

NNLO+PS predictions for diboson processes in MiNNLO_{ps}

Daniele Lombardi



Workshop on Tools for High Precision LHC simulations

Castle Ringberg,

November 1st, 2022

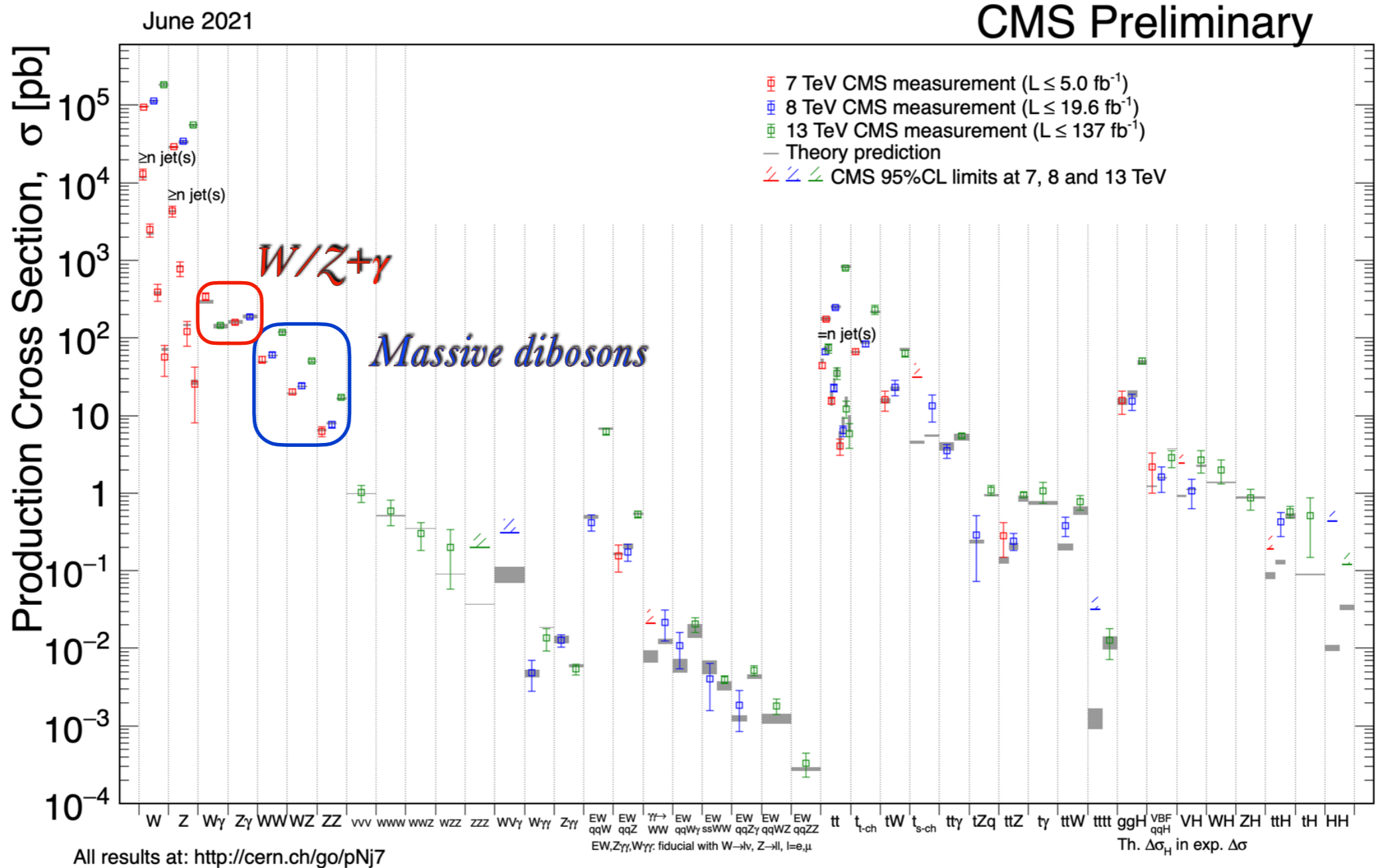
Max-Planck-Institut
für Physik



Importance of diboson processes at LHC



- ❖ They can be measured very accurately (SM candle processes)
- ❖ Background to Higgs measurements and direct new-physics searches
- ❖ Important for indirect BSM searches (anomalous couplings/EFT)

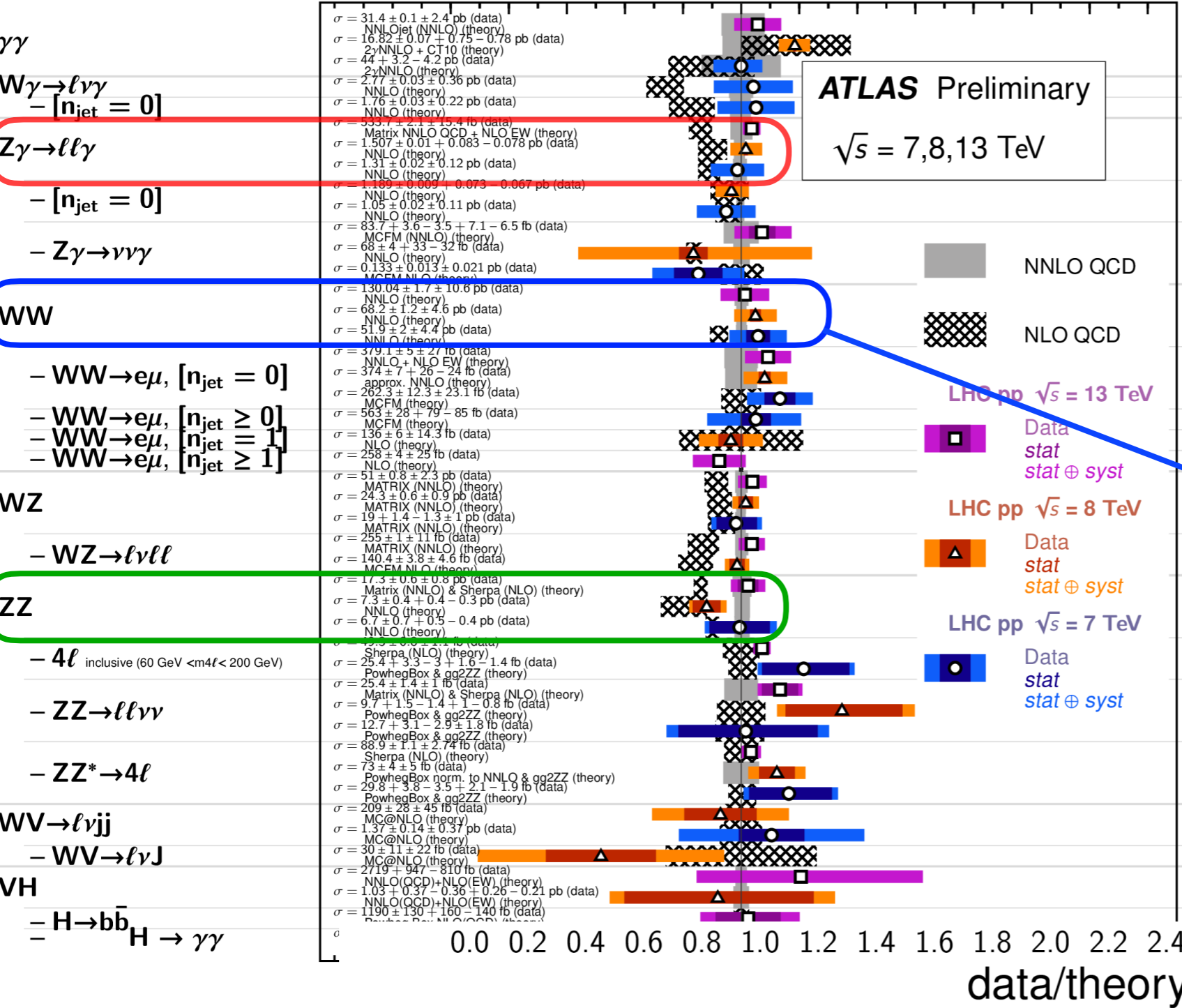


The need for high accuracy ...

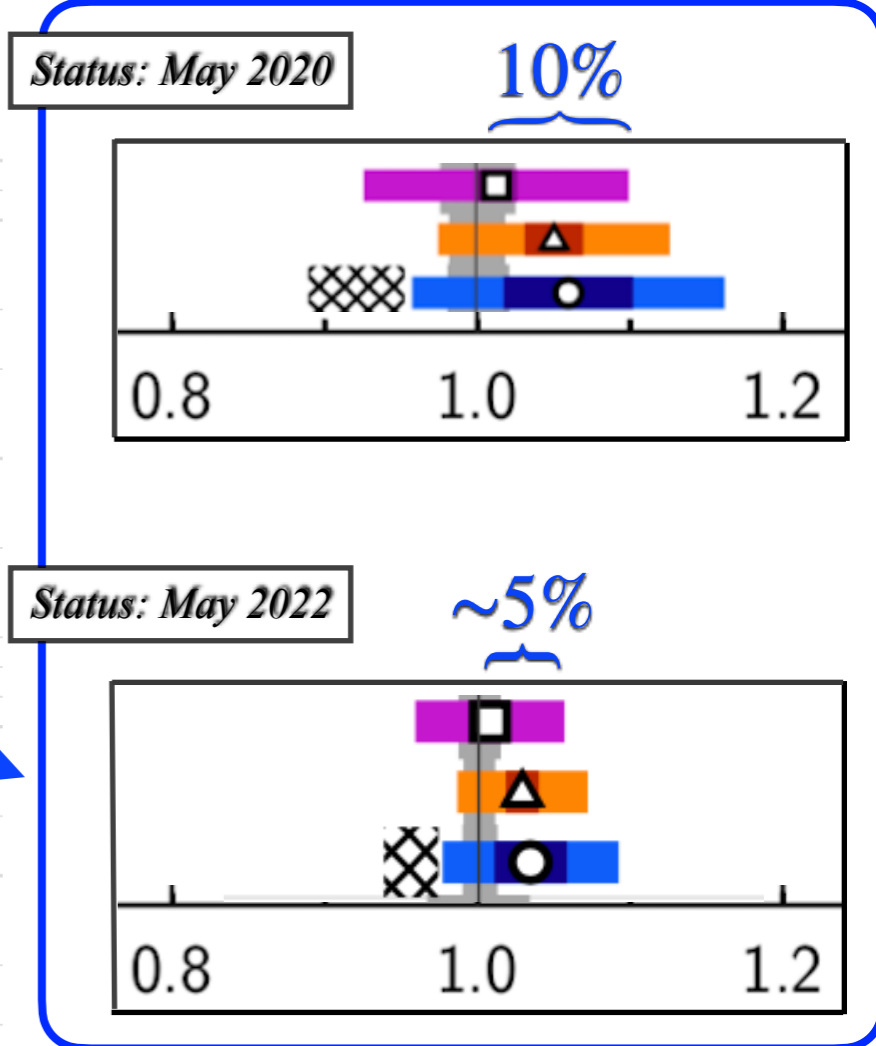
WW

Diboson Cross Section Measurements

Status: February 2022



$\int \mathcal{L} dt$
[fb⁻¹]



Experiments demand $\mathcal{O}(1\%)$ theoretical precision

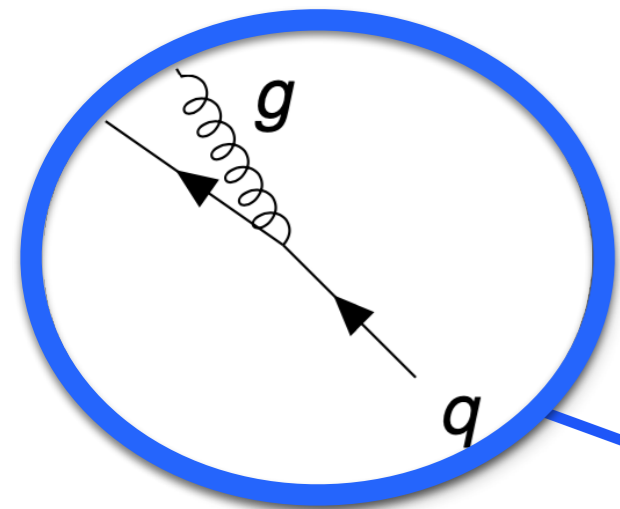
The need for MC event generators



Precise and realistic LHC phenomenology requires full-fledged event simulations.

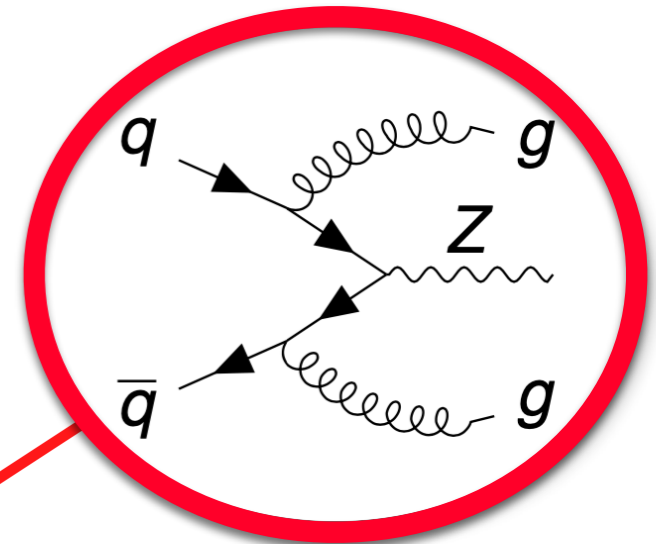
Hadronization + Decays

$$\mu \sim \Lambda_{QCD}$$



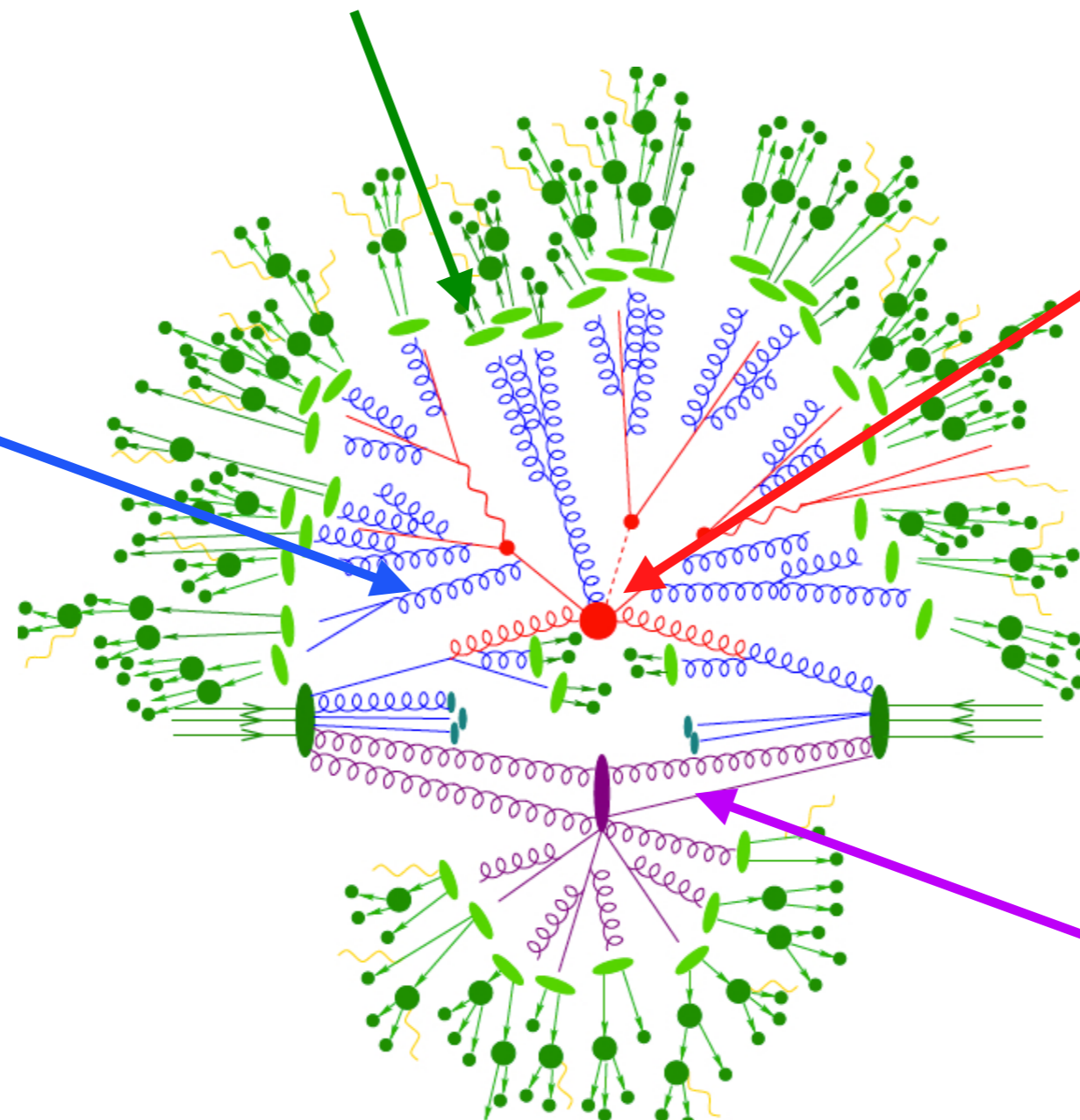
Soft/collinear radiation

$$\Lambda_{QCD} < \mu < Q$$



Hard event

$$\mu \sim Q \gg \Lambda_{QCD}$$

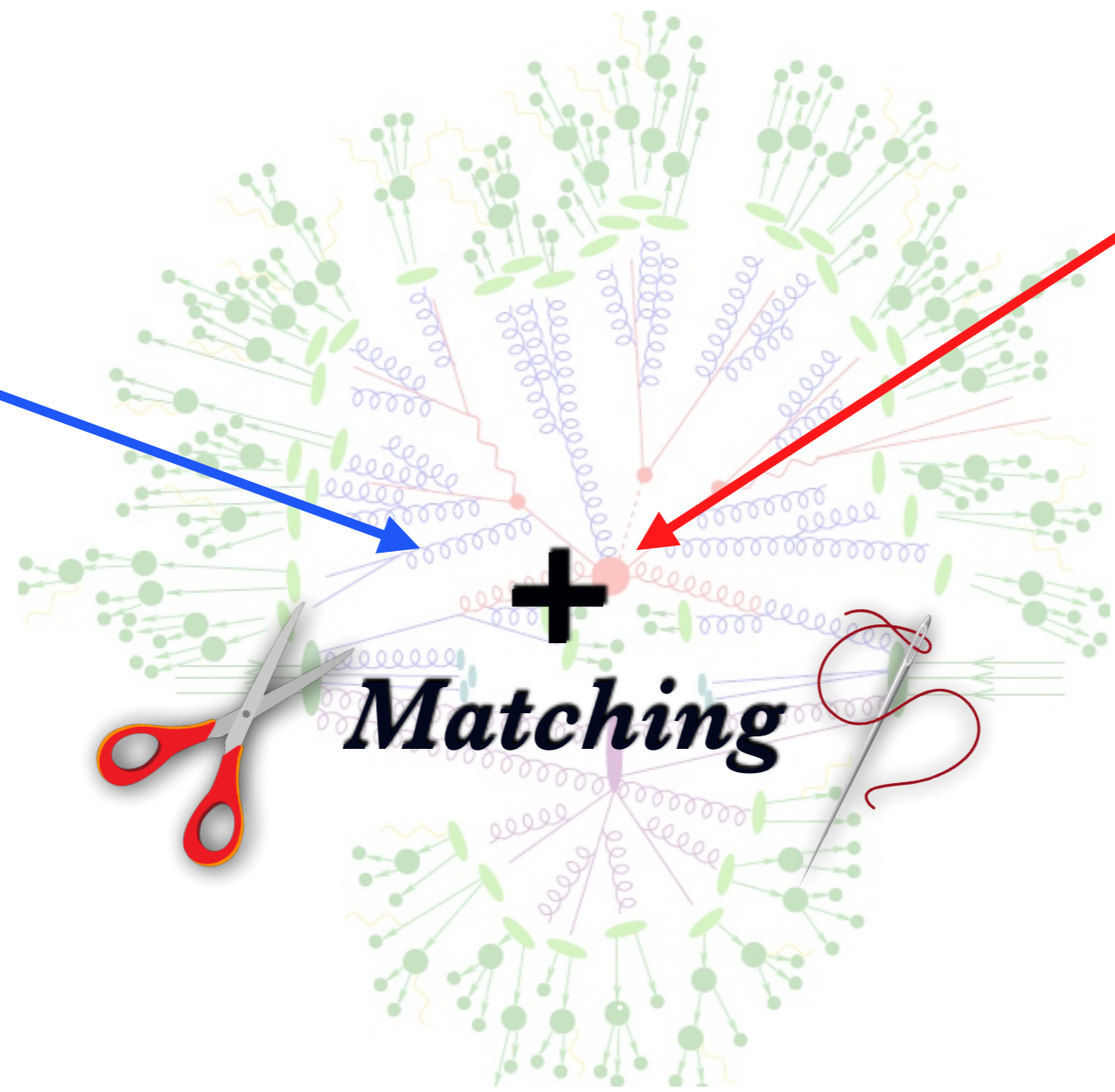
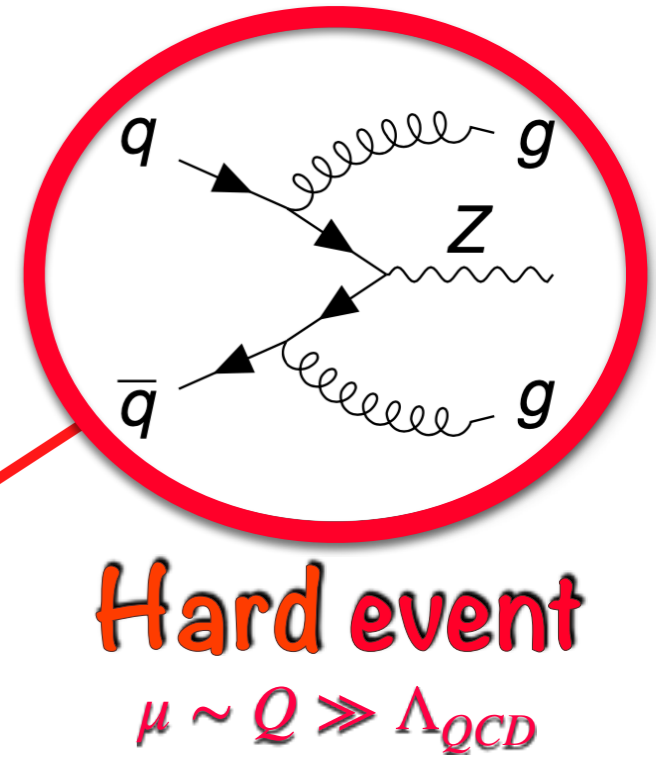
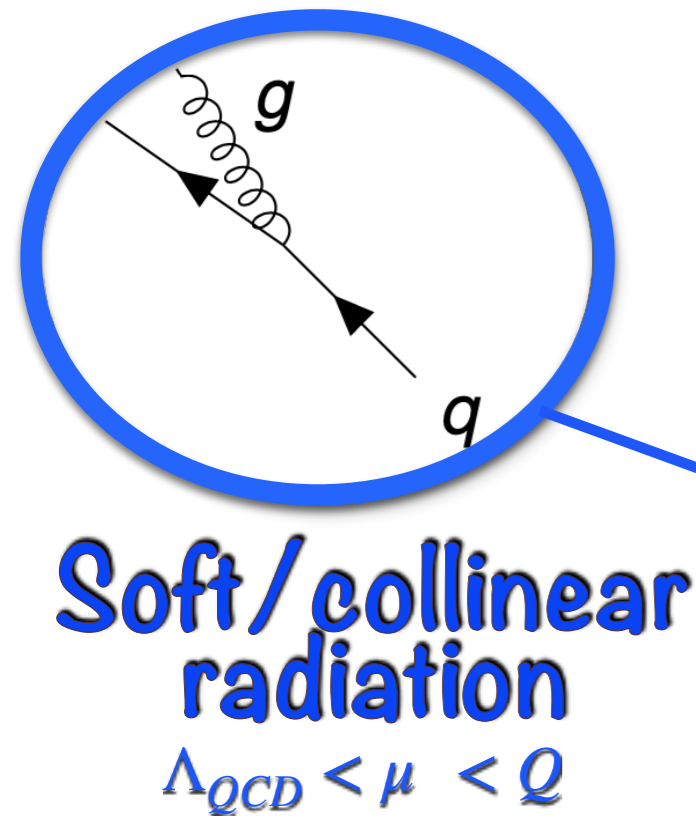


MPI and Underlying events



Fixed-order + PS accuracy

- **NLO+PS** ($\sim 10\%$): long-standing issue \rightarrow standard today
- **NNLO+PS** ($\sim 1\%$): extremely challenging \rightarrow no general approach for involved processes



NNLO+PS generator

- NNLO accuracy for observables inclusive over QCD radiation
- NLO(LO) accuracy for $VV+1(2)$ jets observables
- Resummation obtained by preserving the logarithmic accuracy of the PS

F	$F+1$ jet	$F+2$ jets	$F + \geq 3$ jets
NNLO	NLO	LO	PS

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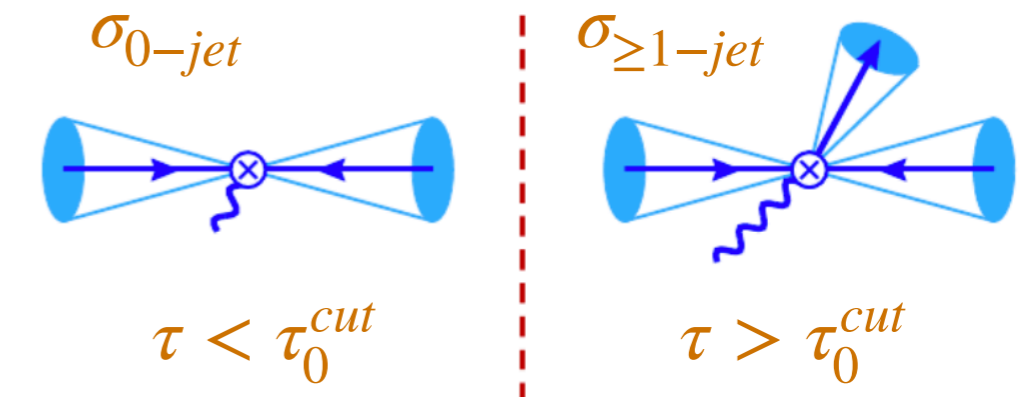
→ MiNLO'

□ Resummation obtained by preserving the logarithmic accuracy of the PS

NLO Merging:

- $FxFx$ [Frixione, Frederix '12]
- MEPS@NLO [Höche, Krauss, Schönherr, Siegert '13]
- UNLOPS [Lönblad, Prestel '12]
- MiNLO' [Hamilton, Nason, Oleari, Zanderighi '12]
- GENEVA [Alioli, Bauer, Berggren, Tackmann, Walsh, Zuberi '13], [Alioli, Bauer, Berggren, Tackmann, Walsh '15]
- ...

Inclusive 0-jet cross section



NNLO+PS generator

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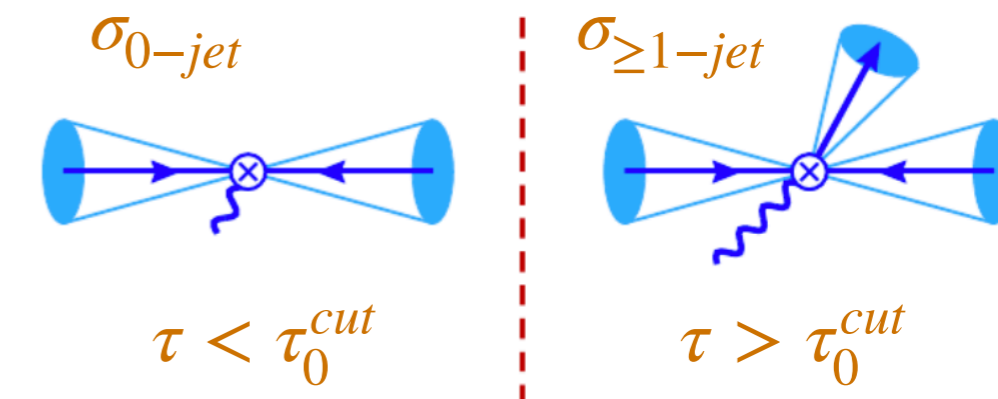
☑ Resummation obtained by preserving the logarithmic accuracy of the PS

→ POWHEG

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- ...

Inclusive 0-jet cross section



NLO Matching:

- MC@NLO [Frixione, Weber '02]
- POWHEG [Nason '04]
- KrkNLO [Jadach, Płaczek, Sapeta, Siódmok, Skrzypek '15]
- MAcNLOPS [Nason, Salam '21]

$$S(t_I) = \Delta(t_I, t_0) \langle \mathbb{I} \rangle + \sum_{l,k=0}^{\infty} \int_{t_l}^{z_l, t_l} \dots \int_{z_k, t_k}^{z, t} \dots \int_{z_1, t_1}^{\tilde{z}_1, \tilde{t}_1} \dots \int_{z_k, \tilde{t}_k}^{t_0}$$

NNLO+PS generator

NNLO accuracy for observables inclusive over QCD radiation

NLO(LO) accuracy for $VV+1(2)$ jets observables

→ MiNLO'

Resummation obtained by preserving the logarithmic accuracy of the PS

→ POWHEG



Issue: interface NNLO calculations to PS by preserving the accuracy of the latter

Solution: different methods available (for colour singlets)

- MiNLO'+reweighting [Hamilton, Nason, Oleari, Zanderighi '12, +Re '13], [Karlberg, Hamilton, Zanderighi '14]
- GENEVA [Alioli, Bauer, Berggren, Tackmann, Walsh, Zuberi '13], [Alioli, Bauer, Berggren, Tackmann, Walsh '15]
- UNNLOPS [Höche, Prestel '14]
- VINCIA sector shower [Campbell, Höche, Li, Preuss, Skands '21]

NNLO+PS generator

NNLO accuracy for observables inclusive over QCD radiation

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→ MiNNLO_{PS}:

- ▶ $2 \rightarrow 1$ [Monni, Nason, Re, Wiesemann, Zanderighi '19], [Monni, Re, Wiesemann '20]
- ▶ $2 \rightarrow 2$ extension [Lombardi, Wiesemann, Zanderighi '20]
- ▶ $\#$ [Mazzitelli, Monni, Nason, Re, Wiesemann, Zanderighi '21, '22]

NNLO accuracy obtained by connecting MiNLO' with transverse-momentum resummation formulae.

$$\underbrace{\frac{d\sigma_F^{\text{res}}}{dp_T d\Phi_B}}_{\text{}} = \frac{d}{dp_T} \{ e^{-S} \mathcal{L} \} = e^{-S} \underbrace{\{ \mathcal{L}' - S' \mathcal{L} \}}_{\equiv D} \quad \rightarrow \mathcal{L} \sim H(C \otimes f) \underbrace{(C \otimes f)}_{\text{luminosity}}$$

$$\bar{B}(\Phi_{FJ}) = d\sigma_F^{\text{res}} + [d\sigma_{FJ}]_{\text{f.o.}} - [d\sigma_F^{\text{res}}]_{\text{f.o.}} = e^{-S} \left\{ D + \underbrace{\frac{[d\sigma_{FJ}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}}}_{1-S^{(1)}\dots} \underbrace{\frac{[d\sigma_F^{\text{res}}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}}}_{-D^{(1)}-D^{(2)}\dots} \right\} \rightarrow \mu_R \sim \mu_F \sim p_T$$

resummation scheme

F : generic colour-singlet system

Expand up to $\mathcal{O}(\alpha_s^3)$: $\int \frac{dp_T}{p_T} e^{-S} \alpha_s^n(p_T) \log^m \frac{Q}{p_T} \sim \mathcal{O}(\alpha_s^{n-\frac{m+1}{2}}(Q))$

$$\bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) \sim e^{-S} \left\{ \underbrace{d\sigma_{FJ}^{(1)}}_{\sim \alpha_s(p_T)} \underbrace{(1 + S^{(1)}) + d\sigma_{FJ}^{(2),\text{sing}}}_{\sim \alpha_s^2(p_T)} + \underbrace{(D - D^{(1)} - D^{(2)})}_{\sim \alpha_s^3(p_T)} + \text{regular} \right\}$$

NNLO accuracy obtained by connecting MiNLO' with transverse-momentum resummation formulae.

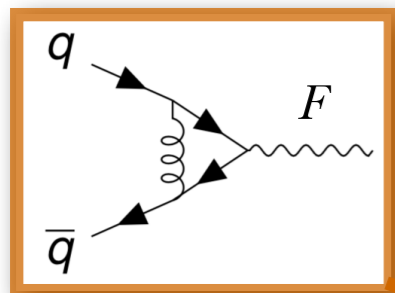
$$\frac{d\sigma_F^{\text{res}}}{dp_T d\Phi_B} = \frac{d}{dp_T} \left\{ e^{-S} \mathcal{L} \right\} = e^{-S} \underbrace{\left\{ \mathcal{L}' - S' \mathcal{L} \right\}}_{\equiv D} \rightarrow \mathcal{L} \sim H(C \otimes f)(C \otimes f)$$

luminosity

$$\bar{B}(\Phi_{FJ}) = d\sigma_F^{\text{res}} + [d\sigma_{FJ}]_{\text{f.o.}} - [d\sigma_F^{\text{res}}]_{\text{f.o.}} = e^{-S} \left\{ D + \underbrace{\frac{[d\sigma_{FJ}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}}}_{1-S^{(1)}\dots} \underbrace{\frac{[d\sigma_F^{\text{res}}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}}}_{-D^{(1)}-D^{(2)}\dots} \right\} \rightarrow \mu_R \sim \mu_F \sim p_T$$

resummation scheme

F: generic colour-singlet system



Include NLO corrections for unresolved radiation

MiNLO'

$$\bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) \sim e^{-S} \left\{ \underbrace{d\sigma_{FJ}^{(1)}}_{\sim \alpha_s(p_T)} \underbrace{(1 + S^{(1)})}_{\sim \alpha_s^2(p_T)} + d\sigma_{FJ}^{(2),\text{sing}} + \underbrace{(D - D^{(1)} - D^{(2)})}_{\sim \alpha_s^3(p_T)} + \text{regular} \right\}$$

Preserve original *F*+1jet accuracy

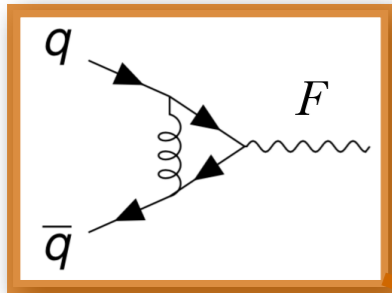
NNLO accuracy obtained by connecting MiNLO' with transverse-momentum resummation formulae.

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resummation scheme

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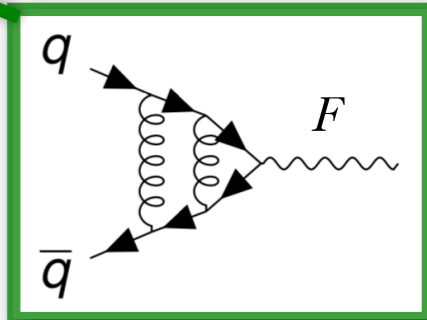
MiNNLO_{PS}

Beyond accuracy

$$\bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) \sim e^{-S} \left\{ \underbrace{d\sigma_{FJ}^{(1)}}_{\sim \alpha_s(p_T)} \underbrace{(1 + S^{(1)})}_{\sim \alpha_s^2(p_T)} + d\sigma_{FJ}^{(2),\text{sing}} + \underbrace{(D - D^{(1)} - D^{(2)})}_{\sim \alpha_s^3(p_T)} + \text{regular} \right\}$$

Preserve original $F+1$ jet accuracy

Include missing NNLO ingredients



MiNNLO_{PS}: 2→2 extension

[Lombardi, Wiesemann, Zanderighi '20]

NNLO accuracy obtained by connecting MiNLO' with transverse-momentum resummation formulae.

$$\frac{d\sigma_F^{\text{res}}}{dp_T d\Phi_B} = \frac{d}{dp_T} \{ e^{-S} \mathcal{L} \} = e^{-S} \underbrace{\{ \mathcal{L}' - S' \mathcal{L} \}}_{\equiv D} \rightarrow \mathcal{L} \sim \underbrace{H}_{\text{luminosity}} (C \otimes f) (C \otimes f)$$

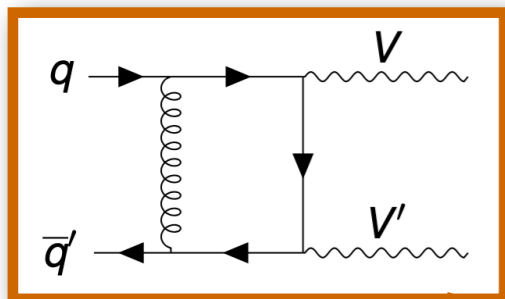
Hard function

$H^{(1)}$ and $H^{(2)}$
generally
flavour and
kinematic
dependent

$$\bar{B}(\Phi_{FJ}) = d\sigma_F^{\text{res}} + [d\sigma_{FJ}]_{\text{f.o.}} - [d\sigma_F^{\text{res}}]_{\text{f.o.}} = e^{-S} \left\{ D + \frac{[d\sigma_{FJ}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}} - \frac{[d\sigma_F^{\text{res}}]_{\text{f.o.}}}{[e^{-S}]_{\text{f.o.}}} \right\}$$

$\underbrace{\hspace{10em}}_{1-S^{(1)}\dots} \quad \underbrace{\hspace{10em}}_{-D^{(1)}-D^{(2)}\dots}$

F : generic colour-singlet system



Include NLO corrections
for unresolved radiation

MiNLO'

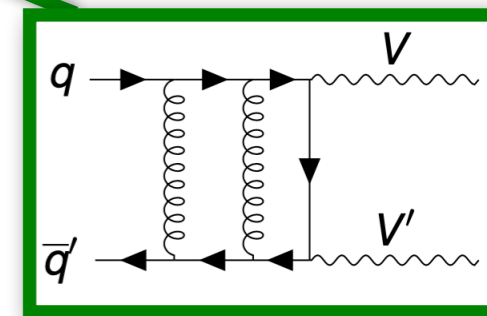
MiNNLO_{PS}

Beyond
accuracy

$$\bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) \sim e^{-S} \left\{ \underbrace{d\sigma_{FJ}^{(1)}}_{\sim \alpha_s(p_T)} (1 + S^{(1)}) + \underbrace{d\sigma_{FJ}^{(2),\text{sing}}}_{\sim \alpha_s^2(p_T)} + \underbrace{(D - D^{(1)} - D^{(2)})}_{\sim \alpha_s^3(p_T)} + \text{regular} \right\}$$

Preserve original
 $F+1$ jet accuracy

Include missing
NNLO ingredients



MiNNLO_{PS}: matching to Parton shower

- **Merging** without unphysical merging scale

$$\bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) \sim e^{-S} \left\{ d\sigma_{FJ}^{(1)}(1 + S^{(1)}) + d\sigma_{FJ}^{(2),\text{sing}} + (D - D^{(1)} - D^{(2)})(p_T) \cdot F^{\text{corr}}(\Phi_{FJ}) + \text{regular} \right\}$$

NNLO corrections are:

❖ spread over the $F+1$ jet phase space



$$F^{\text{corr}}(\Phi_{FJ}) = \frac{J(\Phi_{FJ})}{\int d\Phi'_{FJ} J(\Phi'_{FJ}) \delta(p_T - p'_T) \delta(\Phi_F - \Phi'_F)}$$

❖ switched off at high transverse momentum $p_T \geq Q$



- **Matching** (NNLO+PS accuracy) obtained using POWHEG method

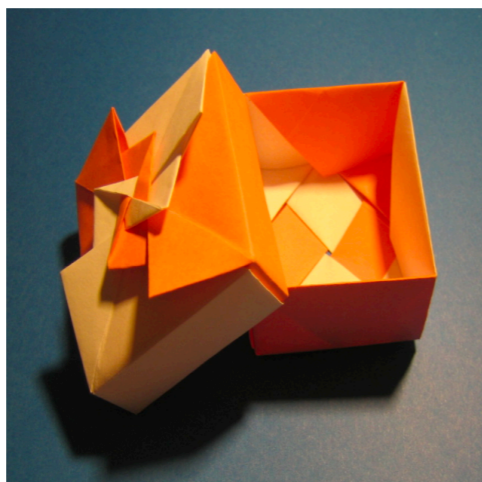
$$d\sigma_{\text{MiNNLO}_{\text{PS}}} = \bar{B}_{\text{MiNNLO}_{\text{PS}}}(\Phi_{FJ}) d\Phi_{FJ} \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{QCD}}) + \Delta_{\text{pwg}}(p_{T,\text{rad}}) \frac{R_s(\Phi_{FJJ})}{B(\Phi_{FJ})} d\Phi_{\text{rad}} \right\}$$

NNLO+PS generators for Dibosons

The POWHEG BOX

Project

The POWHEG BOX is a general computer framework for implementing NLO calculations in shower Monte Carlo programs according to the POWHEG method. It is also a library, where previously included processes are made available to the users. It can be interfaced with all modern shower Monte Carlo programs that support the Les Houches Interface for User Generated Processes.



Available Processes

- Higgs boson production in gluon fusion, P. Monni, P. Nason, E. Re, M. Wiesemann and G. Zanderighi, *JHEP* **05** (2020) 143, arXiv:1908.06987 [\[paper\]](#), P. Monni, E. Re and M. Wiesemann, *Eur. Phys. J. C* **80** (2020), no.11, 1075, arXiv:2006.04133 [\[paper\]](#)

POWHEG-BOX-V2/HJ (inside HJMiNNLO subfolder)

- Next-to-Next-to-Leading Order Event Generation for Top-Quark Pair Production, J. Mazzitelli, P. Monni, P. Nason, E. Re, M. Wiesemann and G. Zanderighi, *Phys.Rev.Lett.* **127** (2021) 062001, arXiv:2012.14267 [\[paper\]](#), Top-pair production at the LHC with MiNNLO_{PS}, J. Mazzitelli, P. Monni, P. Nason, E. Re, M. Wiesemann and G. Zanderighi, arXiv:2112.XXXXX [\[paper\]](#)

POWHEG-BOX-V2/ttJ_MiNNLO

- Advancing MiNNLO_{PS} to diboson processes: $Z\gamma$ production at NNLO+PS, D. Lombardi, M. Wiesemann and G. Zanderighi, *JHEP* **2106** (2021) 095, arXiv:2010.10478 [\[paper\]](#), Anomalous couplings in $Z\gamma$ events at NNLO+PS and improving $\nu\nu\gamma$ backgrounds in dark-matter searches, D. Lombardi, M. Wiesemann and G. Zanderighi, *Phys.Lett. B* **XX** (2021) XXX, arXiv:2108.11315 [\[paper\]](#)

POWHEG-BOX-RES/ZgamJ

- W^+W^- production at NNLO+PS with MiNNLO_{PS}, D. Lombardi, M. Wiesemann and G. Zanderighi, *JHEP* **2111** (2021) 230, 2103.12077 [\[paper\]](#)

POWHEG-BOX-RES/WWJ

- ZZ production at nNNLO+PS with MiNNLO_{PS}, L. Buonocore, G. Koole, D. Lombardi, L. Rottoli, M. Wiesemann, G. Zanderighi, arXiv:2108.05337 [\[paper\]](#)

POWHEG-BOX-RES/ZZJ

MiNNLO_{PS} has been extended to $2 \rightarrow 2$ colour-singlet processes (built in POWHEG-BOX-RES).



First implementation of $Z\gamma$ generator (both $Z \rightarrow \ell^+\ell^-$ and $Z \rightarrow \bar{\nu}\nu + aTGC @NNLO$)

[Lombardi, Wiesemann, Zanderighi '20, '21]



New approach to the existing WW generator [Lombardi, Wiesemann, Zanderighi '21]



ZZ generator with incoherent combination of $\bar{q}q$ and gg channels [Buonocore, Koole, Lombardi, Rottoli, Wiesemann, Zanderighi '21]



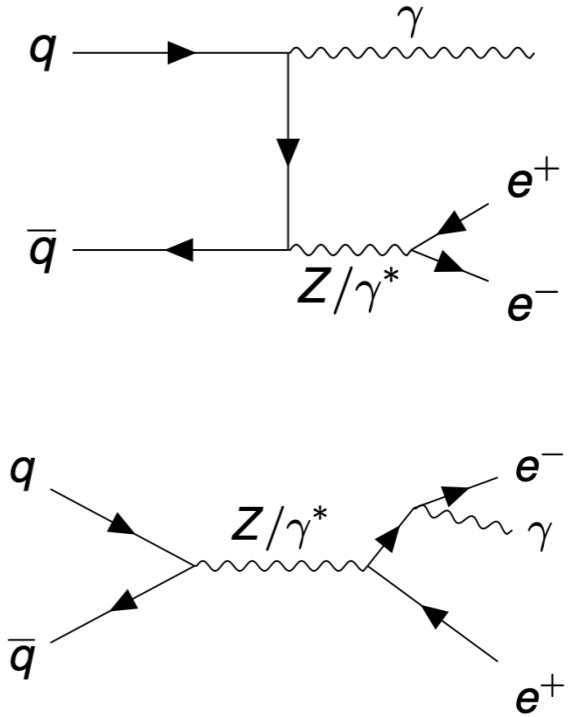
$\gamma\gamma$ generator [Gavardi, Oleari, Re '22]



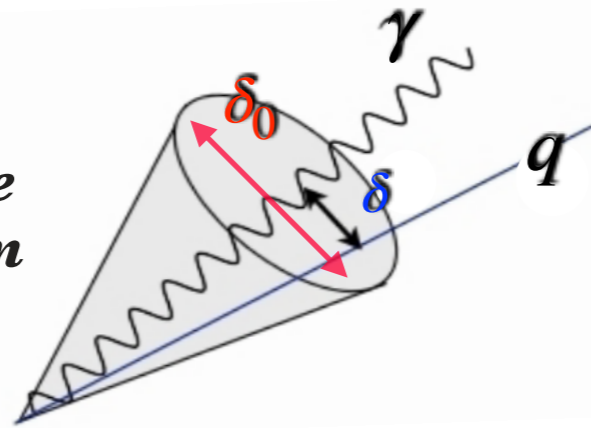
WZ generator including NLO EW corrections and QED PS effects [Lindert, Lombardi, Wiesemann, Zanderighi, Zanolini '22]

Zγ phenomenology at the LHC [Lombardi, Wiesemann, Zanderighi '20]

- ❖ relatively large cross section
- ❖ $Z \rightarrow e^+e^-$ event can be fully reconstructed
- ❖ presence of **isolated photon** \rightarrow theoretically challenging
- ❖ highly relevant as a probe for BSM (especially $Z \rightarrow \nu\bar{\nu}$)



Frixione isolation



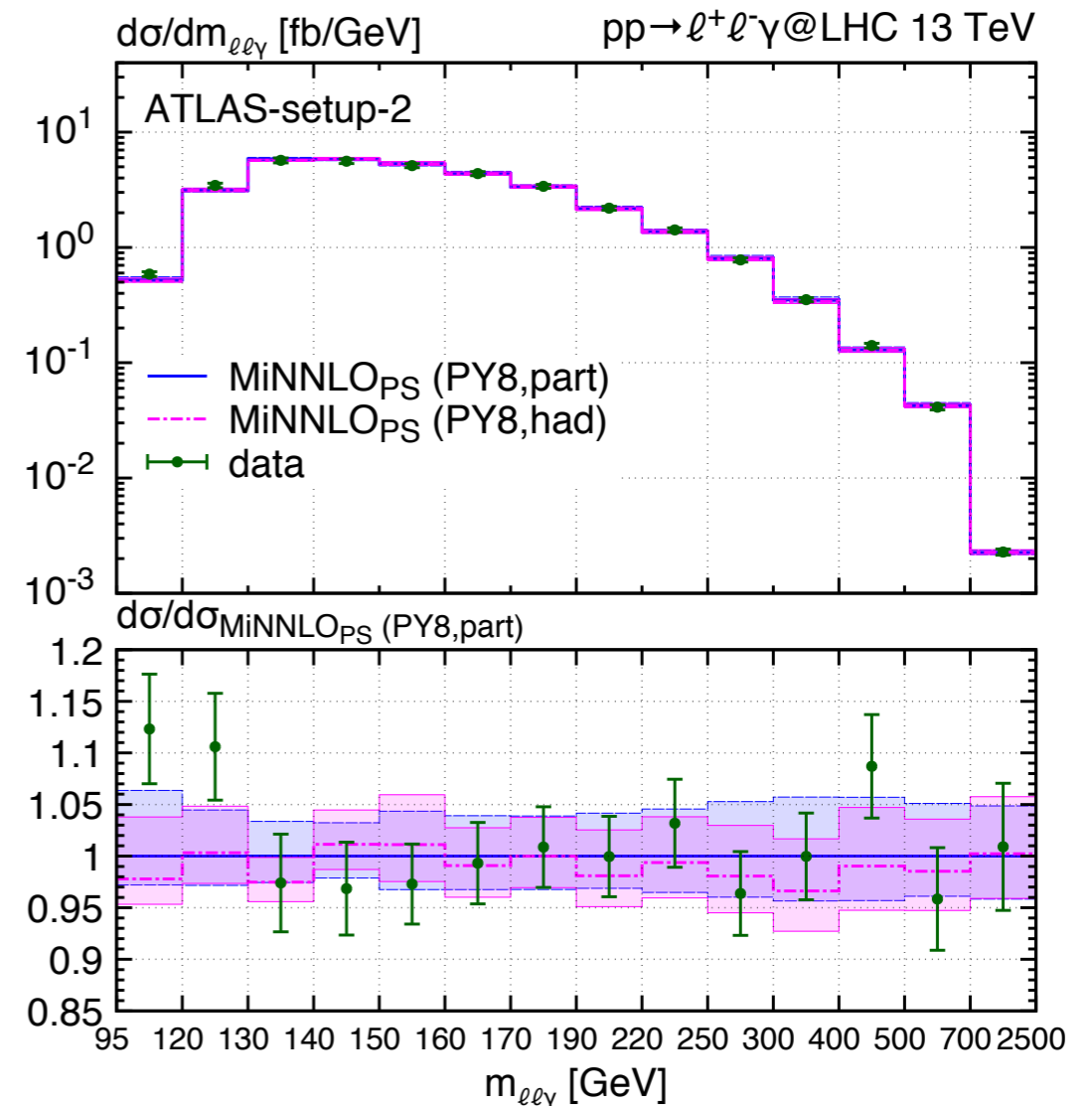
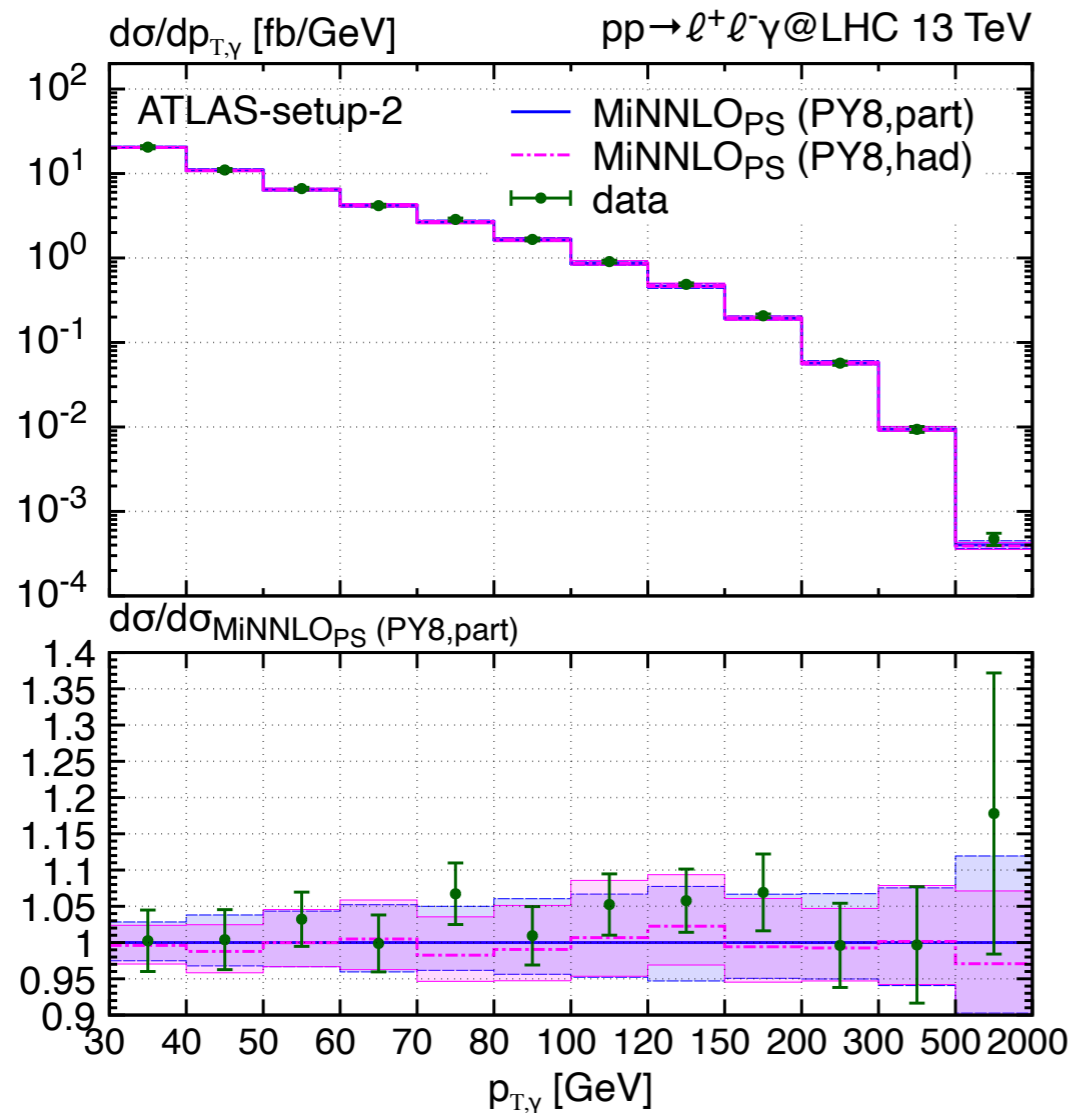
Photon isolation requirement:

- ❖ Experimentally needed to identify hard photons
- ❖ Theoretically delicate definition of an infrared-safe cross section

$$\sum_{\text{hadlpart} \in \delta} E_T^{\text{hadlpart}} \leq E_T^{\text{max}}(\delta) = \epsilon_\gamma p_{T,\gamma} \cdot \left(\frac{1 - \cos\delta}{1 - \cos\delta_0} \right)^n, \quad \forall \delta \leq \delta_0$$

NNLO+PS for $Z\gamma$: Phenomenological results

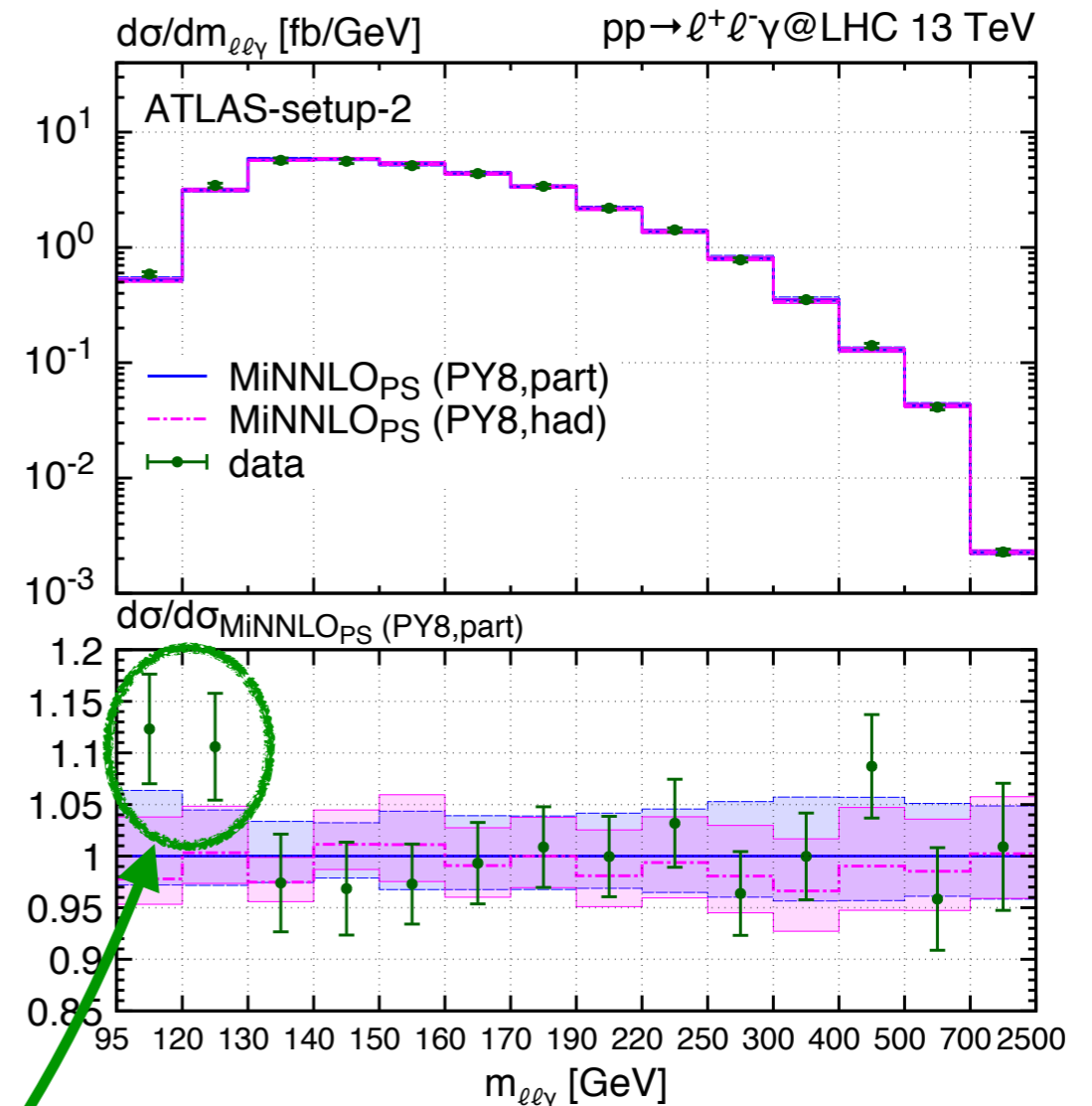
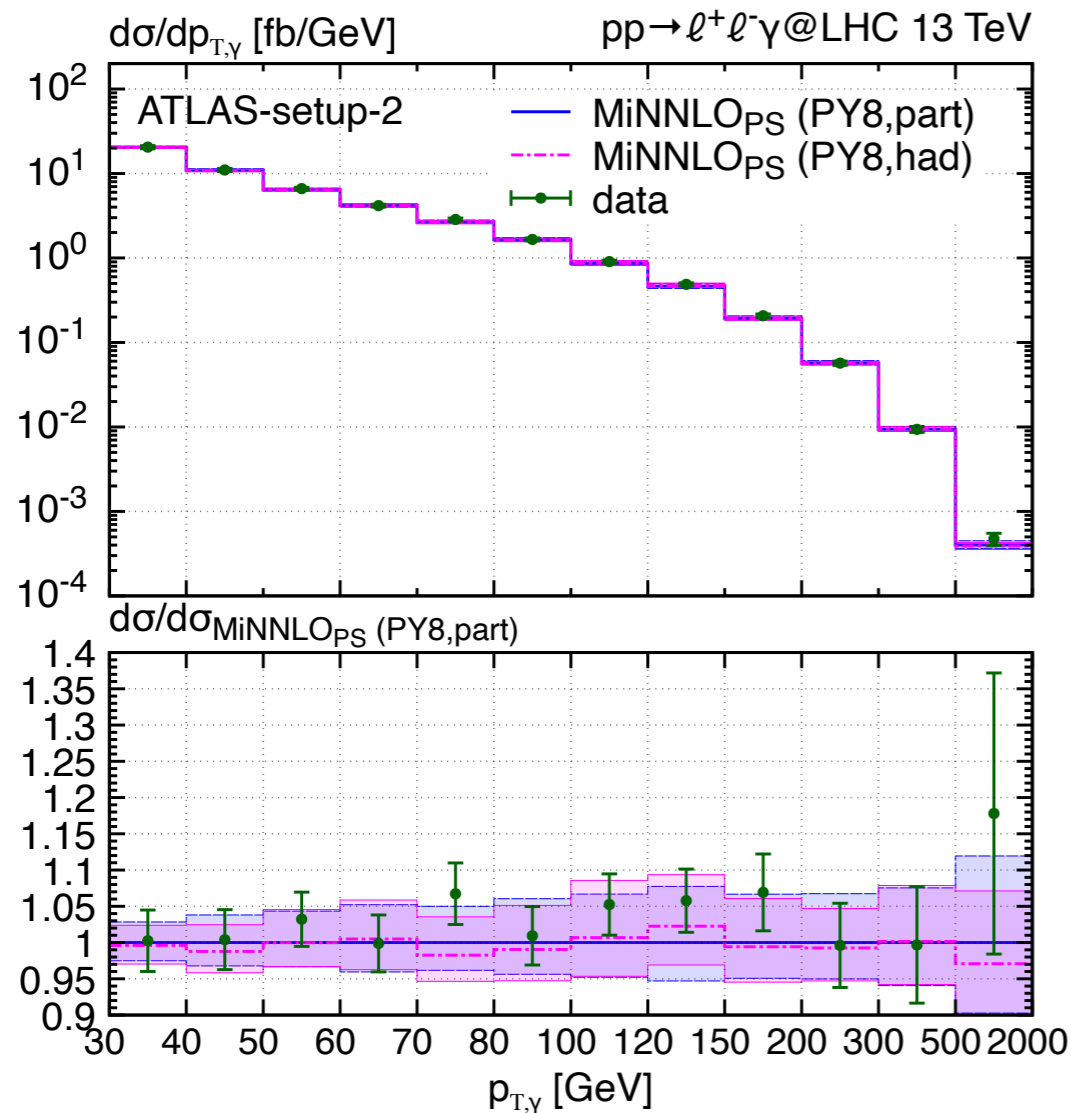
Comparison to data: ATLAS 13TeV analysis



- ✓ *MiNNLO_{PS} predictions in very good agreement with ATLAS 13TeV results, based on the full 139fb⁻¹ Run-2 data ([\[arXiv:1911.04813\]](#))!*

NNLO+PS for $Z\gamma$: Phenomenological results

Comparison to data: ATLAS 13TeV analysis



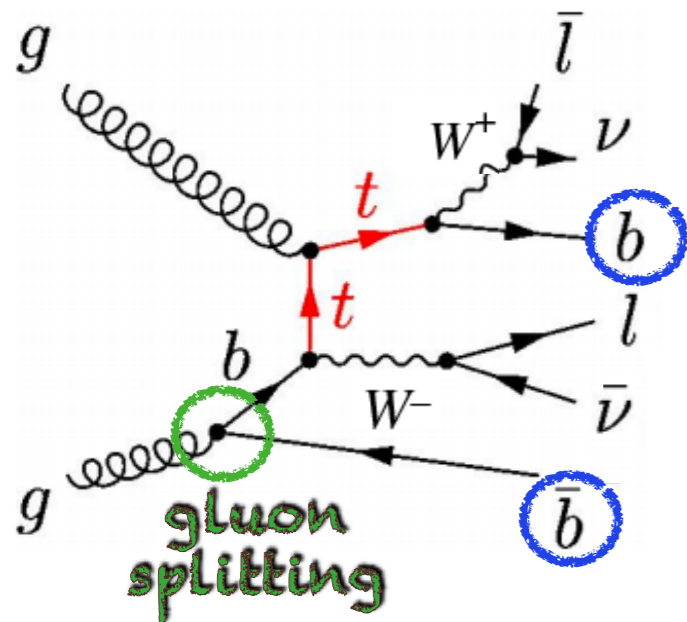
✓ *MiNNLO_{PS} predictions in very good agreement with ATLAS 13TeV results, based on the full 139fb⁻¹ Run-2 data ([\[arXiv:1911.04813\]](#))!*

✓ *Sensitivity to QED radiation*

WW phenomenology at LHC

[Lombardi, Wieseemann, Zanderighi '21]

- ❖ largest cross section among massive diboson processes
- ❖ direct access to anomalous triple gauge couplings
- ❖ no full event reconstruction due to neutrinos \rightarrow high-accurate theoretical predictions required
- ❖ analysis requires a **jet-veto** \rightarrow theoretical modelling important



Jet-veto requirement:

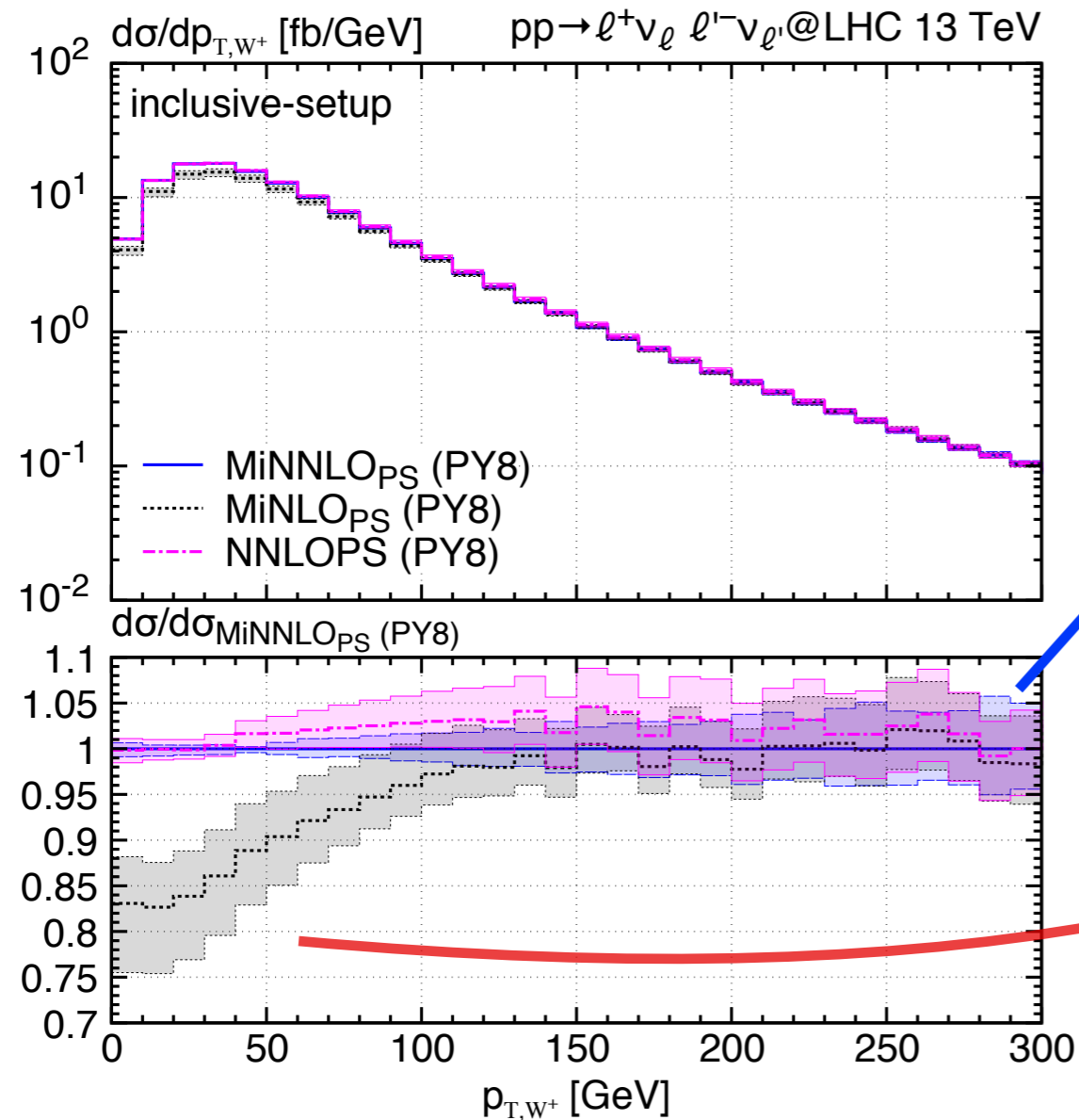
- ❖ Experimentally needed to reduce top background
- ❖ Theoretically involved definition of WW cross section, due to diagrams with resonant top quarks and b final states:
 - ▶ Interference with double-real diagrams
 - ▶ Not separately finite for massless b quarks



4 flavour scheme (4FS) adopted (in good agreement with top-free 5FS [Grazzini, Kallweit, Pozzorini, Rathlev, Wieseemann '17])

NNLO+PS for WW: Phenomenological results

Theoretical studies: inclusive phase space



✓ At high p_{T,W^+} NNLO+PS and MiNLO' predictions agree

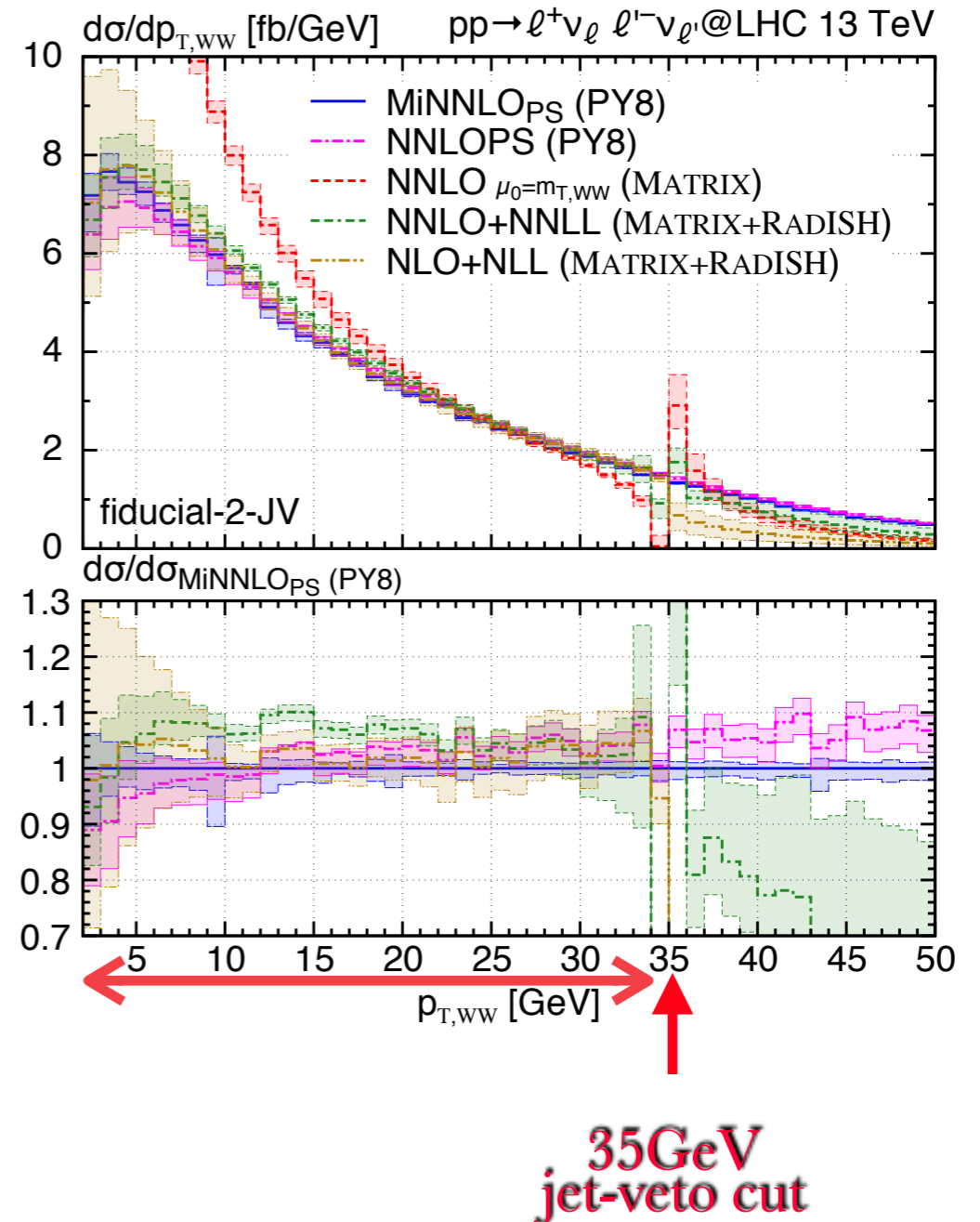
✓ Improved description in the low p_{T,W^+} region (both normalisation and accuracy)

NNLO+PS for WW: Phenomenological results

Comparison with resummation: jet-vetoed $p_{T,WW}$ spectrum

✓ *Fixed-order becomes unphysical at **low p_T** and at the **jet-veto threshold***

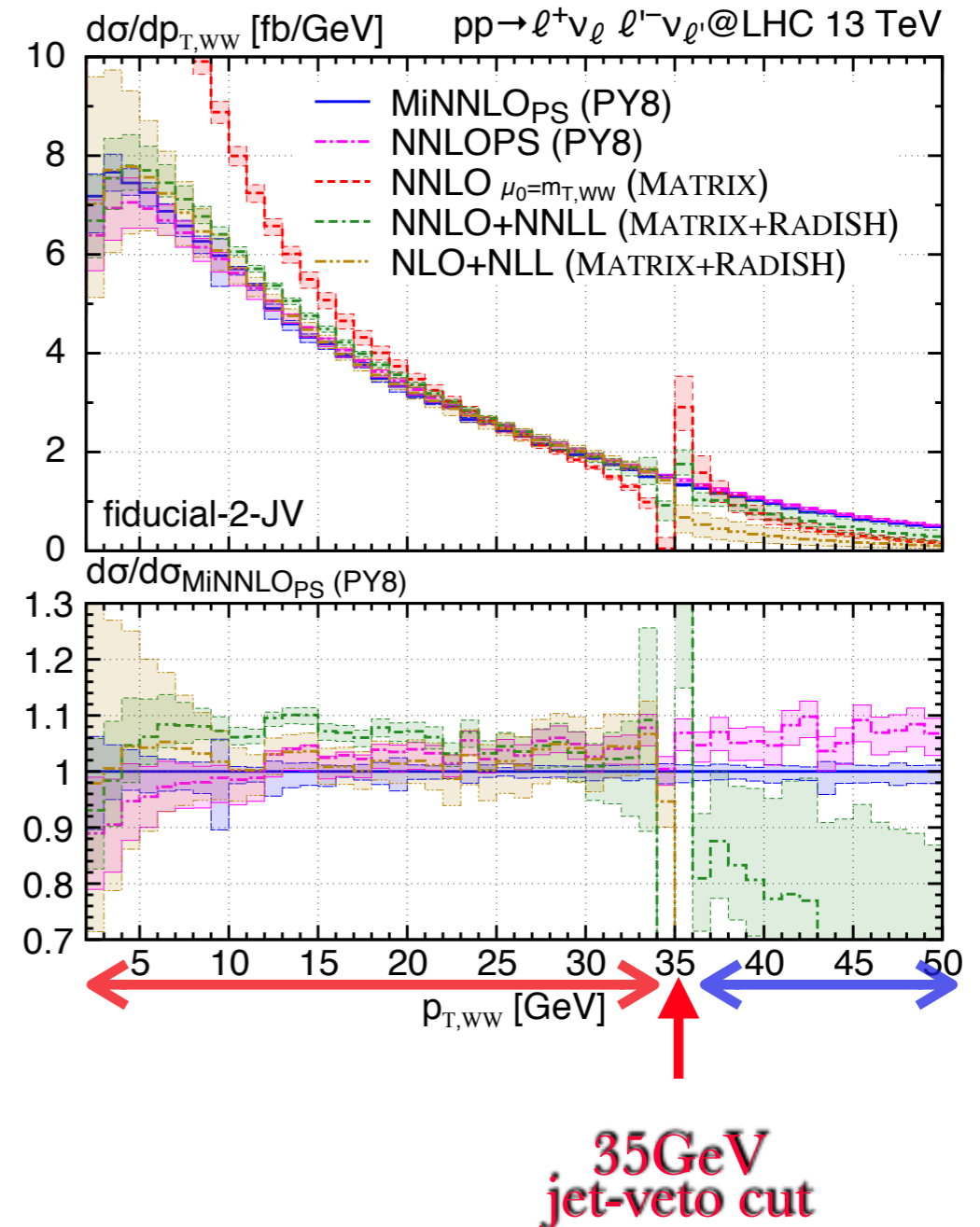
✓ *Analytic resummation only partially cures problems at low p_T , not **above jet-veto cut***



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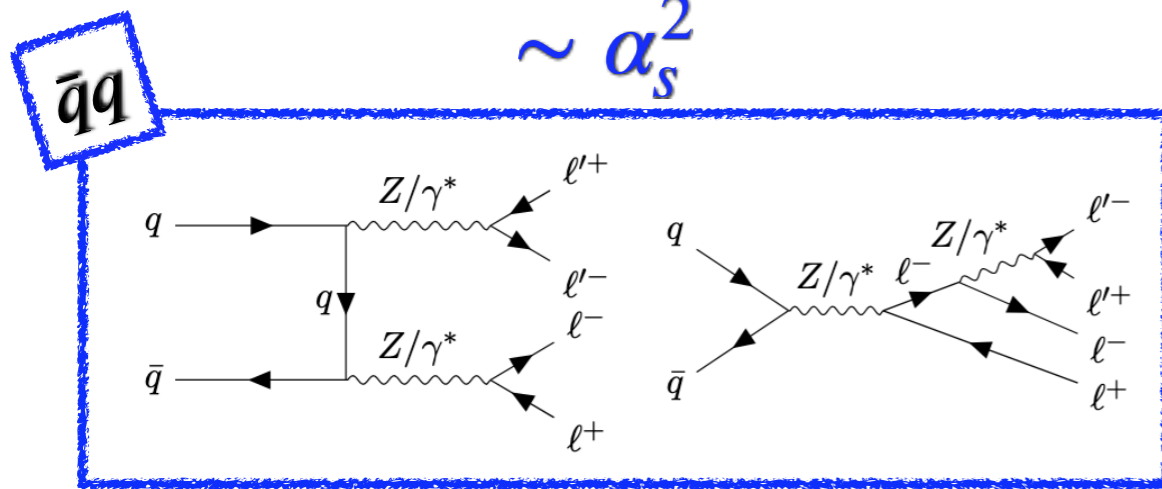
- ✓ *Fixed-order becomes unphysical at **low p_T** and at the **jet-veto threshold***
- ✓ *Analytic resummation only partially cures problems at low p_T , not **above jet-veto cut***
- ✓ *NNLO+PS results provide a more physical description also above jet-veto threshold*



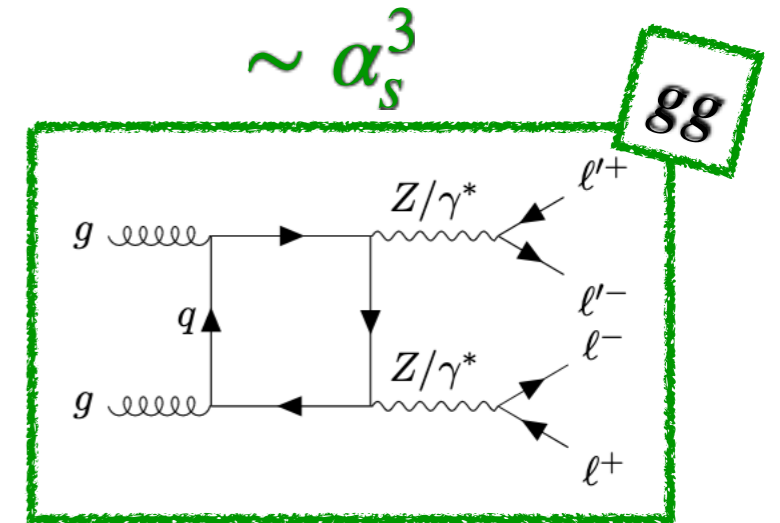
ZZ phenomenology at LHC

[Buonocore, Koole, Lombardi, Rottoli, Wiesemann, Zanderighi '21]

- ❖ smallest cross section among massive diboson processes, but very clean signature
- ❖ relevant for BSM searches
- ❖ important for constraining the Higgs width and Higgs couplings



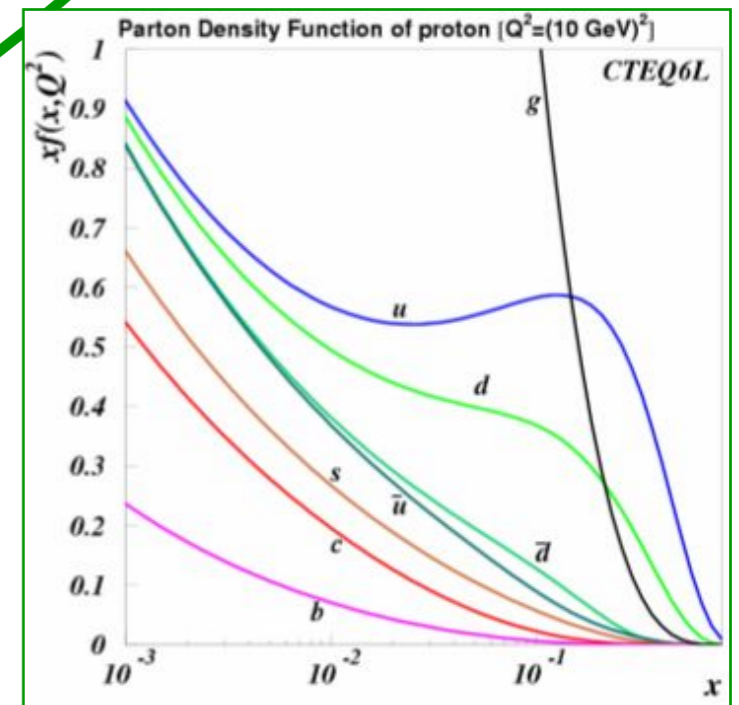
NNLO+PS using MiNNLO_{PS} *



NLO+PS with POWHEG *

$$pp \rightarrow \ell^+ \ell^- \ell^{(\prime)+} \ell^{(\prime)-}$$

Incoherent combination \rightarrow nNNLO+PS

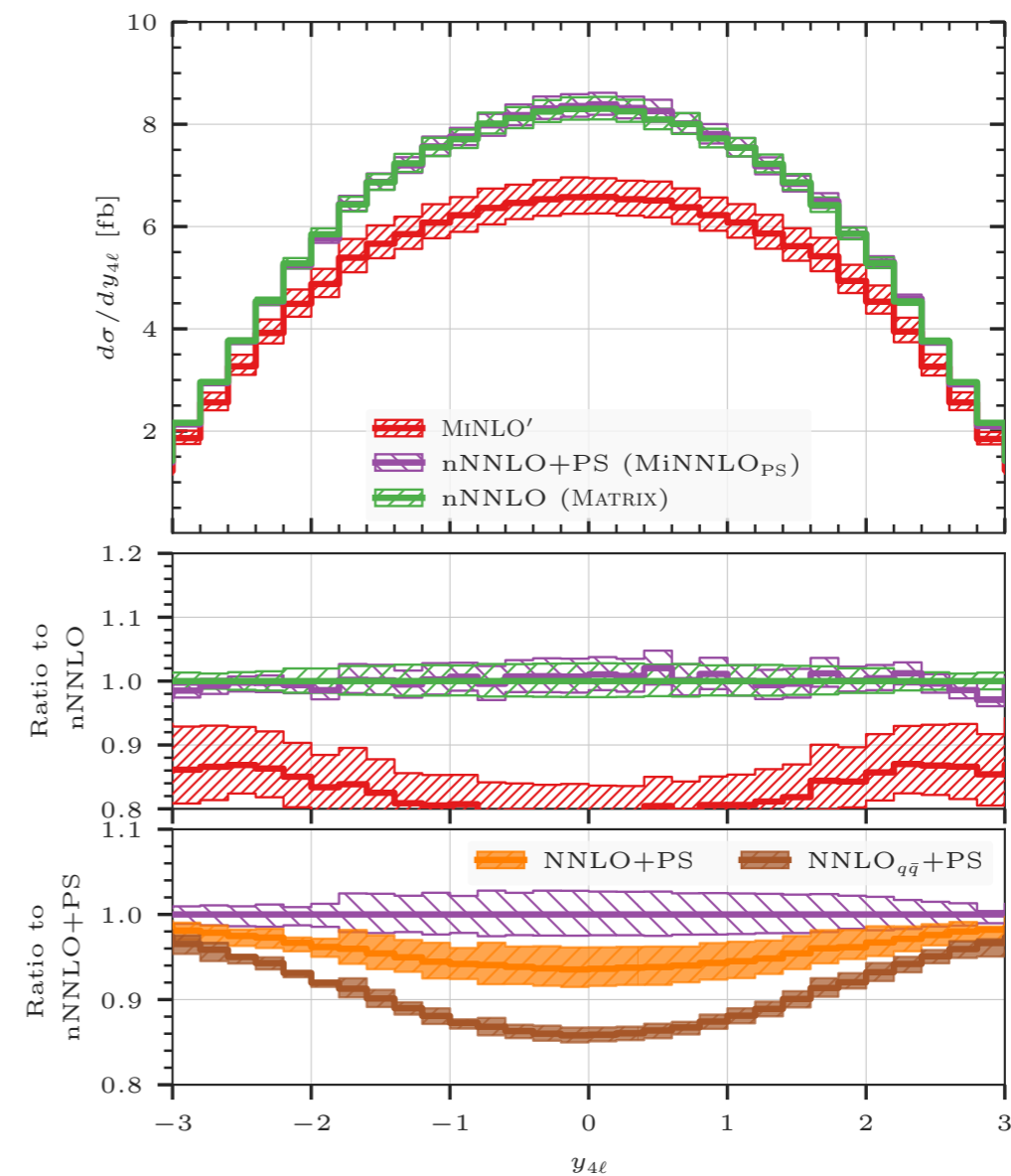
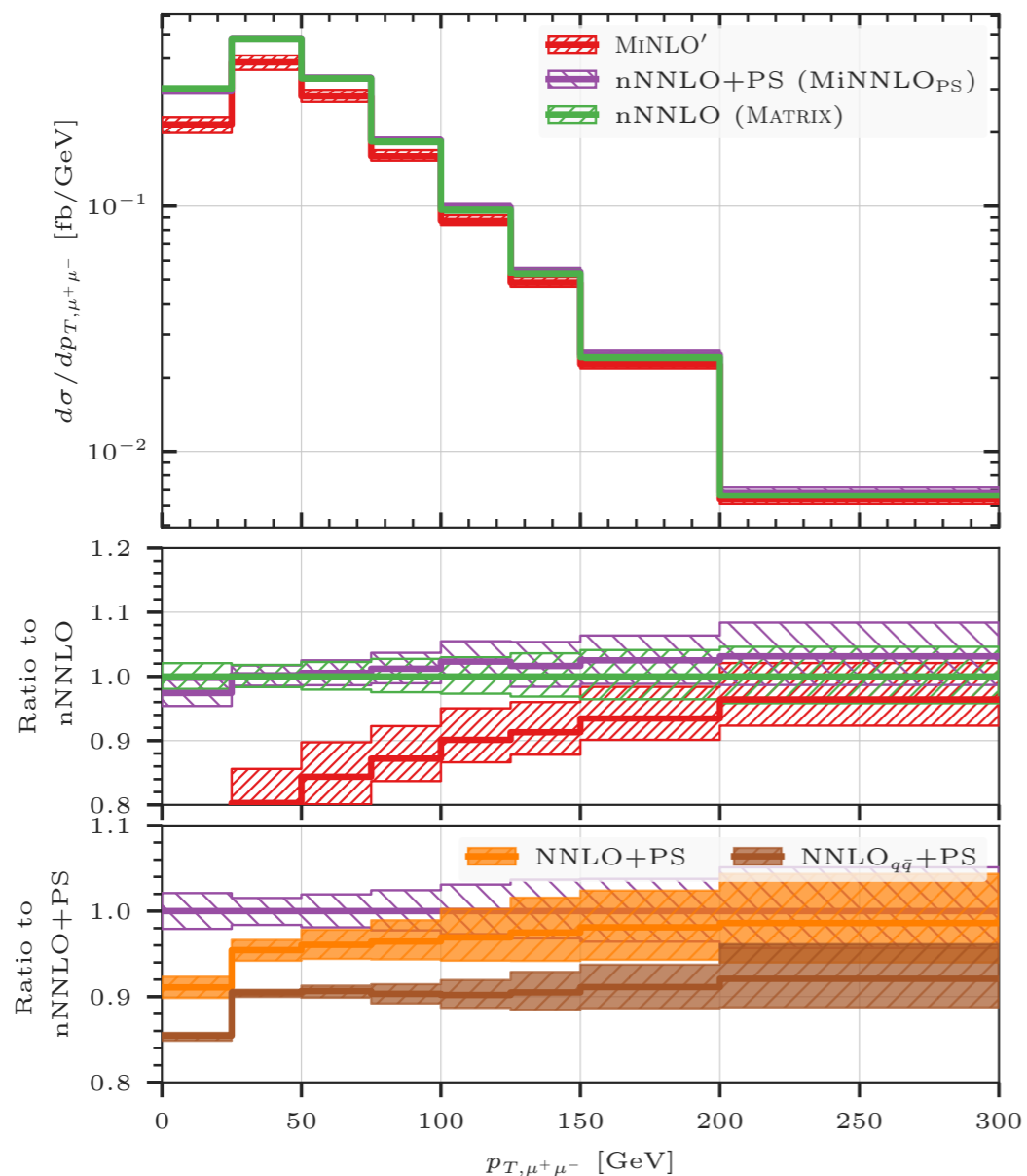


* also done in GENEVA framework [Alioli, Broggio, Gavardi, Kallweit, Lim '21]

* first independent implementation in [Alioli, Ferrario Ravasio, Lindert, Röntsch '21]

nNNLO+PS for ZZ: Phenomenological results

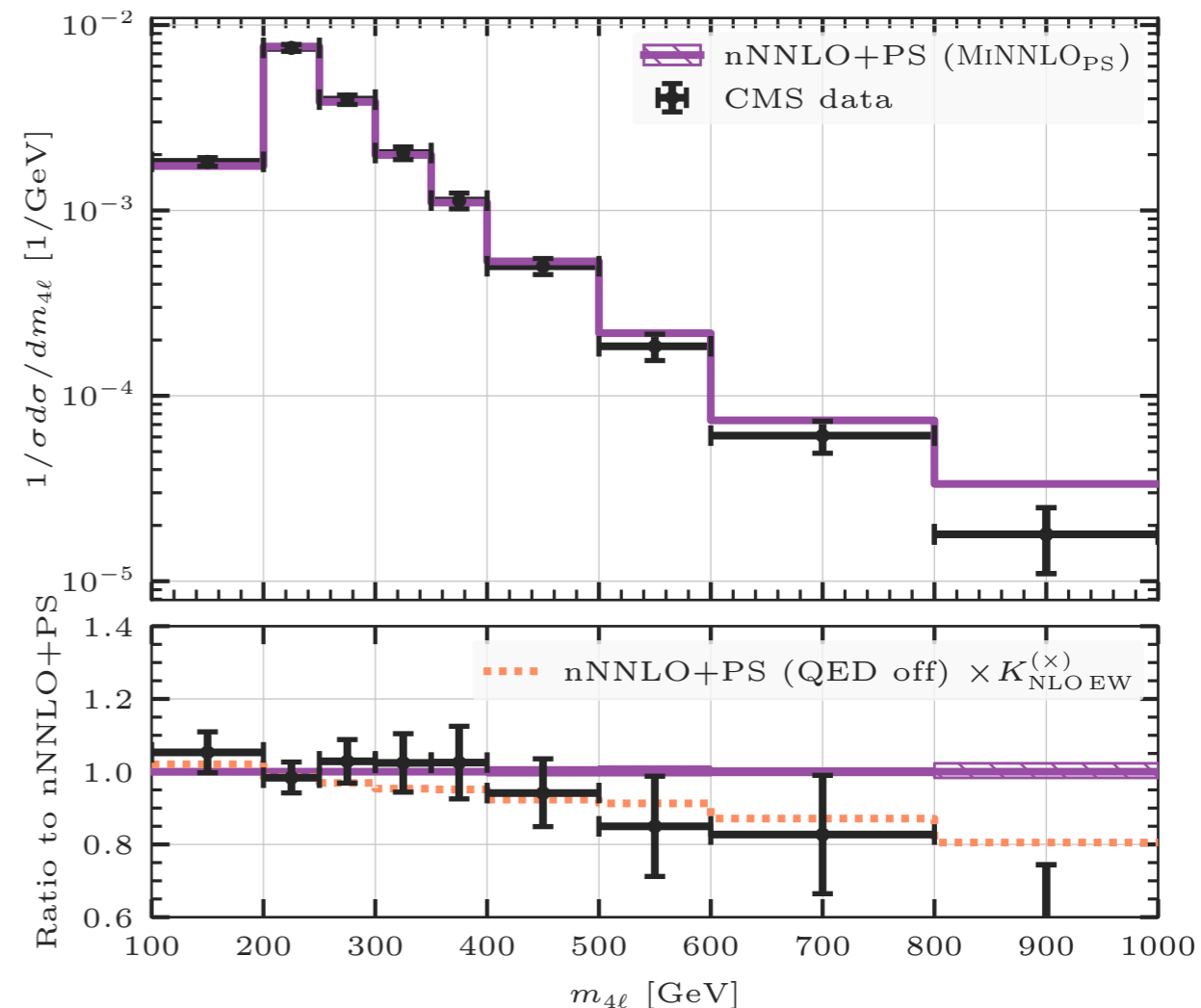
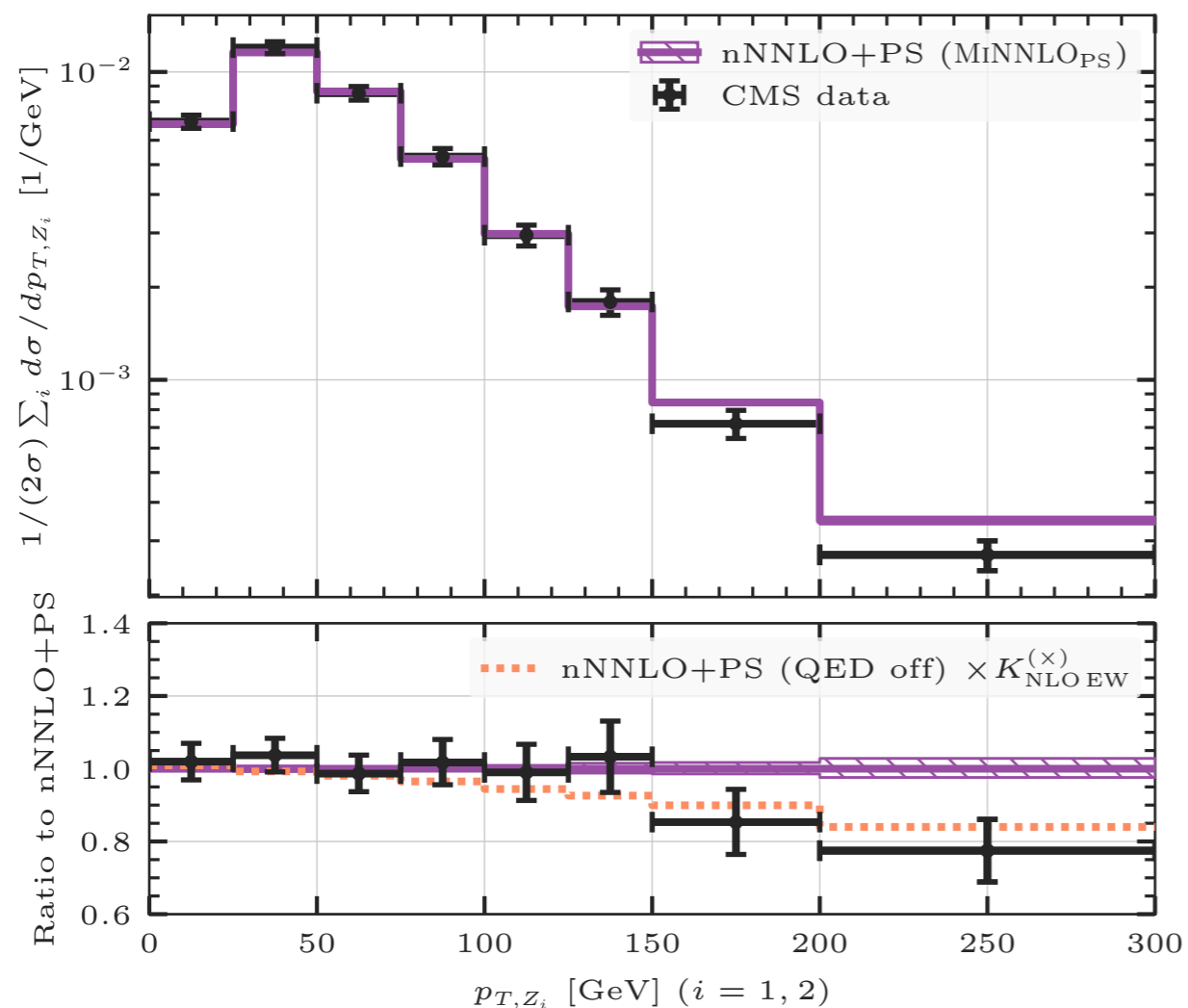
Differential distributions



✓ *Sizeable effects from the inclusion of gg contributions in bulk regions*

nNNLO+PS for ZZ: Phenomenological results

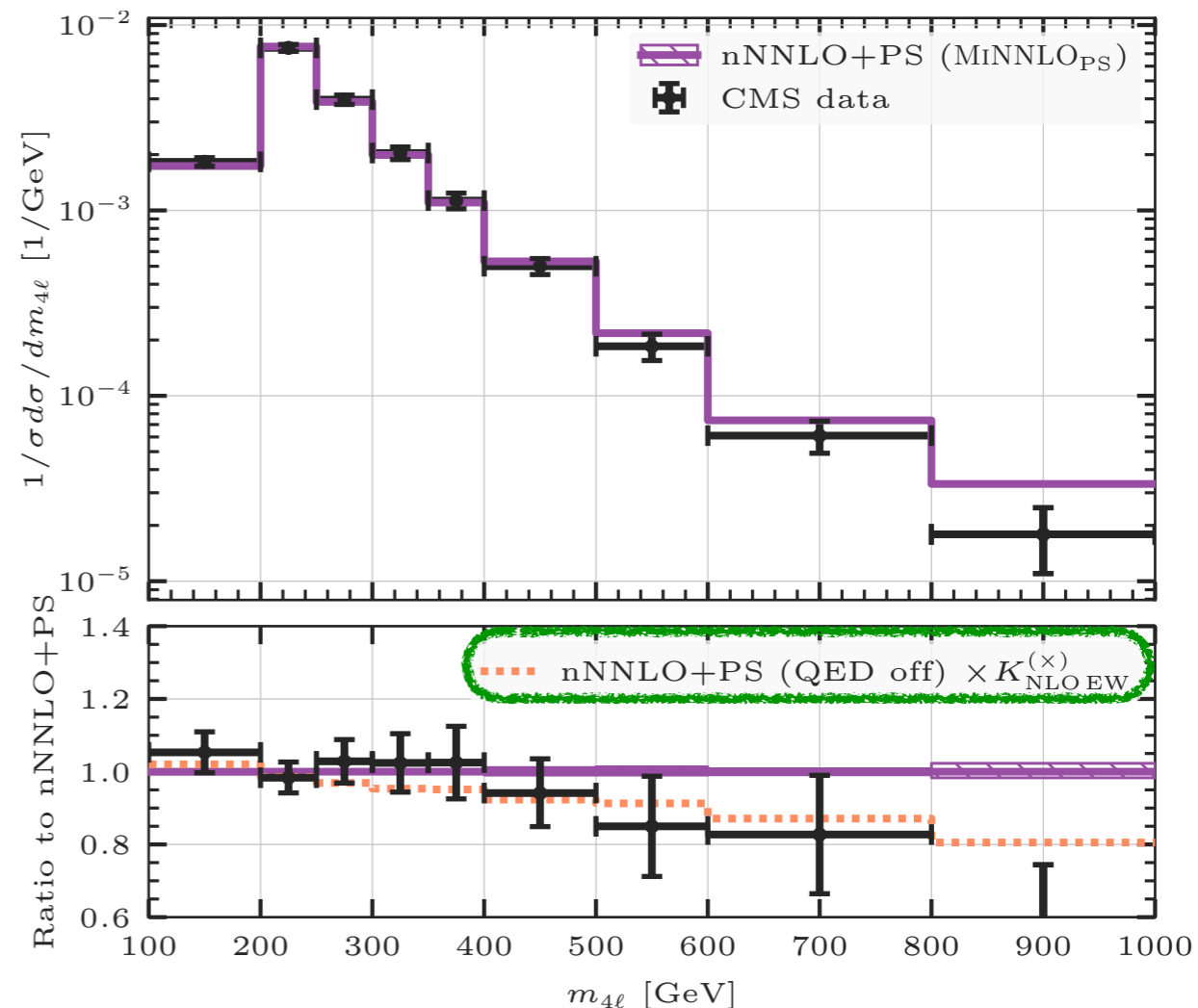
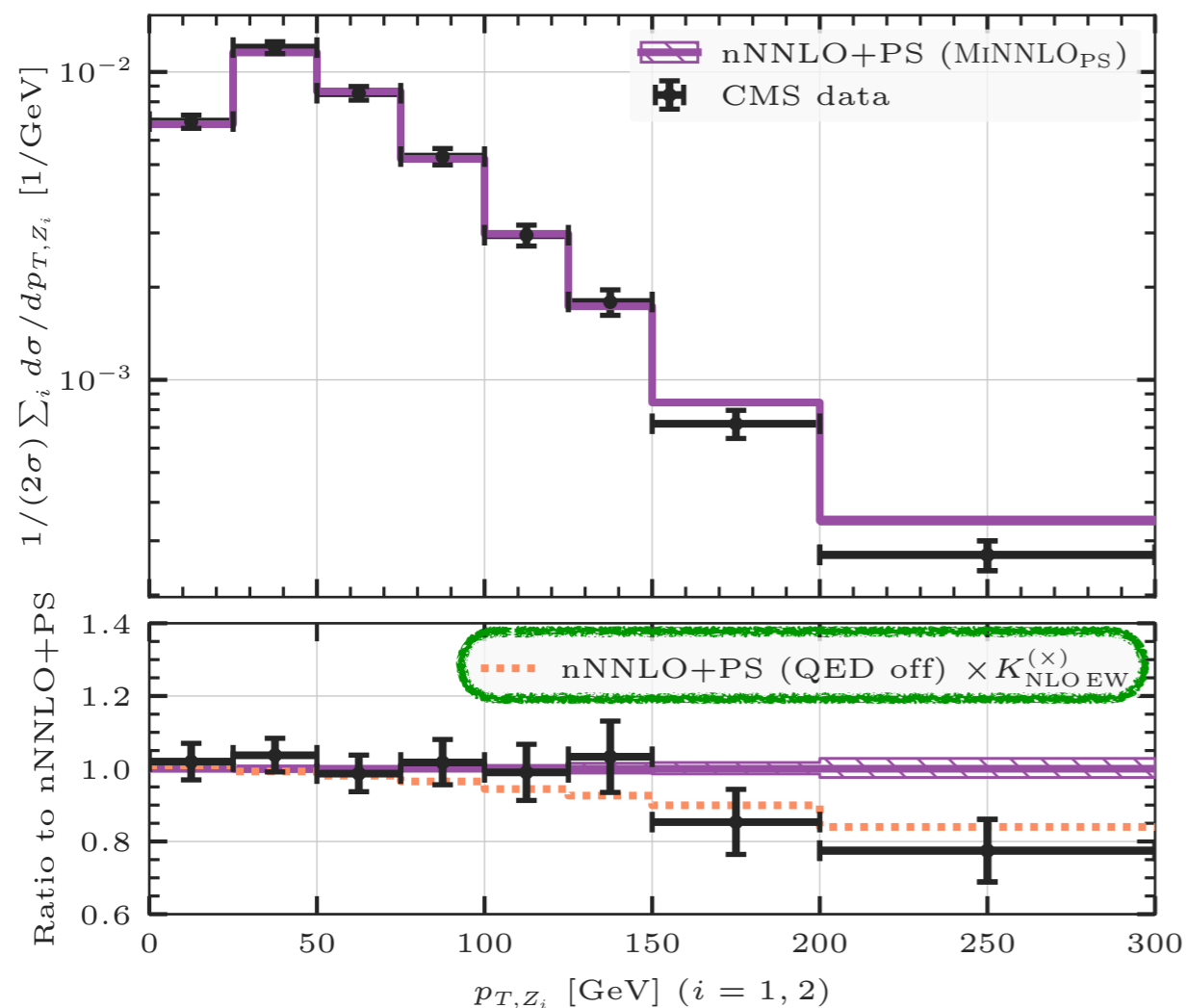
Comparison to data: CMS 13TeV analysis



✓ *nNNLO+PS predictions in good agreement with CMS results, based on the $a137\text{fb}^{-1}$ 13TeV analysis ([arXiv:2009.01186])!*

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Comparison to data: CMS 13TeV analysis



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✓ *Consistent inclusion of EW corrections to describe tails of distributions*

$$K_{NLO EW}^{(x)} = \frac{\sigma_{nNNLO QCD} + \sigma_{LO}^{\bar{q}q} \cdot (\delta_{EW} + \delta_{QCD} \delta_{EW})}{\sigma_{nNNLO QCD}}$$

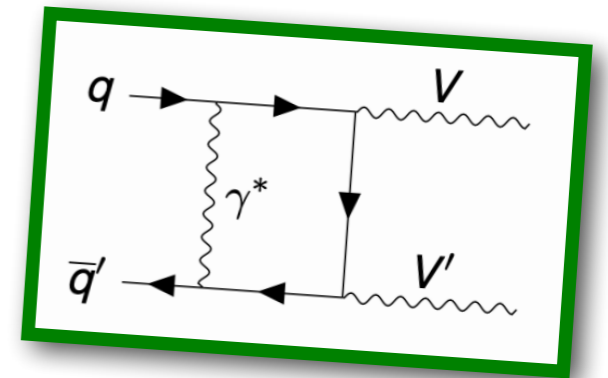
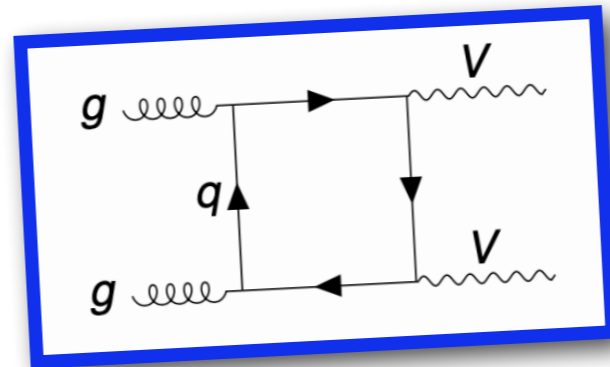
Summary:

- ✓ *NNLO+PS accuracy crucial for realistic and precise LHC phenomenology*
- ✓ *Very good description of data by **Z γ** MiNNLO_{PS} predictions*
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Outlook:

- ✓ *MiNNLO_{PS} to all other diboson processes in the pipeline*
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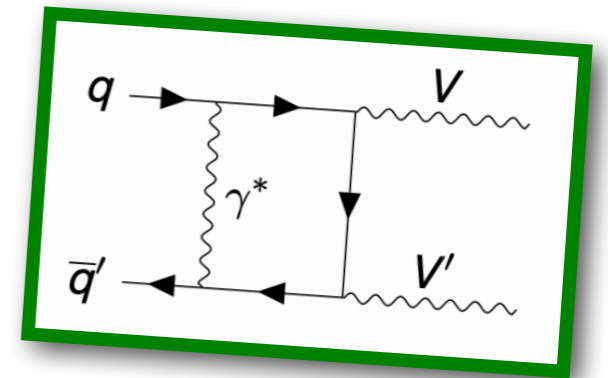
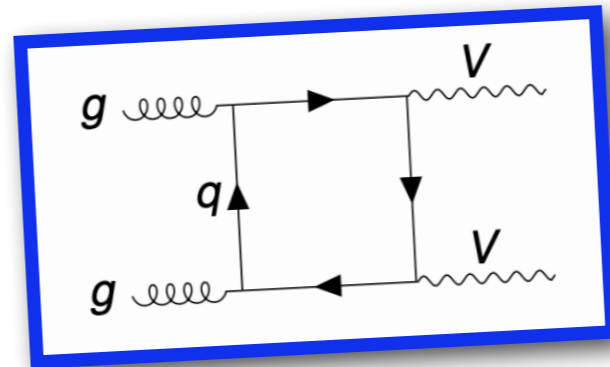
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Thank you for your attention