

Black Holes-Saturons Correspondence

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MPP & LMU

IMPRS Young Scientist Workshop

May, 2022

Based on: Phys. Rev. D 105, 056013 (2022)



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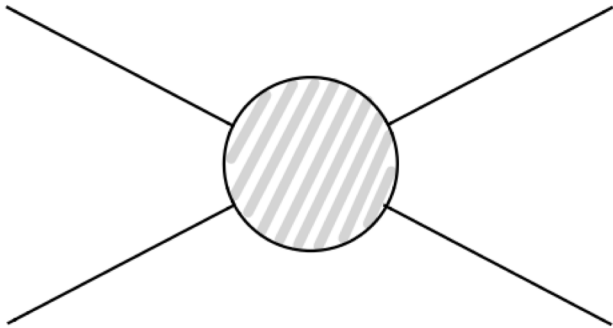
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Recap: Entropy Bound

Unitarity

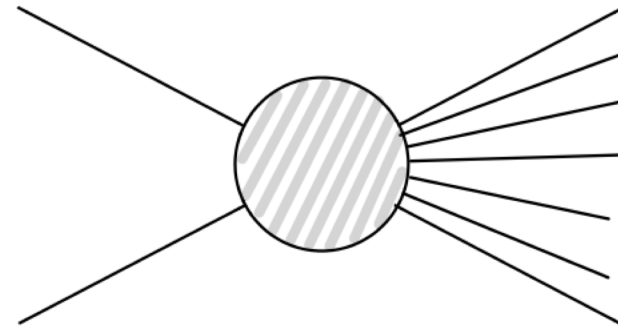
Cross-section:

$$\sigma_{2 \rightarrow 2}$$



Cross-section:

$$\sigma_{2 \rightarrow n} n_{st}$$



[1] G. Dvali, Entropy Bound and Unitarity of Scattering Amplitudes, JHEP03(2021) 126, arXiv:2003.05546.

[2] G. Dvali, Bounds on Quantum Information Storage and Retrieval, Philosophical Transactions of Royal Society A, arXiv:2107.10616

Entropy Bound

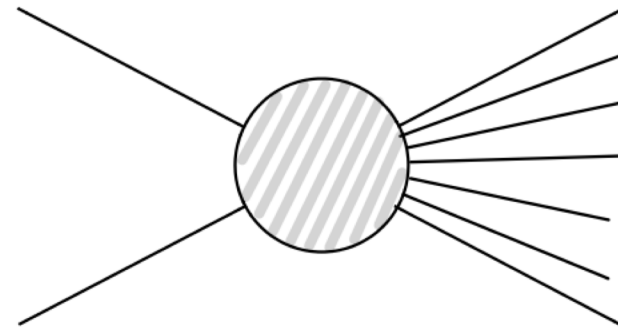
Unitarity

$$\sigma_{2 \rightarrow n} n_{st} \sim e^{-1/\alpha} e^S$$

imposes the following bound on the entropy

$$S_{max} = \frac{1}{\alpha} = \frac{Area}{G_{Gold}}$$

$$\sigma_{2 \rightarrow n} n_{st}$$



Entropy Bound Imposed by Unitarity

- Consider self-sustained object of size R in d space-time dimensions,

$$S \leq \frac{\text{Area}}{G_{Gold}}$$

- $\text{Area} \sim R^{d-2}$
- G_{Gold} is the Goldstone coupling

- Alternatively, the *bound* is

$$S \leq \frac{1}{\alpha}$$

- α is as an effective running coupling of the theory evaluated at the scale $1/R$.

Entropy Bound Imposed by Unitarity

- The maximal entropy compatible with unitarity is

$$S_{\text{Max}} = \frac{1}{\alpha} = \frac{\text{Area}}{f^{-2}},$$

where $f \equiv G_{\text{Gold}}^{-1/2}$ is the canonically normalized *Goldstone decay constant*.

Saturons

- We refer to the objects saturating the entropy bounds as ***Saturons***.
- Different saturons are discussed in [1-3]. These include:
 - Certain solitons, instantons, baryons, oscillons,
 - ***Black Holes***,
 - Lumps of classical fields,
 - ***Vacuum Bubbles***.

[1] G. Dvali, Entropy Bound and Unitarity of Scattering Amplitudes, JHEP03(2021) 126, arXiv:2003.05546.

[2] G. Dvali, Bounds on Quantum Information Storage and Retrieval, PTRS A, arXiv:2107.10616

[3] G. Dvali and O. Sakhelashvili, Black-Hole-Like Saturons in Gross-Neveu, arXiv:2111.03620.

Saturons

Saturons share the following universal properties [1-2]

- Their entropy satisfies the [area law](#):
$$S_{\text{Max}} = \frac{1}{\alpha} = \frac{\text{Area}}{f^{-2}}$$

- If unstable, they decay with a rate which gives the thermal rate of the [temperature](#)

$$T \sim \frac{1}{R},$$

up to $1/S$ corrections.

- In semi-classical treatment, they exhibit an [information horizon](#).
- The minimal time-scale required for the start of [the information retrieval](#) is

$$t_{\min} = \frac{\text{Volume}}{G_{\text{Gold}}} = \frac{R}{\alpha} = S_{\max} R$$

where $\text{Volume} \sim R^d$.

[1] G. Dvali, Entropy Bound and Unitarity of Scattering Amplitudes, JHEP03(2021) 126, arXiv:2003.05546.

[2] G. Dvali, Bounds on Quantum Information Storage and Retrieval, PTRS A, arXiv:2107.10616

Saturons as Black Holes

Black Holes properties

- Their entropy satisfies the [area law](#): [1]

$$S \sim \frac{\text{Area}}{M_p^{-2}} \sim \frac{\text{Area}}{f^{-2}}$$

- $f \sim M_p$

- Decay rate thermal rate of [temperature](#)

$$T \sim \frac{1}{R},$$

up to $1/S$ corrections.

- In semi-classical treatment, they exhibit an [information horizon](#).
- The minimal time-scale required for the start of [the information retrieval](#) is [2]

$$t_{min} = SR \sim \frac{R^3}{M_p^{-2}} \sim \frac{\text{Volume}}{M_p^{-2}}$$

[1] S. W. Hawking, Particle creation by black holes, Commun. Math. Phys. 43, 199 (1975).

[2] D. N. Page, Information in Black Hole Radiation, Phys. Rev. Lett. 71, 3743 (1993).

Saturons

- Saturons share the following universal properties [1-2]

- **Their entropy satisfies the *area law*:**

$$S_{\text{Max}} = \frac{1}{\alpha} = \frac{\text{Area}}{f^{-2}}$$

TODAY

- If unstable, they decay with a rate which gives the thermal rate of the *temperature*

$$T \sim \frac{1}{R},$$

up to $1/S$ corrections.

- **In semi-classical treatment, they exhibit an *information horizon*.**

TODAY

- The minimal time-scale required for the start of the information retrieval is bounded from below by,

$$t_{\text{min}} = \frac{\text{Volume}}{G_{\text{Gold}}} = \frac{R}{\alpha} = S_{\text{max}}R$$

where $\text{Volume} \sim R^d$.

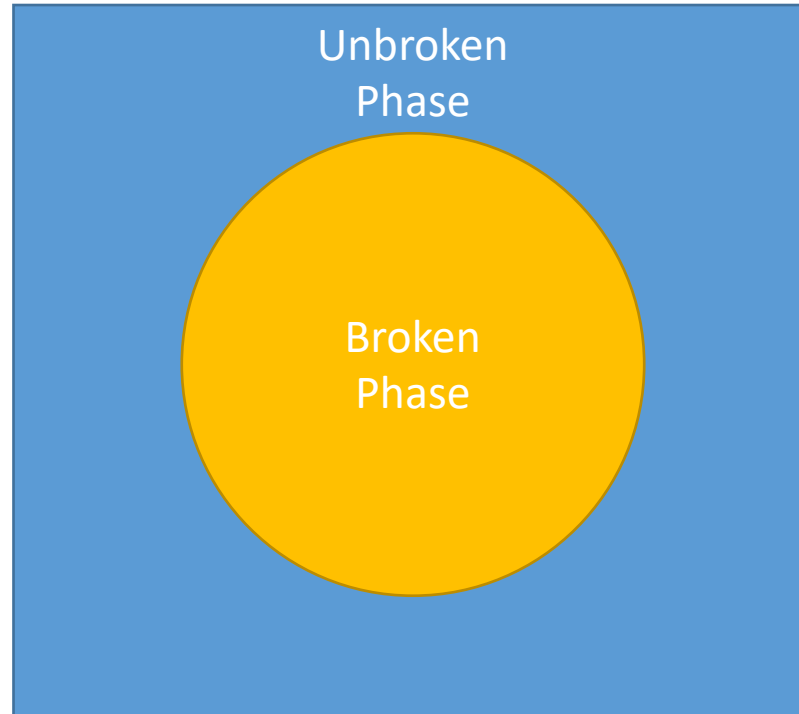
[1] G. Dvali, Entropy Bound and Unitarity of Scattering Amplitudes, JHEP03(2021) 126, arXiv:2003.05546.

[2] G. Dvali, Bounds on Quantum Information Storage and Retrieval, PTRS A, arXiv:2107.10616

Saturn as a Vacuum Bubble

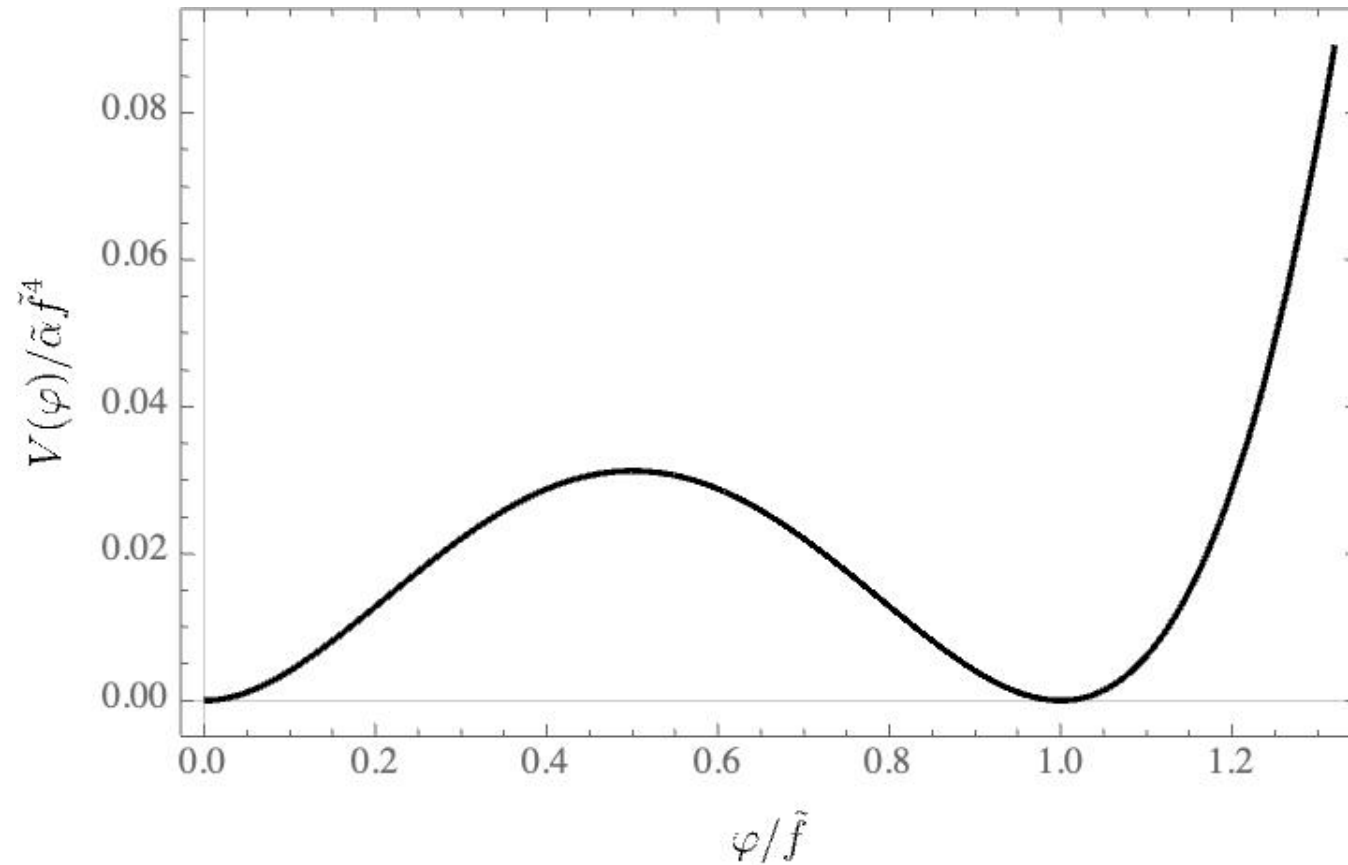
A Model

Vacuum Bubbles



$$SU(N) \rightarrow SU(N - 1) \times U(1)$$

Model of a Saturnon as a Vacuum Bubble

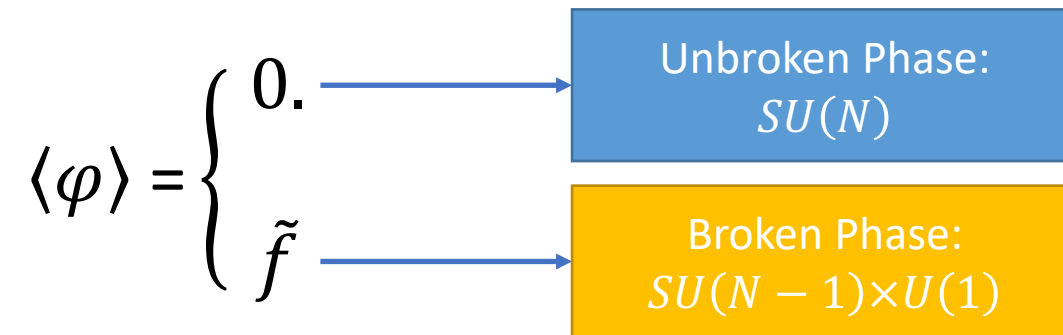


$$V(\varphi) = \frac{\tilde{\alpha}}{2} \varphi^2 (\tilde{f} - \varphi)^2$$

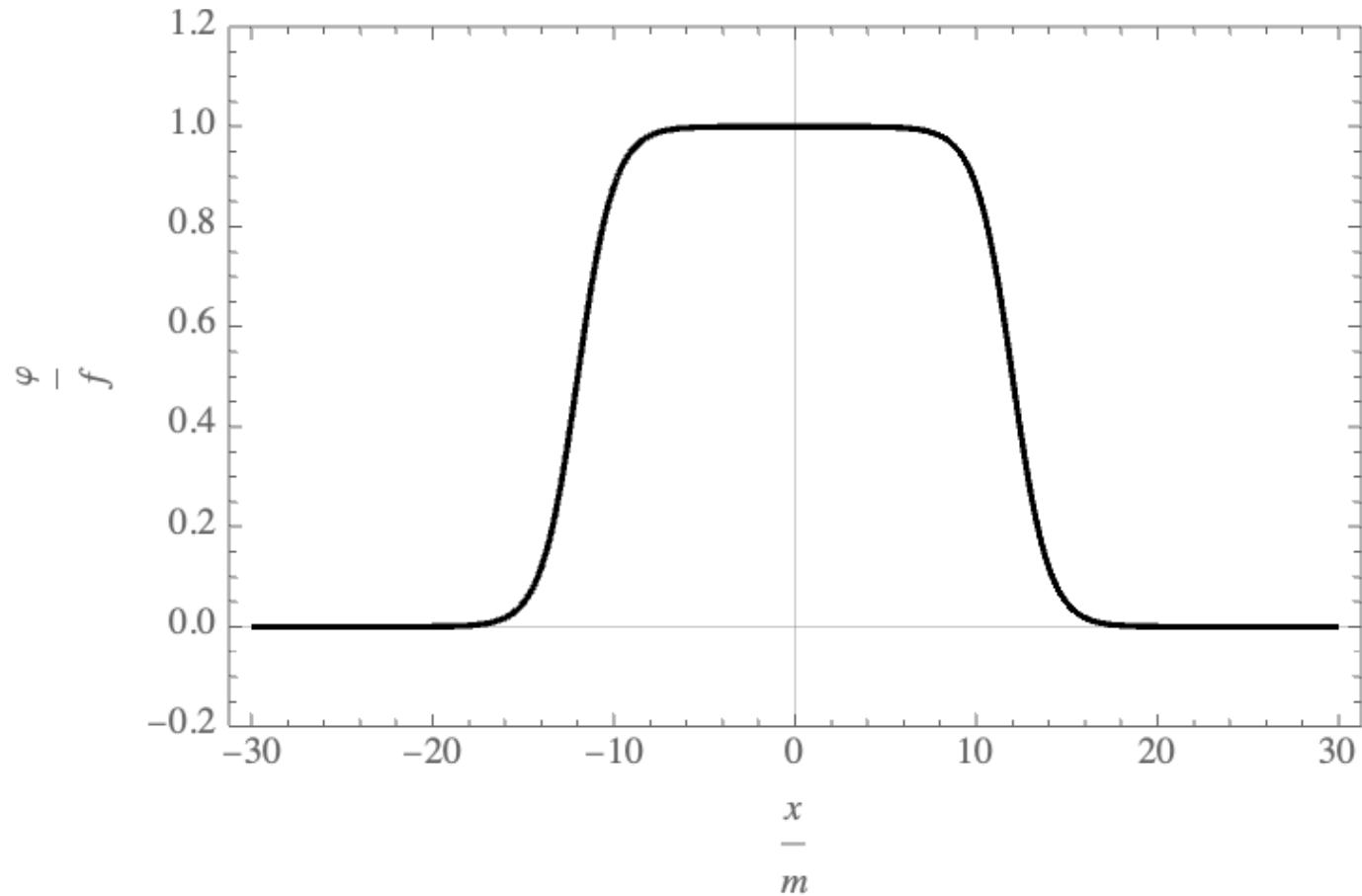
Model of a Saturnon as a Vacuum Bubble

$$V(\varphi) = \frac{\tilde{\alpha}}{2} \varphi^2 (\tilde{f} - \varphi)^2$$

is minimized by



Vacuum Bubbles: Thin Wall Approximation ($R \gg m^{-1}$)



$$\varphi(r) = \frac{\tilde{f}}{2} \left(1 + \tanh \left(\frac{m(R-r)}{2} \right) \right)$$

$$m = \frac{1}{16}$$
$$R = 12$$

Vacuum Bubbles Stabilization

A Memory Burden Effect

Vacuum Bubbles Stabilization

- Consider the ansatz

$$\phi_{\beta}^{\alpha} = (U^{\dagger} \Phi_{\text{D}} U)_{\beta}^{\alpha},$$

where

$$\Phi_{\text{D}} = \frac{\varphi(r)}{\sqrt{N(N-1)}} \text{diag}((N-1), -1, \dots, -1)$$

$$U = \exp[-i\theta^a T^a].$$

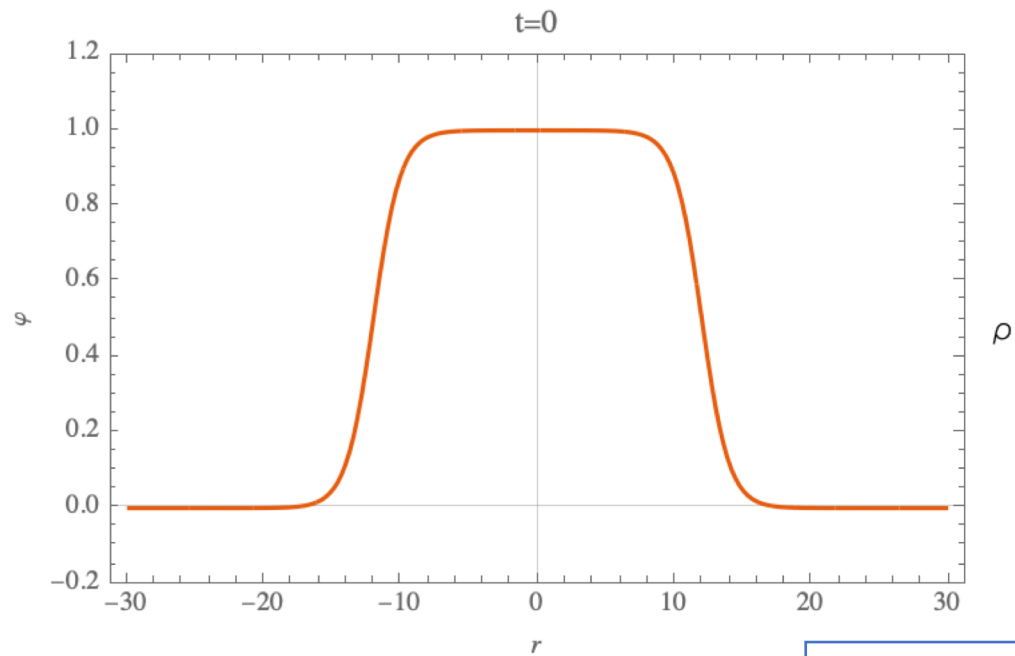
- Here T^a corresponds to the respective broken generators of $SU(N)$ [6].

$$\theta^a = \delta^{a1} \omega t$$

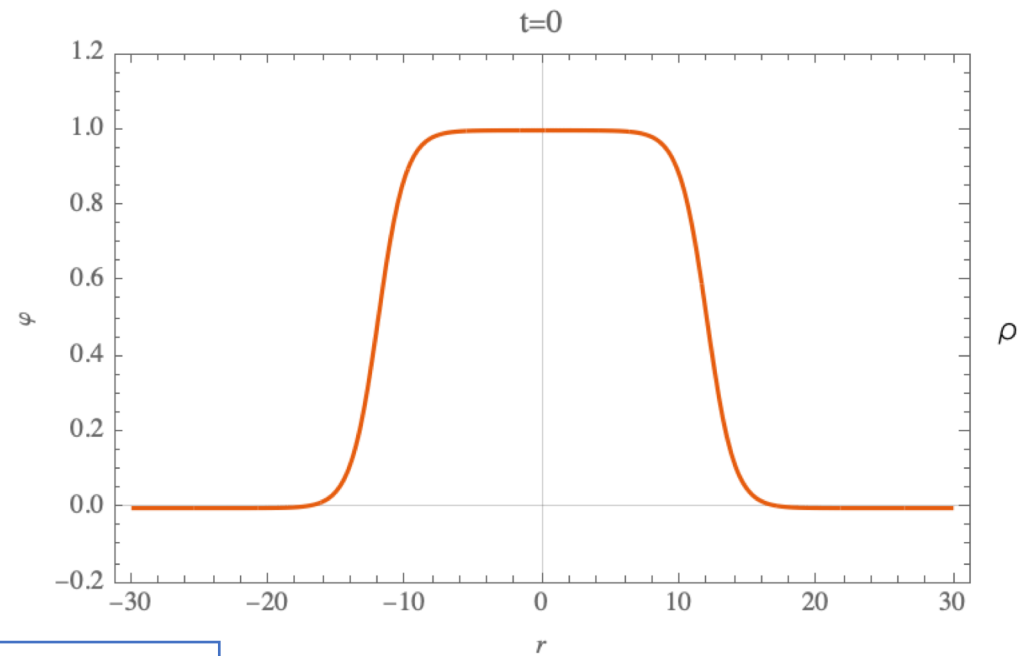
[6] G. Dvali, O. Kaikov, and J. S. V. Bermudez, (2021), arXiv:2112.00551 [hep-th]

Vacuum Bubbles Stabilization:

$$\dot{\theta} = 0$$



$$\dot{\theta} = \omega$$



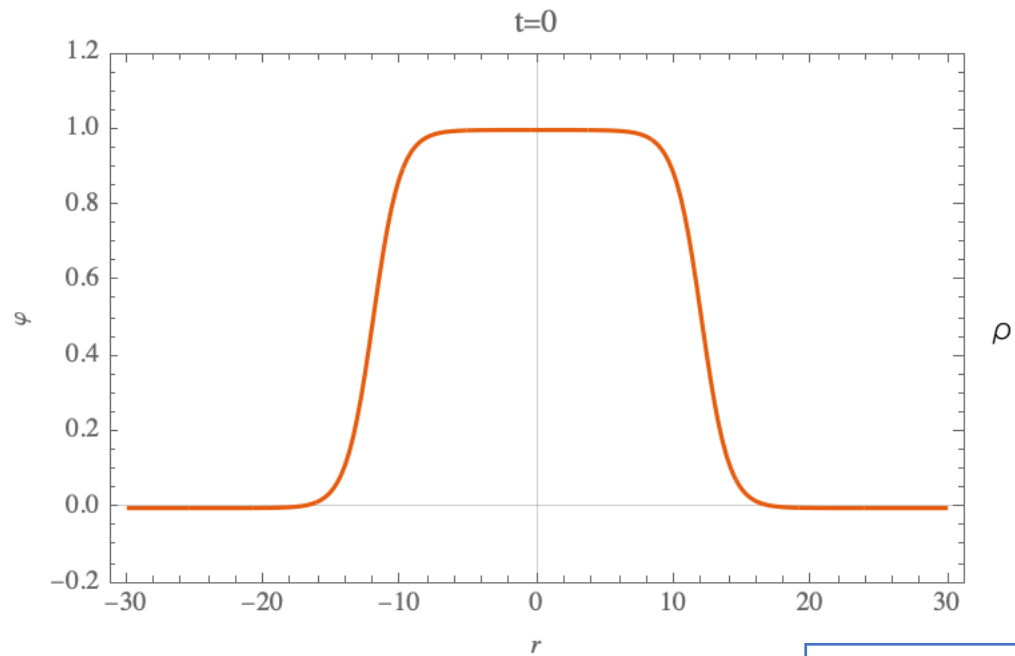
Thin wall approximation for:

$$\omega \approx 0.24 m,$$

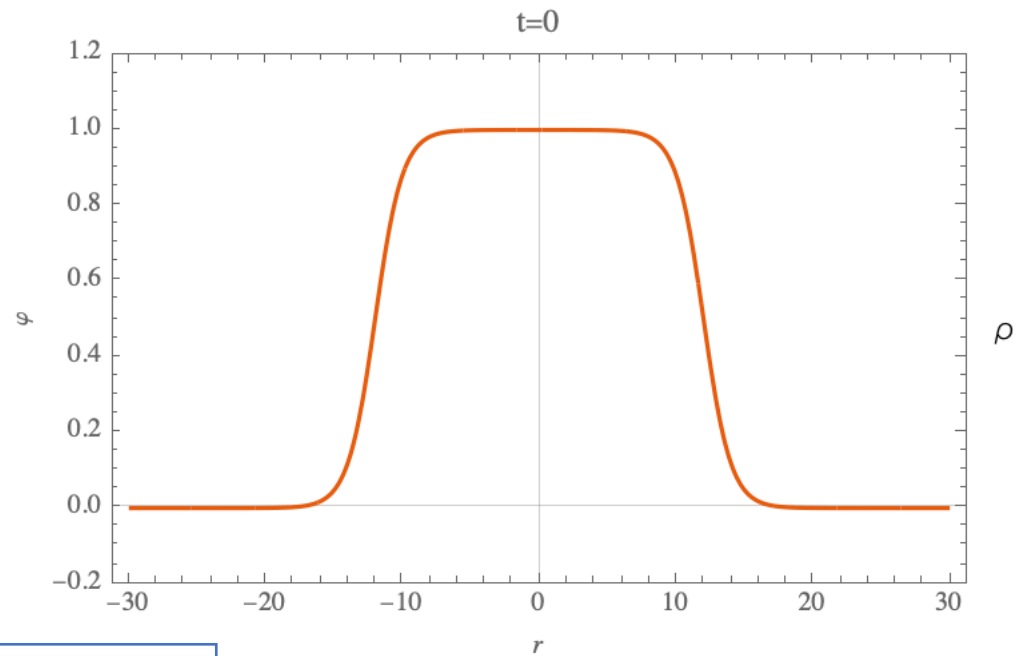
$$R_\omega = \frac{12}{m}$$

Vacuum Bubbles Stabilization:

$$\dot{\theta} = 0$$



$$\dot{\theta} = \omega$$



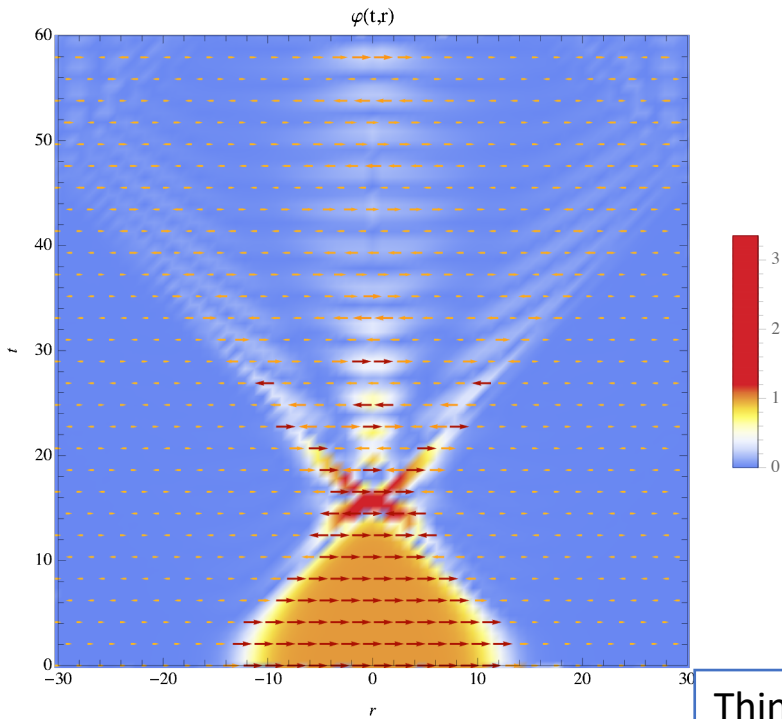
Thin wall approximation for:

$$\omega \approx 0.24 m,$$

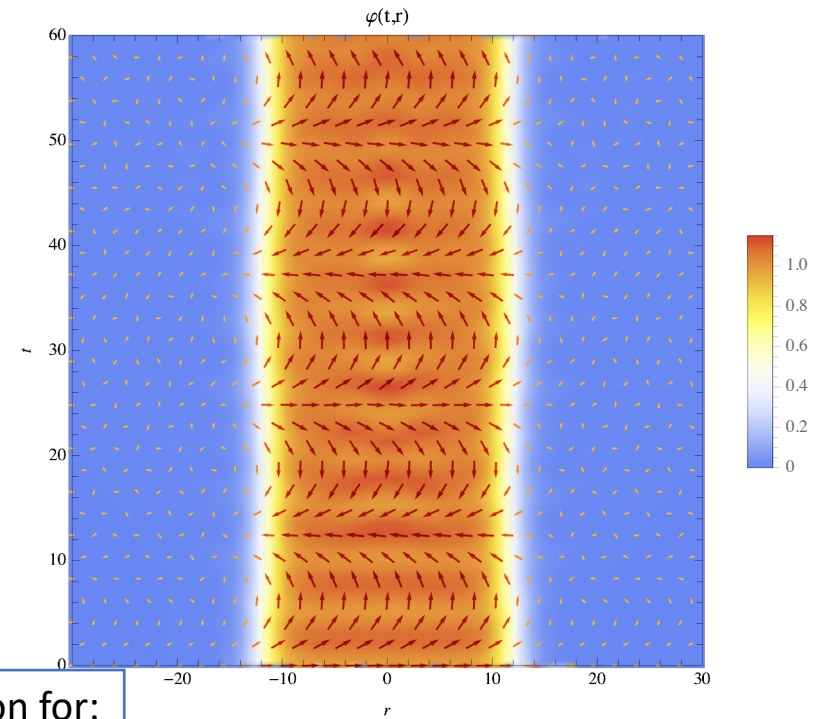
$$R_\omega = \frac{12}{m}$$

Vacuum Bubbles Stabilization:

$$\dot{\theta} = 0$$



$$\dot{\theta} = \omega$$



Thin wall approximation for:

$$\omega \approx 0.24 m,$$

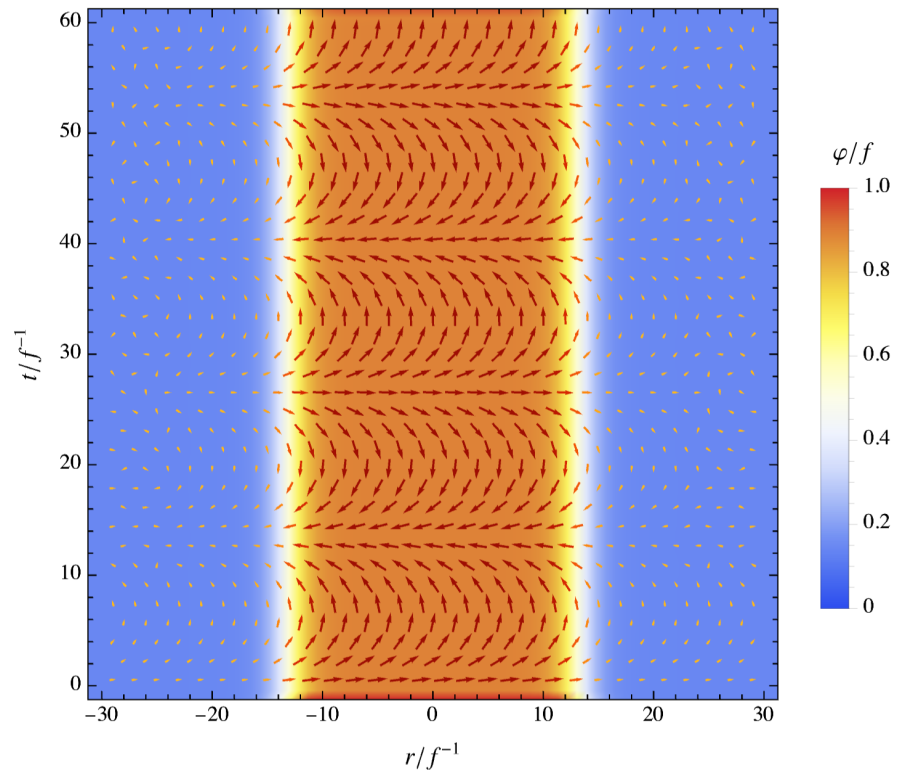
$$R_\omega = \frac{12}{m}$$

Vacuum Bubbles Stabilization

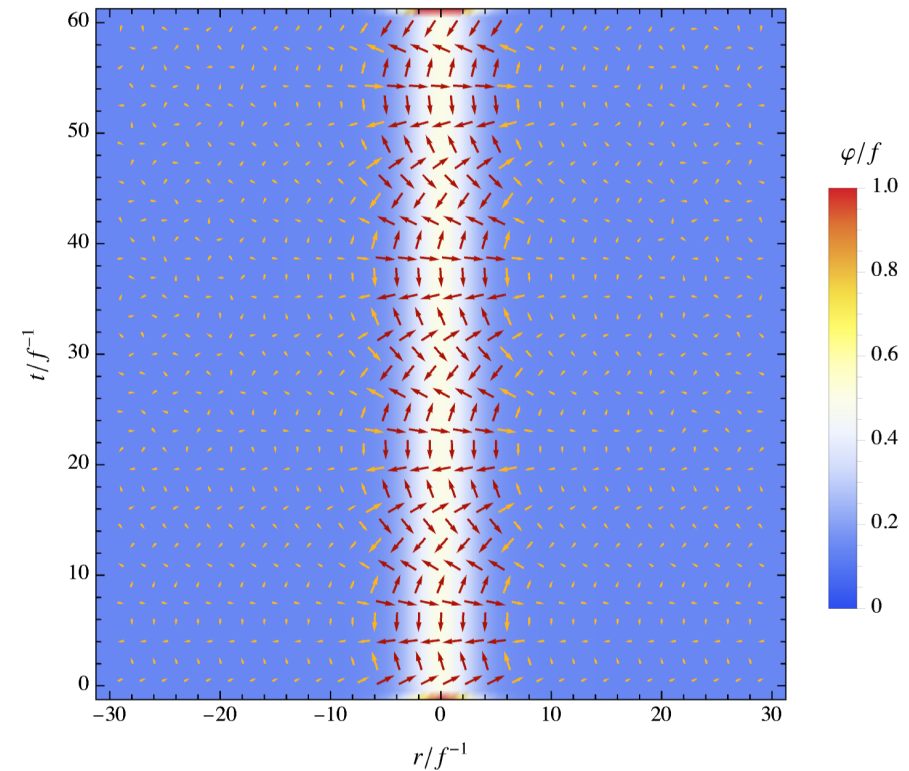
- *A stationary bubble is obtained thanks to the excitations of the) Goldstone mode(s).*
- The bubble is stable because of two factors:
 - 1) The Goldstone $SU(N)$ charge is conserved
 - 2) The same amount of charge in the exterior vacuum would cost much higher energy.

Vacuum Bubbles Stabilization

Large Bubbles
 $R \sim 12m^{-1}$



Small Bubbles
 $R \sim 1.02m^{-1}$



Vacuum Bubbles Micro-state Entropy

Large Bubbles

$$\omega \ll m$$
$$m^{-1} \ll R$$

- $N_G \gg \frac{1}{\alpha} \sim N \sim N_{Gold}$
- $1 \gg \lambda$
- $S \approx 2N \ln \left(\frac{e}{\lambda} \right) \sim \frac{1}{\alpha} \ln \left(\frac{m^{10}}{\omega^{10}} \right)$
- $S \ll S_{max} \sim \frac{1}{\alpha} \frac{m^6}{\omega^6}$

Small Bubbles

$$\omega \sim m$$
$$m^{-1} \sim R$$

- $N_G \sim \frac{1}{\alpha} \sim N \sim N_{Gold}$
- $1 \sim \lambda$
- $S \sim \frac{1}{\alpha} \sim E_{Bubble} R$
- $S \sim S_{max}$

[8] J. D. Bekenstein, Universal Upper Bound on the Entropy-to-Energy Ratio for Bounded Systems, Phys. Rev. D23no. 2 (1981), 287-298.47

Vacuum Bubbles Micro-state Entropy

Large Bubbles

$$\omega \ll m$$
$$m^{-1} \ll R$$

- $N_G \gg \frac{1}{\alpha} \sim N \sim N_{Gold}$
- $1 \gg \lambda$
- $S \approx 2N \ln\left(\frac{e}{\lambda}\right) \sim \frac{1}{\alpha} \ln\left(\frac{m^{10}}{\omega^{10}}\right)$
- $S \ll S_{max} \sim \frac{1}{\alpha} \frac{m^6}{\omega^6}$

Small Bubbles -> **Saturons**

$$\omega \sim m$$
$$m^{-1} \sim R$$

- $N_G \sim \frac{1}{\alpha} \sim N \sim N_{Gold}$
- $1 \sim \lambda$
- $S \sim \frac{1}{\alpha} \sim E_{Bubble}R$
- $S \sim S_{max}$

[8] J. D. Bekenstein, Universal Upper Bound on the Entropy-to-Energy Ratio for Bounded Systems, Phys. Rev. D23no. 2 (1981), 287-298.47

Information Horizon

Saturons in semiclassical limit

Information Horizon

- A universal property of the saturons is that in semi-classical limit, they possess a strict *information horizon*.
- The information stored in the interior of the saturon can not be extracted in any form.
- The general physical reason is that in this limit the *memory modes*, that carry quantum information, *decouple*.

Semi-classical Limit

- The limit in which the classical bubble solution experiences **no back reaction** from quantum fluctuations

$$\alpha \rightarrow 0, \quad R = \text{finite}, \quad \omega = \text{finite}, \quad \alpha N = \text{finite}.$$

- Simultaneously

$$f \rightarrow \infty, \quad m = \text{finite}.$$

Semi-classical Limit

- The limit in which the classical bubble solution experiences **no back reaction** from quantum fluctuations

$$\alpha \rightarrow 0, \quad R = \text{finite}, \quad \omega = \text{finite}, \quad \alpha N = \text{finite}.$$

- Simultaneously

$$f \rightarrow \infty, \quad m = \text{finite}.$$

- Recall: For BH $f \sim M_p$

Semi-classical Limit

- The limit in which the classical bubble solution experiences **no back reaction** from quantum fluctuations

$$\begin{array}{llll} \alpha \rightarrow 0, & R = \text{finite}, & \omega = \text{finite}, & \alpha N = \text{finite}. \\ f \rightarrow \infty, & m = \text{finite}. & & \end{array}$$

- In this limit, the effective coupling of a Goldstone mode of frequency ε

$$\alpha_G = \frac{\varepsilon^2}{f^2}$$

goes to zero for any finite ε .

Goldstone Horizon

- **At finite f** , a Goldstone mode of frequency $\varepsilon \ll m$ cannot propagate outside the bubble, even though the coupling α_G is finite.

- Two cases

- The energy of an internal perturbation ε is such

$$\varepsilon < m,$$

propagation is impossible due to the finite energy gap.

- $\varepsilon \ll m$, the perturbation energy can exceed the mass gap at the expense of a large occupation number n_ε of Goldstone quanta.

$$n_\varepsilon \rightarrow 1,$$

such a process is exponentially suppressed by a factor e^{-n_ε}

Goldstone Horizon: An Example

- Lets consider a perturbation on a stable vacuum bubble, ϕ_{VB} ,

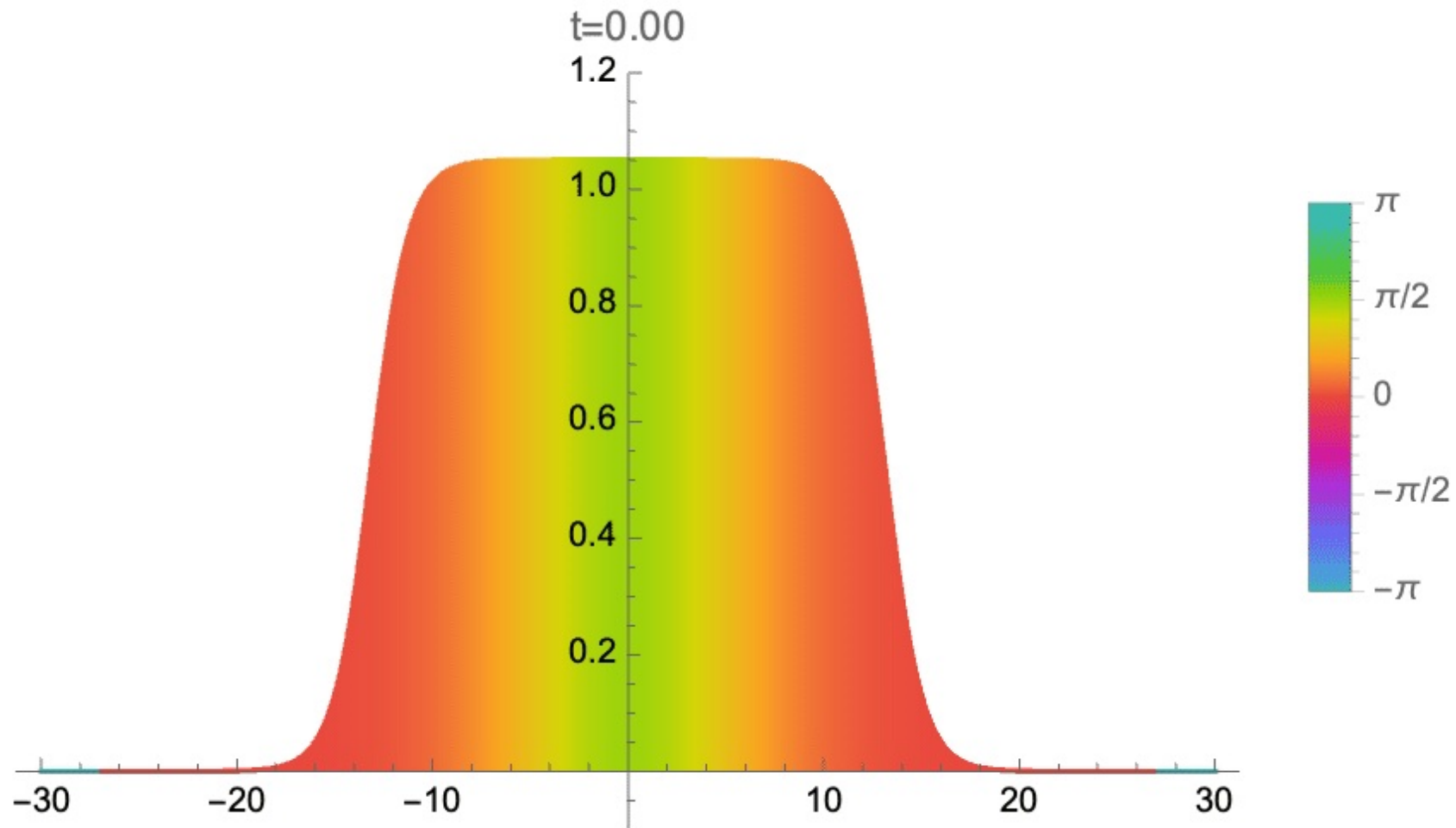
$$\phi_{\beta}^{\alpha} \Big|_{t=0} = p(r) \phi_{VB}^{\alpha}(r)$$

$$p(r) = \exp \left[\frac{i\pi f(r)}{2 f(0)} T^1 \right]$$

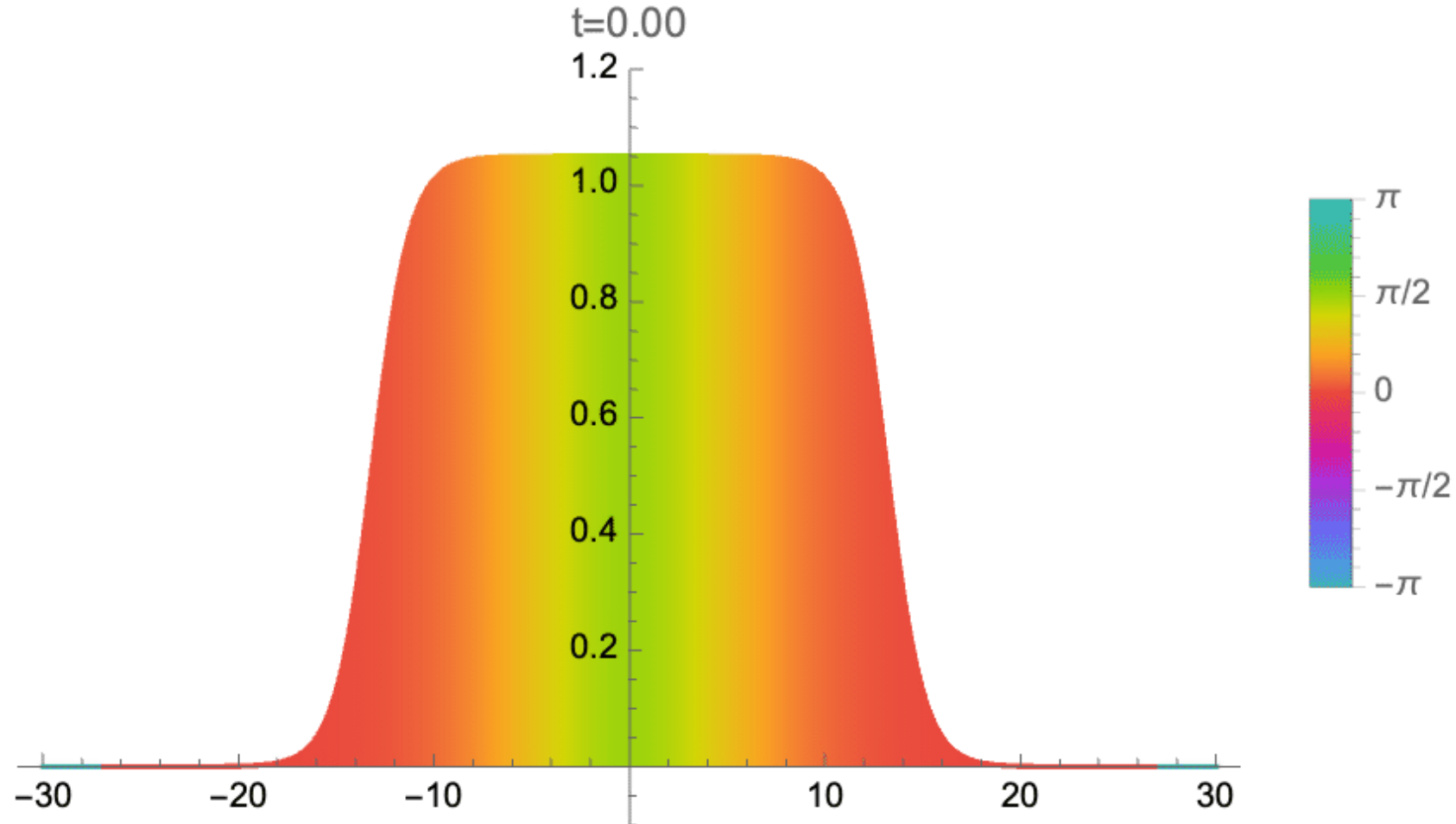
where

$$f(r) = \frac{1}{\sqrt{2\pi\sigma}} e^{-\frac{r^2}{2\sigma^2}}$$

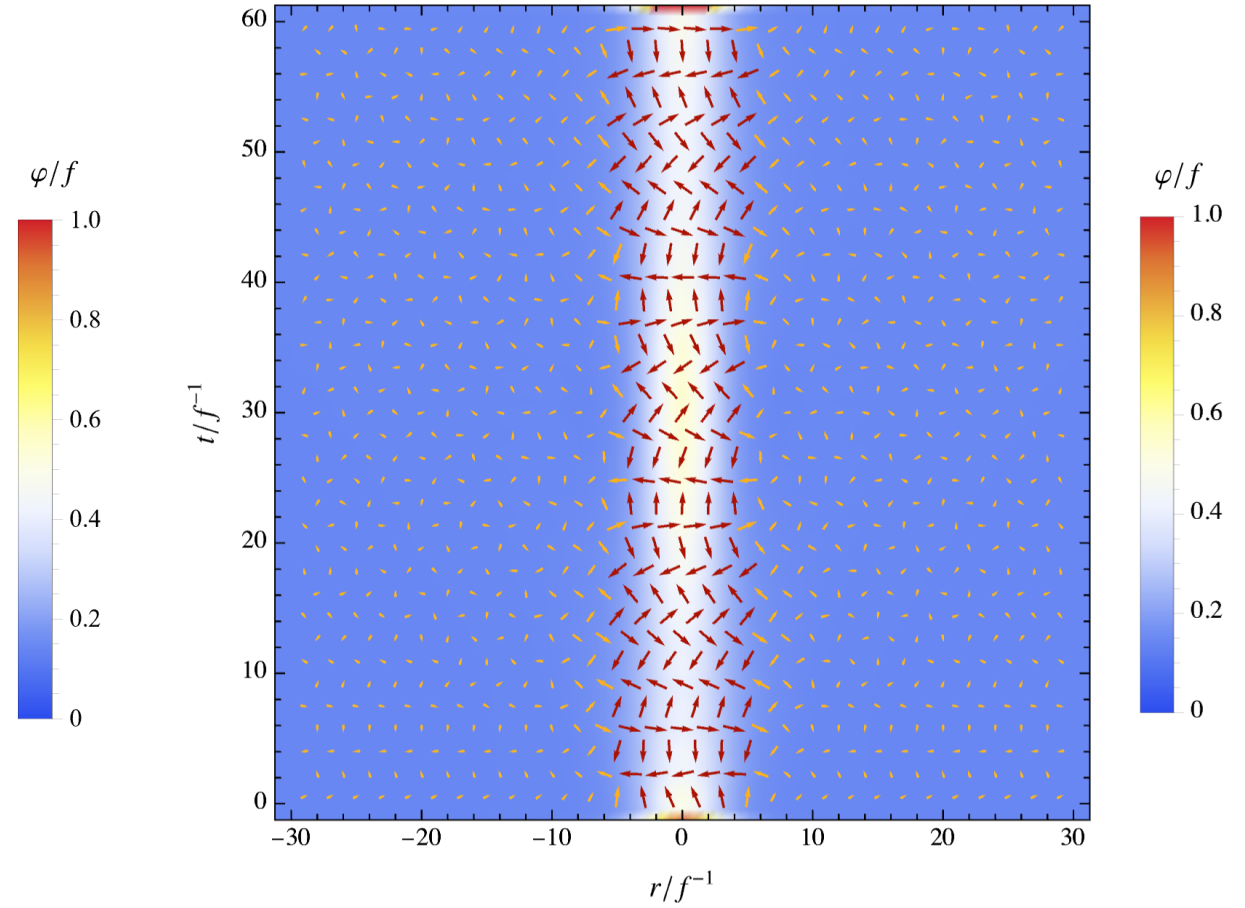
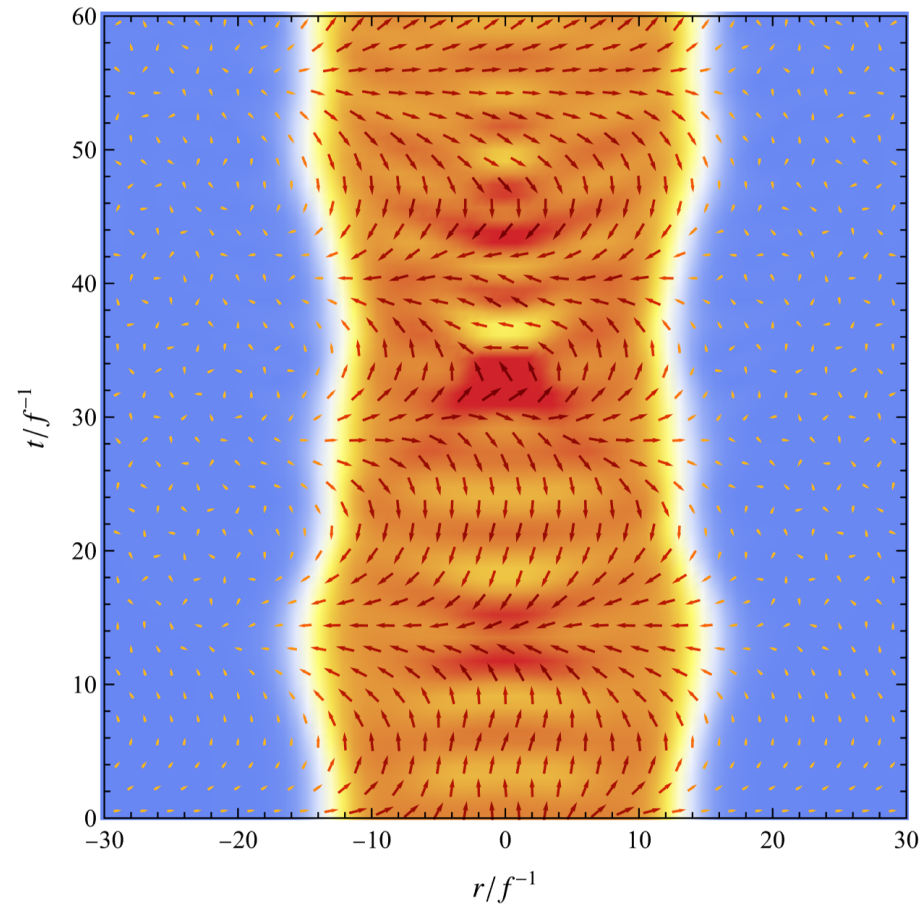
Goldstone Horizon: An Example



Goldstone Horizon: An Example



Goldstone Horizon



$$\sigma \equiv 5/m$$

Correspondence to Black Holes

Correspondence to Black Holes

Saturons

- $S = (fR)^2 = \alpha^{-1}$
- $T = R^{-1}$
- $t_{\min} = R^3 f^2 = SR$
- Information/Goldstone Horizon

Black Holes

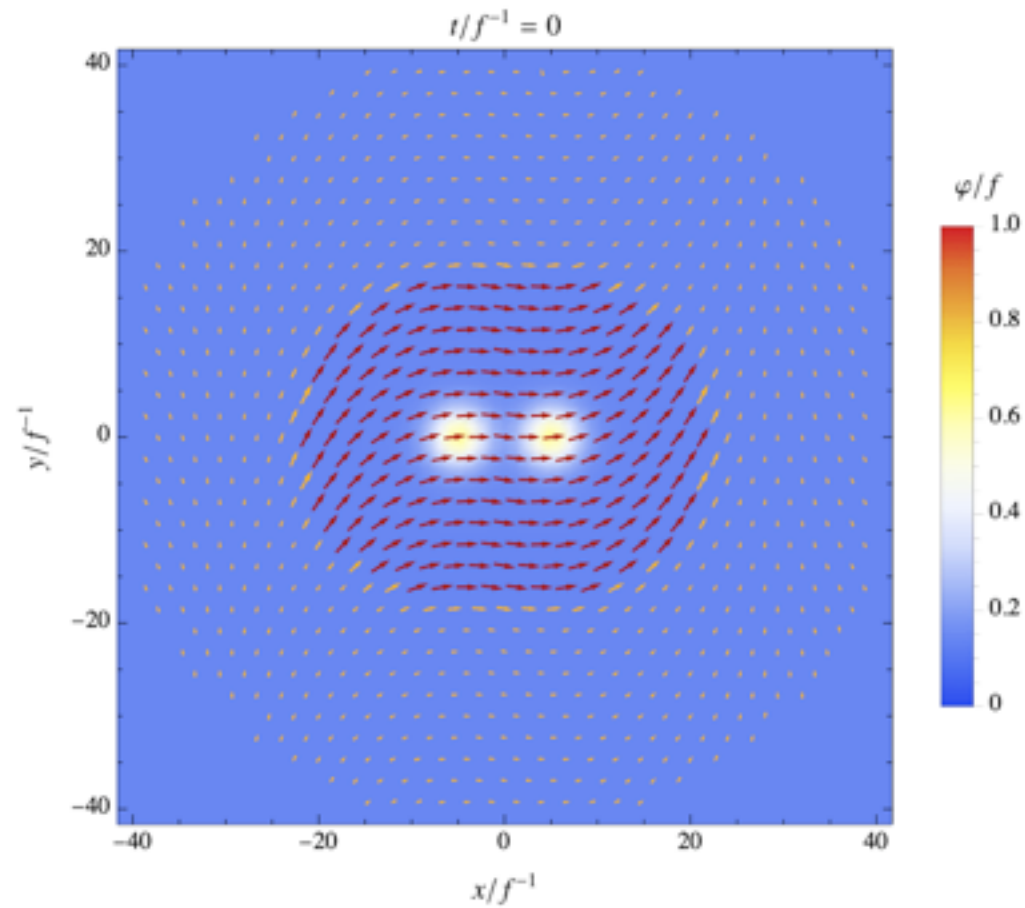
- $S = (M_P R)^2$
- $T = R^{-1}$
- $t_{\min} = R^3 M_P^2 = SR$
- Information Horizon

Outlook

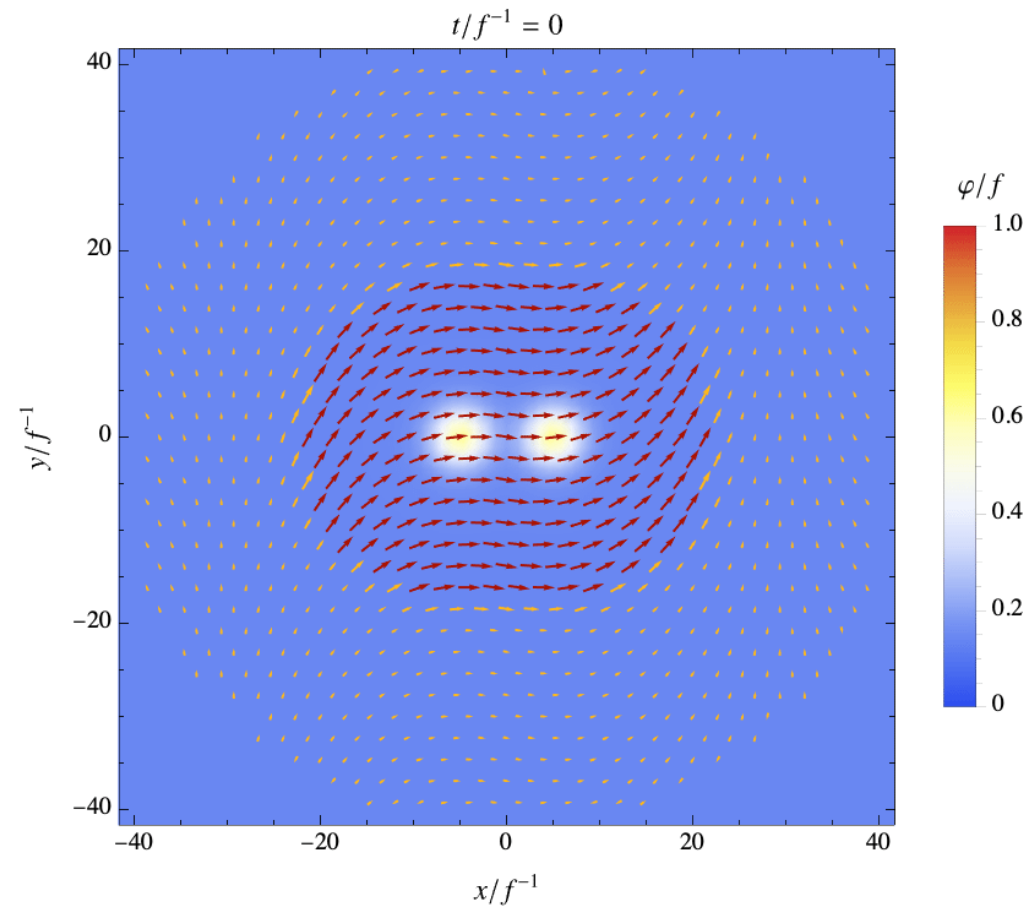
- Generalizations of our study are currently researched.
- Departures from semi-classical behavior can become observable for the black holes that are relatively old and close to their half-decay time.
- The light primordial black holes, provided they exist, can be within a potentially interesting window.
- Very recently some other possible observational consequences for rotating black holes were discussed in [7]

[7] G. Dvali, F. Kühnel and M. Zantedeschi, Vortexes in Black Holes, arXiv:hep-th/2112.08354.

Outlook



Outlook



Conclusions

- We have shown an explicit example of a **Saturon as a Vacuum Bubble**. Small bubbles saturate the entropy bound and correspond to the saturons of the theory.
- A large (macroscopic) occupation number of the Goldstone modes, parametrized by ω , stabilizes the Vacuum Bubbles. This phenomenon is due to the **memory burden effect and could be relevant for PBH**
- Saturons (Vacuum Bubbles) exhibit a **goldstone horizon**, analog to the information horizon of Black holes.

Conclusions

- Area law:

$$S_{\max} = \frac{1}{\alpha} = \frac{\text{Area}}{f^{-2}}$$

- Temperature

$$T \sim \frac{1}{R},$$

up to $1/S$ corrections.

- In semi-classical treatment, they exhibit an information horizon.
- Information retrieval after

$$t_{\min} = \frac{\text{Volume}}{G_{\text{Gold}}} = \frac{R}{\alpha} = S_{\max} R$$

Thank you

Appendix

Vacuum Bubbles Stabilization: Quantum picture of classical stability

- The energies in terms of the occupation numbers of the corresponding quanta are

$$E_{\text{int}} = \omega N_G, \quad \text{where,} \quad N_G \equiv \frac{1}{\alpha} \frac{m^5}{\omega^5} \left(\frac{16\pi}{81} \right),$$

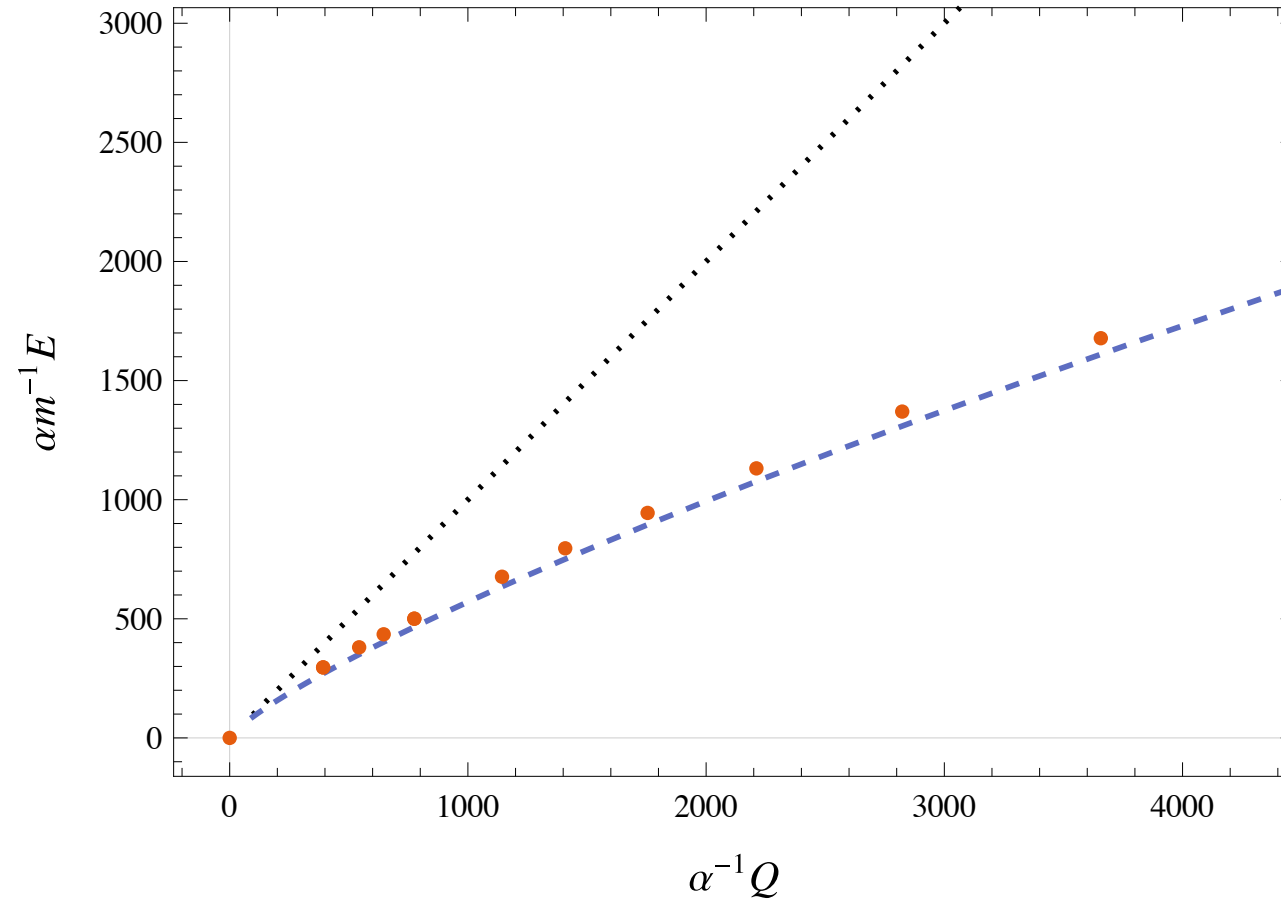
$$E_{\text{wall}} = m N_\varphi, \quad \text{where,} \quad N_\varphi \equiv \frac{1}{\alpha} \frac{m^4}{\omega^4} \left(\frac{8\pi}{27} \right).$$

- Observe that $\frac{N_G}{N_\varphi} = \frac{2}{3} \frac{m}{\omega} \gg 1$ for the thin walls.

Vacuum Bubbles Stabilization: Quantum picture of classical stability

- Conclusion: *A stationary bubble is obtained thanks to the excitations of the) Goldstone mode(s).*
- The bubble is stable because of two factors:
 - 1) The fact that the Goldstone $SU(N)$ charge is conserved; and
 - 2) The fact that the same amount of charge in the exterior vacuum would cost much higher energy.

Spectrum



Memory Burden Effect

Large amount of
Memory patterns

Stored quantum
information.

Slowdown of the
system's evolution

Vaccum Bubbles Stabilization

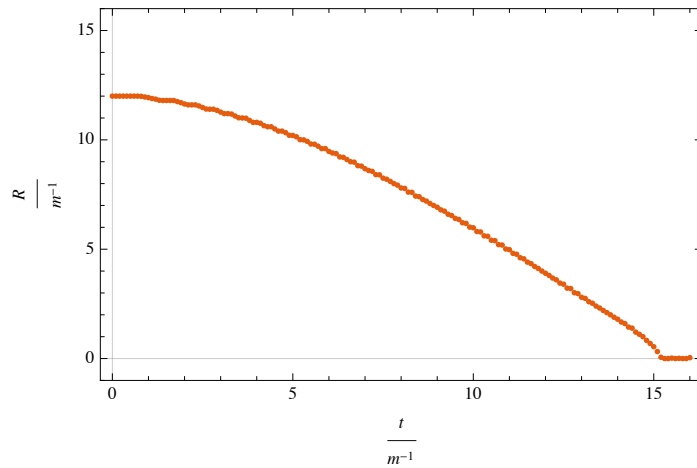
Large amount of
Bubble micro-states

Excitations of the
Goldstone modes

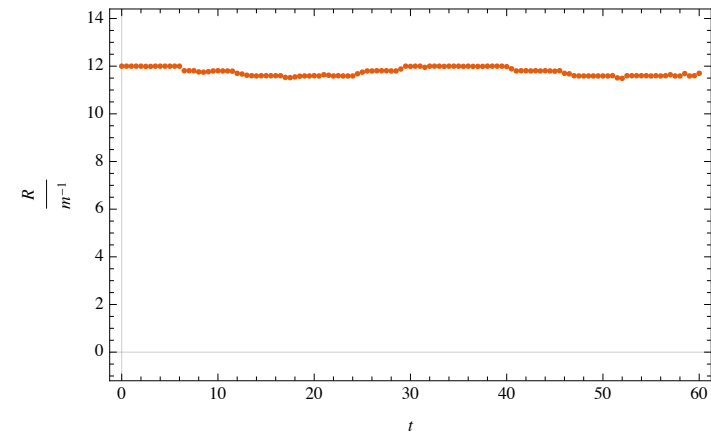
Slowdown of
bubble's decay

Vacuum Bubble

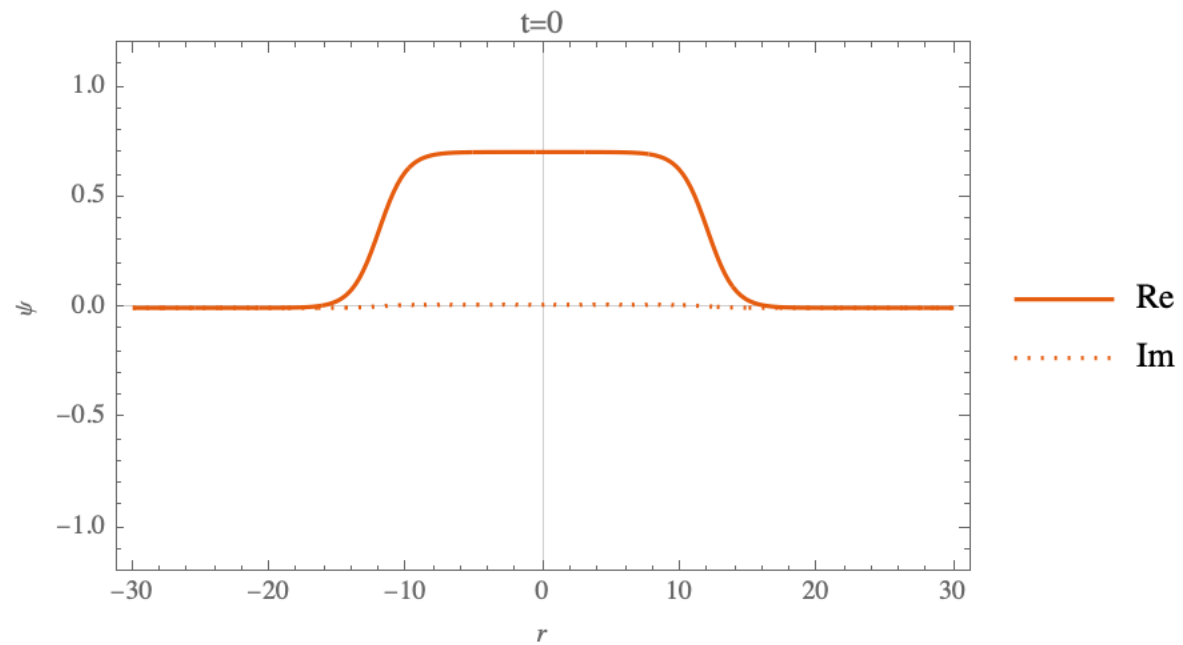
Decay
 $(\dot{\theta} = 0)$



Stabilization
 $(\dot{\theta} \neq 0)$

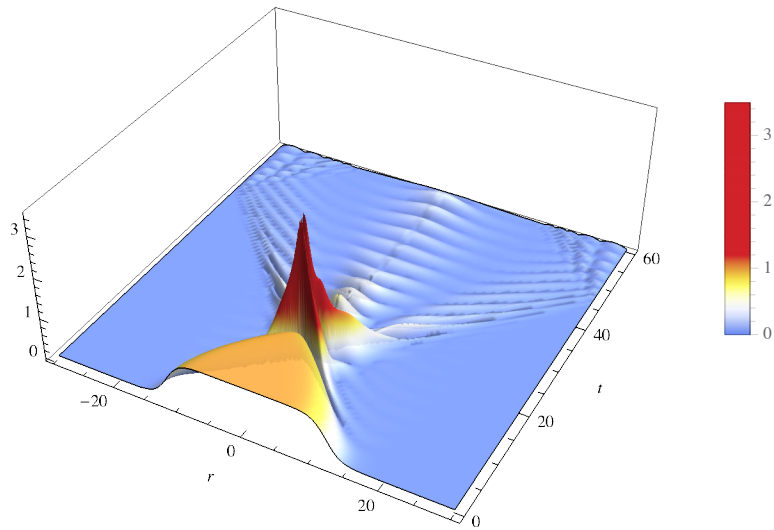


Stabilization

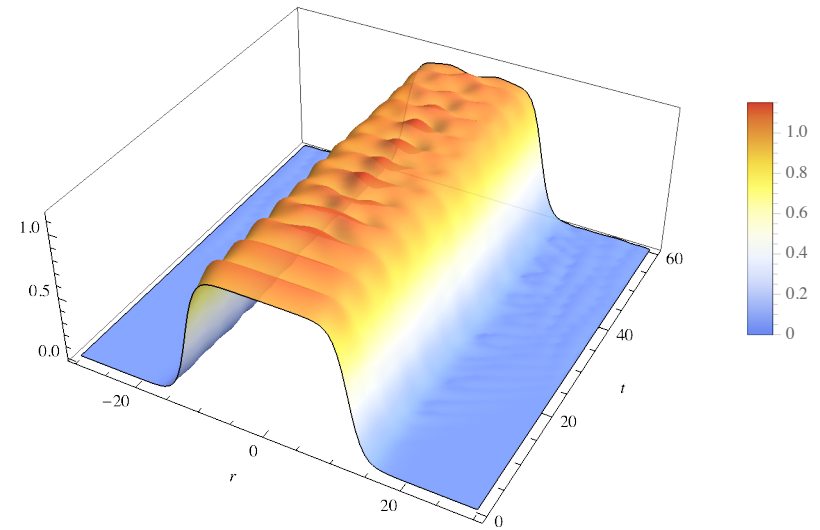


Vacuum Bubble Stabilization

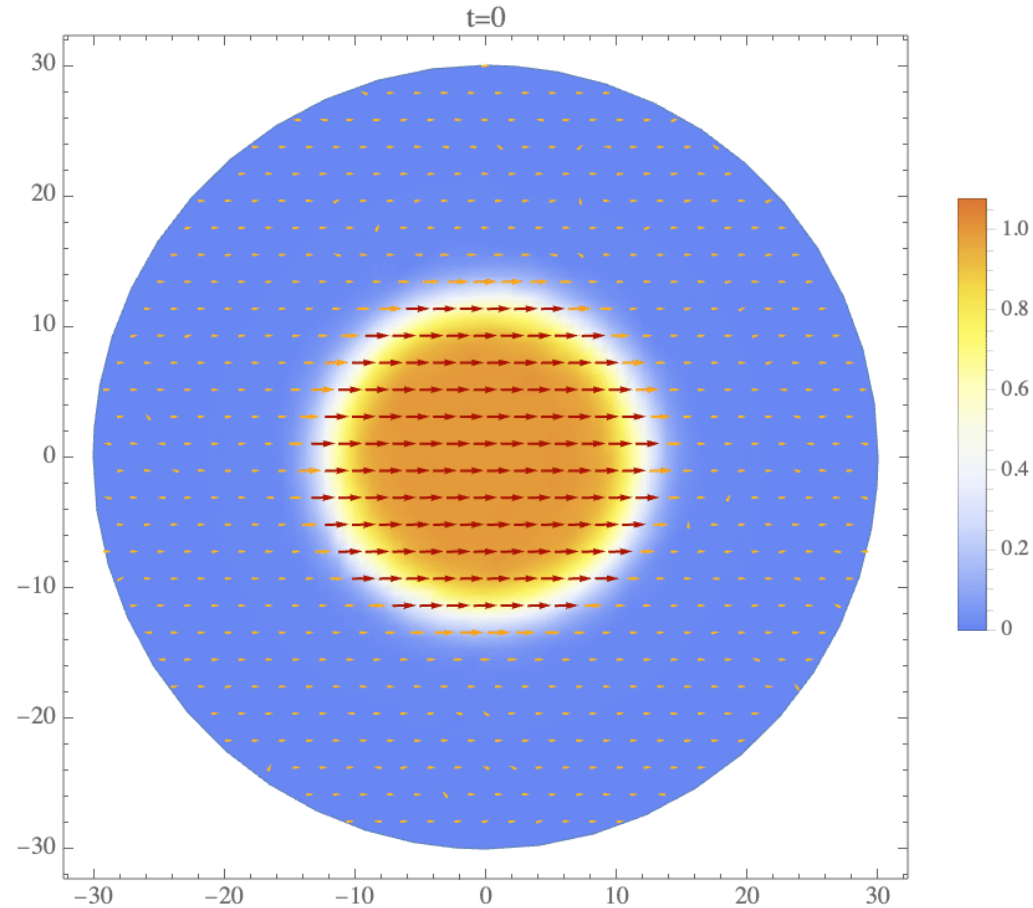
Decay
 $(\dot{\theta} = 0)$



Stabilization
 $(\dot{\theta} \neq 0)$



Vacuum Bubble Stabilization in $d=3$



Vacuum bubble Decay

