

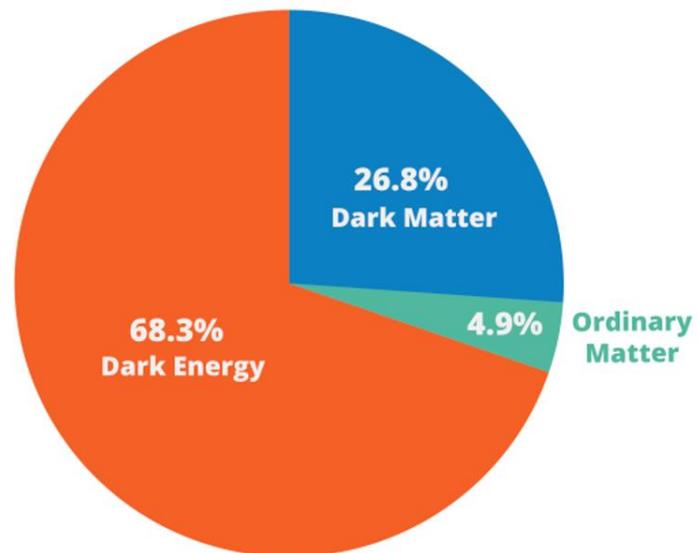
A photograph of a sunset over the ocean. The sun is a bright, glowing orb in the upper left quadrant, casting a warm, golden light across the sky and the water. The sky is filled with soft, wispy clouds. In the distance, the silhouette of the IceCube detector structure is visible on the horizon. The water in the foreground is dark and calm.

Constraining generalized Dark Matter-Nucleon Interactions based on Dark Matter Capture in the Sun with IceCube

Lilly Peters

Dark Matter

Estimated matter-energy content of the Universe



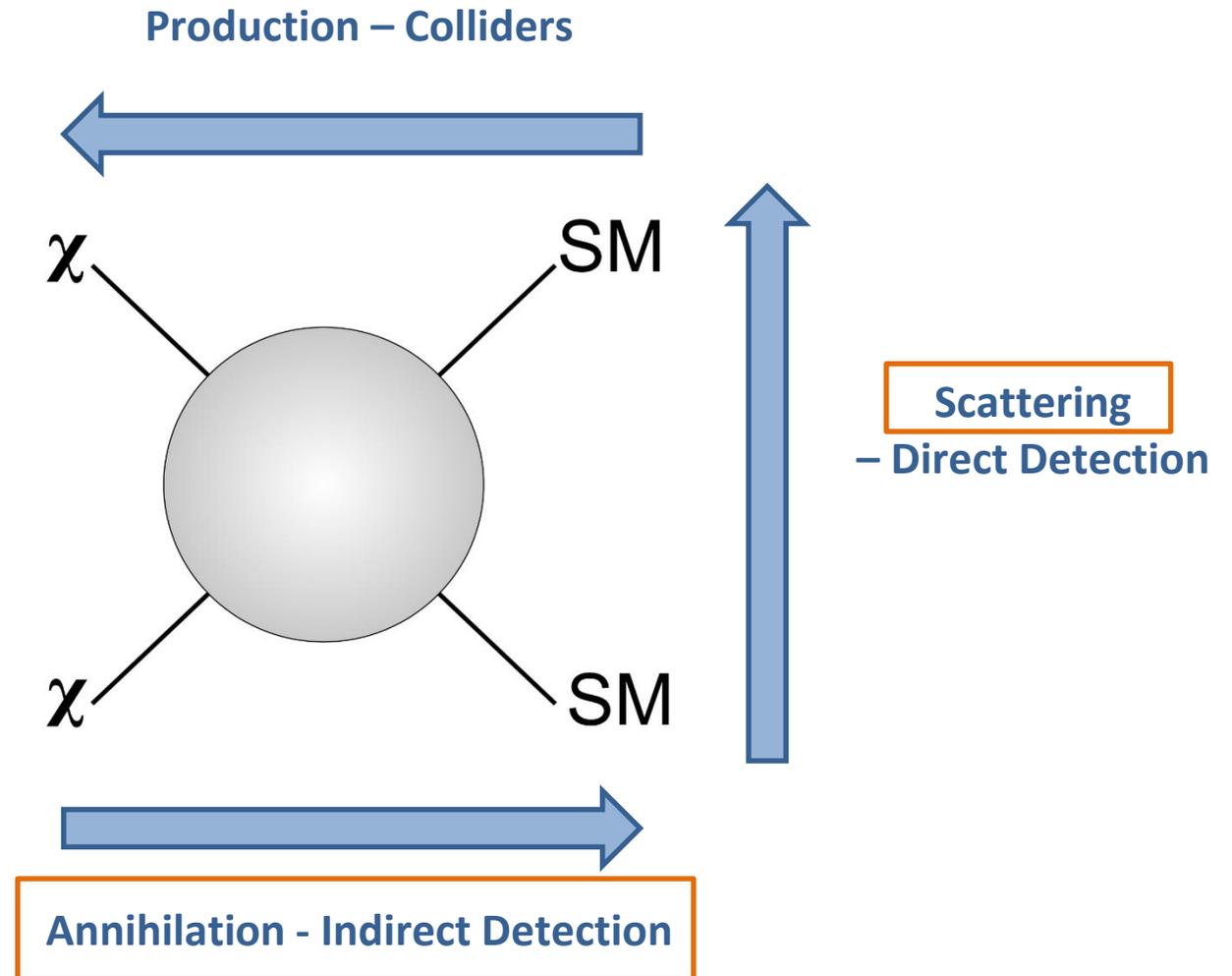
- Non-luminous
- Compelling observational evidence
- Its nature remains unknown

How to detect dark matter?

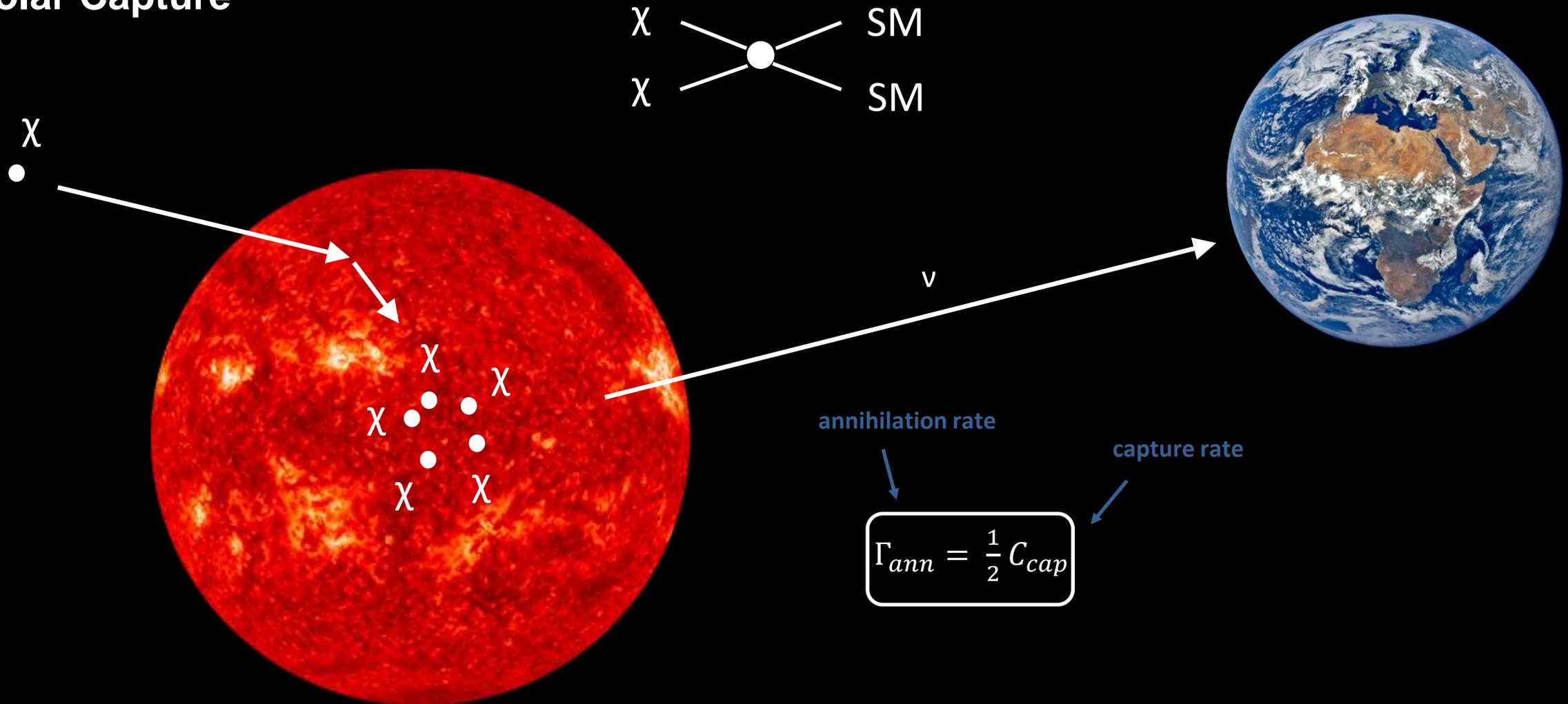
Dark Matter Detection

- Different interaction channels/search strategies for dark matter particles
- This work: **Solar capture**

→ special: **indirect** detection, but sensitive to **scattering** cross section



Solar Capture



Effective Theory

Motivation: standard spin-dependent (SD) and spin-independent (SI) interactions can be suppressed or forbidden
 → include **velocity** and **momentum dependent** interactions

„standard assumption“

SI

SD

$$O_1 = \mathbb{1}_{\chi N}$$

$$O_3 = i\hat{\mathbf{S}}_N \cdot \left(\frac{\hat{\mathbf{q}}}{m_N} \times \hat{\mathbf{v}}^\perp \right) \mathbb{1}_\chi$$

$$O_4 = \hat{\mathbf{S}}_\chi \cdot \hat{\mathbf{S}}_N$$

$$O_5 = i\hat{\mathbf{S}}_\chi \cdot \left(\frac{\hat{\mathbf{q}}}{m_N} \times \hat{\mathbf{v}}^\perp \right) \mathbb{1}_N$$

$$O_6 = \left(\hat{\mathbf{S}}_\chi \cdot \frac{\hat{\mathbf{q}}}{m_N} \right) \left(\hat{\mathbf{S}}_N \cdot \frac{\hat{\mathbf{q}}}{m_N} \right)$$

$$O_7 = \hat{\mathbf{S}}_N \cdot \hat{\mathbf{v}}^\perp \mathbb{1}_\chi$$

$$O_8 = \hat{\mathbf{S}}_\chi \cdot \hat{\mathbf{v}}^\perp \mathbb{1}_N$$

$$O_9 = i\hat{\mathbf{S}}_\chi \cdot \left(\hat{\mathbf{S}}_N \times \frac{\hat{\mathbf{q}}}{m_N} \right)$$

$$O_{10} = i\hat{\mathbf{S}}_N \cdot \frac{\hat{\mathbf{q}}}{m_N} \mathbb{1}_\chi$$

$$O_{11} = i\hat{\mathbf{S}}_\chi \cdot \frac{\hat{\mathbf{q}}}{m_N} \mathbb{1}_N$$

$$O_{12} = \hat{\mathbf{S}}_\chi \cdot \left(\hat{\mathbf{S}}_N \times \hat{\mathbf{v}}^\perp \right)$$

$$O_{13} = i \left(\hat{\mathbf{S}}_\chi \cdot \hat{\mathbf{v}}^\perp \right) \left(\hat{\mathbf{S}}_N \cdot \frac{\hat{\mathbf{q}}}{m_N} \right)$$

$$O_{14} = i \left(\hat{\mathbf{S}}_\chi \cdot \frac{\hat{\mathbf{q}}}{m_N} \right) \left(\hat{\mathbf{S}}_N \cdot \hat{\mathbf{v}}^\perp \right)$$

$$O_{15} = - \left(\hat{\mathbf{S}}_\chi \cdot \frac{\hat{\mathbf{q}}}{m_N} \right) \left[\left(\hat{\mathbf{S}}_N \times \hat{\mathbf{v}}^\perp \right) \cdot \frac{\hat{\mathbf{q}}}{m_N} \right]$$

$$O_{17} = i \frac{\hat{\mathbf{q}}}{m_N} \cdot \mathcal{S} \cdot \hat{\mathbf{v}}^\perp \mathbb{1}_N$$

$$O_{18} = i \frac{\hat{\mathbf{q}}}{m_N} \cdot \mathcal{S} \cdot \hat{\mathbf{S}}_N$$

$$O_{19} = \frac{\hat{\mathbf{q}}}{m_N} \cdot \mathcal{S} \cdot \frac{\hat{\mathbf{q}}}{m_N}$$

$$O_{20} = \left(\hat{\mathbf{S}}_N \times \frac{\hat{\mathbf{q}}}{m_N} \right) \cdot \mathcal{S} \cdot \frac{\hat{\mathbf{q}}}{m_N}$$

$$\mathcal{L}_{\text{int}} = \sum_{N=n,p} \sum_i c_i^N \mathcal{O}_i \chi^+ \chi^- N^+ N^-$$

isoscalar:

$$c^0 = c^p + c^n$$

isovector:

$$c^1 = c^p - c^n$$

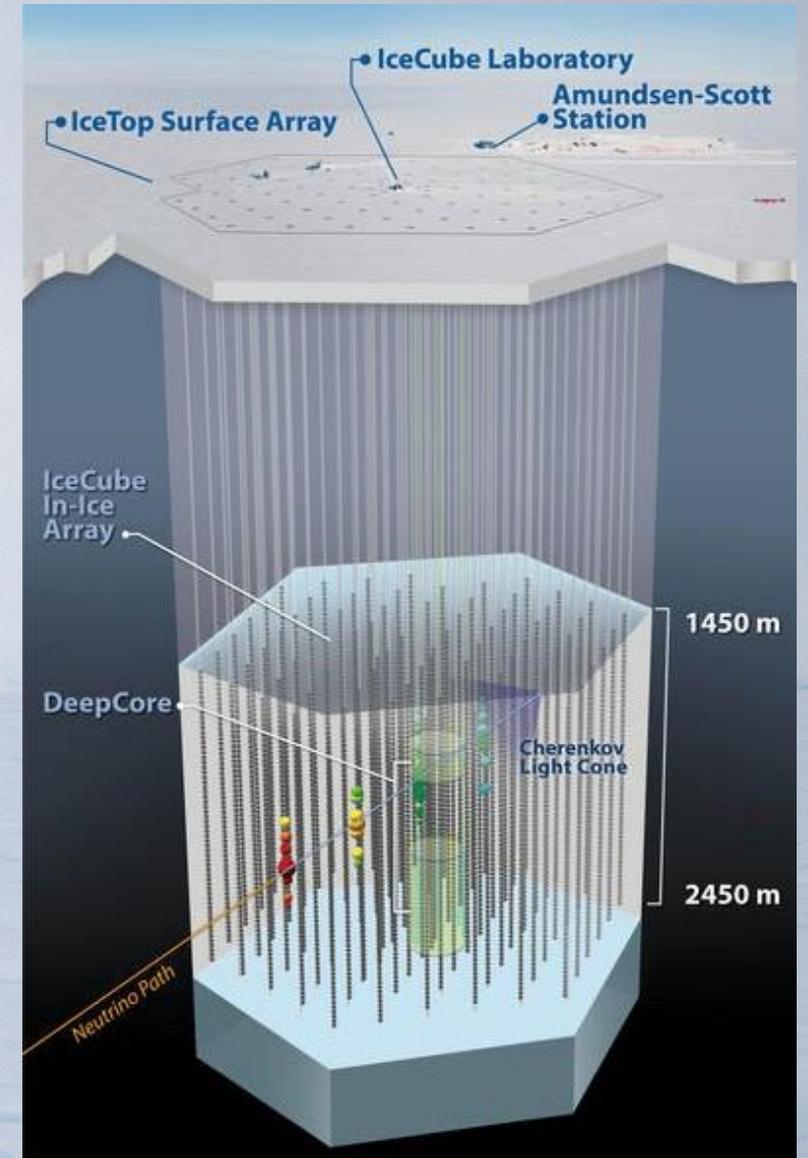
Goal: Constrain isoscalar and isovector **coupling constants** of all interaction operators (4, 14 and 18 for spin-0, spin-1/2 and spin-1 DM)

The IceCube Neutrino Observatory



Image: Yuya Makino, IceCube/NSF

- Neutrinos undergo **deep-inelastic scattering** in the Antarctic ice
- Optical detection of **Cherenkov light** emitted by secondary charged particles



Method

Goal: Constrain isoscalar and isovector coupling constants of all interaction operators

- Capture rate is **proportional** to coupling constant squared: $C_{cap} \propto c^2$
- Use proportionality to **convert limits** from IceCube solar dark matter searches

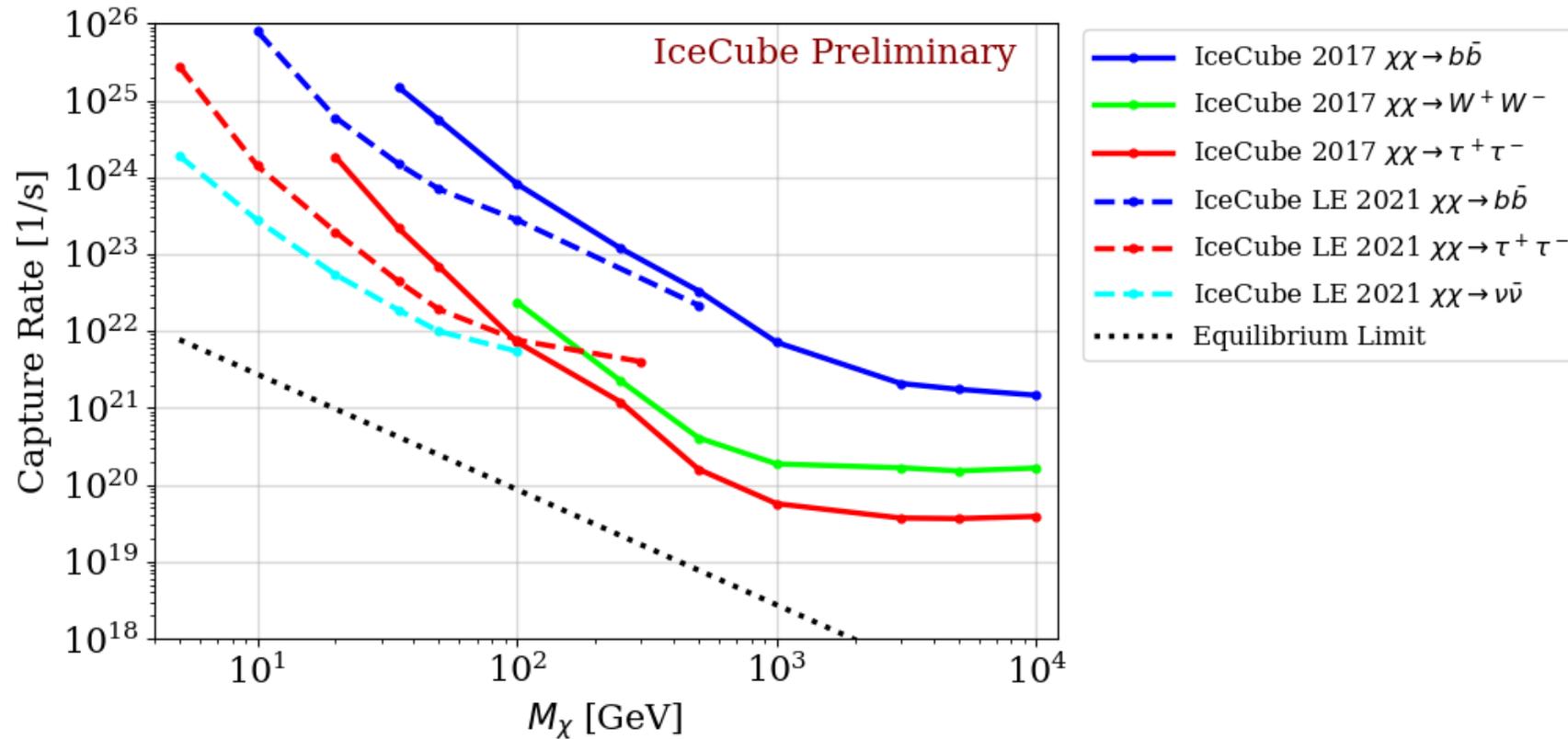
$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$

Calculate **capture rates** in the Sun for each interaction operator assuming $c_i = c_0$
(All other coupling constants are set to zero)

- Choose $c_0 = 10^{-3} m_v^{-2}$ $m_v = 246.2 \text{ GeV}$.

Capture Rate Limits

$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$



2017:

- 3 years of data

LE 2021:

- 7 years of data
- low energy (DeepCore)
- also annihilation directly into neutrinos
- not published yet

new!

Capture Rate Calculation

$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$



Capture Rate Calculation

$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$



$$C_{cap}^T = n_\chi \int_0^{R_\odot} dr 4\pi r^2 n_T(r) \int_0^\infty du 4\pi u^2 f_\odot(u) \frac{u^2 + v_\odot^{esc}(r)^2}{u} \cdot \int_{E_{min}}^{E_{max}} dE_R \frac{d\sigma_T}{dE_R} \theta(\Delta E)$$

↑ Sum over 16 most abundant elements
 ↑ DM density
 ↑ Solar Model
 ↑ DM velocity distribution

Capture Rate Calculation

$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$

$$C_{cap}^T = n_\chi \int_0^{R_\odot} dr 4\pi r^2 n_T(r) \int_0^\infty du 4\pi u^2 f_\odot(u) \frac{u^2 + v_\odot^{esc}(r)^2}{u} \cdot \int_{E_{min}}^{E_{max}} dE_R \frac{d\sigma_T}{dE_R} \theta(\Delta E)$$

Sum over 16
most abundant
elements

DM density

Solar Model

DM velocity distribution

DM response functions → Contain the coupling constants

$$\frac{d\sigma_T}{dE_R}(\omega^2, q^2) = \frac{2m_T}{\omega^2} \frac{1}{2J+1} \sum_{\tau, \tau'} \left[\sum_{k=M, \Sigma', \Sigma''} R_k^{\tau\tau'} \left(v_T^{\perp 2}, \frac{q^2}{m_N^2} \right) W_k^{\tau\tau'}(q^2) + \frac{q^2}{m_N^2} \sum_{k=\Phi'', \Phi''M, \tilde{\Phi}', \Delta, \Delta\Sigma'} R_k^{\tau\tau'} \left(v_T^{\perp 2}, \frac{q^2}{m_N^2} \right) W_k^{\tau\tau'}(q^2) \right]$$

Nuclear response functions

Nuclear response operators

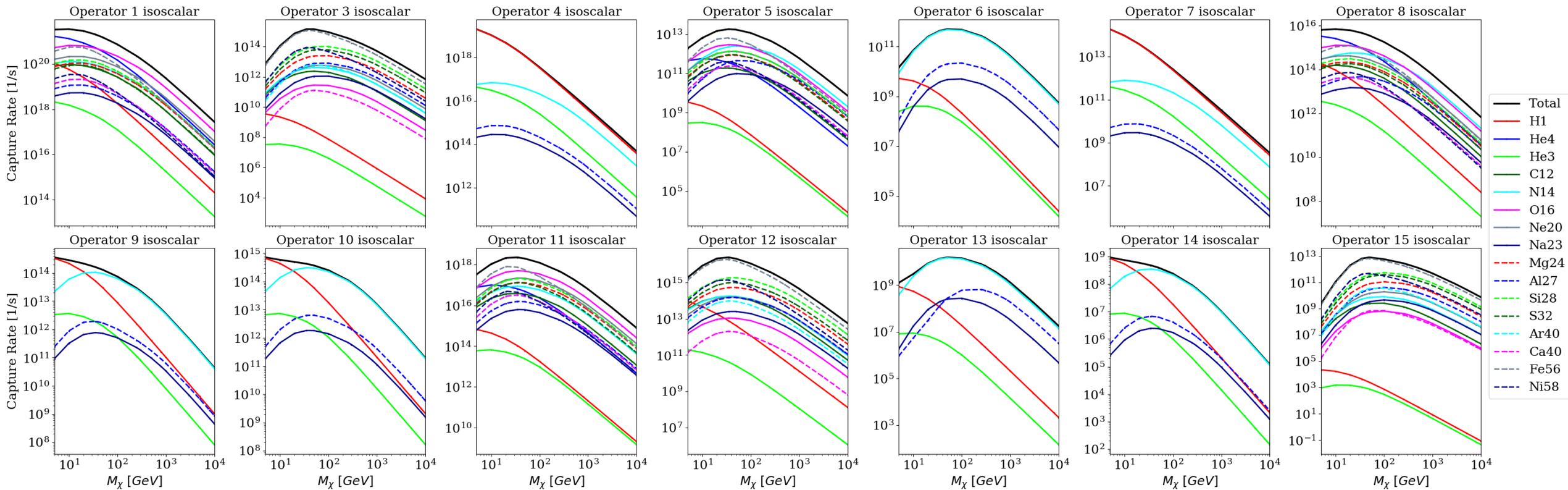
- Multiple expansion of nuclear charges and currents gives six nuclear response operators: $M, \Sigma', \Sigma'', \Delta, \Phi', \Phi''$
- Measure different quantities (spin, mass, ...)
- Favor different elements

Results

- Capture Rates
- Systematic Uncertainties
- Limits on the Coupling Constants

Capture Rates

spin $\frac{1}{2}$ isoscalar



- Most relevant elements can be **H1, He4, N14, O16, Al27 or Fe56**
- Dependence on **momentum transfer** \rightarrow suppression of lighter elements

Systematics

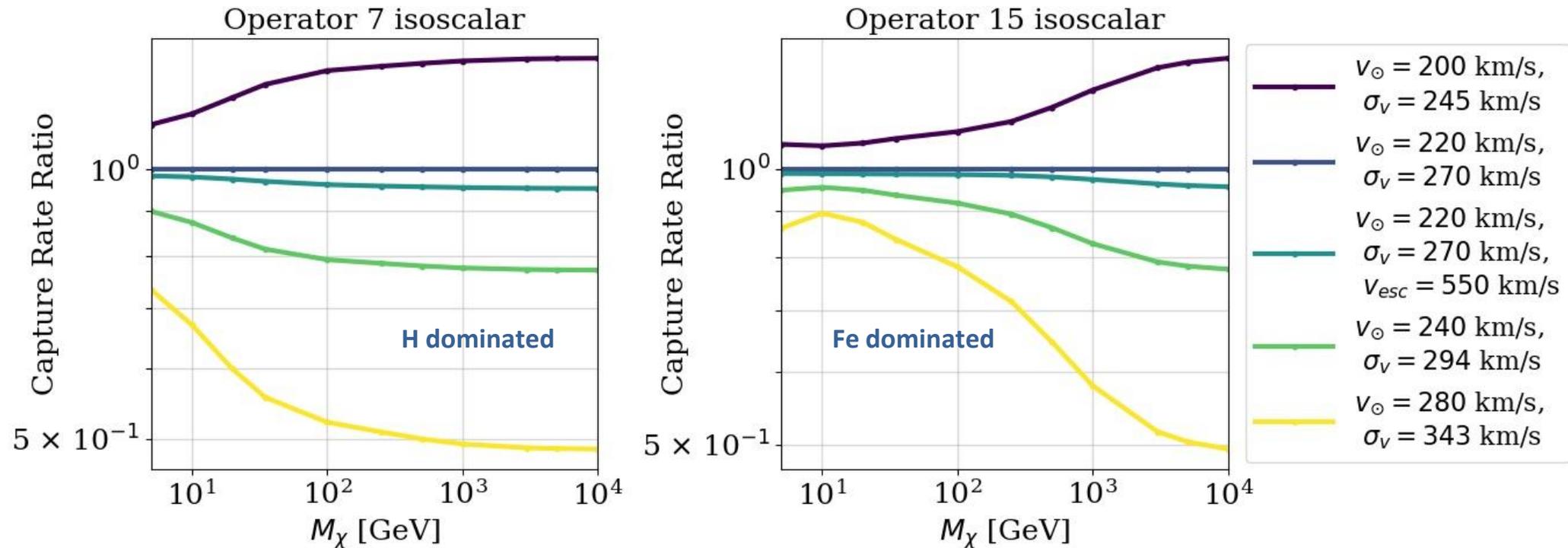
$$C_{\text{cap}}^{\text{T}} = n_{\chi} \int_0^{R_{\odot}} dr 4\pi r^2 n_{\text{T}}(r) \int_0^{\infty} du 4\pi u^2 f_{\odot}(u) \frac{u^2 + v_{\odot}^{\text{esc}}(r)^2}{u} \cdot \int_{E_{\text{min}}}^{E_{\text{max}}} dE_{\text{R}} \frac{d\sigma_{\text{T}}}{dE_{\text{R}}} \theta(\Delta E)$$

↑
DM density
↑
Solar Model
↑
DM velocity distribution

Systematics – Velocity Distribution

spin $\frac{1}{2}$ isoscalar

- Standard Assumptions: **Maxwellian** distribution with rotational velocity of 220 km/s, dispersion of 270 km/s and no galactic escape velocity

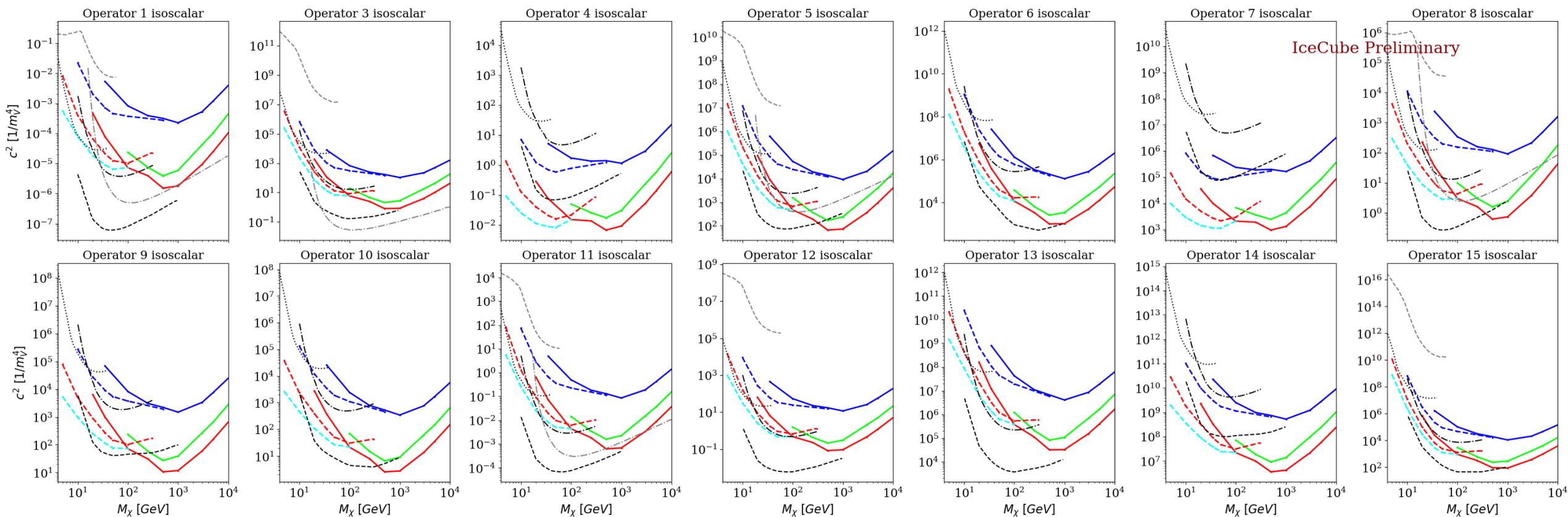


→ Operators which are dominated by **hydrogen** are affected at **smaller dark matter masses**

Limits

$$(c_i^{limit})^2 = \frac{C_{cap}^{limit}}{C_{cap,i}} c_0^2$$

spin 1/2 isoscalar

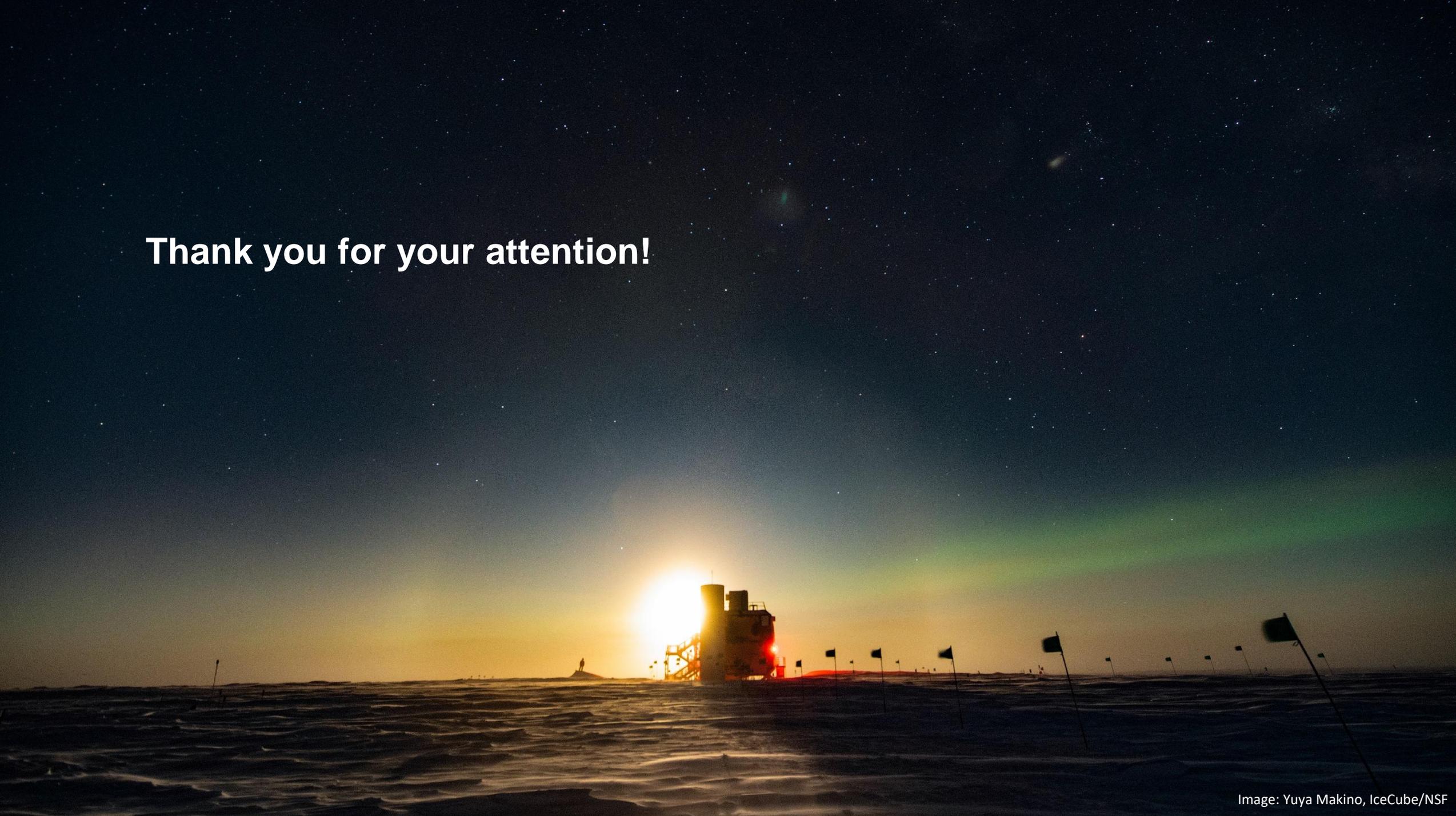


Summary

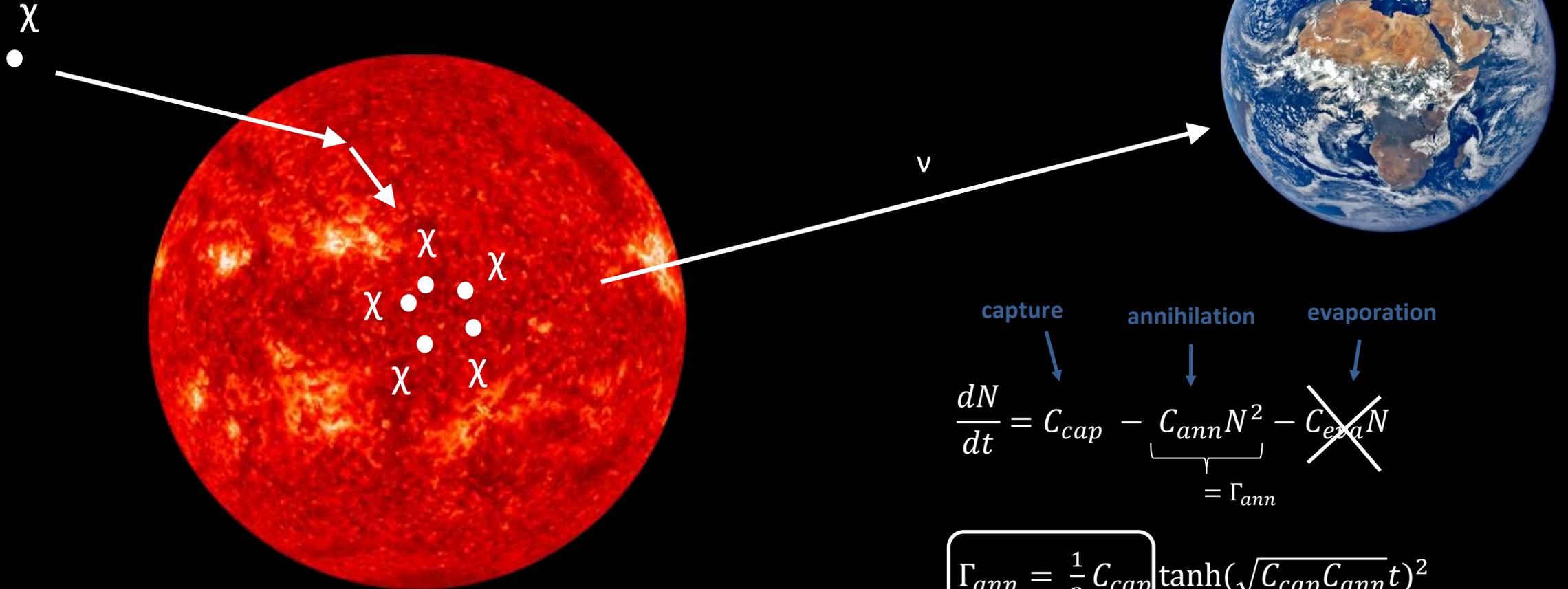


- Calculated **capture rates** of dark matter in the Sun within the **general effective theory** of isoscalar and isovector **dark matter-nucleon interactions** in the non-relativistic regime
- The theory predicts **four, 14 or 18 interaction operators** for dark matter particles with **spin 0, spin 1/2 or spin 1** with a non-trivial **dependence on velocity and momentum transfer**
- Studied impact of **systematic uncertainties** on the capture rate, in particular the **velocity distribution** of dark matter particles in the Galaxy and the **elemental composition** of the Sun
- Used the computed capture rates **to convert exclusion limits** on the capture rate by IceCube into **limits on the coupling constants** of the theory
- Leading limits for various interaction types, dark matter masses and annihilation channels

Thank you for your attention!



Solar Capture



$$\frac{dN}{dt} = \overset{\text{capture}}{\downarrow} C_{cap} - \underbrace{C_{ann} N^2}_{= \Gamma_{ann}} - \overset{\text{evaporation}}{\downarrow} \cancel{C_{eva} N}$$

$$\boxed{\Gamma_{ann} = \frac{1}{2} C_{cap}} \underbrace{\tanh(\sqrt{C_{cap} C_{ann} t})^2}_{= 1 \text{ for } t \gg 1/\sqrt{C_{cap} C_{ann}}}$$

DM Response Functions

$$R_M^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = c_1^\tau c_1^{\tau'} + \frac{j_X(j_X + 1)}{3} \left(\frac{q^2}{m_N^2} v_T^{\perp 2} c_5^\tau c_5^{\tau'} + v_T^{\perp 2} c_8^\tau c_8^{\tau'} + \frac{q^2}{m_N^2} c_{11}^\tau c_{11}^{\tau'} \right) \quad (4.3)$$

$$R_{\Phi''}^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = \frac{q^2}{4m_N^2} c_3^\tau c_3^{\tau'} + \frac{j_X(j_X + 1)}{12} \left(c_{12}^\tau - \frac{q^2}{m_N^2} c_{15}^\tau \right) \cdot \left(c_{12}^{\tau'} - \frac{q^2}{m_N^2} c_{15}^{\tau'} \right) \quad (4.4)$$

$$R_{\Phi'}^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = \frac{j_X(j_X + 1)}{12} \left(c_{12}^\tau c_{12}^{\tau'} + \frac{q^2}{m_N^2} c_{13}^\tau c_{13}^{\tau'} \right) \quad (4.5)$$

$$R_{\Sigma''}^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = \frac{q^2}{4m_N^2} c_{10}^\tau c_{10}^{\tau'} + \frac{j_X(j_X + 1)}{12} \left[c_4^\tau c_4^{\tau'} + \frac{q^2}{m_N^2} \left(c_4^\tau c_6^{\tau'} + c_6^\tau c_4^{\tau'} \right) + \frac{q^4}{m_N^4} c_6^\tau c_6^{\tau'} + v_T^{\perp 2} c_{12}^\tau c_{12}^{\tau'} + \frac{q^2}{m_N^2} v_T^{\perp 2} c_{13}^\tau c_{13}^{\tau'} \right] \quad (4.6)$$

$$R_{\Sigma'}^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = \frac{1}{8} \left(\frac{q^2}{m_N^2} v_T^{\perp 2} c_3^\tau c_3^{\tau'} + v_T^{\perp 2} c_7^\tau c_7^{\tau'} \right) + \frac{j_X(j_X + 1)}{12} \cdot \left[c_4^\tau c_4^{\tau'} + \frac{q^2}{m_N^2} c_9^\tau c_9^{\tau'} + \frac{v_T^{\perp 2}}{2} \left(c_{12}^\tau - \frac{q^2}{m_N^2} c_{15}^\tau \right) \cdot \left(c_{12}^{\tau'} - \frac{q^2}{m_N^2} c_{15}^{\tau'} \right) + \frac{q^2}{2m_N^2} v_T^{\perp 2} c_{14}^\tau c_{14}^{\tau'} \right] \quad (4.7)$$

$$R_{\Delta}^{\tau\tau'}\left(v_T^{\perp 2}, \frac{q^2}{m_N^2}\right) = \frac{j_X(j_X + 1)}{3} \left(\frac{q^2}{m_N^2} c_5^\tau c_5^{\tau'} + c_8^\tau c_8^{\tau'} \right). \quad (4.8)$$

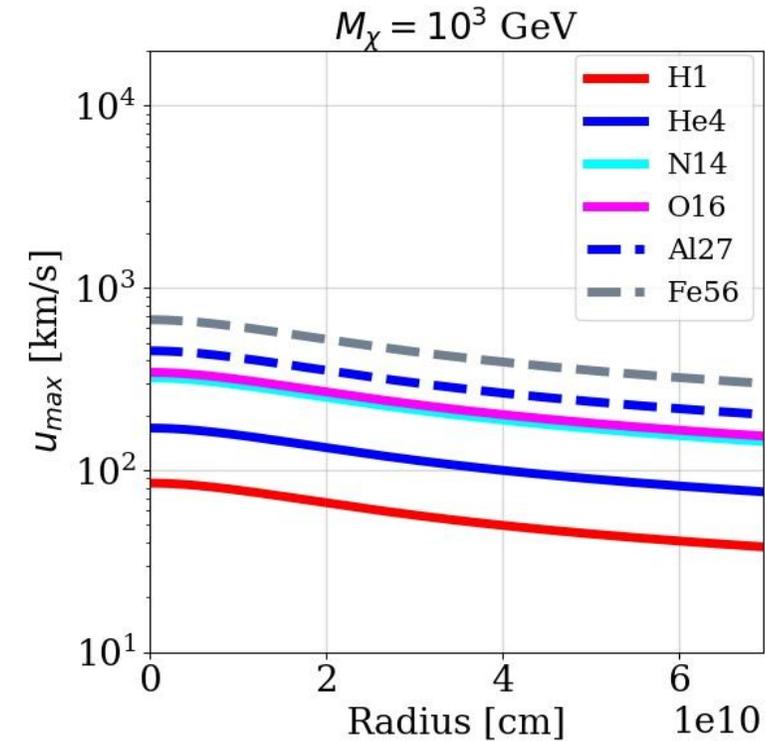
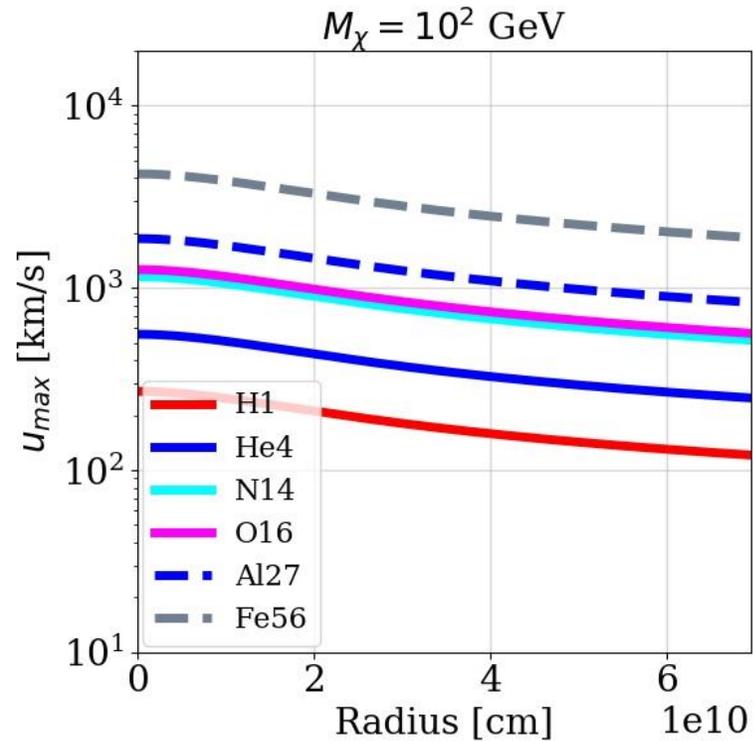
Capture Rate Calculation

$$C_{\text{cap}}^{\text{T}} = n_{\chi} \int_0^{R_{\odot}} dr 4\pi r^2 n_{\text{T}}(r) \int_0^{\infty} du 4\pi u^2 f_{\odot}(u) \frac{u^2 + v_{\odot}^{\text{esc}}(r)^2}{u} \cdot \int_{E_{\text{min}}}^{E_{\text{max}}} dE_{\text{R}} \frac{d\sigma_{\text{T}}}{dE_{\text{R}}} \theta(\Delta E)$$

$$E_{\text{min}} = \frac{1}{2} M_{\chi} u^2, \quad E_{\text{max}} = \frac{2\mu_{\text{T}}^2}{m_{\text{T}}} (u^2 + v_{\odot}^{\text{esc}}(r)^2), \quad \Delta E = E_{\text{max}} - E_{\text{min}}$$

- Integration over u has effectively an **upper bound** u_{max} where $E_{\text{min}} = E_{\text{max}}$
- u_{max} is **smaller** for **heavier** dark matter masses and **lighter** target elements

Capture Rate Calculation

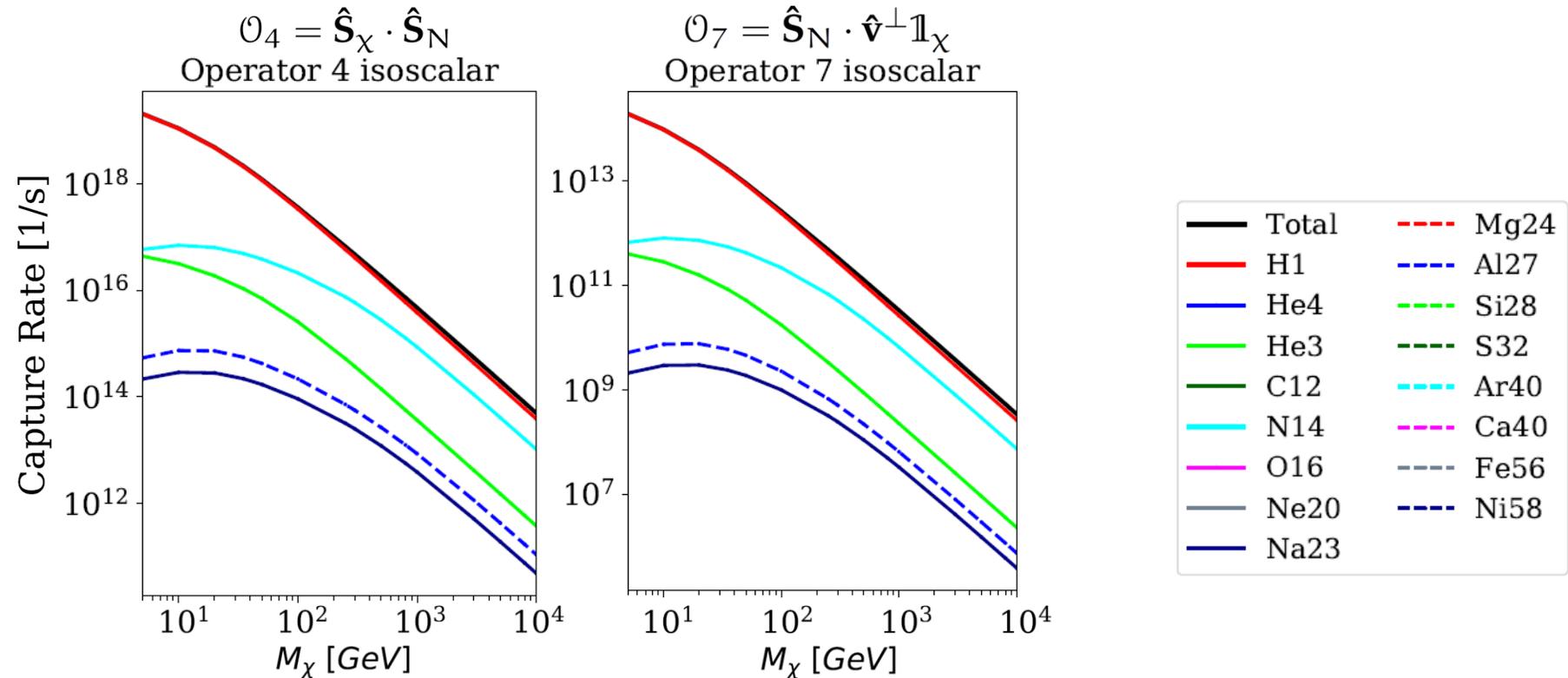


→ Integration over u has effectively an **upper bound** u_{max} where $E_{min} = E_{max}$

→ u_{max} is **smaller** for **heavier** dark matter masses and **lighter** target elements

Capture Rates

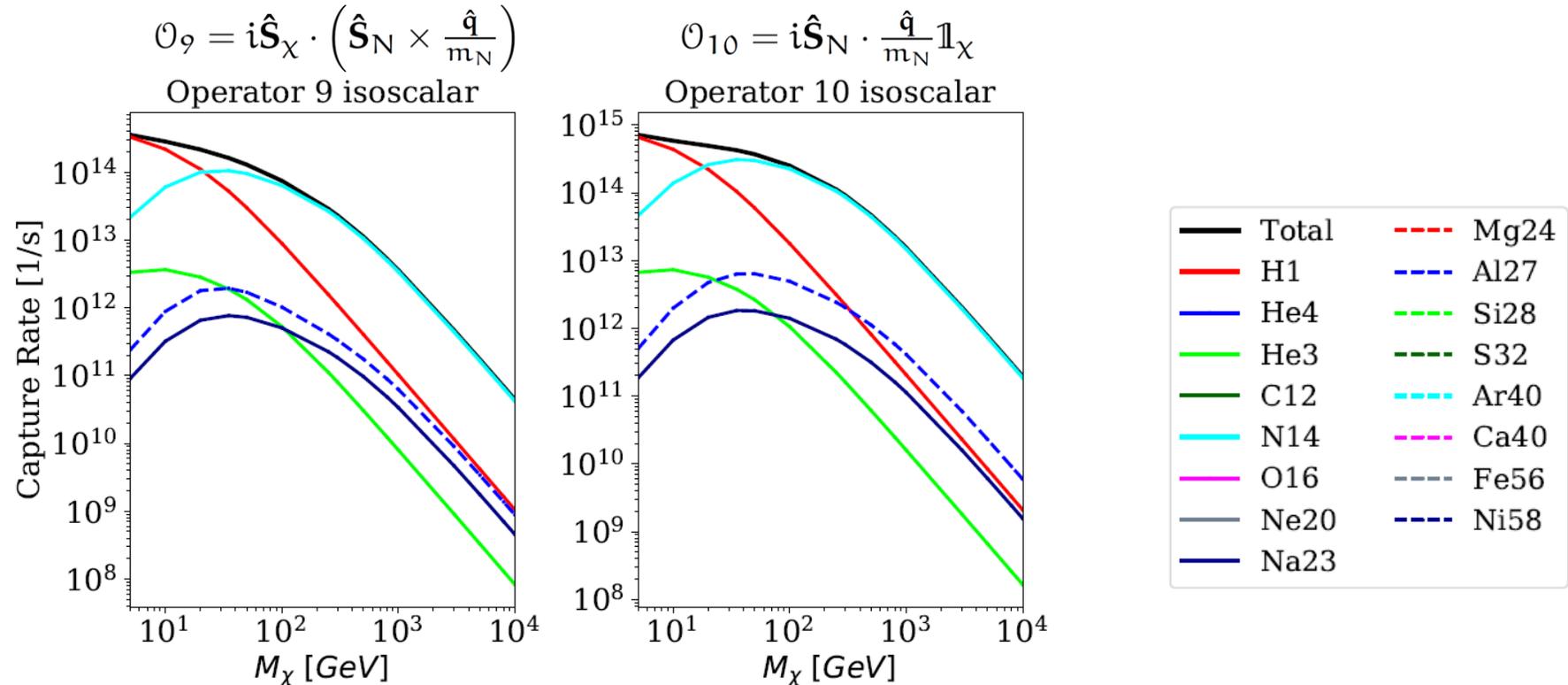
spin $\frac{1}{2}$ isoscalar



- Two **hydrogen** dominated operators
- They contain the **spin dependent** nuclear response operators Σ' and Σ''
- Capture rate generated by operator 7 is smaller due to **velocity suppression**

Capture Rates

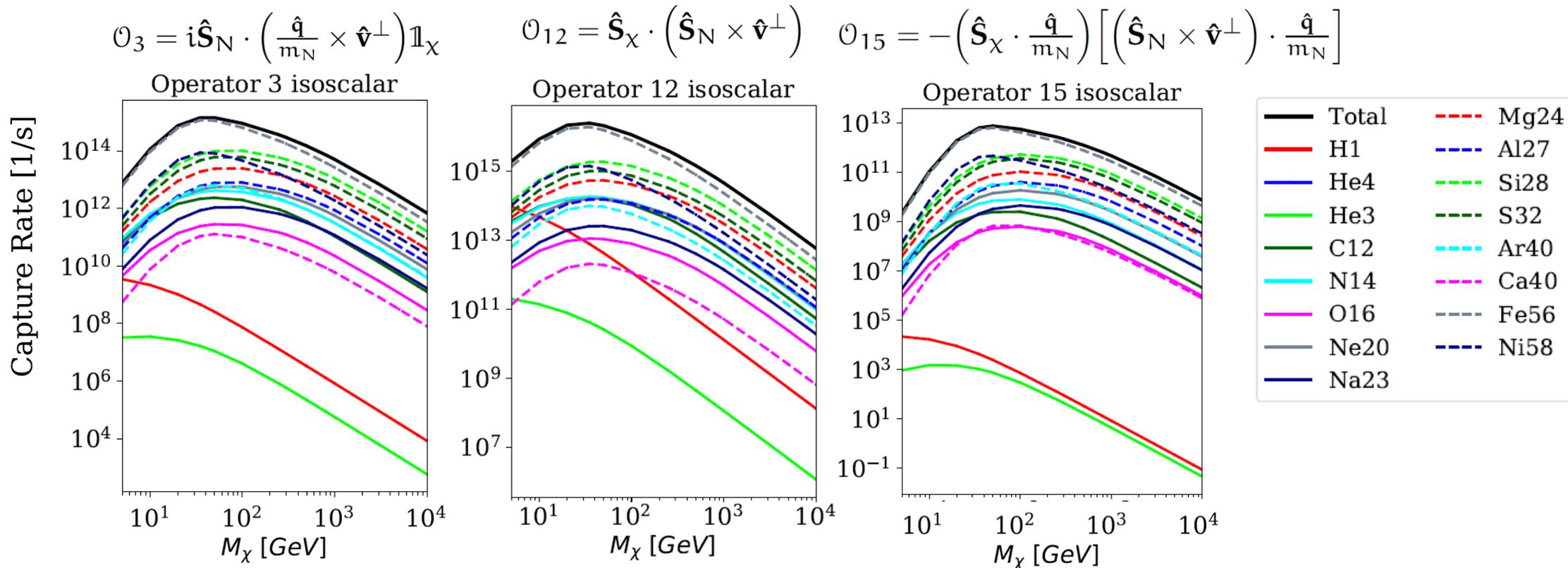
spin ½ isoscalar



- Operator 6, 9, 10, 13 and 14 also contain only Σ' and Σ''
- Depend on **momentum transfer** \rightarrow suppression of lighter elements
- In addition to the suppression of lighter elements due to a lower maximum accessible velocity u_{max} this leads to a **domination of nitrogen** at large DM masses

Capture Rates

spin ½ isoscalar



- **Iron** dominated operators
- Dominant response Φ that favors **heavy elements** with orbits of **large angular momentum**

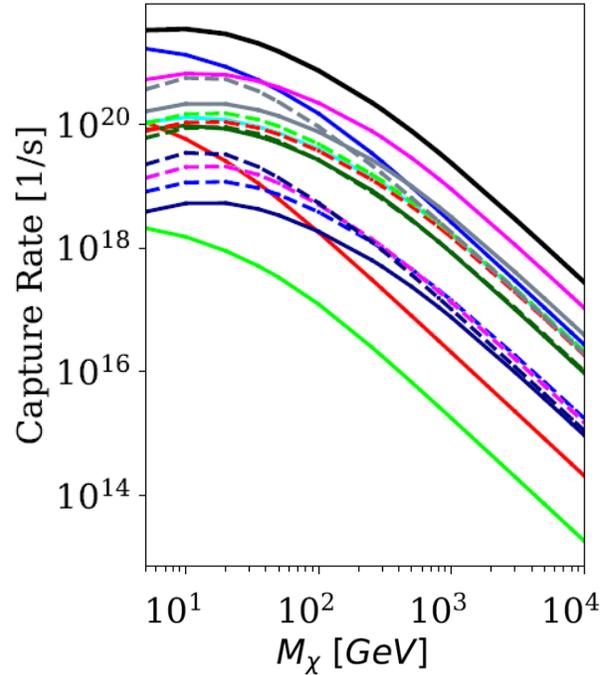
Capture Rates

spin $\frac{1}{2}$ isoscalar



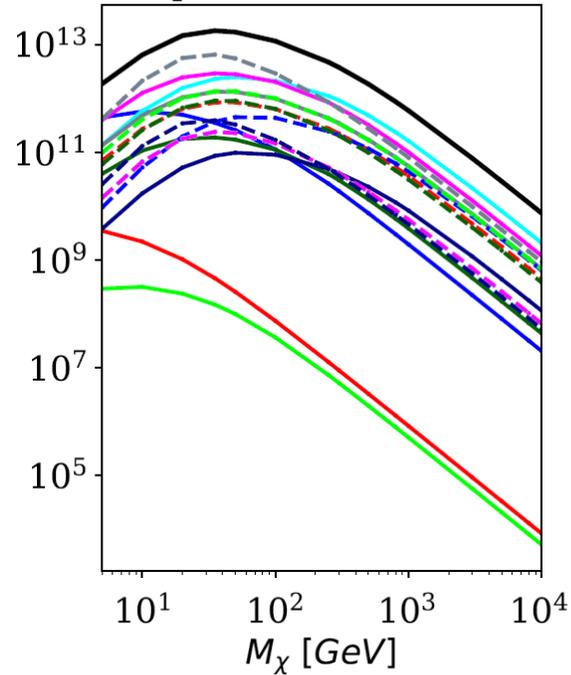
$$\mathcal{O}_1 = \mathbb{1}_{\chi N}$$

Operator 1 isoscalar



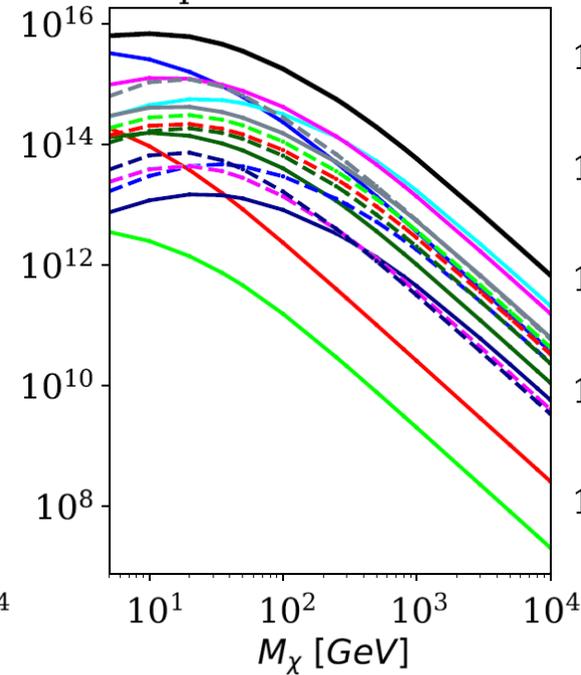
$$\mathcal{O}_5 = i\hat{\mathbf{S}}_{\chi} \cdot \left(\frac{\hat{\mathbf{q}}}{m_N} \times \hat{\mathbf{v}}^{\perp} \right) \mathbb{1}_N$$

Operator 5 isoscalar



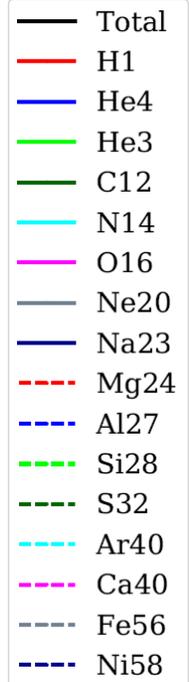
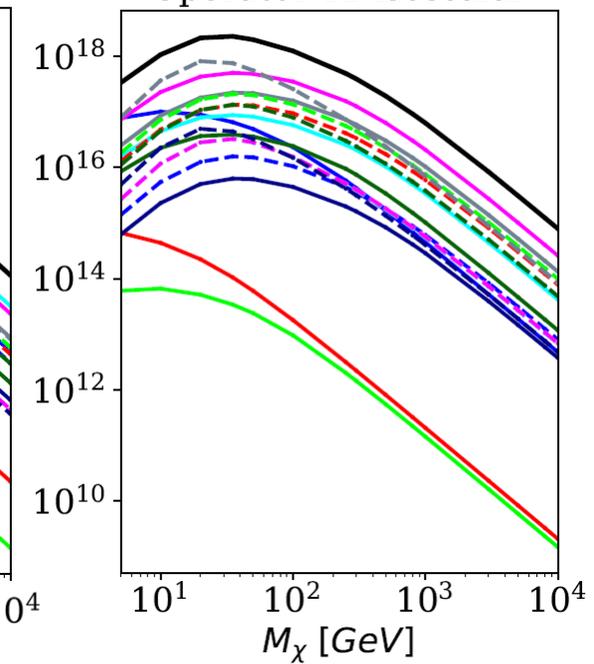
$$\mathcal{O}_8 = \hat{\mathbf{S}}_{\chi} \cdot \hat{\mathbf{v}}^{\perp} \mathbb{1}_N$$

Operator 8 isoscalar



$$\mathcal{O}_{11} = i\hat{\mathbf{S}}_{\chi} \cdot \frac{\hat{\mathbf{q}}}{m_N} \mathbb{1}_N$$

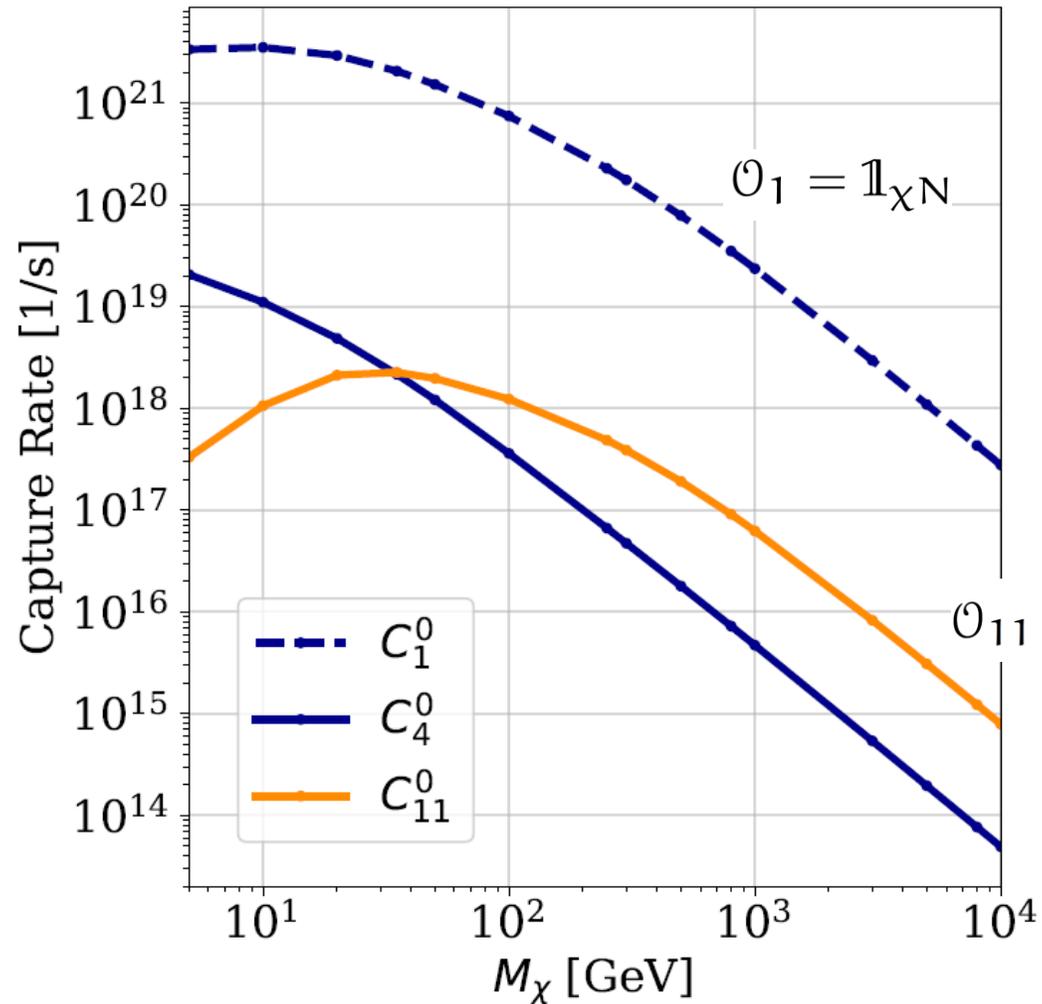
Operator 11 isoscalar



- The standard SI response \mathbf{M} favors heavy elements
- For Operator 5 and 8 additionally Δ contributes (\rightarrow Nitrogen)

Capture Rates

spin $\frac{1}{2}$ isoscalar

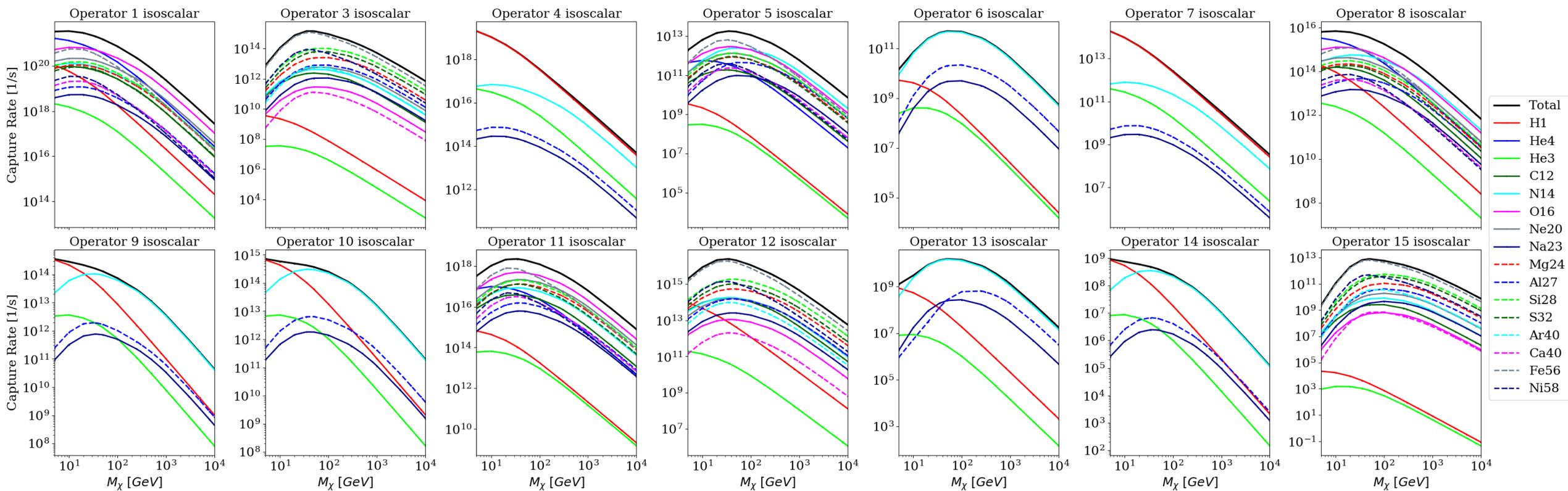


- Operator 1 and 4 are the ,standard' SI and SD interactions
- Operator 11 generates a **larger capture rate** for DM masses > 35 GeV for isoscalar couplings

→ Should be **included** in ,standard analyses'

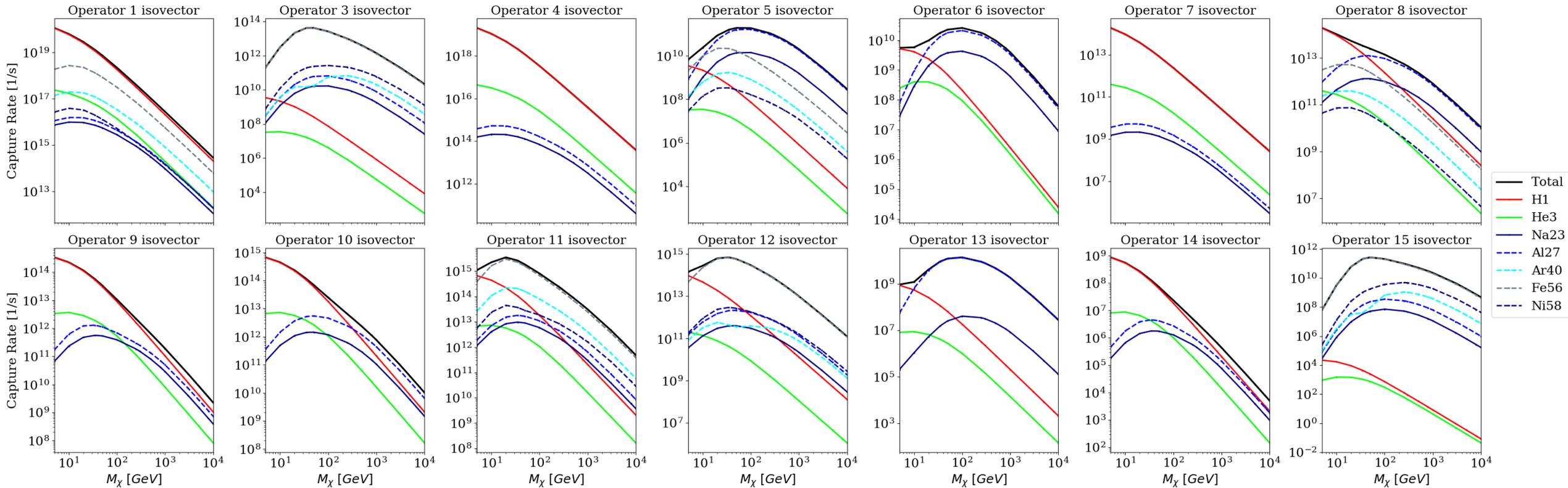
Capture Rates

spin $\frac{1}{2}$ isoscalar



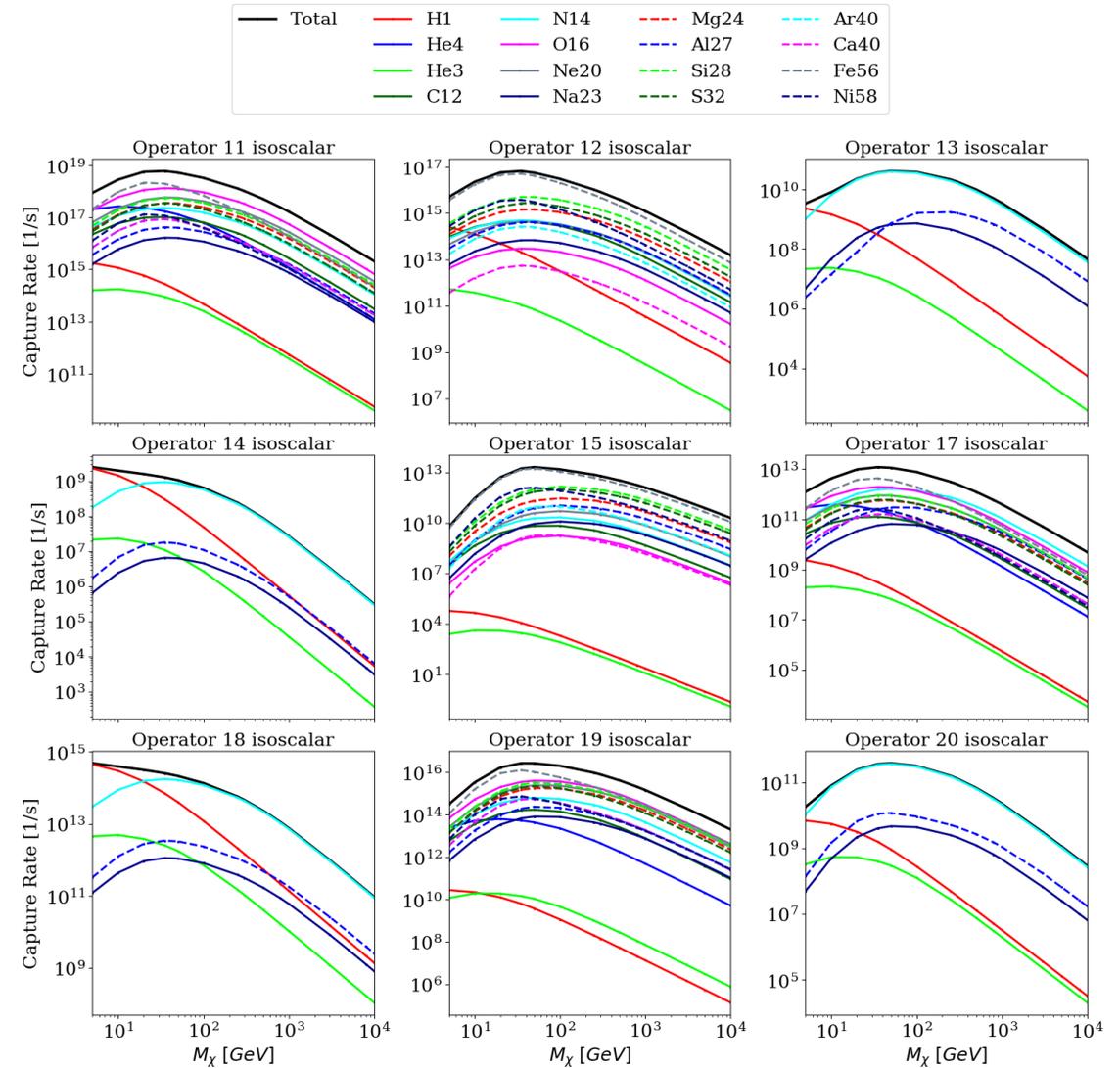
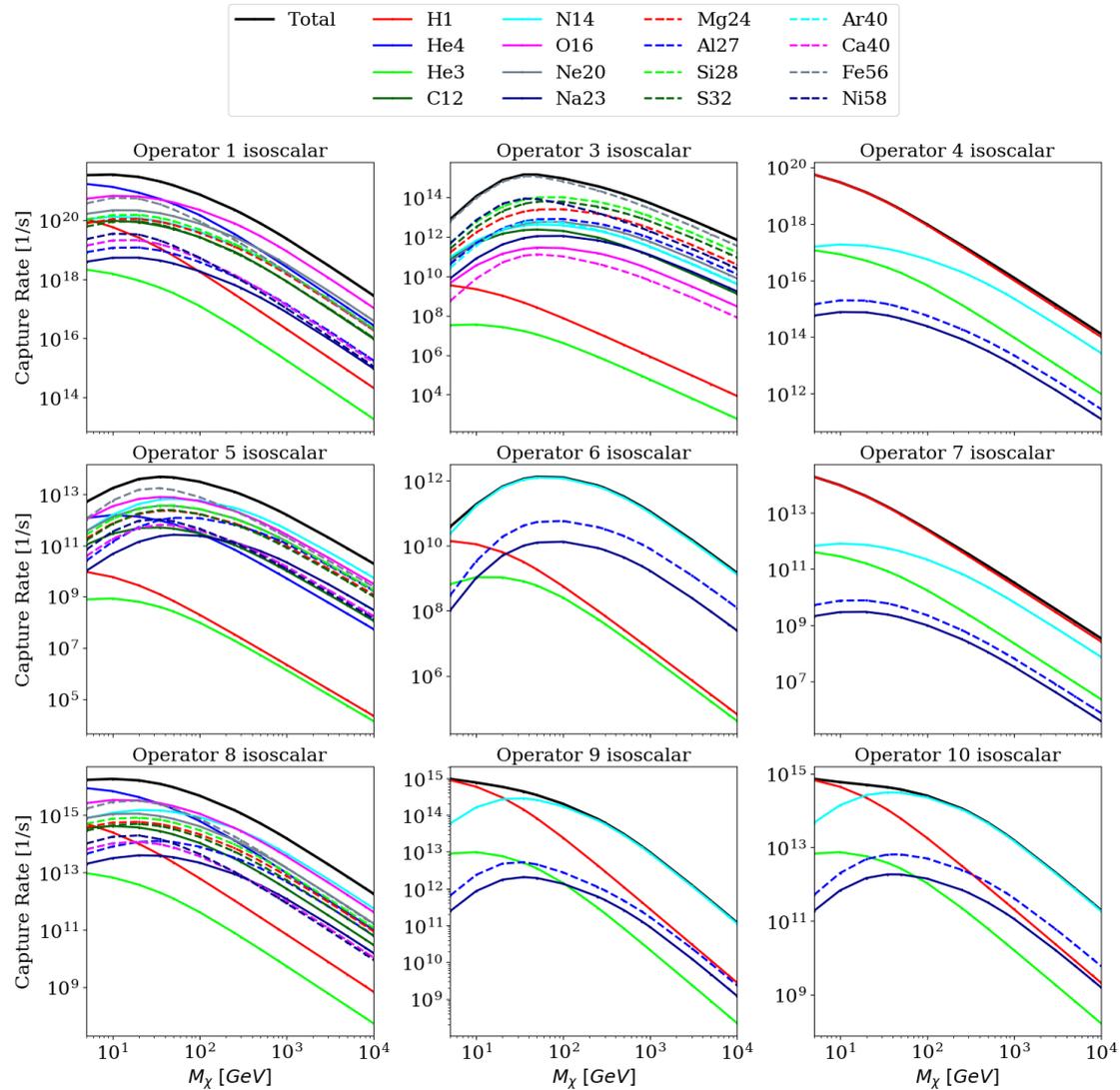
Capture Rates

spin $\frac{1}{2}$ isovector



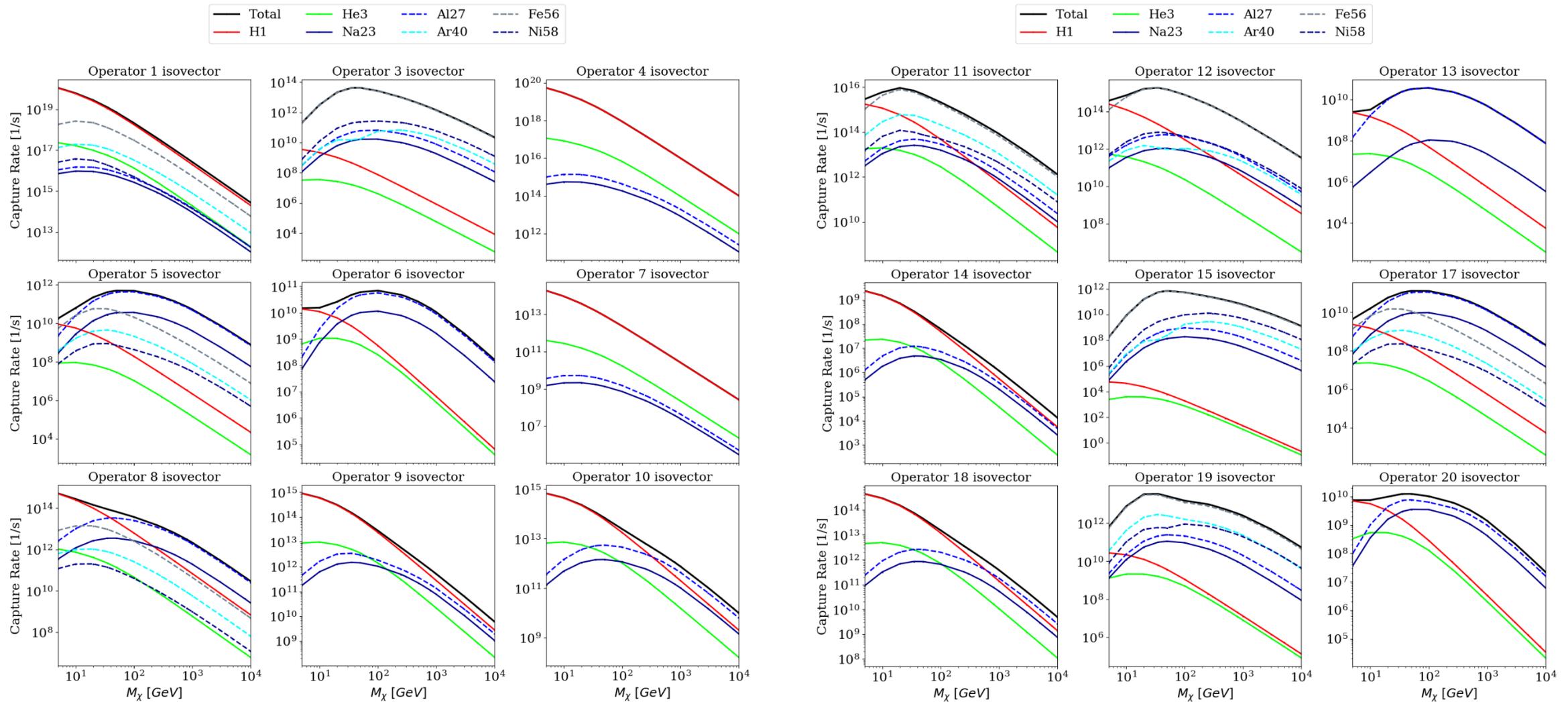
Capture Rates

spin 1 isoscalar



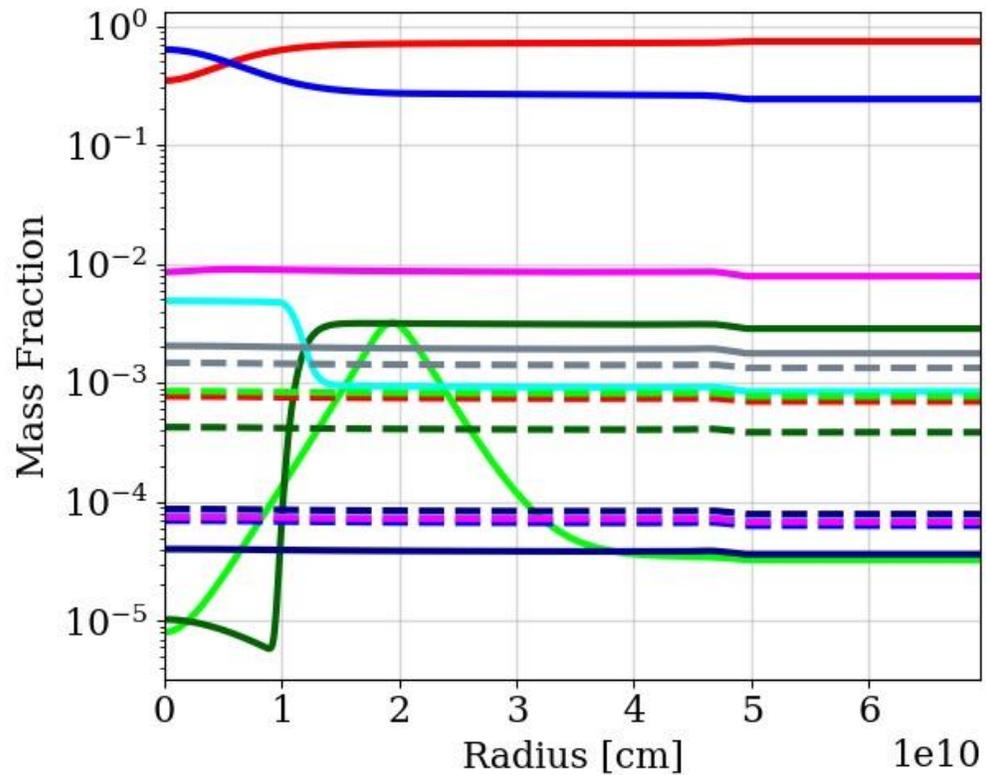
Capture Rates

spin 1 isovector

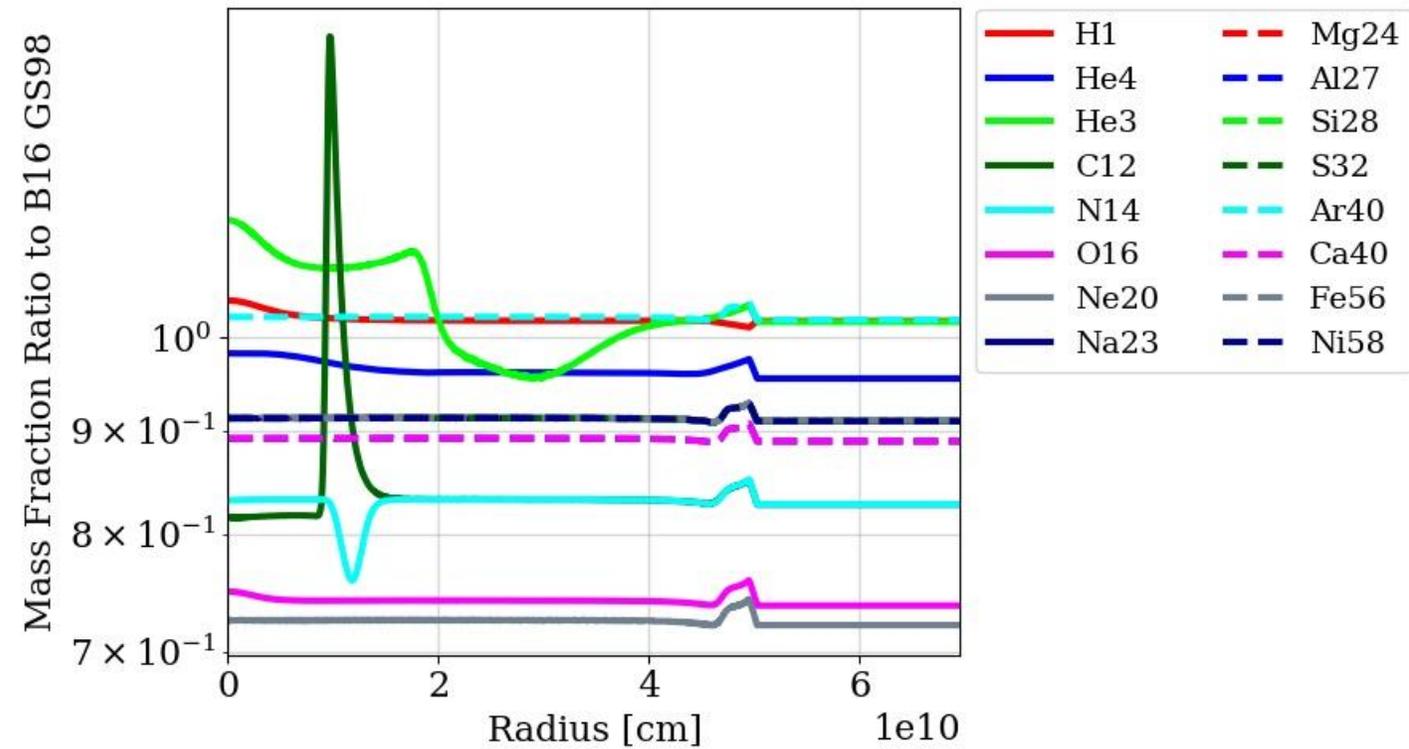


Systematics – Solar Model

B16 GS98



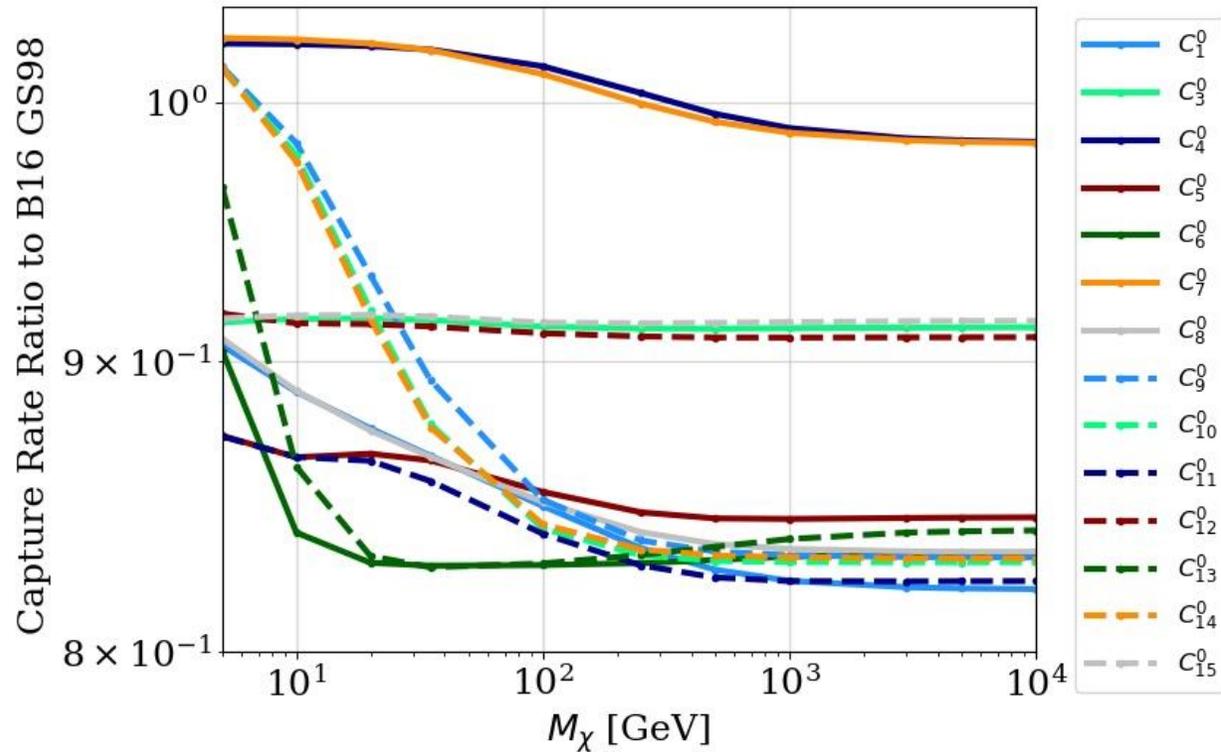
B16 AGSS09



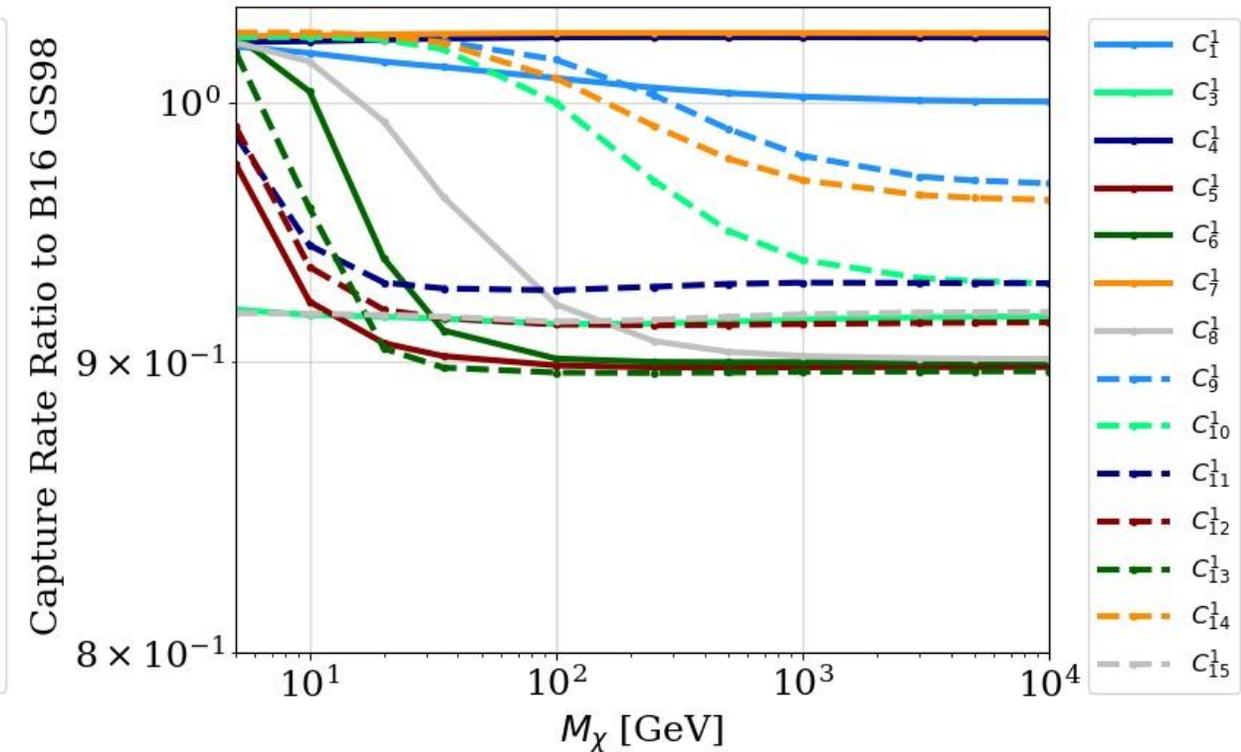
Systematics – Solar Model

spin $\frac{1}{2}$

B16 AGSS09 - isoscalar



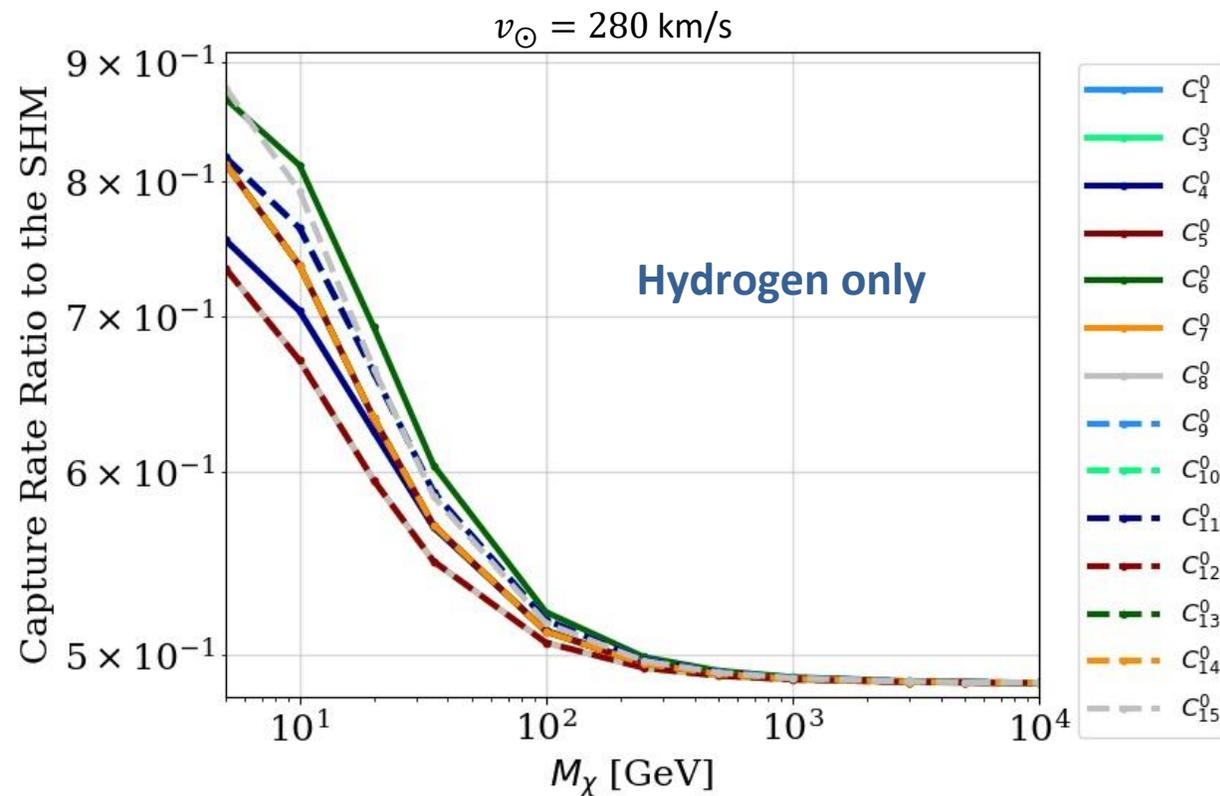
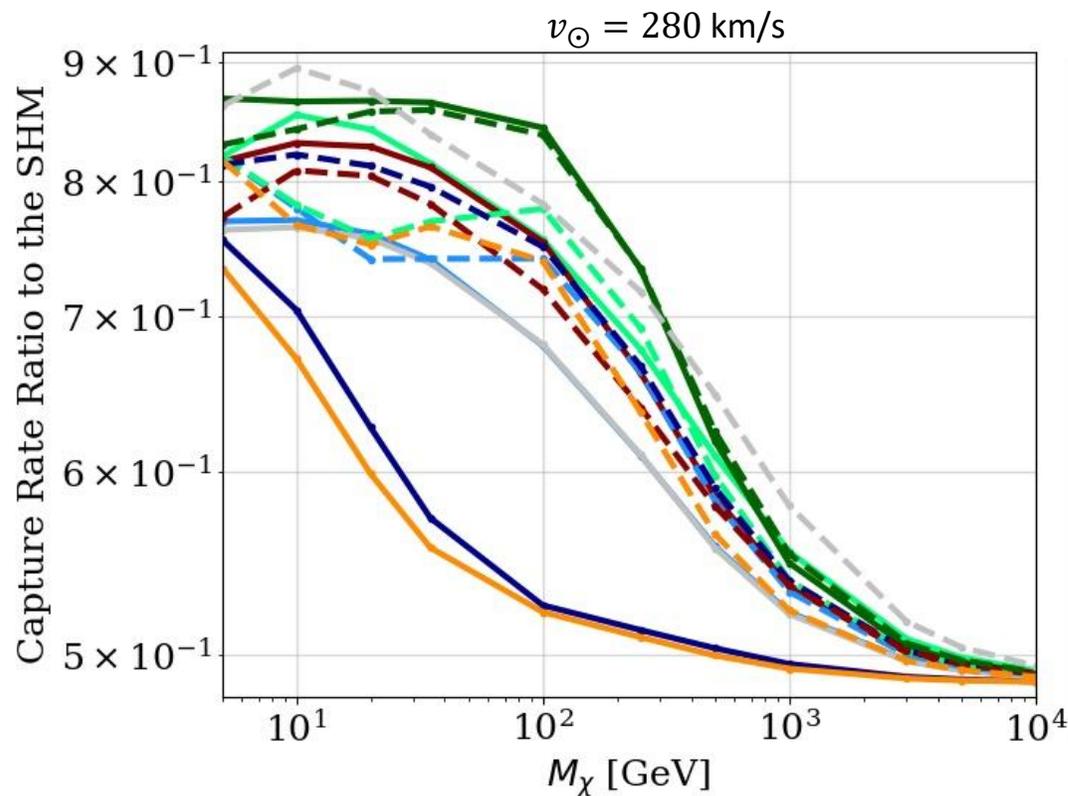
B16 AGSS09 - isovector



- Hydrogen dominated O4 and O7 least affected
- Largest effect for O16 dominated interaction types (up to 18%)
- Less deviation for isovector interactions because N14 and O16 do not contribute

Systematics – Velocity Distribution

spin $\frac{1}{2}$ isoscalar



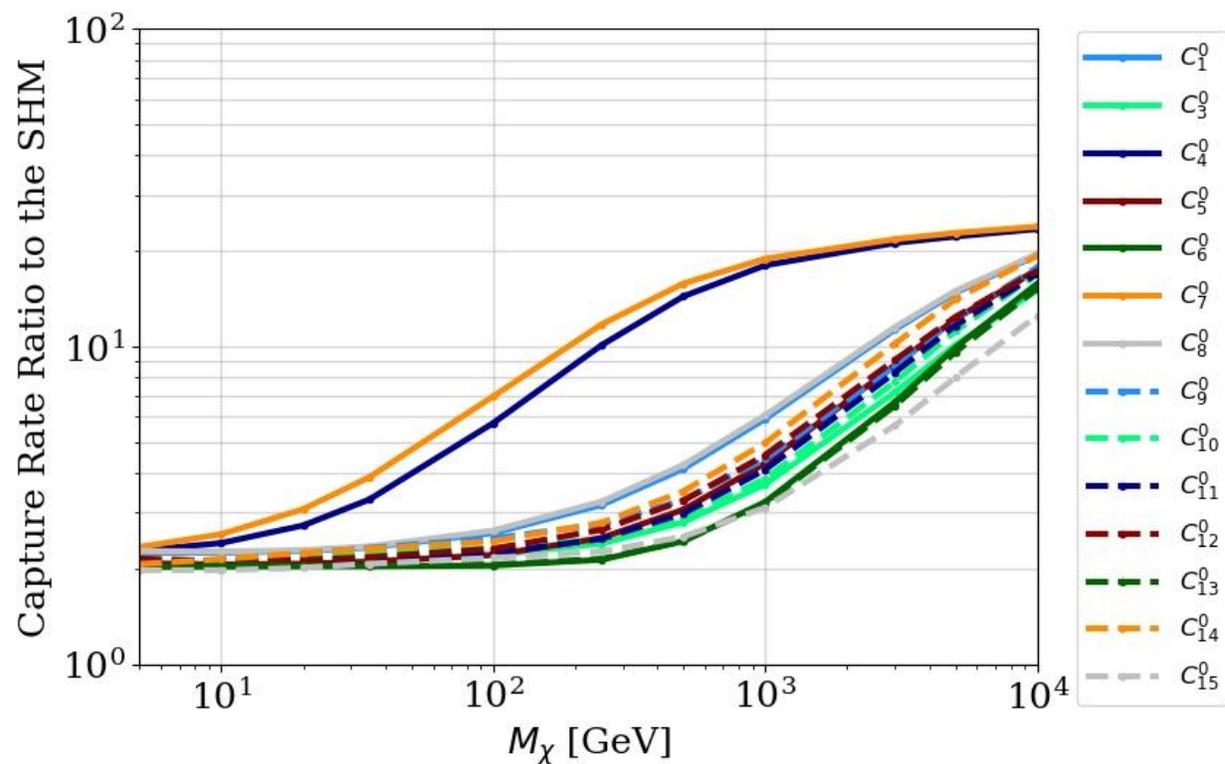
→ Velocity and momentum dependence in operators is **subdominant**

→ Capture rate is **less sensitive** to changes in the velocity distribution with **increasing** order of q and **decreasing** order of \mathbf{v}

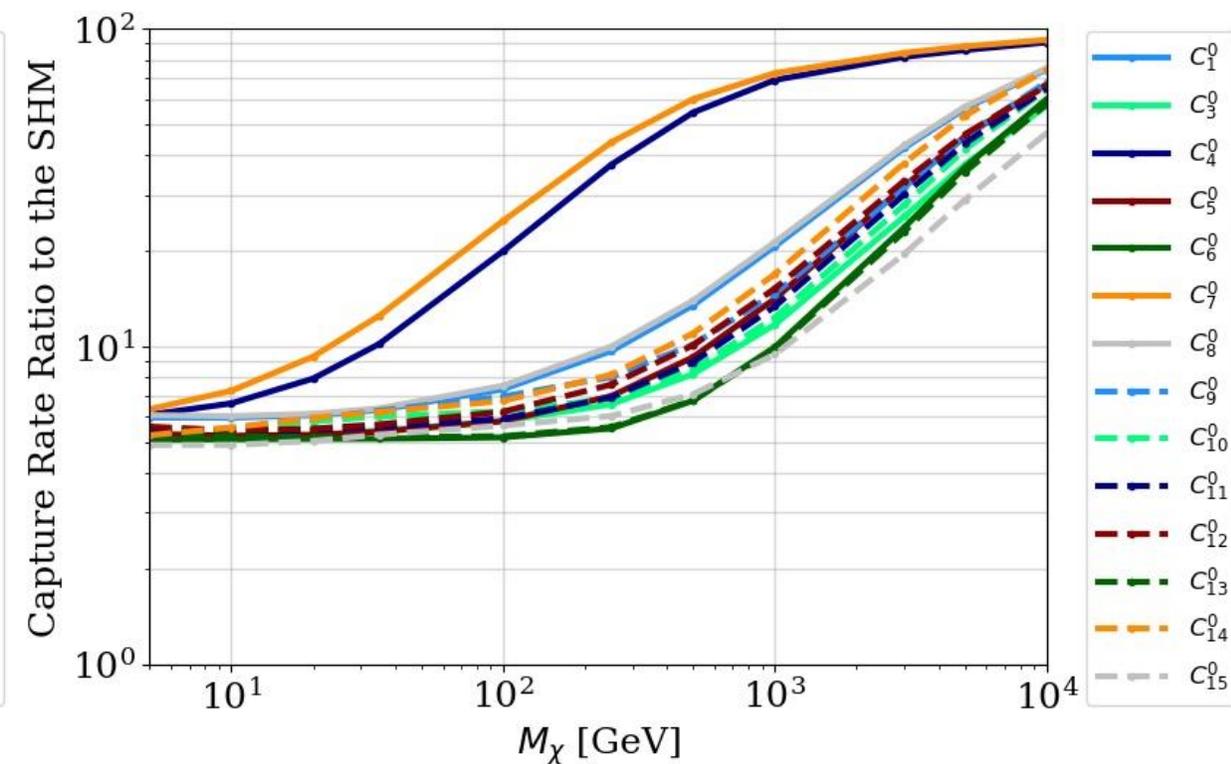
Systematics – Dark Disk

spin $\frac{1}{2}$ isoscalar

$$\rho_{\text{dd}}/\rho_{\text{h}} = 0.25$$

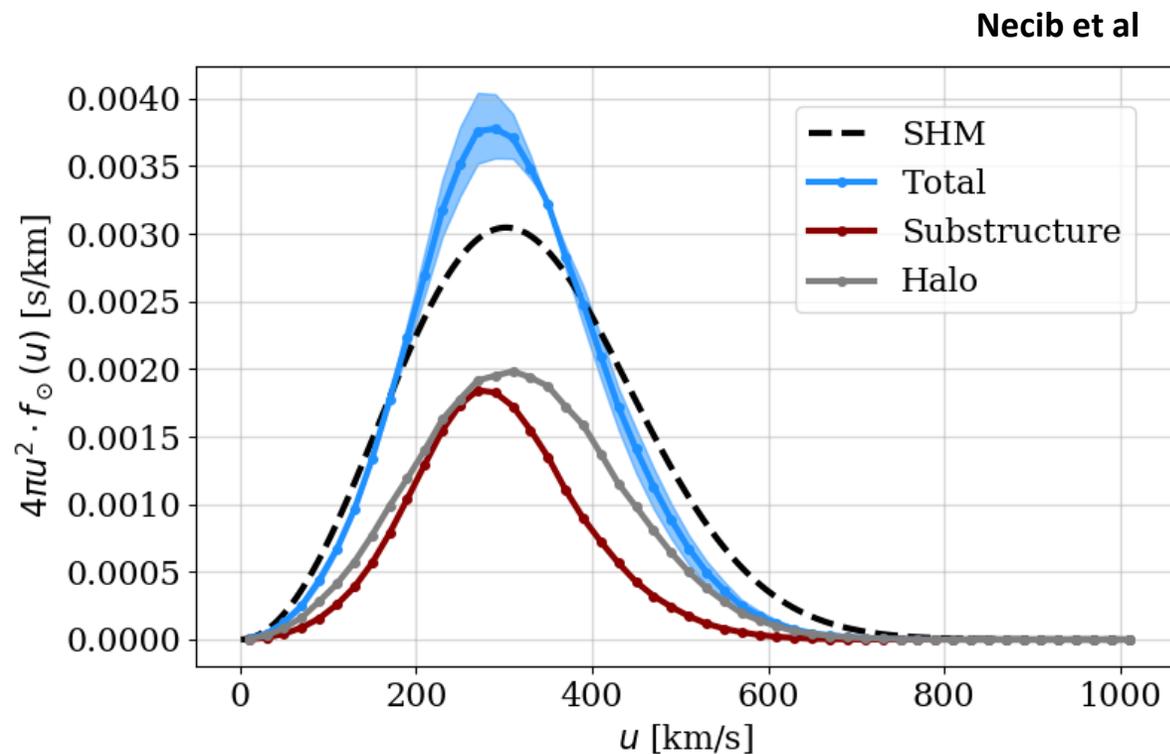


$$\rho_{\text{dd}}/\rho_{\text{h}} = 1$$

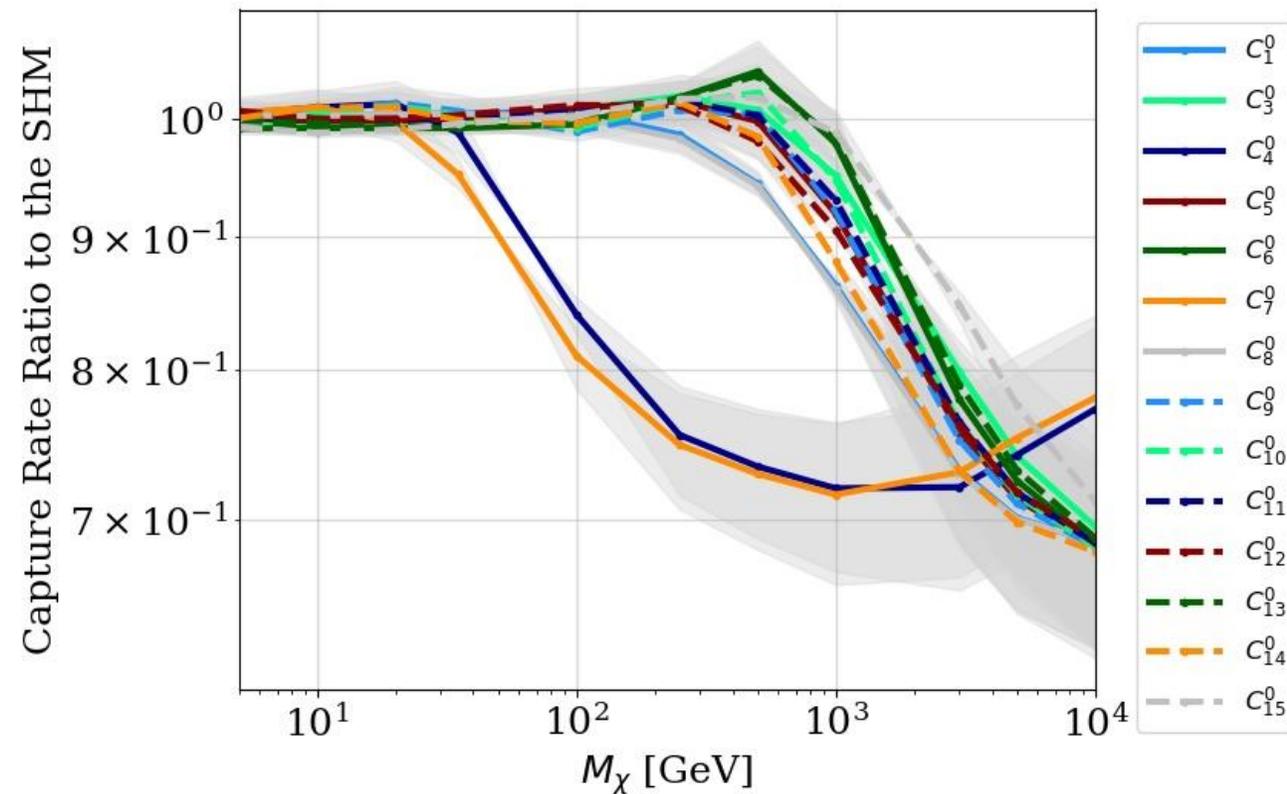


Systematics – Gaia Sausage

spin $\frac{1}{2}$ isoscalar



$$f(u) = (1 - \eta_{\text{sub}}) f_{\text{halo}}(u) + \eta_{\text{sub}} f_{\text{sub}}(u).$$



Limits

spin $\frac{1}{2}$ isovector