Theoretical Physics of the Early Universe Department of Physics Technical University of Munich



# False Vacuum Decay and Functional Determinants

## The Fubini Lipatov instanton and the lifetime of the SM

#### Matthias Carosi

Theoretical Physics of the Early Universe Department of Physics Technical University of Munich

IMPRS Young Scientists Workshop 2022



Tur Uhronturm



#### False Vacuum Decay

- Metastability in the quantum world
- The Higgs false vacuum decay
- How do we evaluate the vacuum decay rate?

#### 2 The Problem with Functional Determinants



In a quantum theory, tunneling from one minimum of the potential to another is possible due to quantum fluctuations.



- In a quantum theory, tunneling from one minimum of the potential to another is possible due to quantum fluctuations.
- If a local minimum is not the absolute minimum of the potential, it is therefore meta-stable and we refer to it as a **false vacuum**.



- In a quantum theory, tunneling from one minimum of the potential to another is possible due to quantum fluctuations.
- If a local minimum is not the absolute minimum of the potential, it is therefore meta-stable and we refer to it as a **false vacuum**. Example: scalar  $\varphi^4$  theory in euclidean spacetime

$$S[\varphi] = \int d^D x \left(\frac{1}{2}\partial^\mu \varphi(x)\partial_\mu \varphi(x) + V(\varphi(x))\right) \text{ with } V(\varphi) = -\frac{m^2}{2}\varphi^2 + \frac{g}{3}\varphi^3 + \frac{\lambda}{4}\varphi^4$$
 (1)



- In a quantum theory, tunneling from one minimum of the potential to another is possible due to quantum fluctuations.
- If a local minimum is not the absolute minimum of the potential, it is therefore meta-stable and we refer to it as a false vacuum.
   Example: scalar φ<sup>4</sup> theory in euclidean spacetime

$$S[\varphi] = \int d^D x \left(\frac{1}{2}\partial^\mu \varphi(x)\partial_\mu \varphi(x) + V(\varphi(x))\right) \text{ with } V(\varphi) = -\frac{m^2}{2}\varphi^2 + \frac{g}{3}\varphi^3 + \frac{\lambda}{4}\varphi^4$$
 (1)

The field in the false vacuum will eventually tunnel through the potential barrier and land on a more stable minimum. We call this phenomenon **false vacuum decay**.



Plot from https://arxiv.org/pdf/astro-ph/0005003.pdf



#### The Higgs potential



Plot from Bass, S.D., De Roeck, A. & Kado, M. The Higgs boson implications and prospects for future discoveries. Nat Rev Phys 3, 608–624 (2021). https://doi.org/10.1038/s42254-021-00341-2





#### The Higgs false vacuum decay Is the Standard Model unstable?

Plot from https://arxiv.org/pdf/1205.6497.pdf



#### ТШ

#### The Higgs false vacuum decay Is the Standard Model unstable?

Current measurements suggest that the Higgs field sits in a metastable vacuum. Plot from https://arxiv.org/pdf/1205.6497.pdf



#### ТШ

#### The Higgs false vacuum decay Is the Standard Model unstable?

- Current measurements suggest that the Higgs field sits in a metastable vacuum.
- An estimate for the lifetime of the Standard Model, namely for the decay rate of the Higgs field, can be obtained through known methods.



### ТШ

#### The Higgs false vacuum decay Is the Standard Model unstable?

- Current measurements suggest that the Higgs field sits in a metastable vacuum.
- An estimate for the lifetime of the Standard Model, namely for the decay rate of the Higgs field, can be obtained through known methods.
- However, these rely on approximations not entirely valid in this case.





The decay rate is given by the imaginary part of the ground state energy in the false vacuum.



The decay rate is given by the imaginary part of the ground state energy in the false vacuum.

$$Z = \sum_{n} e^{-\beta E_{n}} \longrightarrow E_{0} = -\lim_{\beta \to \infty} \frac{1}{\beta} \log Z , \quad \beta = \text{Euclidean time}$$
(2)

To evaluate it, we can deploy a technique developed by Callan & Coleman: the Coleman bounce  $\varphi_b(x)$ , a non homogeneous solution to the equations of motion that interpolates between the false and the true vacuum and then bounces back.



The decay rate is given by the imaginary part of the ground state energy in the false vacuum.

$$Z = \sum_{n} e^{-\beta E_{n}} \longrightarrow E_{0} = -\lim_{\beta \to \infty} \frac{1}{\beta} \log Z , \quad \beta = \text{Euclidean time}$$
(2)

- To evaluate it, we can deploy a technique developed by Callan & Coleman: the Coleman bounce  $\varphi_b(x)$ , a non homogeneous solution to the equations of motion that interpolates between the false and the true vacuum and then bounces back.
- We evaluate Z using a saddle point approximation around the bounce.



The decay rate is given by the imaginary part of the ground state energy in the false vacuum.

$$Z = \sum_{n} e^{-\beta E_{n}} \longrightarrow E_{0} = -\lim_{\beta \to \infty} \frac{1}{\beta} \log Z , \quad \beta = \text{Euclidean time}$$
(2)

- To evaluate it, we can deploy a technique developed by Callan & Coleman: the Coleman bounce  $\varphi_b(x)$ , a non homogeneous solution to the equations of motion that interpolates between the false and the true vacuum and then bounces back.
- We evaluate Z using a saddle point approximation around the bounce.

$$\Gamma \propto \left(\frac{\det S''[\varphi_b]}{\det S''[\varphi_{FV}]}\right)^{-\frac{1}{2}} e^{-S[\varphi_b] + S[\varphi_{FV}]} \tag{3}$$

**Outline** 



False Vacuum Decay

#### 2 The Problem with Functional Determinants

- The Fubini-Lipatov instanton
- Zero and negative modes
- The Resolvent Method
- Challenges and outlook



(4)

A decent approximation for the Higgs potential is

$$V(arphi) = -rac{\lambda}{4}arphi^4 \quad ext{with} \quad \lambda > 0$$

ТШ

A decent approximation for the Higgs potential is

$$V(\varphi) = -\frac{\lambda}{4}\varphi^4 \quad \text{with} \quad \lambda > 0$$
 (4)

This theory admits a Coleman bounce: the Fubini-Lipatov instanton.

$$\varphi_{\rm FL}(x) = \sqrt{\frac{8}{\lambda}} \frac{R}{R^2 + (r - r_0)^2}, \qquad r^2 = x_i x_i$$

Matthias Carosi | False Vacuum Decay and Functional Determinants | 11/05/2022

(5)



A decent approximation for the Higgs potential is

$$V(\varphi) = -\frac{\lambda}{4}\varphi^4 \quad \text{with} \quad \lambda > 0$$
 (4)

This theory admits a Coleman bounce: the Fubini-Lipatov instanton.

$$\varphi_{\rm FL}(x) = \sqrt{\frac{8}{\lambda}} \frac{R}{R^2 + (r - r_0)^2}, \qquad r^2 = x_i x_i$$
(5)

We can identify two invariances of the theory



A decent approximation for the Higgs potential is

$$V(\varphi) = -\frac{\lambda}{4}\varphi^4 \quad \text{with} \quad \lambda > 0$$
 (4)

This theory admits a Coleman bounce: the Fubini-Lipatov instanton.

$$\varphi_{\rm FL}(x) = \sqrt{\frac{8}{\lambda}} \frac{R}{R^2 + (r - r_0)^2}, \qquad r^2 = x_i x_i$$
 (5)

We can identify two invariances of the theory

 $\Box$  Position of the centre of mass  $r_0 \to$  translations  $\to$  translational zero modes  $\varphi_{tr} \propto \partial_i \varphi_{FL}$ 

ТШ

A decent approximation for the Higgs potential is

$$V(\varphi) = -\frac{\lambda}{4}\varphi^4 \quad \text{with} \quad \lambda > 0$$
 (4)

This theory admits a Coleman bounce: the Fubini-Lipatov instanton.

$$\varphi_{\rm FL}(x) = \sqrt{\frac{8}{\lambda}} \frac{R}{R^2 + (r - r_0)^2}, \qquad r^2 = x_i x_i$$
(5)

We can identify two invariances of the theory

- $\Box$  Position of the centre of mass  $r_0 \to$  translations  $\to$  translational zero modes  $\varphi_{
  m tr} \propto \partial_i \varphi_{
  m FL}$
- $\Box$  Instanton size R o dilatation o dilatational zero mode  $arphi_{
  m dil} \propto \partial_R arphi_{
  m FL}$



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty \tag{6}$$



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty$$
 (6)

We would like to define a procedure for extracting the zero modes from the evaluation of the determinant, and obtain a final result.



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty$$
(6)

We would like to define a procedure for extracting the zero modes from the evaluation of the determinant, and obtain a final result.

Strategy:



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty$$
(6)

We would like to define a procedure for extracting the zero modes from the evaluation of the determinant, and obtain a final result.

Strategy:

1. Subtract the translational zero modes from the definition of the determinant.



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty$$
(6)

We would like to define a procedure for extracting the zero modes from the evaluation of the determinant, and obtain a final result.

Strategy:

- 1. Subtract the translational zero modes from the definition of the determinant.
- 2. Add 1-loop quantum corrections and renormalize. This will break scale invariance, eliminating the remaining dilatational zero mode.



The decay rate (3) is unfortunately divergent

$$(\det S''[\varphi_b])^{-1/2} \sim \left(\prod_n \lambda_n\right)^{-1/2} \xrightarrow{\text{zero modes}} "\left(0^{-1/2}\right)" \to \infty$$
(6)

We would like to define a procedure for extracting the zero modes from the evaluation of the determinant, and obtain a final result.

Strategy:

- 1. Subtract the translational zero modes from the definition of the determinant.
- 2. Add 1-loop quantum corrections and renormalize. This will break scale invariance, eliminating the remaining dilatational zero mode.
- The operator  $S''[\varphi_b]$  also contains a negative mode, which is responsible for the imaginary part of the ground state energy.

Matthias Carosi | False Vacuum Decay and Functional Determinants | 11/05/2022

#### **Resolvent Method**

Green function's problem

$$S''[\varphi]G(x,x') = \delta(x-x'), \quad \text{where} \quad S''[\varphi] = -\Box_x + V''(\varphi(x)) \tag{7}$$

Spectral decomposition

$$G(x, x') = \int d\lambda \frac{f_{\lambda}(x)f_{\lambda}(x')}{\lambda}$$
(8)

Define

$$(S''[\varphi] + s)G(s; x, x') = \delta(x - x')$$
(9)

then

$$\log \frac{\det S''[\varphi_b]}{\det S''[\varphi_{\mathsf{FV}}]} = -\int dx \int_0^\infty ds \left[G_{\mathsf{bounce}}(s;x,x) - G_{\mathsf{FV}}(s;x,x)\right]$$
(10)





#### Subtracting the zero modes



Define the subtracted Green's function as orthogonal to the space of zero modes

$$\int dx' G^{\perp}(x, x')\phi_0(x') = 0$$
(11)

Differential equation for the subtracted Green's function

.

$$S''[\varphi]G^{\perp}(x,x') = \delta(x-x') - \sum \phi_0(x)\phi_0^*(x')$$
(12)

#### Matthias Carosi | False Vacuum Decay and Functional Determinants | 11/05/2022

#### **Quantum corrections**

1-loop quantum corrections to the bounce equation

$$-\Box \varphi_b^{(1)}(x) + V'(\varphi_b^{(1)}(x)) + \varphi_b^{(1)}(x) \Pi_{\mathrm{ren}}^{(1)}(x(\varphi_b^{(1)})) = 0$$

Quantum corrected equation for the 2-point function

$$\left[ -\Box + V''(\varphi_b^{(1)}(x)) + \Pi_{\text{ren}}^{(1)}(x(\varphi_b^{(1)})) \right] G_{1-\text{loop}}(x, x') + \int dy \varphi_b^{(1)}(x) \Sigma(x, y) \varphi_b^{(1)}(y) G_{1-\text{loop}}(y, x') = \delta(x - x')$$
(14)

and



(13)

#### Challenges and outlook What do we learn about functional determinants?



To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:

ПΠ

- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!



- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.



- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.
  - 3. Quantum corrections to the Green's function equation include non-local terms which make it impossible to evaluate without defining some cut-off procedure.



- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.
  - 3. Quantum corrections to the Green's function equation include non-local terms which make it impossible to evaluate without defining some cut-off procedure.
  - The next steps would be
    - 1. Obtain a numerical result for the decay rate at 1-loop.



- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.
  - 3. Quantum corrections to the Green's function equation include non-local terms which make it impossible to evaluate without defining some cut-off procedure.
  - The next steps would be
    - 1. Obtain a numerical result for the decay rate at 1-loop.
    - 2. Include the effects of additional fields (fermionic and bosonic).



- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.
  - 3. Quantum corrections to the Green's function equation include non-local terms which make it impossible to evaluate without defining some cut-off procedure.
  - The next steps would be
    - 1. Obtain a numerical result for the decay rate at 1-loop.
    - 2. Include the effects of additional fields (fermionic and bosonic).
    - 3. Study the behavious of functional determinants when transforming the operators.



#### What do we learn about functional determinants?

- To evaluate functional determinants in a proper way we use the resolvent method. This can be nicely extended to include quantum corrections but presents a number of challenges:
  - 1. Finding the Green's function in the space orthogonal to the zero modes is hard!
  - 2. Quantum corrections to the bounce and to the Green's function can only be evaluated self-consistently.
  - 3. Quantum corrections to the Green's function equation include non-local terms which make it impossible to evaluate without defining some cut-off procedure.

#### The next steps would be

- 1. Obtain a numerical result for the decay rate at 1-loop.
- 2. Include the effects of additional fields (fermionic and bosonic).
- 3. Study the behavious of functional determinants when transforming the operators.
- 4. Eventually yield a more precise evaluation of the SM lifetime.

#### **THANK YOU FOR THE ATTENTION!**





#### **THANK YOU FOR THE ATTENTION!**





#### **THANK YOU FOR THE ATTENTION!**



