

Theory Predictions in the Precision Era

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Theory Predictions in the Precision Era

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New Physics and Precision Calculations

- **SM** thoroughly tested at the **Large Hadron Collider**
- Higgs boson discovered in 2012
- No other new particles or new physics signals observed so far...
- Main focus is finding small new physics effects
- **Precise predictions** for the SM are always relevant...
- But they are **indispensable** to detect small deviations
- Theory needs at least to match experimental precision

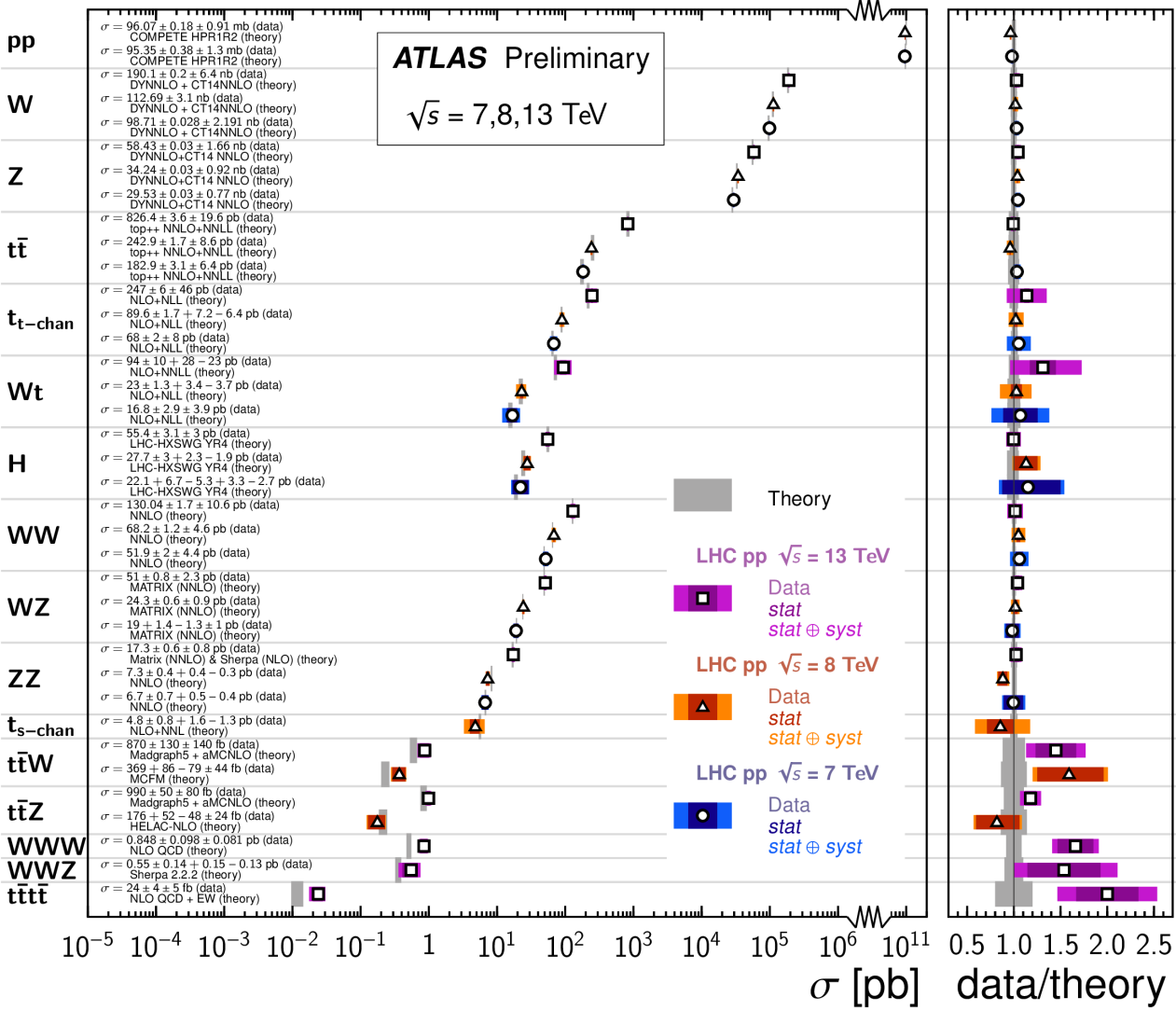


Smoking gun
new physics signal



Small deviations
from the SM

LHC: Present and Future



- Experimental accuracy really impressive!
- W, Z, $t\bar{t}$ at the few percent level
- H and diboson at the 5% level
- Only a small fraction of total integrated luminosity

Current $\sim 150 \text{ fb}^{-1}$
 Run3: 300 fb^{-1} (2024)
 HL-LHC: 3000 fb^{-1} (~ 2040)

- Experimental precision will only get better!



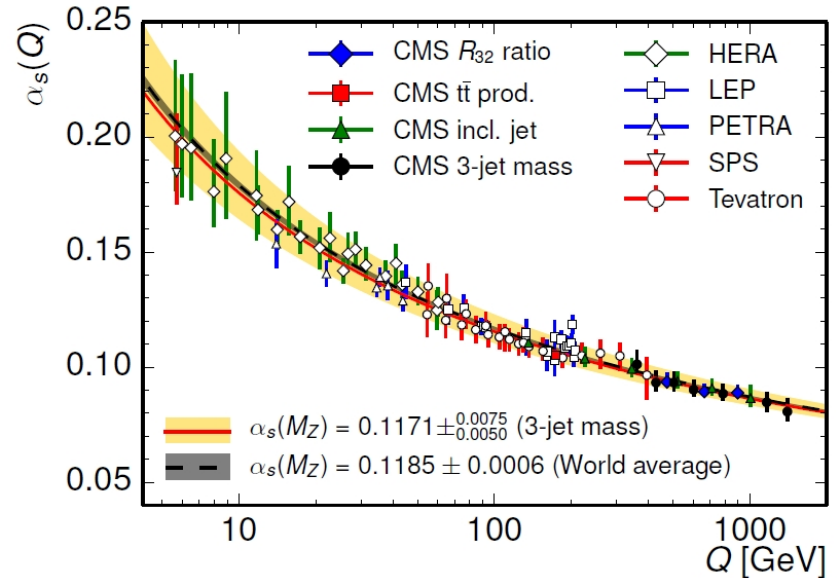
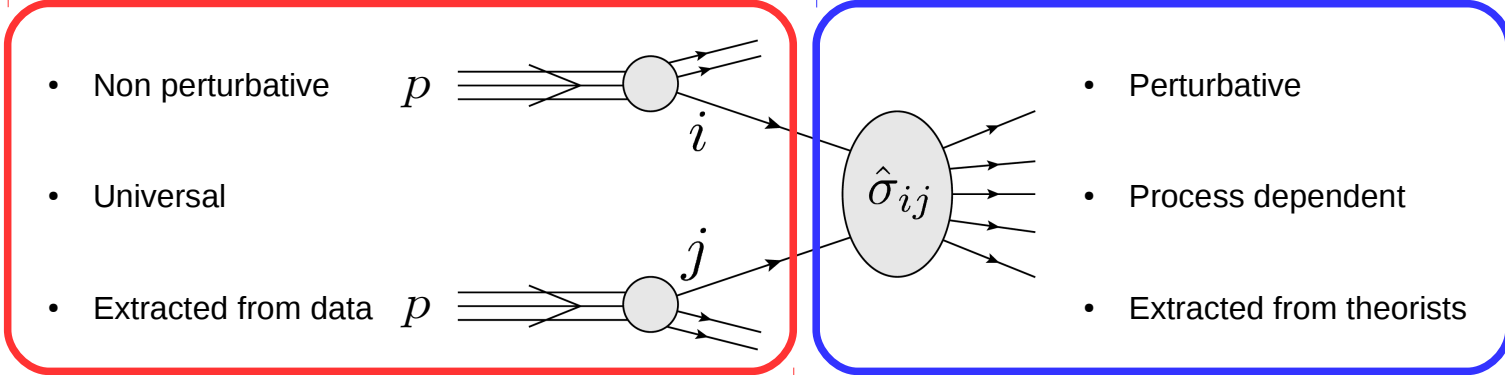
Demand for precision calculations

Quantum Chromodynamics

- QCD corrections fundamental in hadron collisions, very large corrections due to large value of strong coupling
- Running of $\alpha_s \rightarrow$ grows at low energies (large distances)
- LHC hard collisions in perturbative regime, but... 😎
- The physics of the proton ($\sim 1\text{GeV}$) are not! 😓

Factorization of long and short distances:

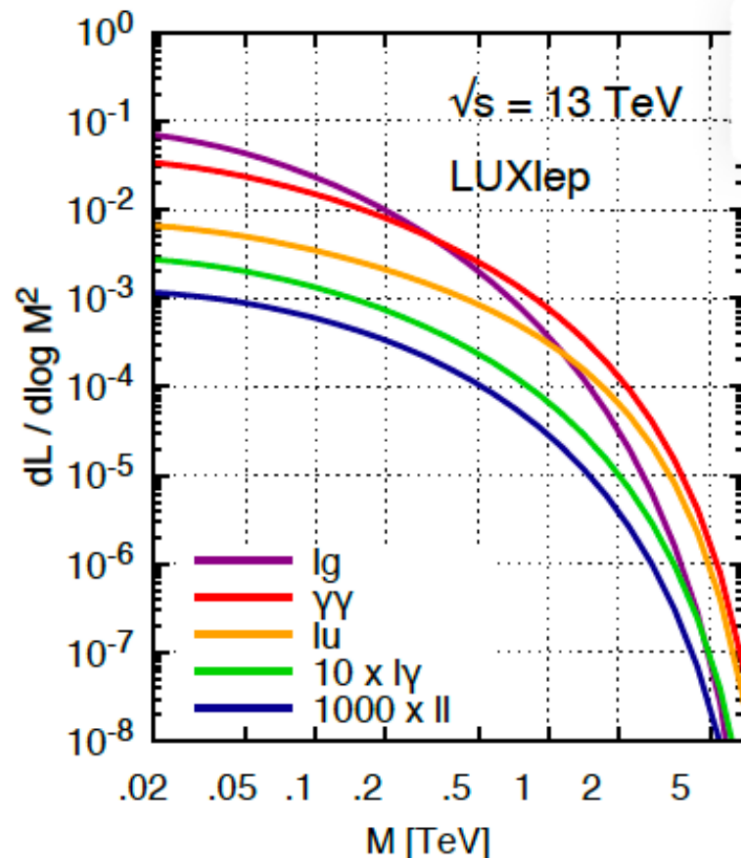
$$\sigma_{pp \rightarrow X} = \sum_{ij} \int dx_1 dx_2 \underbrace{f_i(x_1) f_j(x_2)}_{\text{Parton Distribution Functions}} \underbrace{\hat{\sigma}_{ij \rightarrow X}}_{\text{Partonic cross section (hard process)}} + \mathcal{O}(\Lambda/Q)$$



Photon and Lepton Induced Processes at the LHC

- Quarks and gluons inside the proton... and much more!
- At this level of precision or for some rare processes photon and lepton content is relevant

Lepton luminosities



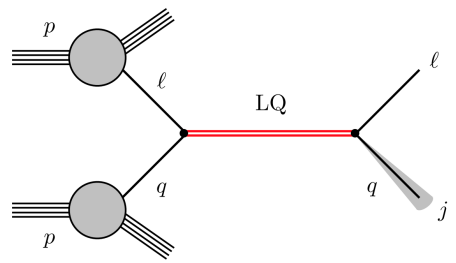
$$\mathcal{L}_{ij} \equiv M^2 \int_0^1 dz dy (f_i(z, M^2) f_j(y, M^2) \delta(M^2 - szy))$$

Probability to pick a parton "i" in the proton with momentum fraction z

First precise predictions for photon and lepton induced processes at the LHC:

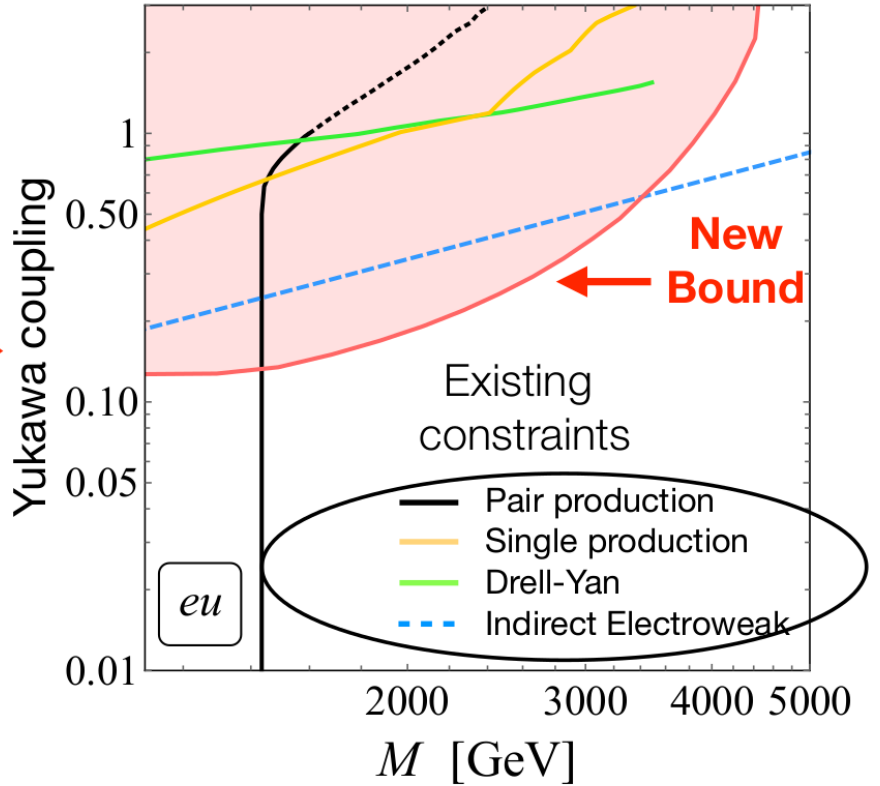
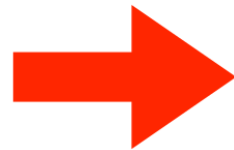
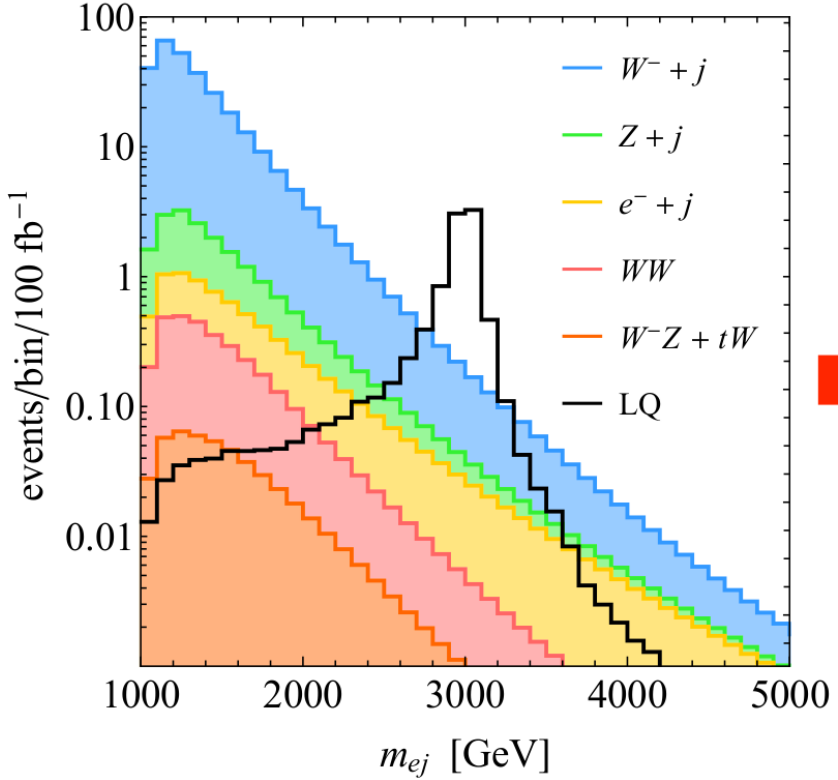
- opens up possibility to new measurements
- backgrounds for BSM searches for rare processes involving leptons/photons

Resonant Leptoquark Production



- Application: set stringent constraints to Leptoquark production

LHC, $\sqrt{s} = 13$ TeV



Fixed Order Calculations

$$\sigma_{pp \rightarrow X} = \sum_{ij} \int dx_1 dx_2 f_i(x_1) f_j(x_2) \hat{\sigma}_{ij \rightarrow X} + \mathcal{O}(\Lambda/Q)$$

Uncertainties (indicative!)

LO $\rightarrow \mathcal{O}(100\%)$

NLO $\rightarrow \mathcal{O}(10\%)$

NNLO $\rightarrow \mathcal{O}(1\%)$

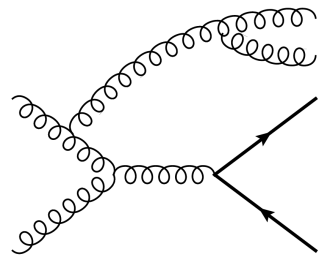


$$\hat{\sigma} = \underbrace{\hat{\sigma}_{\text{LO}}}_{\text{LO}} \underbrace{(1 + \alpha_S + \alpha_S^2 + \dots)}_{\text{NLO}} \underbrace{\quad}_{\text{NNLO}}$$

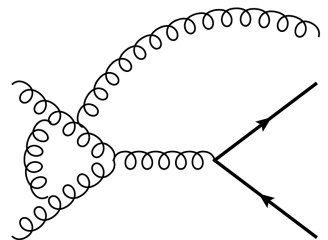
Higher-order calculations are very complex: large number of Feynman diagrams, intermediate divergencies, complicated phase-space integrations, etc... But this is not even the end of the story!

Fixed-order calculations look like this:

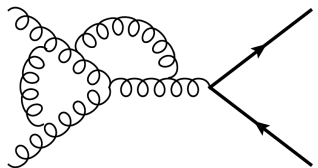
E.g.: NNLO $t\bar{t}$



Double real



Real-virtual



Two-loop virtual

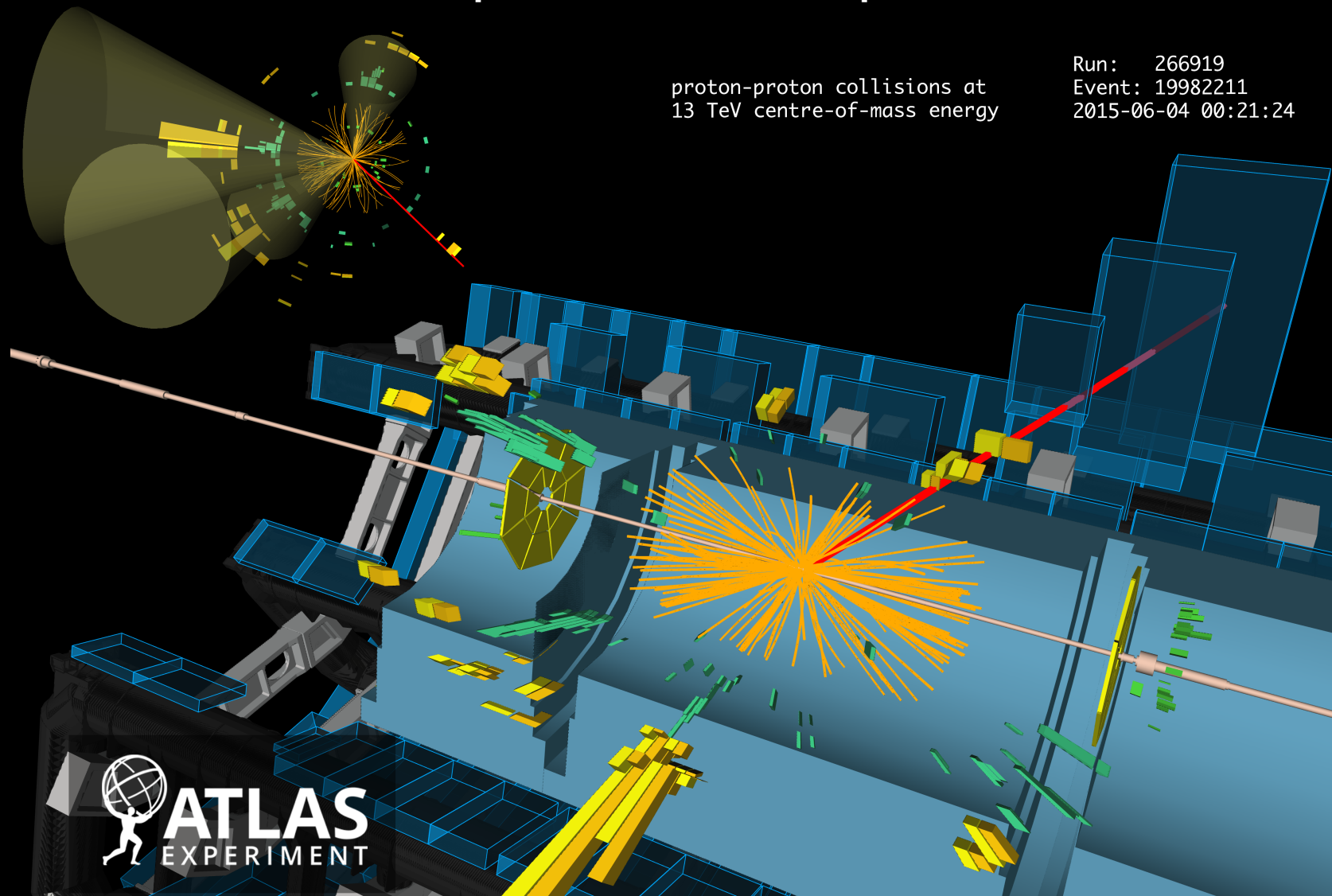
Small number of final-state particles

Free partons in the final state

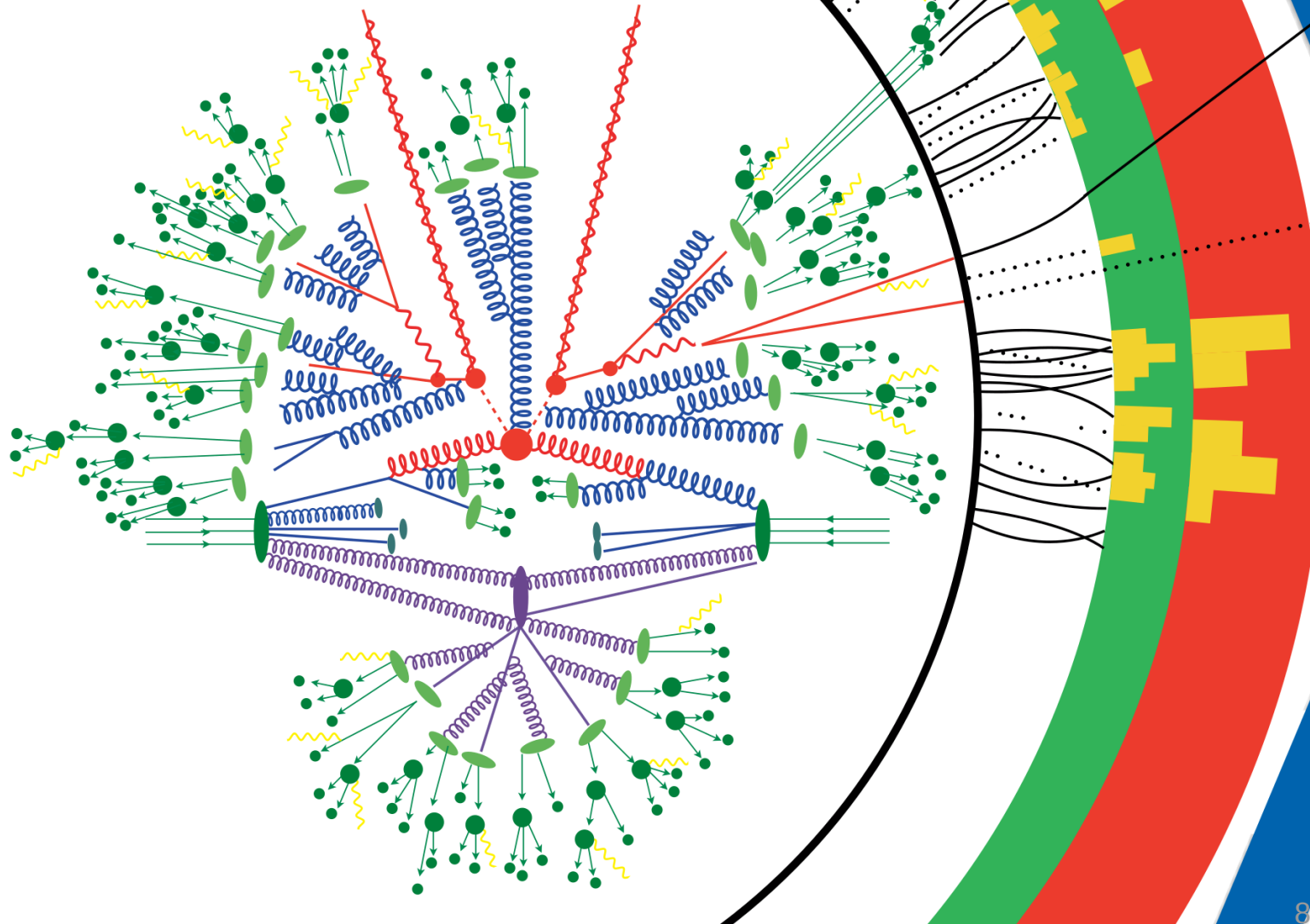
Actual events are more complicated → more complex simulations are needed!

proton-proton collisions at
13 TeV centre-of-mass energy

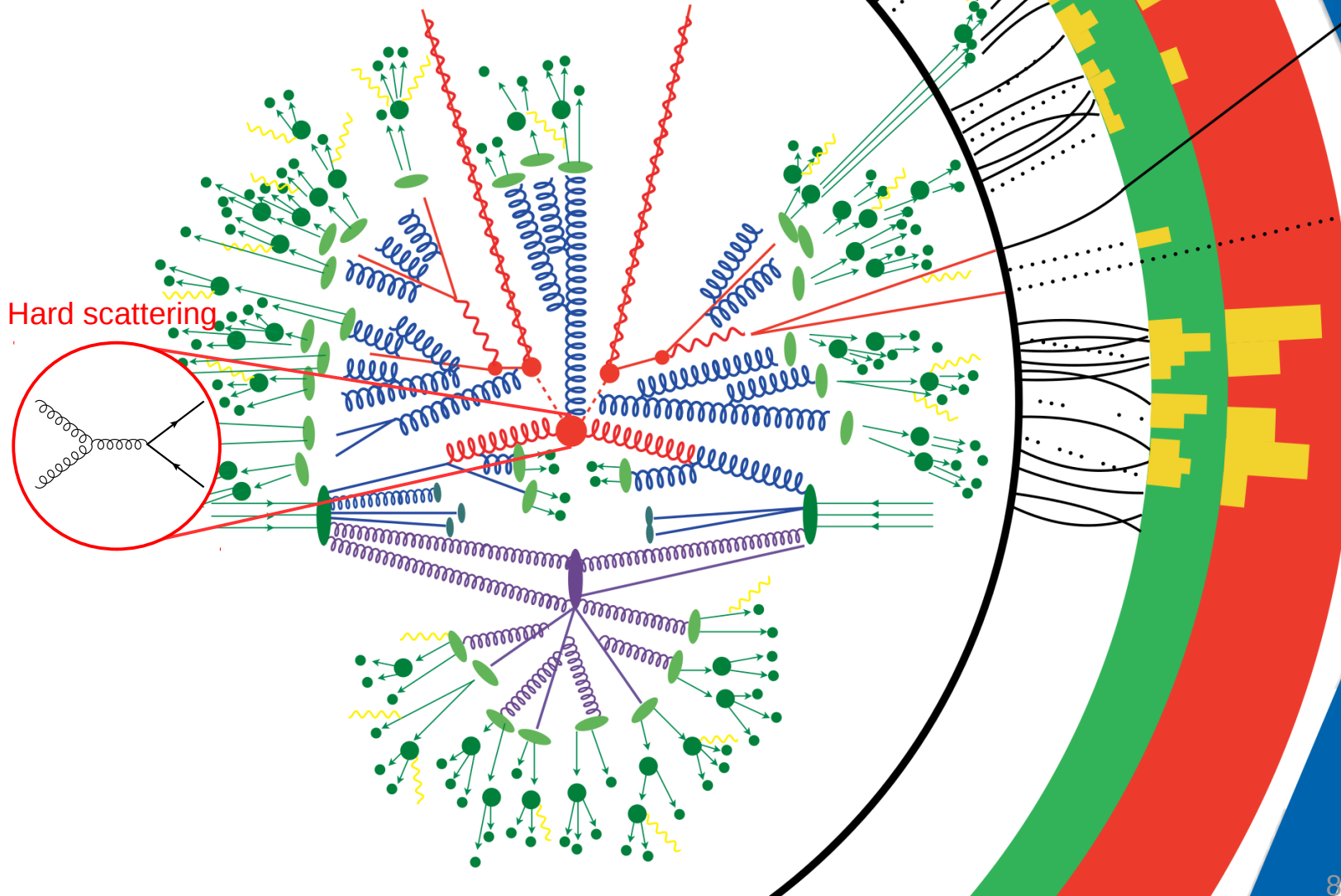
Run: 266919
Event: 19982211
2015-06-04 00:21:24



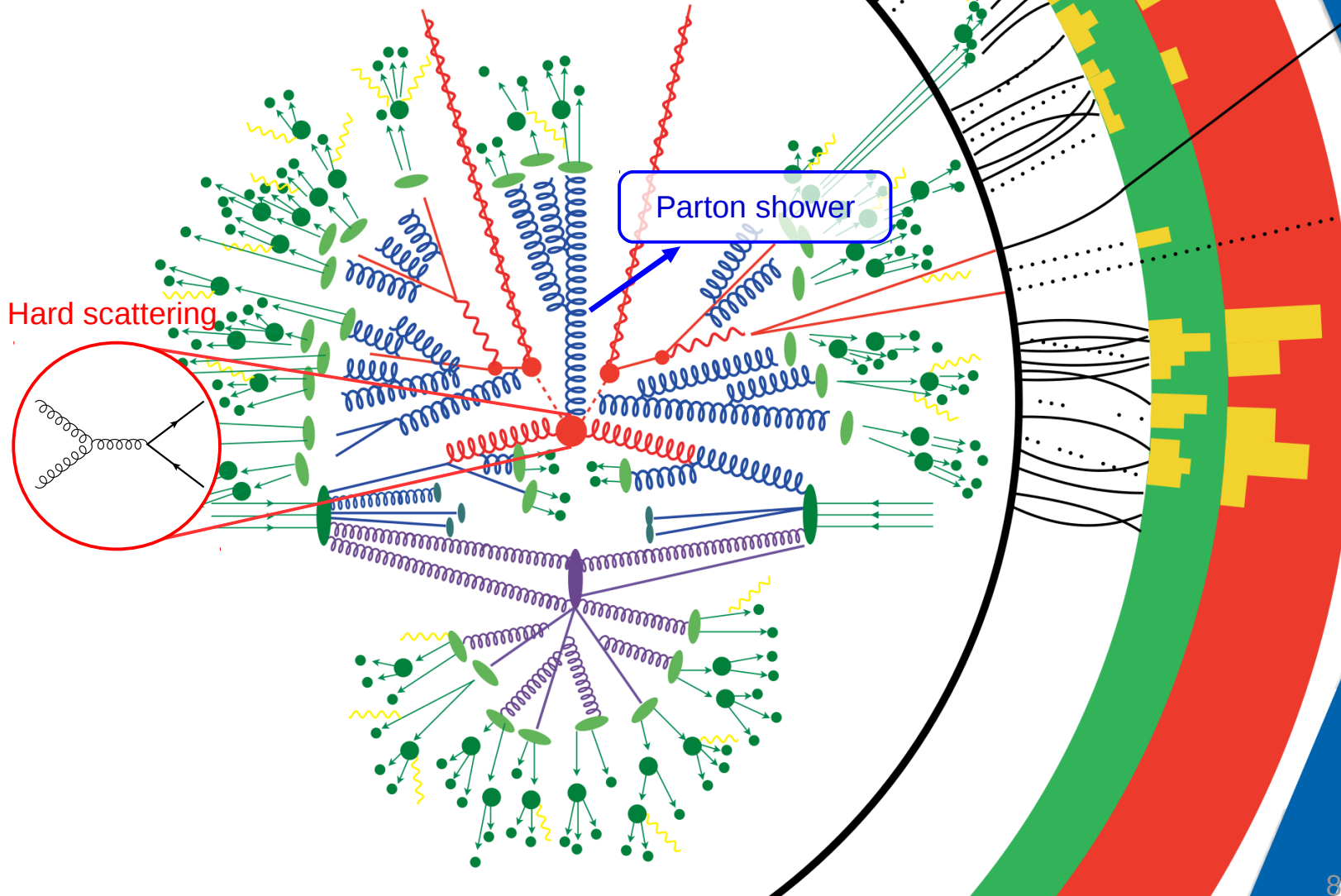
Event Generators



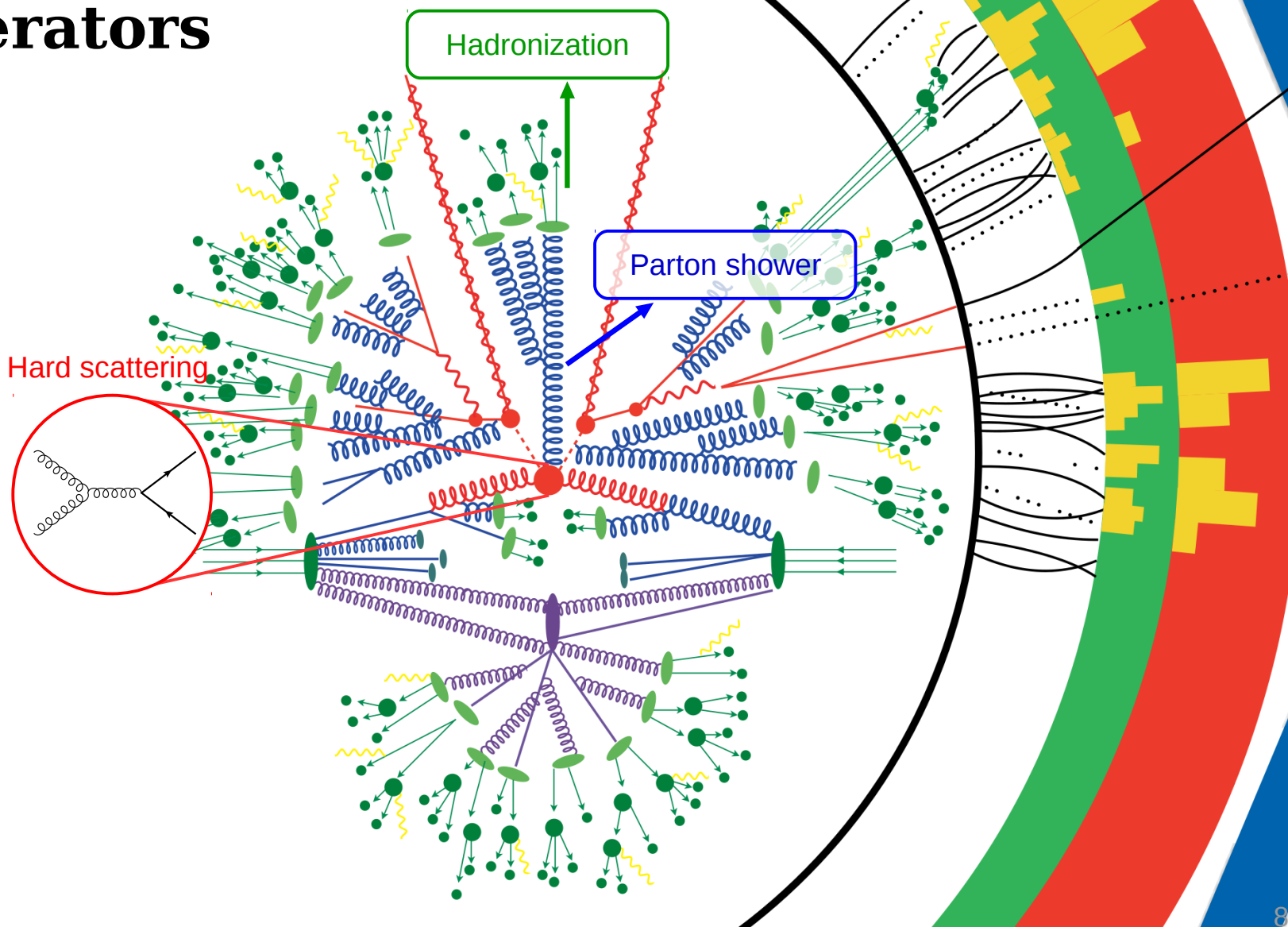
Event Generators



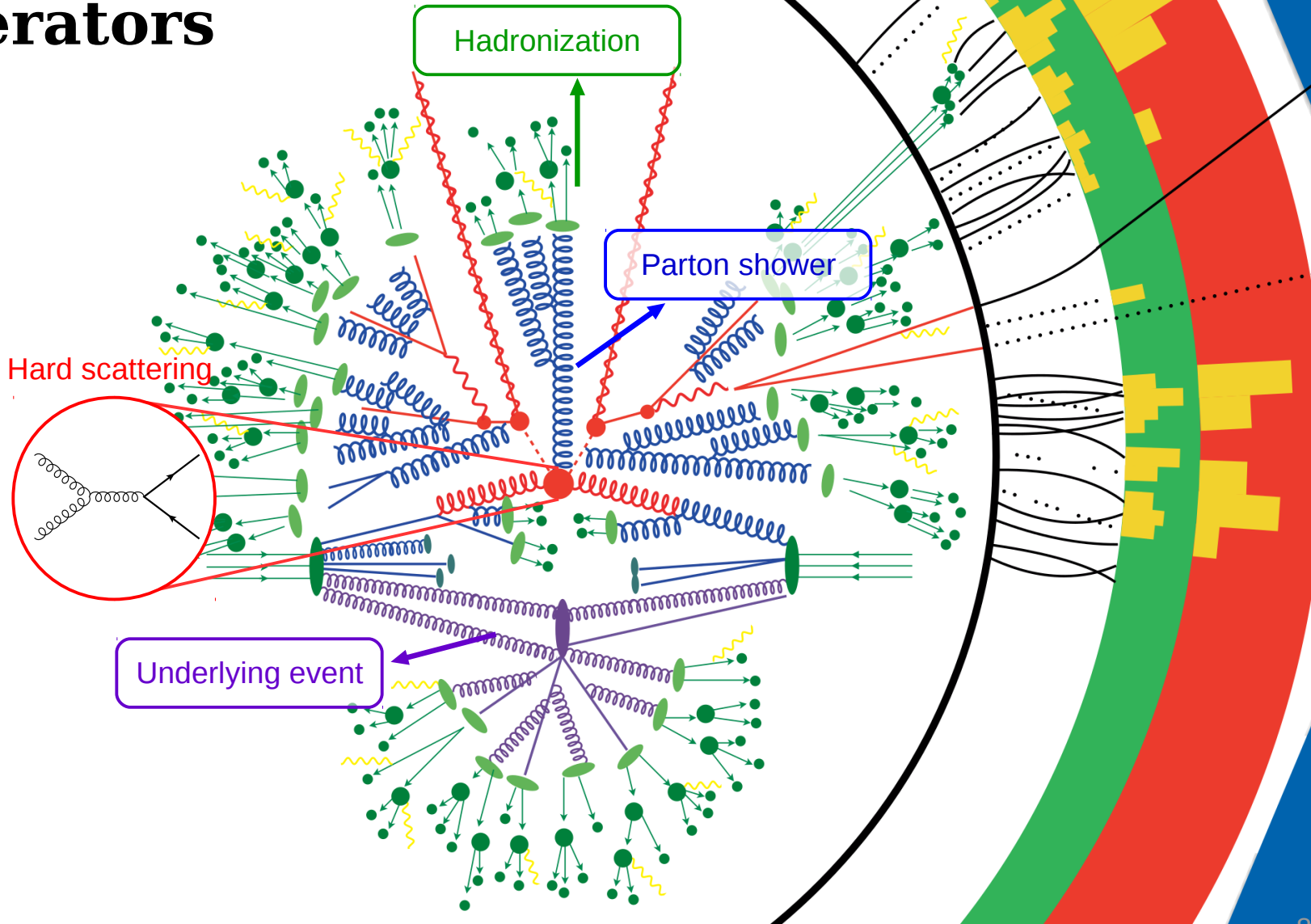
Event Generators



Event Generators

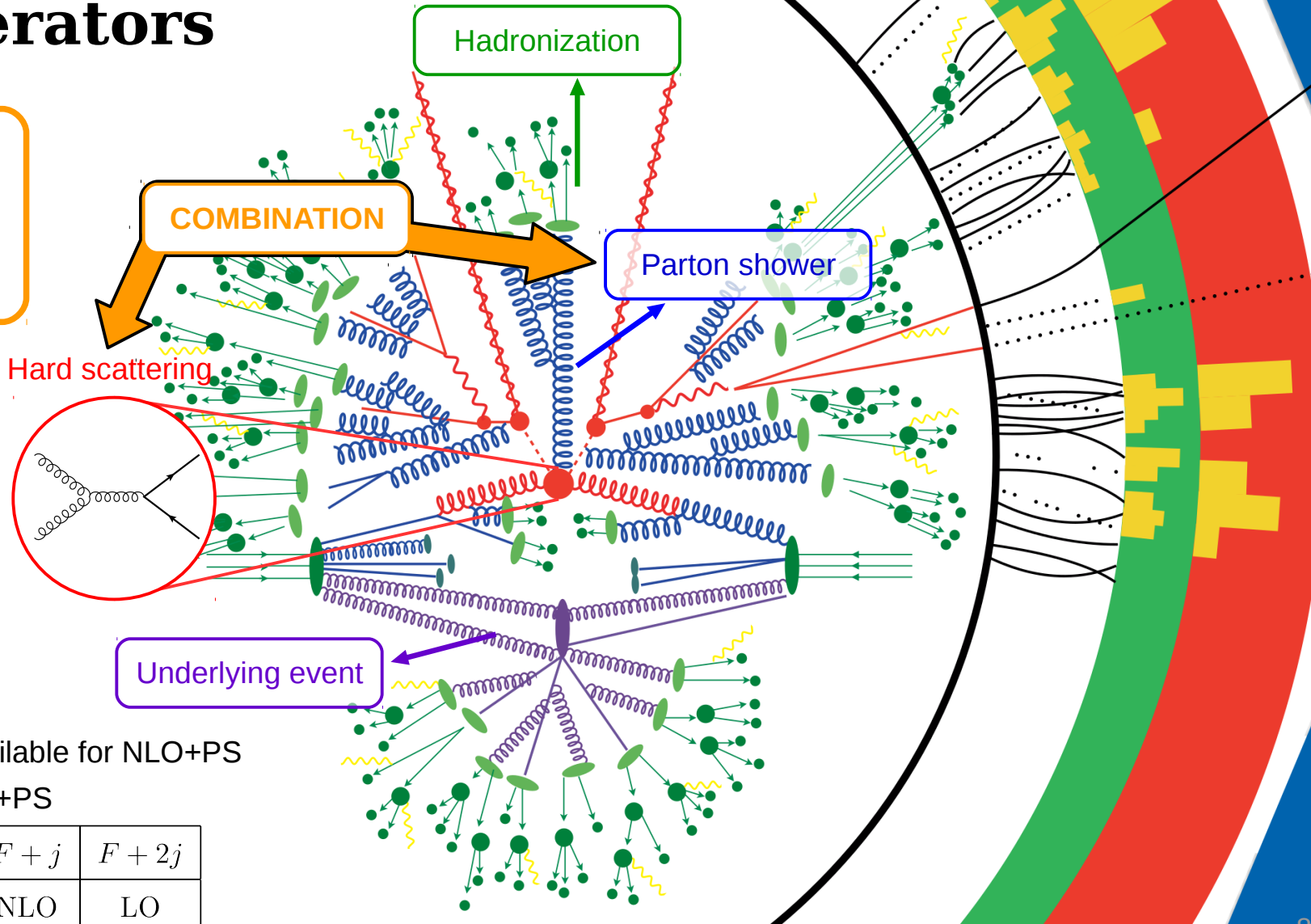


Event Generators



Event Generators

We want to keep the fixed-order accuracy when computing inclusive observables



- General approaches available for NLO+PS
- Current frontier is NNLO+PS

	F	$F + j$	$F + 2j$
$F@NNLO_{PS}$	NNLO	NLO	LO

NNLO+PS

◆ seminal approaches for NLO+PS many years ago (POWHEG, MC@NLO)

◆ first NNLO+PS for simple $2 \rightarrow 1$ processes

❖ **MiNLO+reweighting** [Hamilton, Nason, Zanderighi '12, + Re '13],
[Karlberg, Hamilton, Zanderighi '14]

❖ **Geneva** [Alioli, Bauer, Berggren, Tackmann, Walsh, Zuberi '13],
[Alioli, Bauer, Berggren, Tackmann, Walsh '15]

❖ **UNNLOPS** [Höche, Prestel '14]

◆ **MiNNLO_{PS}**: new approach with enormous potential

[Monni, Nason, Re, MW, Zanderighi '19], [Monni, Re, MW '20]

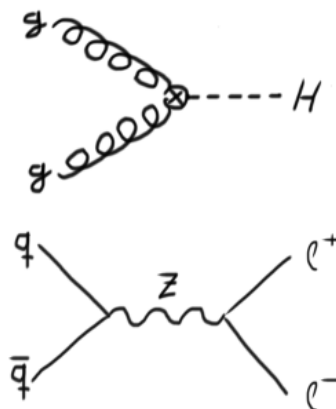
❖ NNLO corrections extracted from analytic resummation formula

$$d\sigma^{(\text{sing})} \sim d\sigma_{c\bar{c}}^{(0)} \times \exp[-S_c(b)] \times [HC_1 C_2]_{c\bar{c}; a_1 a_2} \times f_{a_1} f_{a_2}$$

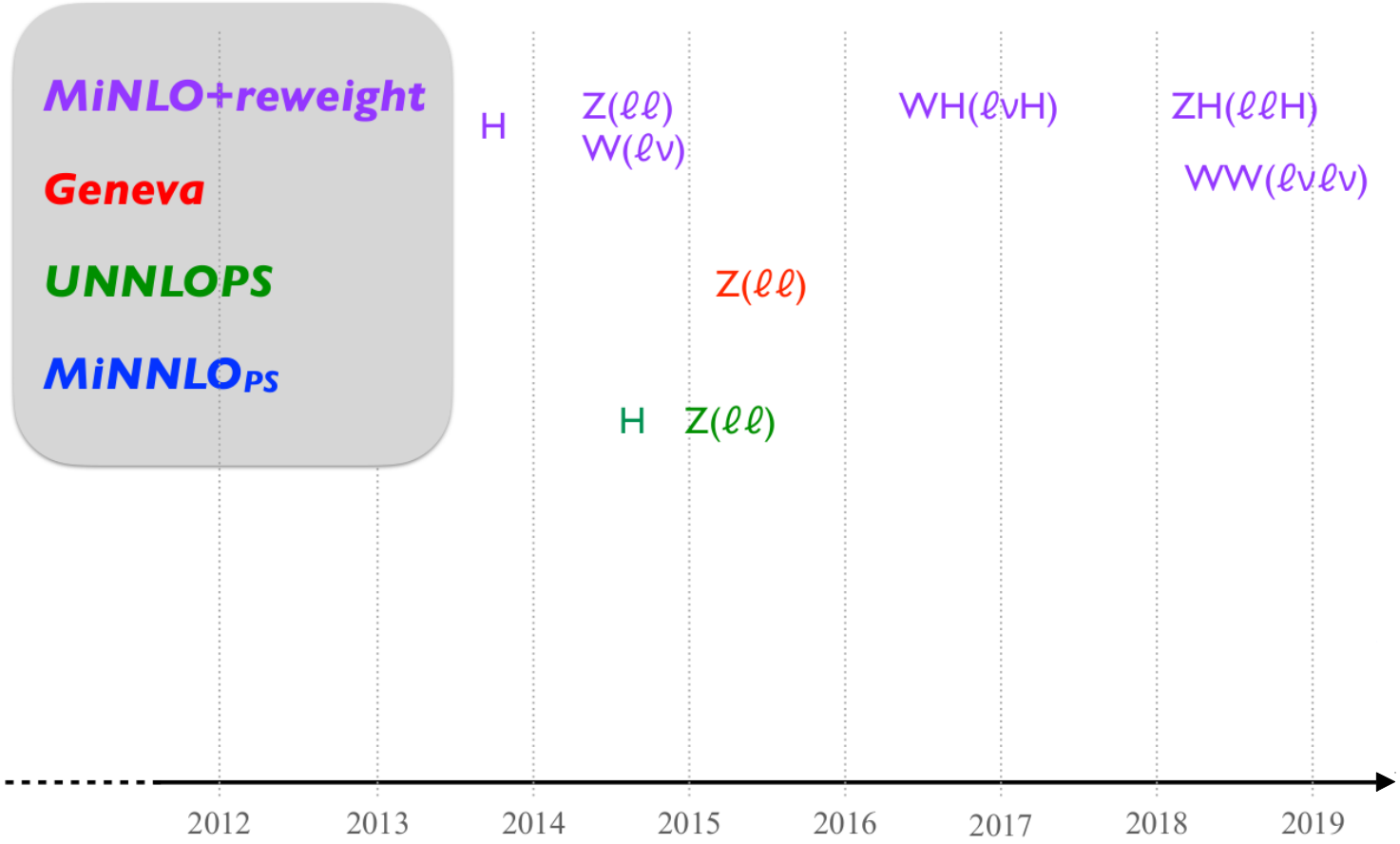
❖ physically sound (no new unphysical scale)

❖ applicable beyond $2 \rightarrow 1$ processes (even beyond colour-singlet)

❖ numerically efficient

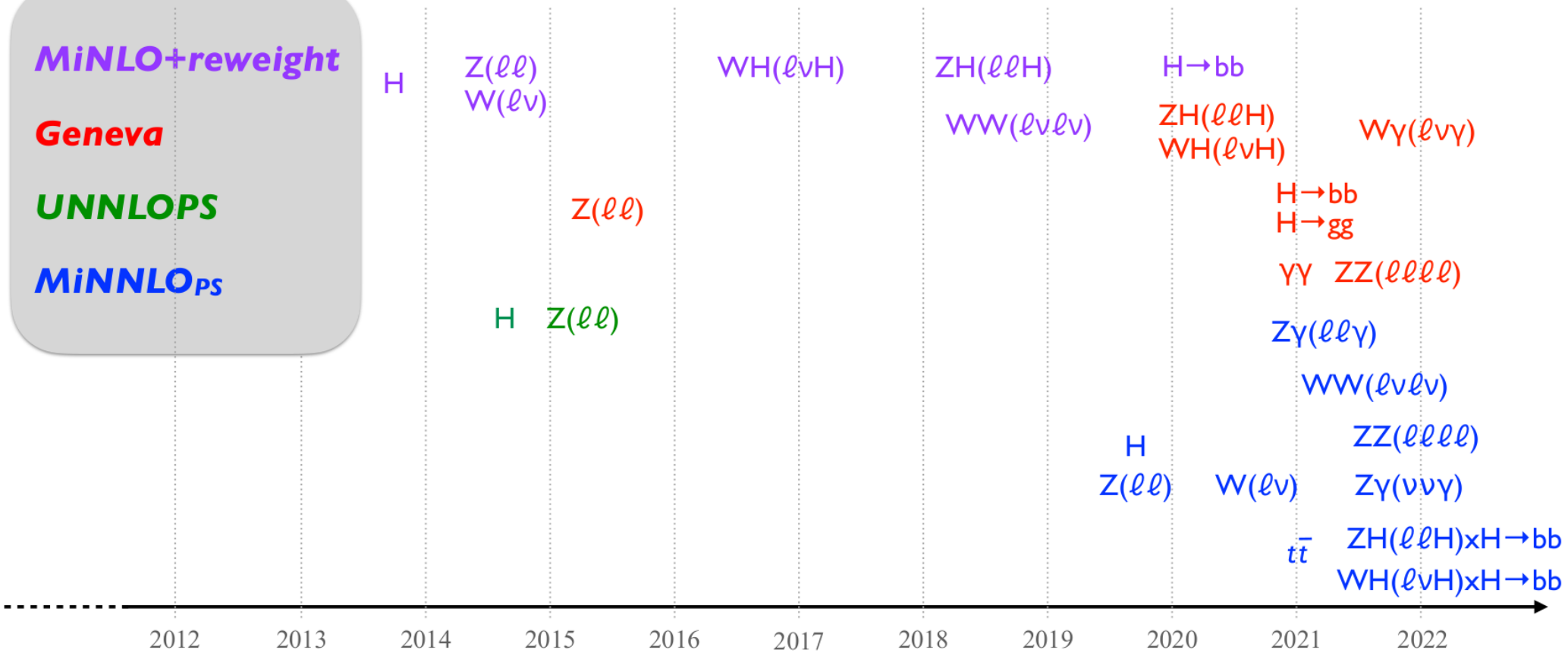


NNLO+PS Timeline

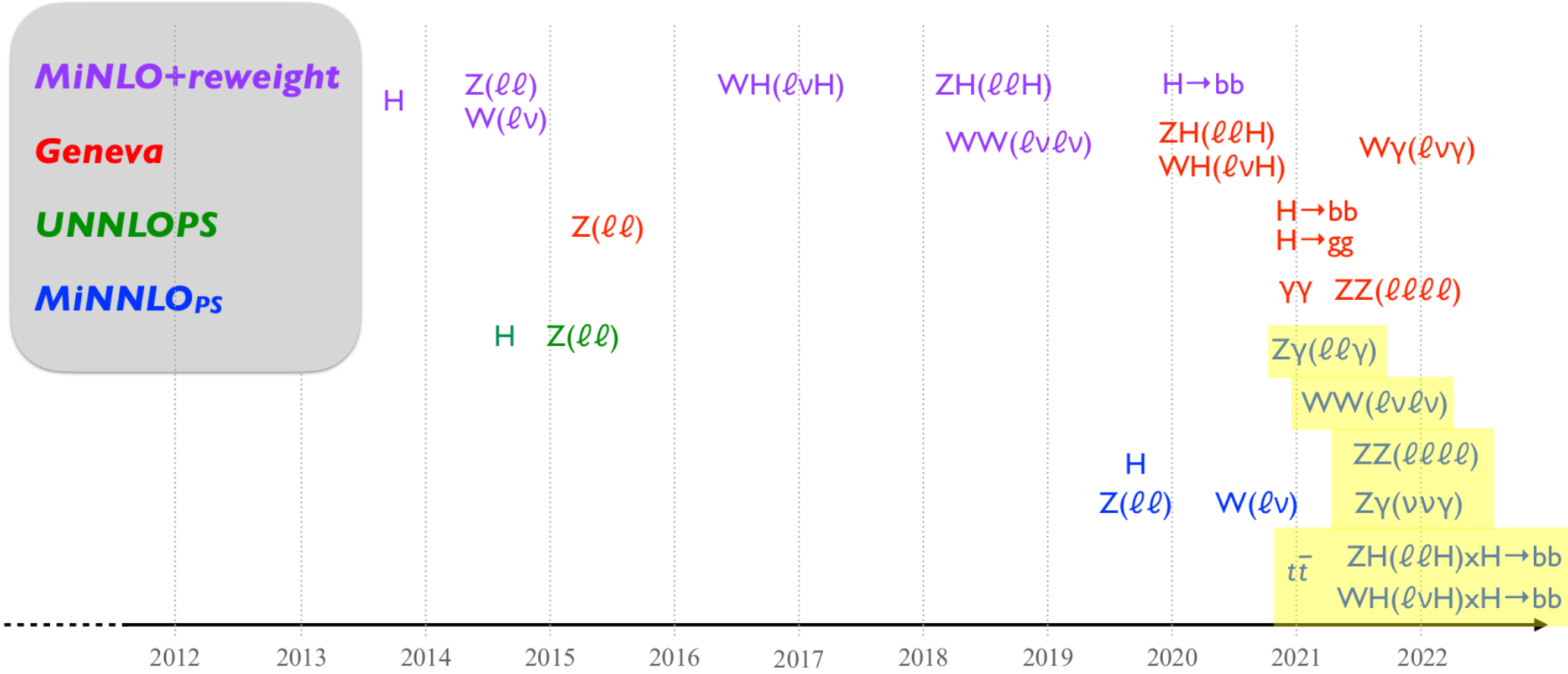


NNLO+PS Timeline

MiNLO+reweight
Geneva
UNNLOPS
MiNNLO_{PS}

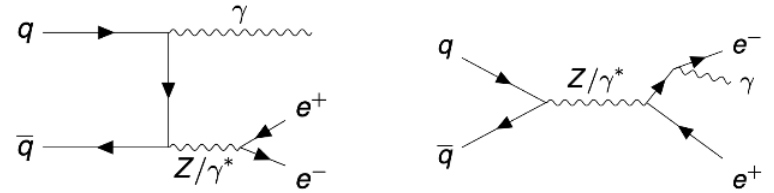


NNLO+PS Timeline

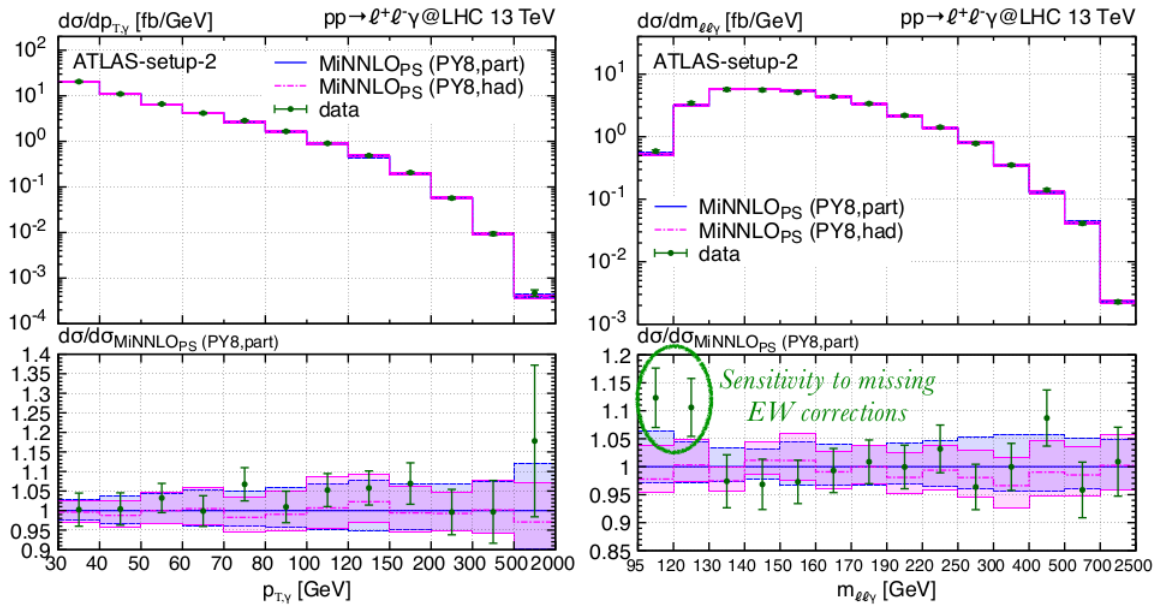


Zγ Phenomenology at the LHC

- ❖ $Z \rightarrow e^+e^-$ event can be fully reconstructed
- ❖ presence of **isolated photon** → theoretically challenging
- ❖ highly relevant as a probe for BSM (especially $Z \rightarrow \nu\bar{\nu}$)



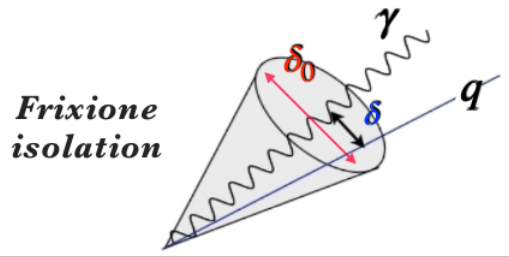
Very good description of data from ATLAS 139fb⁻¹ analysis!



Photon isolation requirement:

- ❖ Experimentally needed to identify hard photons
- ❖ Theoretically delicate definition of an infrared-safe cross section

$$\sum_{had/part \in \delta} E_T^{had/part} \leq E_T^{max}(\delta) = E_T^{ref} \cdot \left(\frac{1 - \cos\delta}{1 - \cos\delta_0} \right)^n, \quad \forall \delta \leq \delta_0$$



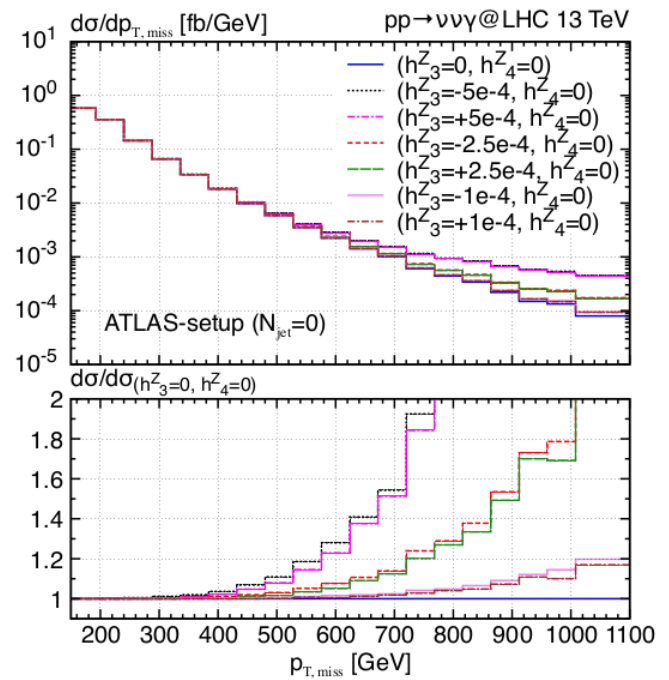
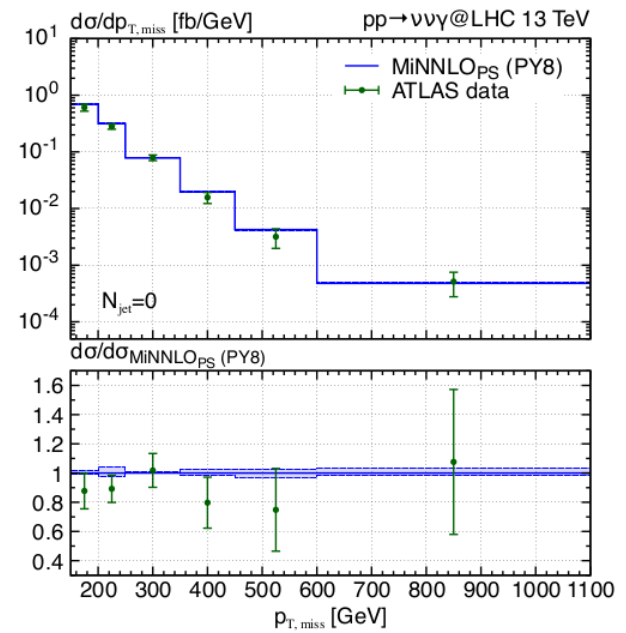
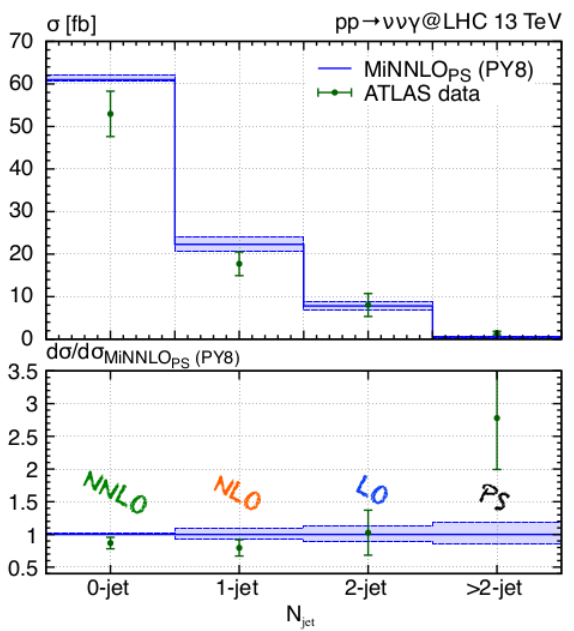
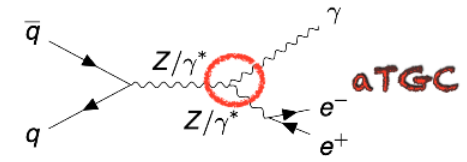
Zγ Phenomenology at the LHC

[Lombardi, Wieseemann, Zanderighi '21] Z → νν̄

- ❖ Z → e⁺e⁻ event can be fully reconstructed
- ❖ presence of isolated photon → theoretically challenging
- ❖ highly relevant as a **probe for BSM** (especially Z → νν̄)

Good agreement with experimental data from ATLAS 36.1fb⁻¹ analysis!

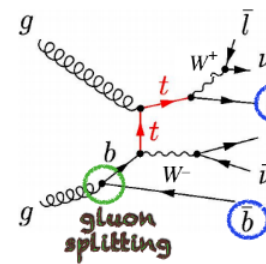
$$\Gamma_{Z\gamma V}^{\alpha\beta\mu}(q_1, q_2, p) = \frac{i(p^2 - m_V^2)}{\Lambda^2} \left(h_1^V (q_2^\mu g^{\alpha\beta} - q_2^\alpha g^{\mu\beta}) + \right. \\ \left. + \frac{h_2^V}{\Lambda^2} p^\alpha (p \cdot q_2 g^{\mu\beta} - q_2^\mu p^\beta) - h_3^V \varepsilon^{\mu\alpha\beta\nu} q_{2\nu} - \frac{h_4^V}{\Lambda^2} \varepsilon^{\mu\beta\nu\sigma} p^\alpha p_\nu q_{2\sigma} \right)$$



WW Phenomenology at the LHC

[Lombardi, Wieseemann, Zanderighi '21]

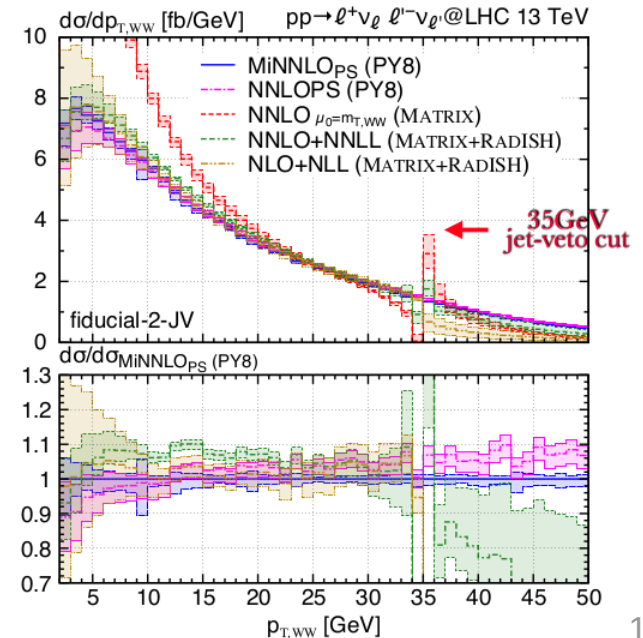
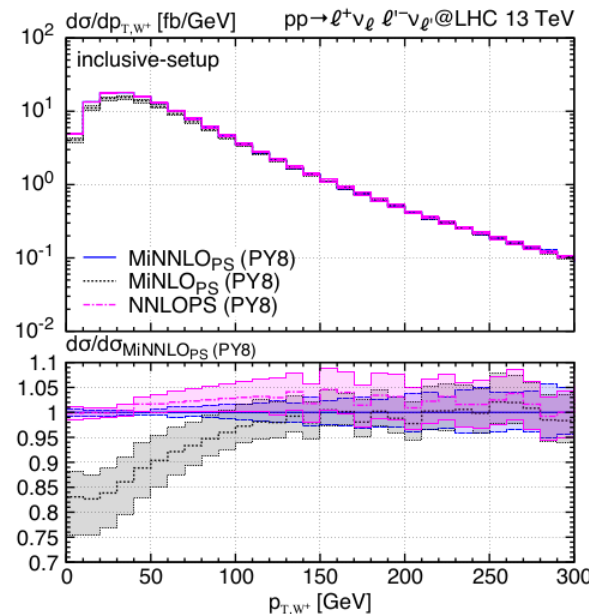
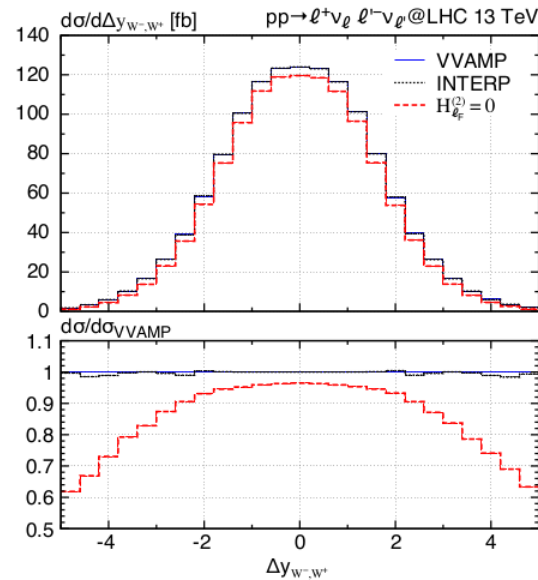
- ❖ largest cross section among massive diboson processes
- ❖ direct access to anomalous triple gauge couplings
- ❖ no full event reconstruction due to neutrinos \rightarrow high-accurate theoretical predictions required
- ❖ analysis requires a **jet-veto** \rightarrow theoretical modelling important



Jet-veto requirement:

- ❖ Experimentally needed to reduce top background
- ❖ Theoretically involved definition of WW cross section, due to diagrams with resonant top quarks and b final states:
 - ▶ Interference with double-real diagrams
 - ▶ Not separately finite for massless b quarks

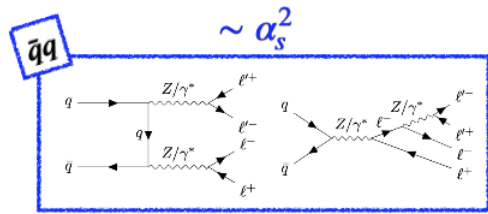
Two-loop contributions evaluated with 4D cubic spline interpolation



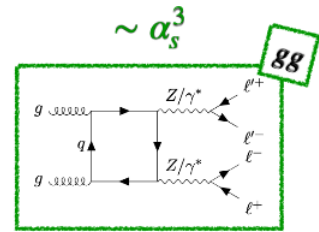
ZZ Phenomenology at the LHC

[Buonocore, Koole, Lombardi, Rottoli, Wiesemann, Zanderighi '21]

- ❖ smallest cross section among massive diboson processes, but very clean signature
- ❖ relevant for BSM searches
- ❖ important for constraining the Higgs width and Higgs couplings



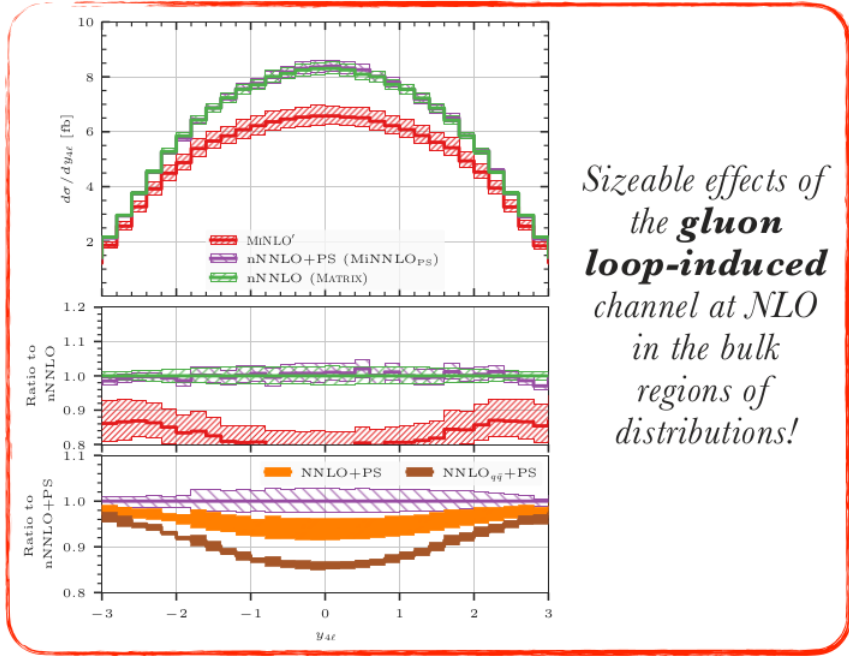
NNLO+PS using MiNNLO_{PS}



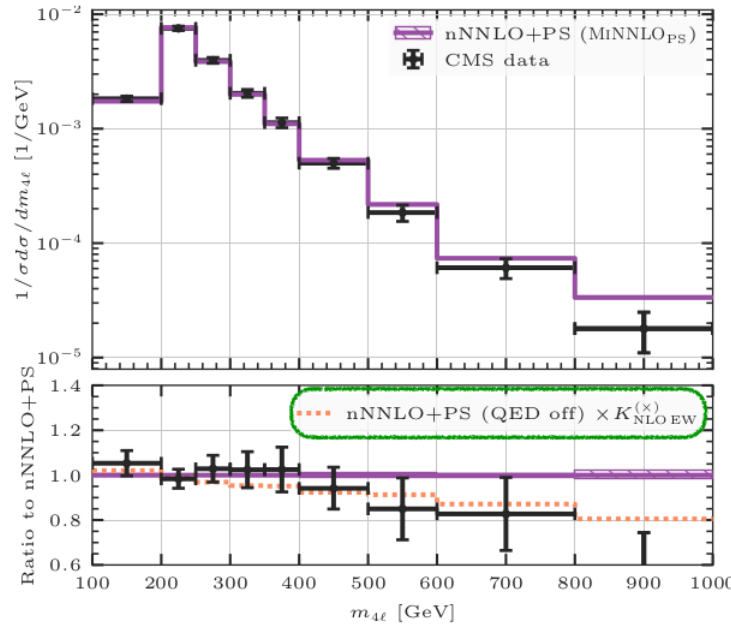
NLO+PS with POWHEG

$pp \rightarrow \ell^+ \ell^- \ell^{(\prime)+} \ell^{(\prime)-}$

Incoherent combination \rightarrow nNNLO+PS



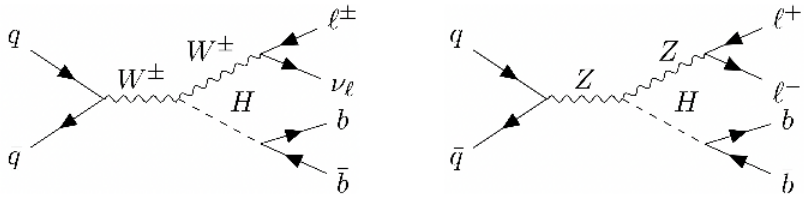
Sizeable effects of the **gluon loop-induced** channel at NLO in the bulk regions of distributions!



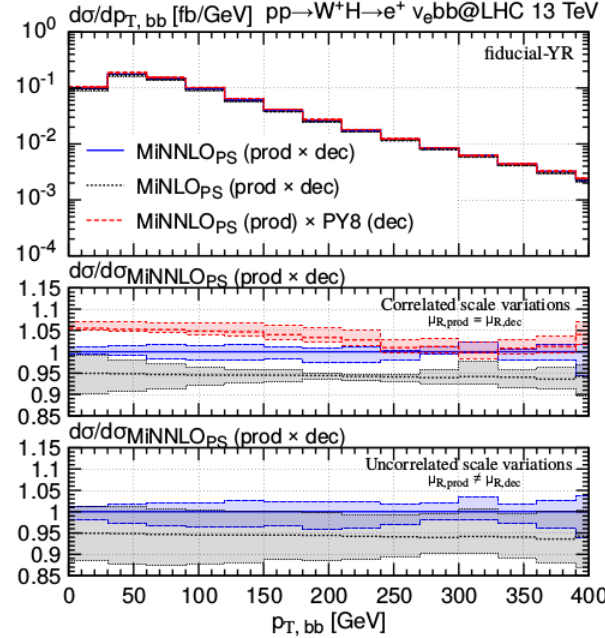
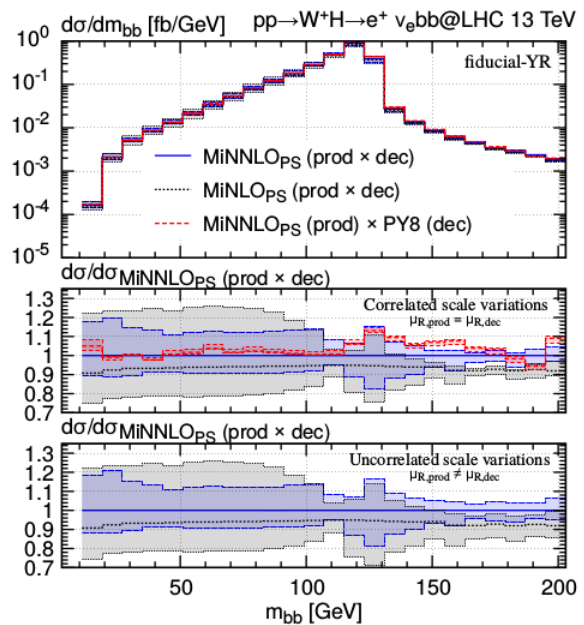
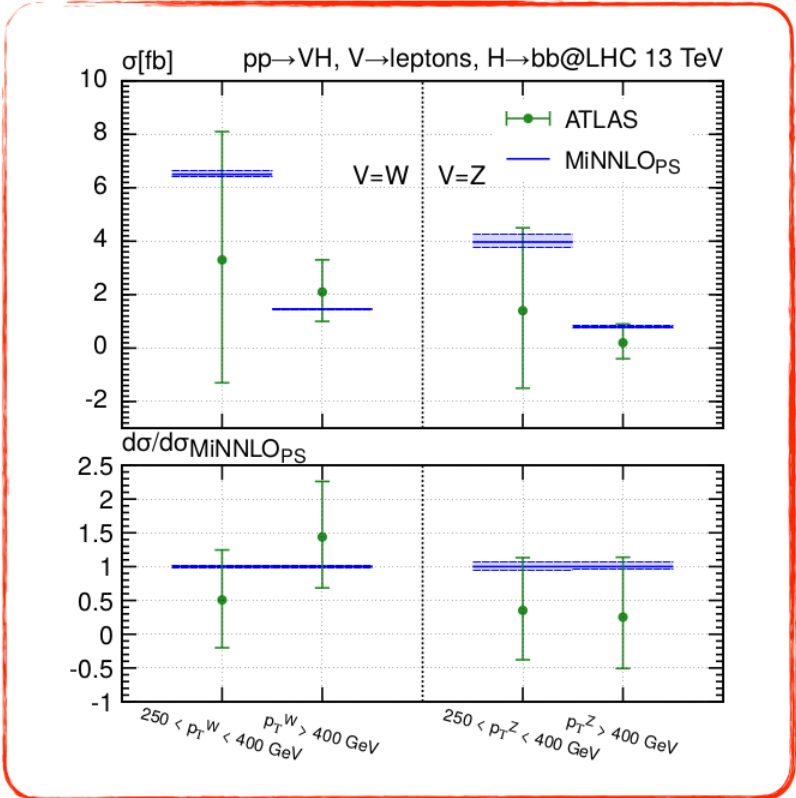
Comparison with CMS 137fb⁻¹ analysis!

NLO EW corrections needed for a reliable descriptions of tails of distributions

Higgsstrahlung with $H \rightarrow bb$ decay

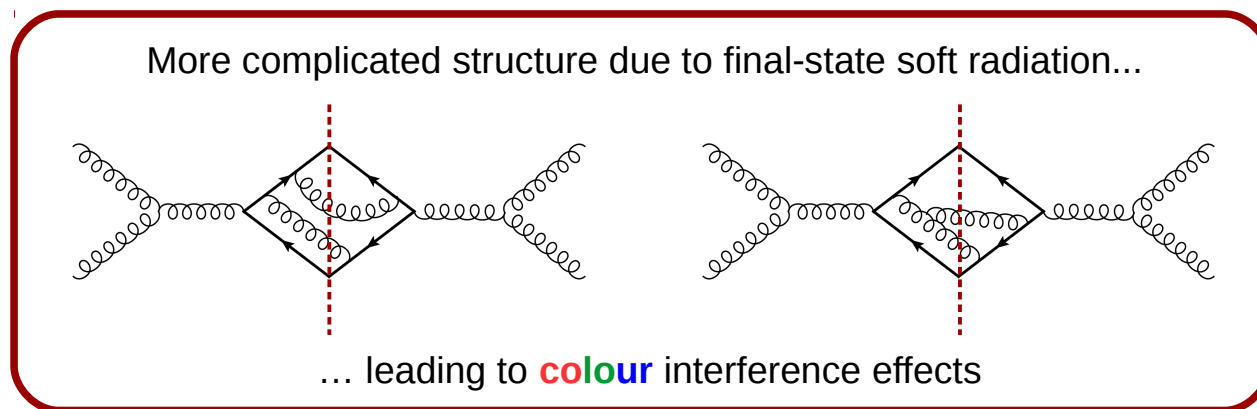


- ❖ Needed for **precision measurement** in the Higgs sector
- ❖ One of the main production channels + largest branching fraction in decay
- ❖ **NNLO+PS** accuracy in both **production and decay**



NNLO+PS for Heavy Quarks

- MiNNLO_{PS} method recently extended to top-quark pair production
- Qualitative advancement in the NNLO_{PS} field: first method beyond colour singlet
- Paving the way to more complicated processes
- Incredibly relevant for LHC phenomenology → ~40% of analyses using $t\bar{t}$ predictions



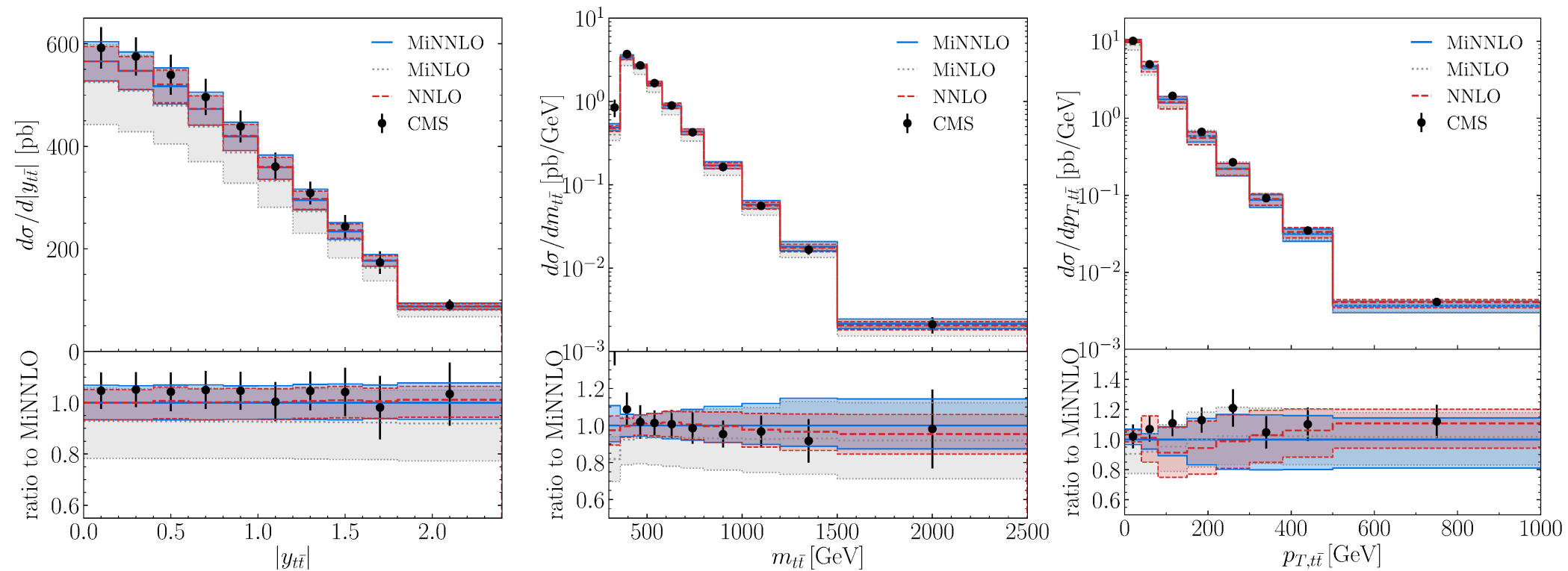
Colour singlet: $d\sigma^{(\text{sing})} \sim d\sigma_{c\bar{c}}^{(0)} \times \exp(-S_c) \times [HC_1C_2]_{c\bar{c};a_1a_2} \times f_{a_1}f_{a_2}$

Top-pair production: $d\sigma^{(\text{sing})} \sim d\sigma_{c\bar{c}}^{(0)} \times \exp(-S_c) \times [\text{Tr}(\mathbf{H}\Delta)C_1C_2]_{c\bar{c};a_1a_2} \times f_{a_1}f_{a_2}$

Effects coming from soft emissions from the FS contained in operator Δ

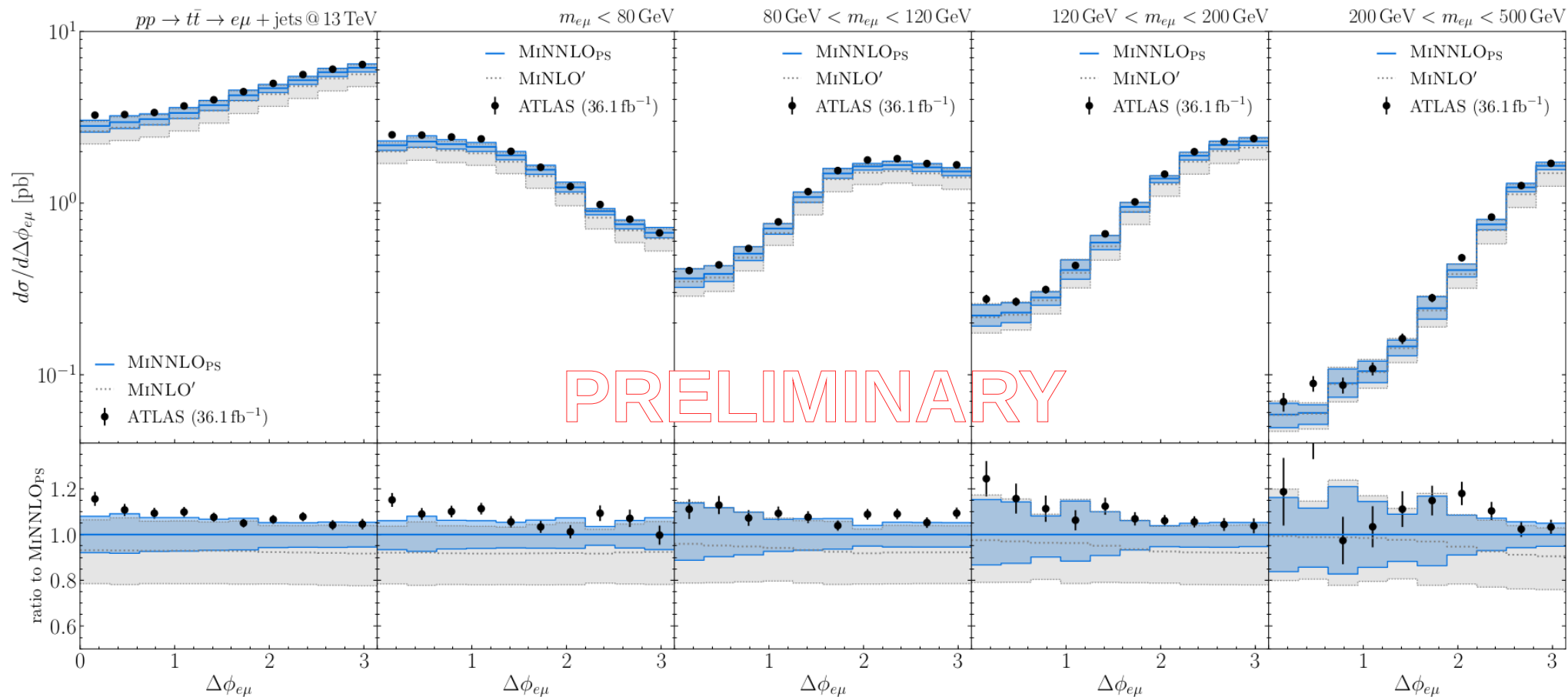
Top-pair Production at NNLO_{PS}

- Results for stable top quarks:



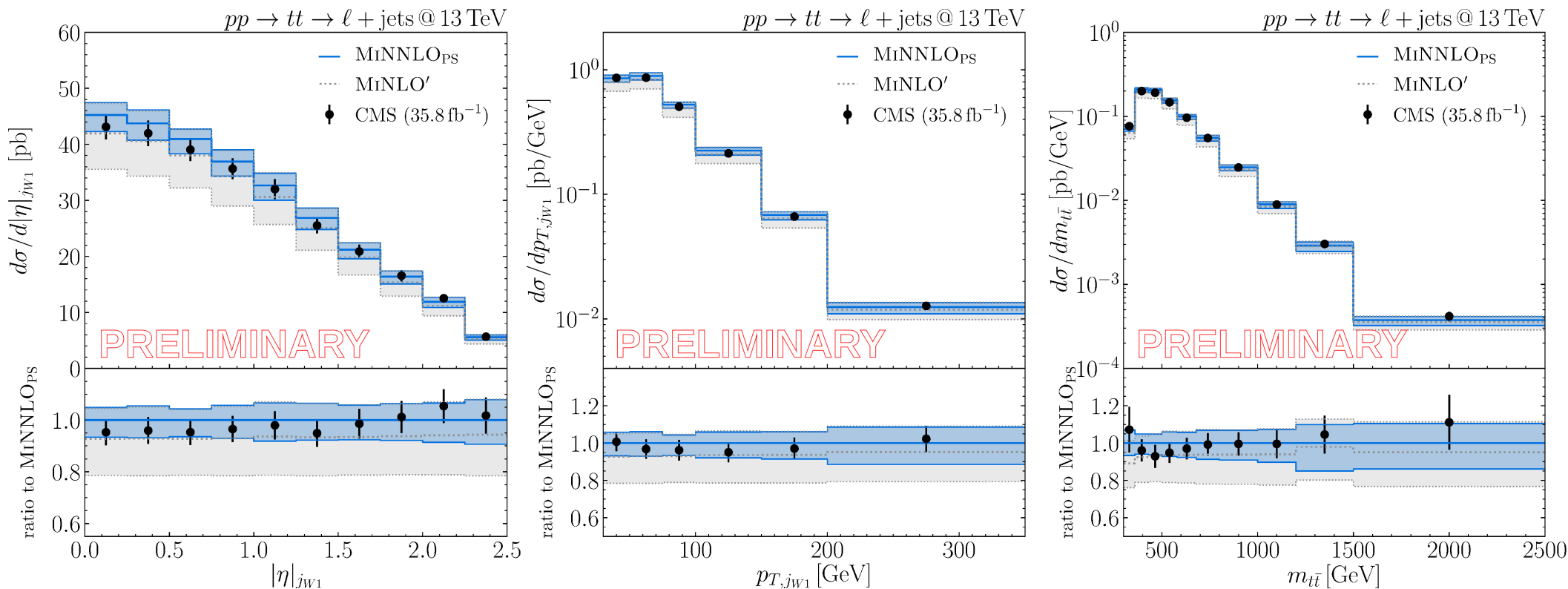
- Excellent agreement with NNLO fixed-order, both in normalization and shape!

- Including top decays using ratio of tree-level decayed and undecayed MEs
- Both W bosons decaying leptonically, comparison to ATLAS analysis:



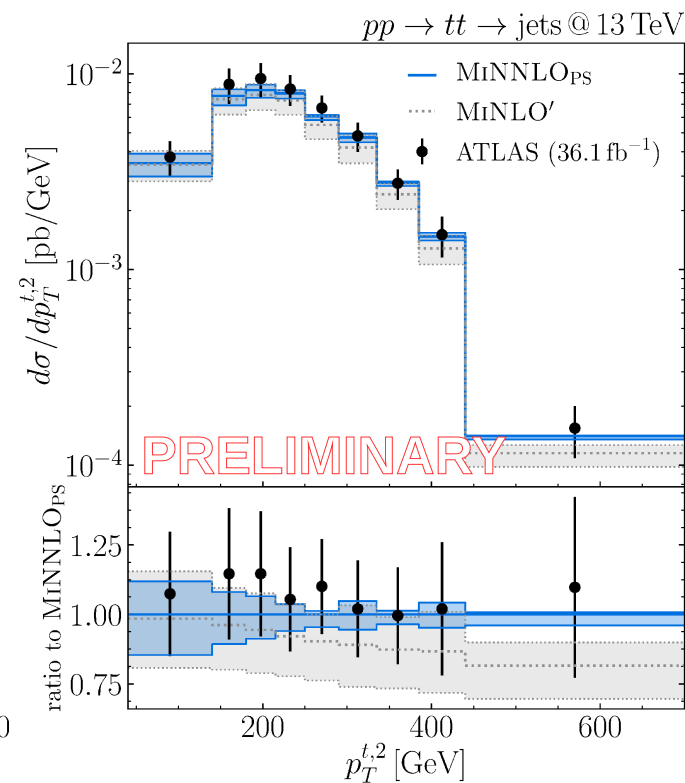
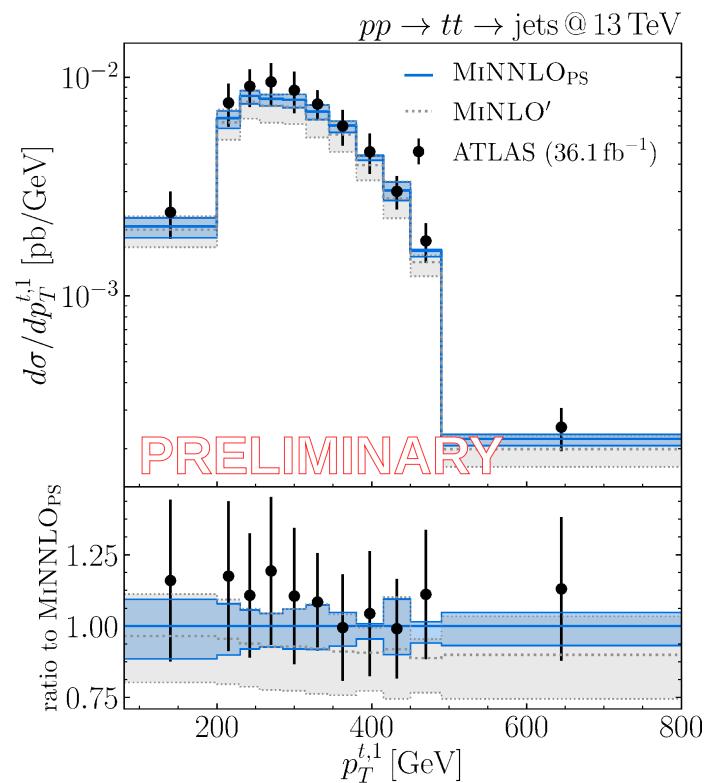
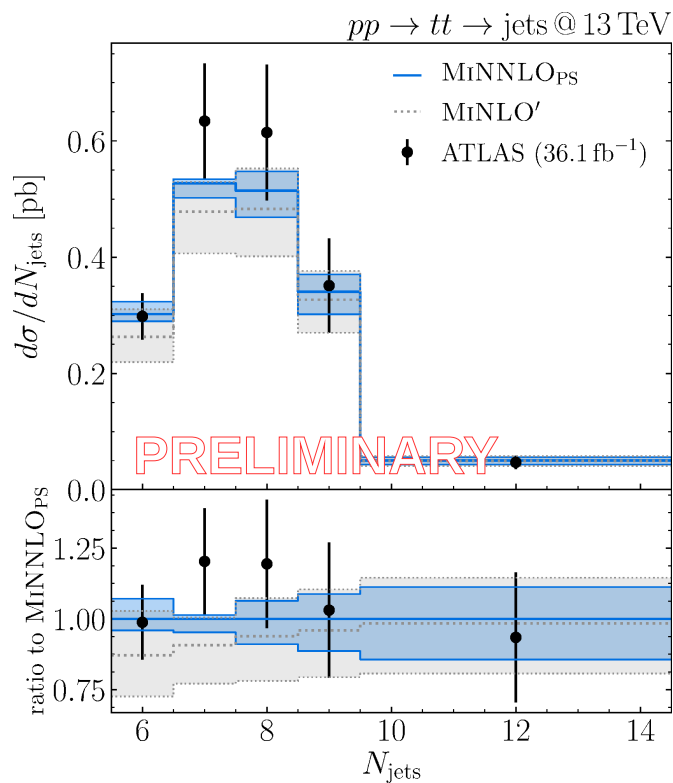
- Azimuthal angle between electron and muon → spin correlations

- Including top decays using ratio of tree-level decayed and undecayed MEs
- W bosons decaying semi-leptonically, comparison to CMS analysis:



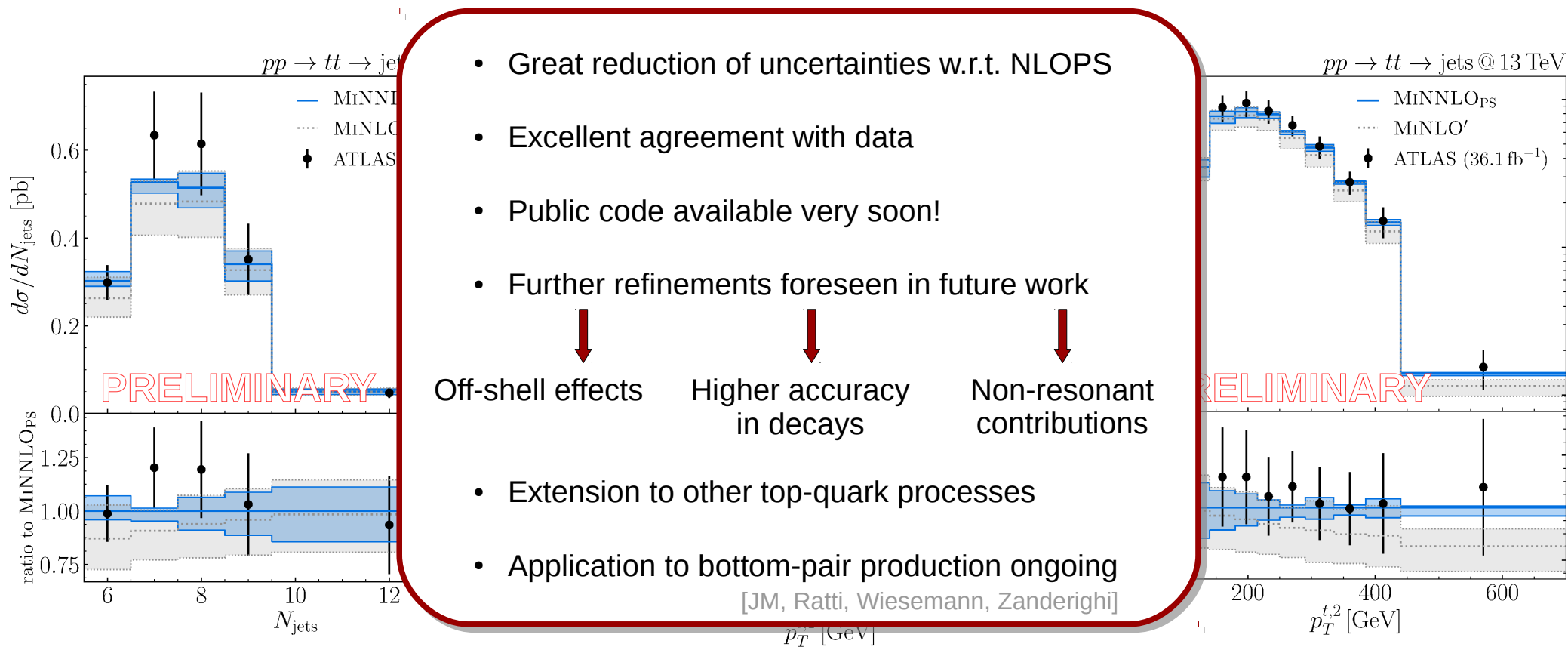
- Rapidity and p_T of leading jet coming from reconstructed W, and reconstructed $t\bar{t}$ invariant mass

- Including top decays using ratio of tree-level decayed and undecayed MEs
- Both W bosons decaying hadronically, comparison to ATLAS analysis:



- Rapidity and p_T of leading jet coming from reconstructed W, and reconstructed $t\bar{t}$ invariant mass

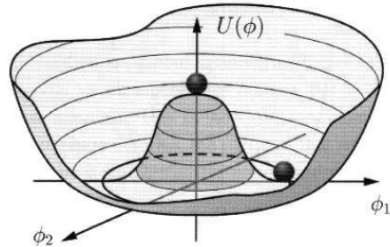
- Including top decays using ratio of tree-level decayed and undecayed MEs
- Both W bosons decaying hadronically, comparison to ATLAS analysis:



- Rapidity and p_T of leading jet coming from reconstructed W, and reconstructed $t\bar{t}$ invariant mass

The Higgs Potential

- Most of the SM already under scrutiny from LHC searches
- One sector barely explored so far: Higgs boson potential, responsible for EWSB

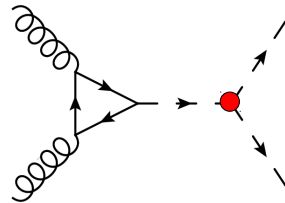


After EWSB, the potential leads to Higgs self interactions

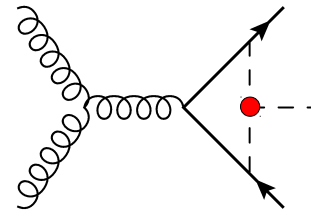
$$V(H) = \frac{1}{2}M_H^2 H^2 + \lambda v H^3 + \frac{1}{4}\lambda' H^4$$

$$\text{In the SM: } \lambda = \lambda' = M_H^2/(2v^2)$$

- Measuring or bounding the trilinear coupling is one of the main tasks of future LHC runs
- Two ways to measure λ



Higgs pair production

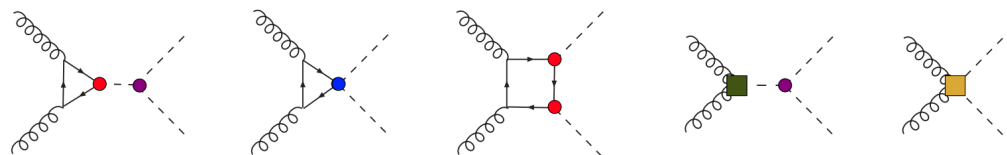


Loop-induced effects
in other observables

HH at NNLO with Anomalous Couplings

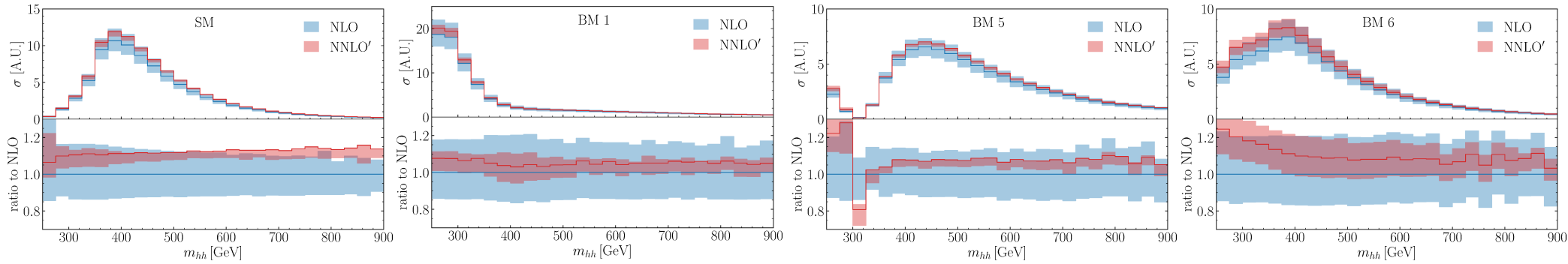
- We consider the HEFT operators affecting HH production

$$\mathcal{L} \supset -m_t \left(c_t \frac{h}{v} + c_{tt} \frac{h^2}{v^2} \right) - c_{hhh} \frac{m_h^2}{2v} h^3 + \frac{\alpha_s}{8\pi} \left(c_{ggh} \frac{h}{v} + c_{gghh} \frac{h^2}{v^2} \right) G_{\mu\nu}^a G^{a,\mu\nu}$$



- NLO corrections → full top mass dependence
- NNLO piece → Born-improved Heavy Top Limit

} NNLO'

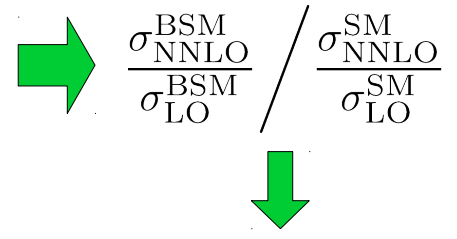
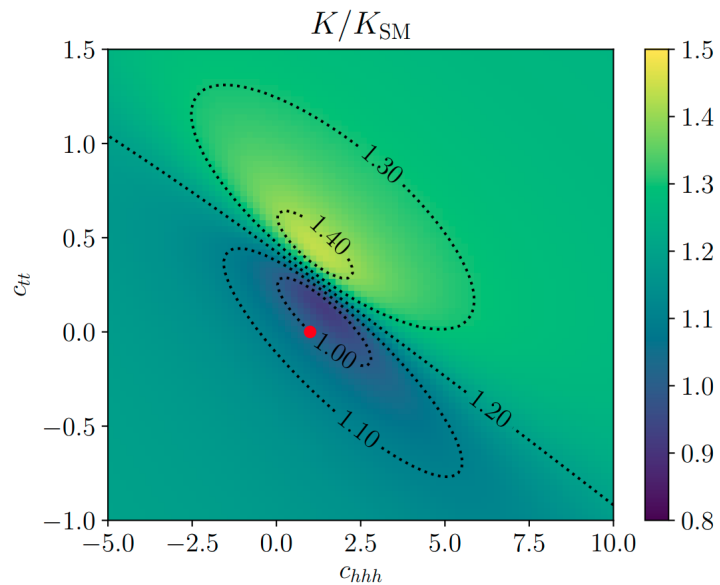
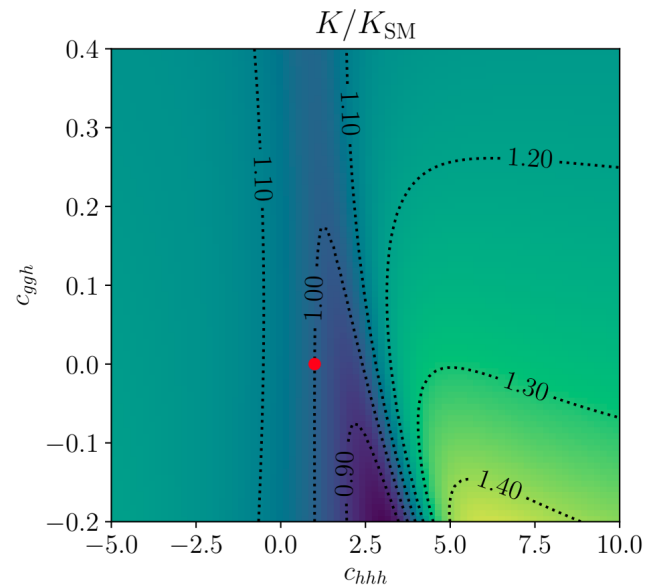


HH at NNLO with Anomalous Couplings

- NNLO cross section can be parameterized in terms of anomalous couplings

$$\begin{aligned}
 \sigma_{\text{BSM}}/\sigma_{\text{SM}} = & a_1 c_t^4 + a_2 c_{tt}^2 + a_3 c_t^2 c_{hhh}^2 + a_4 c_{ggh}^2 c_{hhh}^2 + a_5 c_{gghh}^2 + a_6 c_{tt} c_t^2 + a_7 c_t^3 c_{hhh} \\
 & + a_8 c_{tt} c_t c_{hhh} + a_9 c_{tt} c_{ggh} c_{hhh} + a_{10} c_{tt} c_{gghh} + a_{11} c_t^2 c_{ggh} c_{hhh} + a_{12} c_t^2 c_{gghh} \\
 & + a_{13} c_t c_{hhh}^2 c_{ggh} + a_{14} c_t c_{hhh} c_{gghh} + a_{15} c_{ggh} c_{hhh} c_{gghh} + a_{16} c_t^3 c_{ggh} \\
 & + a_{17} c_t c_{tt} c_{ggh} + a_{18} c_t c_{ggh}^2 c_{hhh} + a_{19} c_t c_{ggh} c_{gghh} + a_{20} c_t^2 c_{ggh}^2 \\
 & + a_{21} c_{tt} c_{ggh}^2 + a_{22} c_{ggh}^3 c_{hhh} + a_{23} c_{ggh}^2 c_{gghh} + a_{24} c_{ggh}^4 + a_{25} c_{ggh}^3 c_t
 \end{aligned}$$

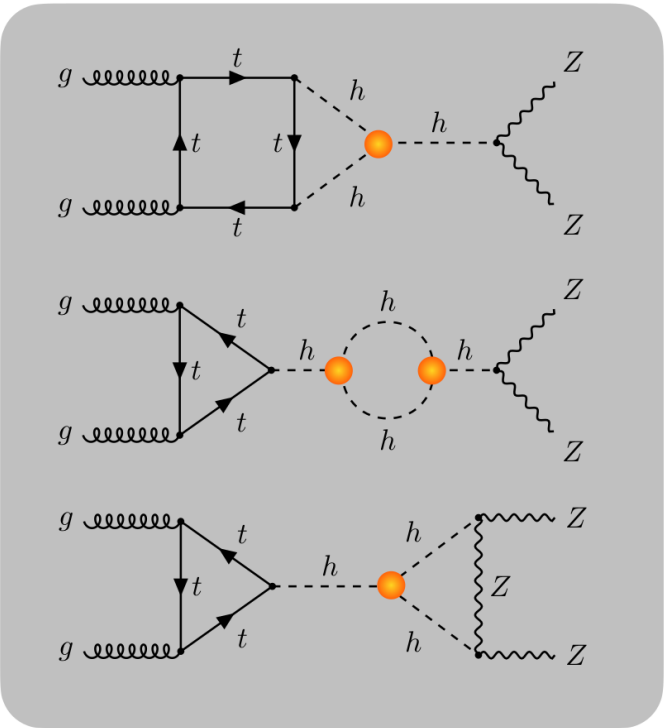
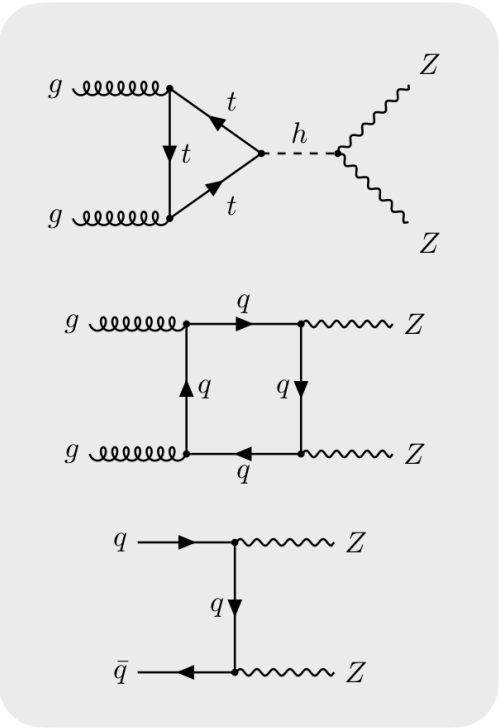
- We performed a fit and obtained the values of the a_i coefficients



Can present ~40% variations depending on the couplings

Importance of computing QCD corrections for BSM couplings

Limits from off-shell H production



Calculation in SMEFT



Two operators affecting the trilinear coupling



$$\mathcal{O}_6 = -\lambda|H|^6,$$

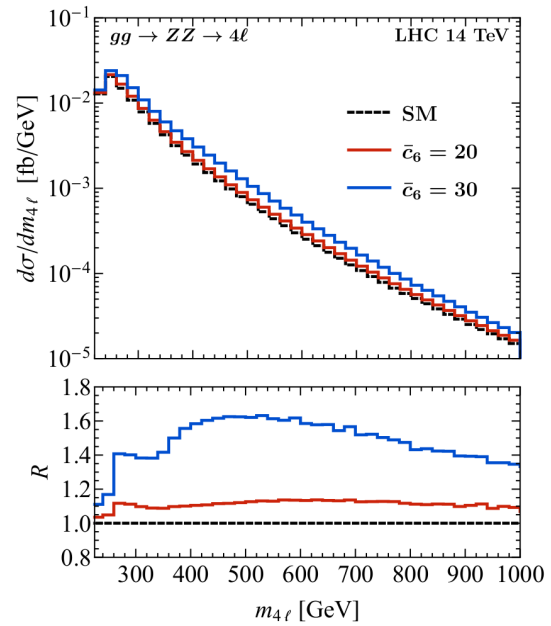
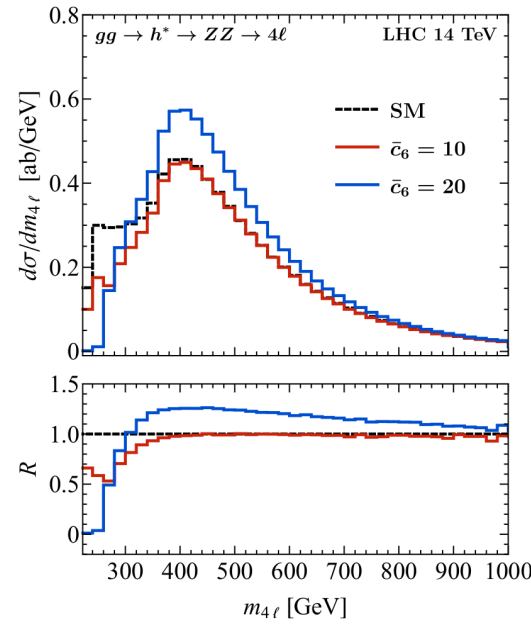


$$\mathcal{O}_H = \frac{1}{2}(\partial_\mu|H|)^2$$

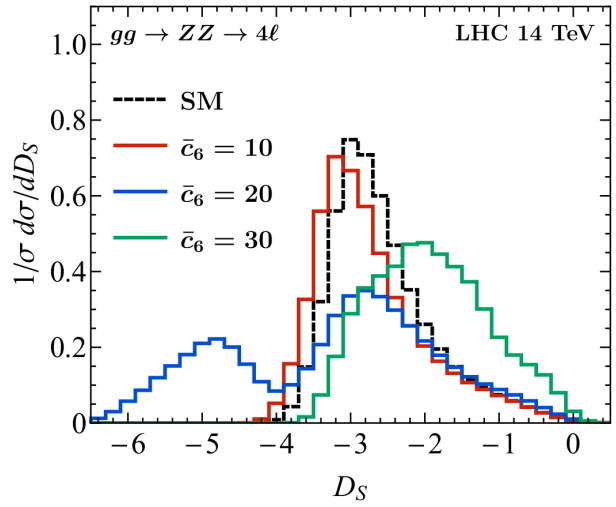
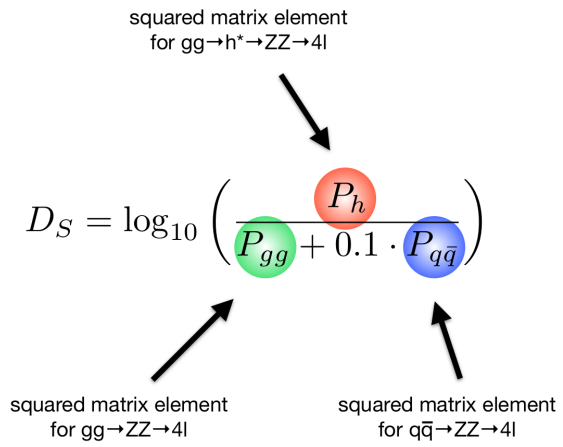
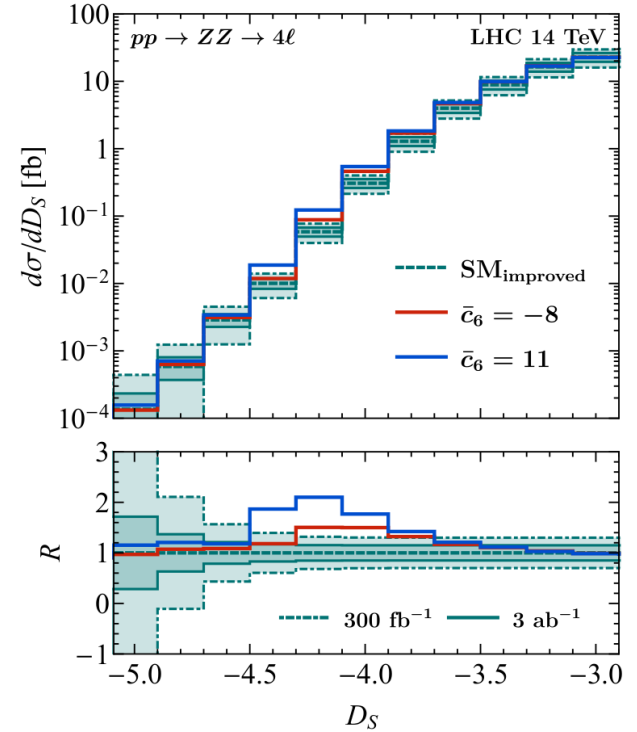


$$c_3 = 1 + \bar{c}_6 - \frac{3}{2}\bar{c}_H$$

- Both resonant and non-resonant contributions included
- Effect from SMEFT operators \mathcal{O}_6 and \mathcal{O}_H on self coupling (c_3) via loop corrections

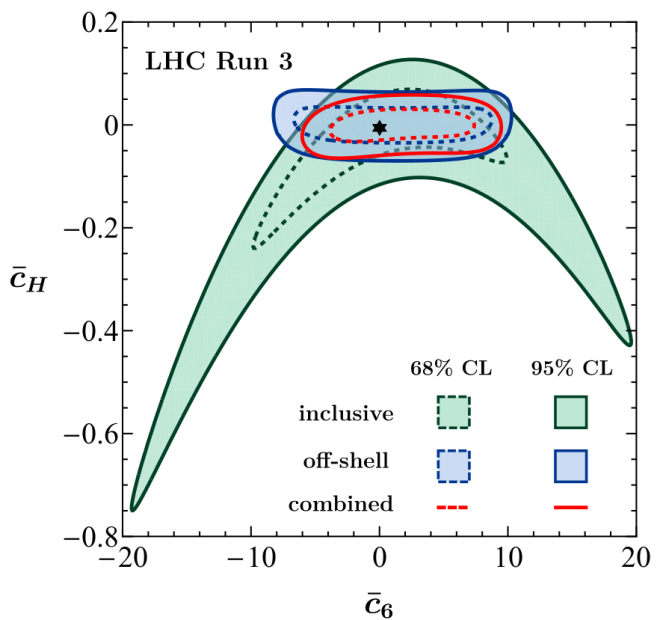


- Four lepton invariant mass and ME-based discriminant D_S have been studied
- Phenomenologically relevant kinematic effects found in both spectra



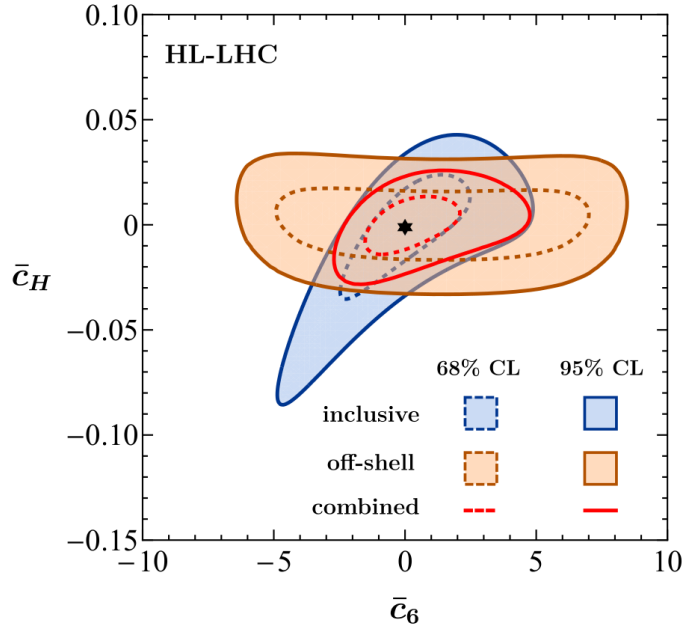
Projected Limits on c_3

- Projected limits for Run3 and HL-LHC have been derived
- Compared to the projected limits from inclusive single-H production
- Bounds found to be competitive and complementary to the inclusive ones → remove flat directions



Combined 95% C.L. limits:

$$c_3 \in [-4.0, 6.1]$$



$$c_3 \in [-1.7, 5.7]$$

- Projections from double-Higgs production will strengthen the constraints

Summary

- Without clear NP signals, quest for BSM is focusing more in small deviations
- Precise theory predictions are in this context indispensable
- Big contributions to the field coming from our group in the last year:
 - NNLO_{PS} for diboson production
 - NNLO_{PS} for top-pair production
 - Precise determination of photon and lepton PDFs
 - Studies related to the Higgs self coupling
 - ...
- Promising prospects for the future research activities!

Thanks!