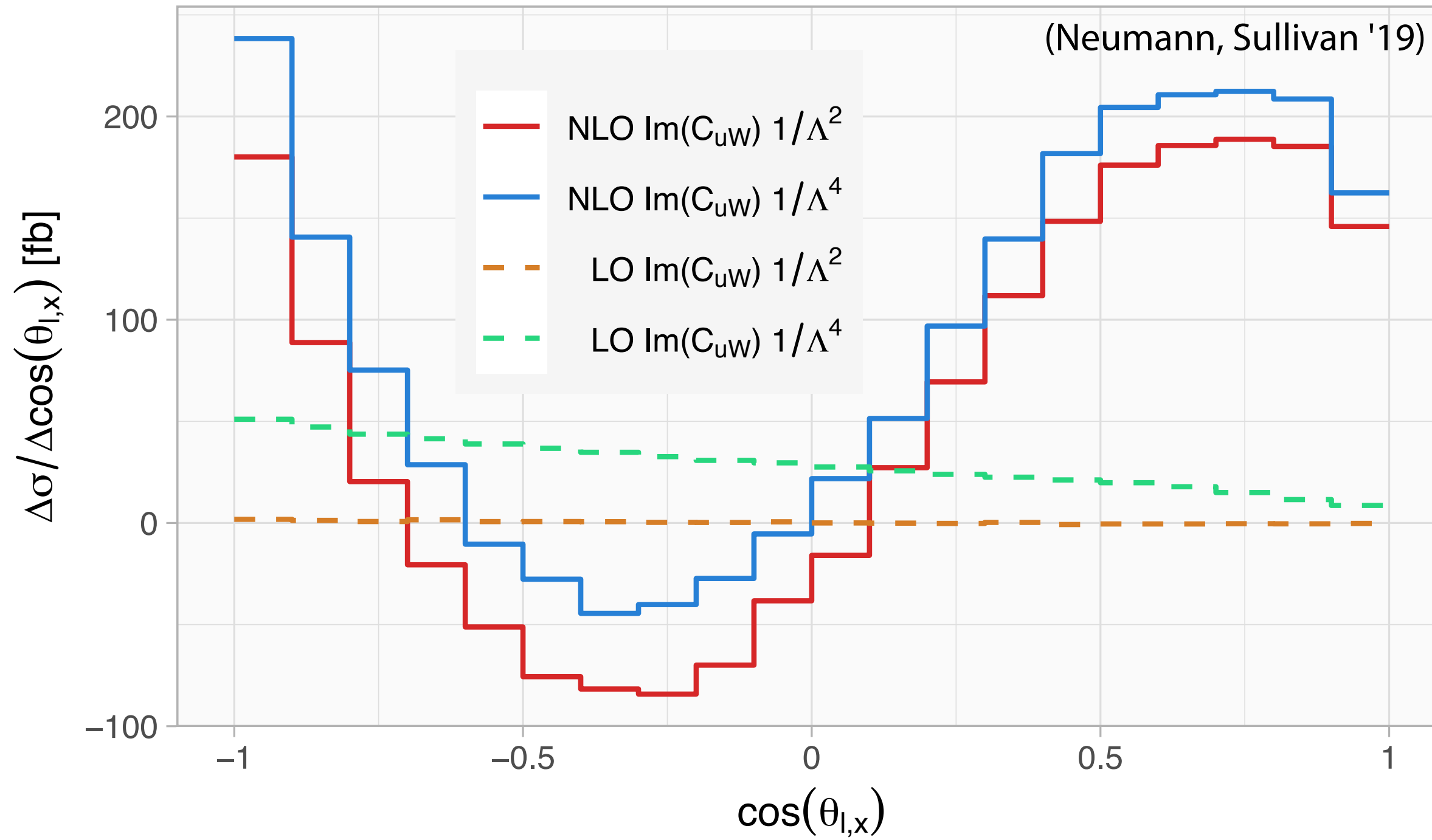
t-channel single-top-quark off-shell + SMEFT

Inclusively: Off-shell effects $\mathcal{O}(\Gamma_t/m_t)$



(for a decay study see also (Boughezal, Chen, Petriello, Wiegand '19))

Importance of NLO effects in the SMEFT



MCFM 9

Improvements in MCFM 9 (Neumann, Campbell '19)

automatic taucut fitting, power corrections, improved jettiness definition

Color singlets at NNLO, public in MCFM	
H	1605.08011: Boughezal, Campbell, Ellis, Focke, Giele, Liu, Petriello, Williams
Z (w. NLO EW)	
W^\pm	
ZH	
$W^\pm H$	1601.00658: Campbell, Ellis, Williams
γ	1612.04333: Campbell, Ellis, Williams
$\gamma\gamma$	1603.02663: Campbell, Ellis, Li, Williams
$Z\gamma$	1708.02925: Campbell, Neumann, Williams

NNLO benchmark results in MCFM-9.0

Process	nproc	τ_{cut} [GeV]	σ^{NLO}	σ^{NNLO}	fitted corr.	CPU time [h]
W^+	1	$6 \cdot 10^{-3} m_W$	4.221 nb	4.209 ± 0.005 nb	-27 ± 15 pb	7.6
W^-	6	$6 \cdot 10^{-3} m_W$	3.315 nb	3.275 ± 0.004 nb	-25 ± 10 pb	7.8
Z	31	$6 \cdot 10^{-3} m_Z$	885.3 pb	875.8 ± 0.9 nb	-3.5 ± 2.0 fb	13.0
H	112	$4 \cdot 10^{-3} m_H$	1.396 pb	1.872 ± 0.002 pb	7 ± 6 fb	9.7
$\gamma\gamma$	285	$1 \cdot 10^{-4} m_{\gamma\gamma}$	27.91 pb	43.54 ± 0.08 pb	0.36 ± 0.10 pb	83.2
W^+H	91	$3 \cdot 10^{-3} m_{W^+H}$	2.204 fb	2.262 ± 0.004 fb	0.002 ± 0.008 fb	16.0
W^-H	96	$3 \cdot 10^{-3} m_{W^-H}$	1.491 fb	1.526 ± 0.003 fb	-0.005 ± 0.007 fb	13.0
ZH	110	$3 \cdot 10^{-3} m_{ZH}$	0.753 fb	0.842 ± 0.001 fb	-0.005 ± 0.003 fb	12.5
$Z\gamma$	300	$3 \cdot 10^{-4} m_{Z\gamma}$	434 fb	525.5 ± 1.0 fb	4.5 ± 1.7 fb	202.5

Large performance improvements through implementation of subleading terms in τ -factorization

(Moult, Rothen, Stewart, Tackmann, Zhu '17; Ebert, Moult, Stewart, Tackmann, Vita, Zhu '18;) and (Boughezal, Liu, Petriello '16;

Boughezal, Isgro, Petriello '18)