

Workshop on Tools for High Precision LHC Simulations

Inclusion of NLO EW corrections in NNLO+PS event generation with MiNNLO_{PS}

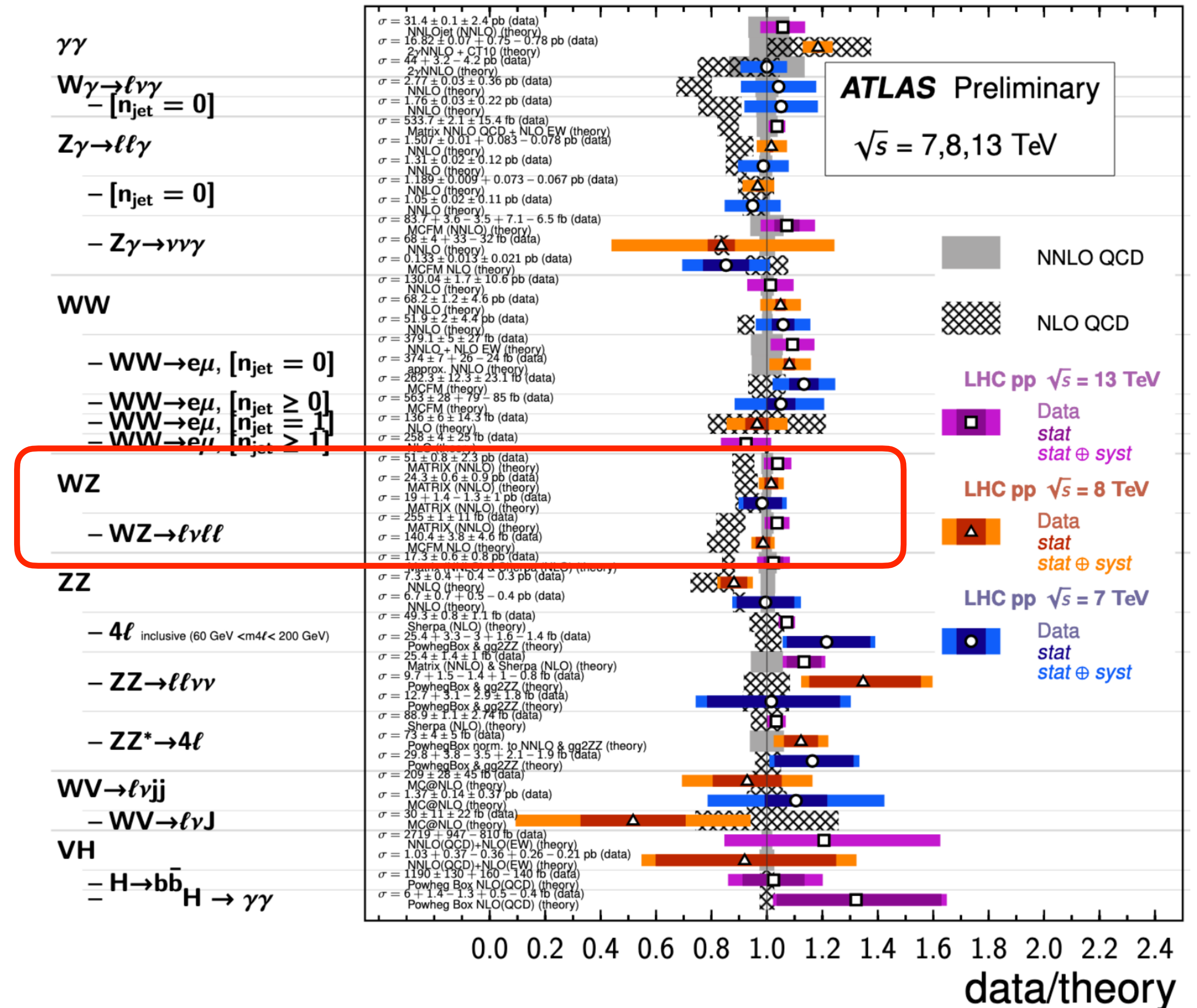
MAX PLANCK INSTITUTE
FOR PHYSICS



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Based on [2208.12660] J. Lindert, D. Lombardi, M. Wiesemann, G. Zanderighi, S.Z.

- Precision physics is a promising path for the **observation of effects beyond the Standard Model**.
- NNLO (QCD) computations are **crucial** for an accurate description of data, but they need to be supplemented by **EW corrections** for reaching the required experimental accuracy.
- The **matching** of a fixed-order calculation **with parton showers** is required for a realistic description of an event at a collider.



WHY WZ PRODUCTION?

- **The production of a pair of vector bosons is highly relevant**, as it provides access to trilinear gauge couplings and to the gauge symmetry structure of the EW sector.
- **WZ production is particularly interesting** both for the large cross section and the clean experimental signature (we consider the purely leptonic decay with one neutrino).

CURRENT STATE OF THE ART:

NLO EW calculation

[Bierweiler, Kasprzik, Kühn (2013), Baglio, Ninh, Weber (2013)]

[Biedermann, Denner, Hofer (2017)]

NNLO QCD calculation

[Grazzini, Kallweit, Rathlev, Wiesemann (2016), (2017)]

NLO QCD + NLO EW matched to Parton Showers

[Chiesa, Oleari, Re (2020)]

NNLO QCD + NLO EW combination

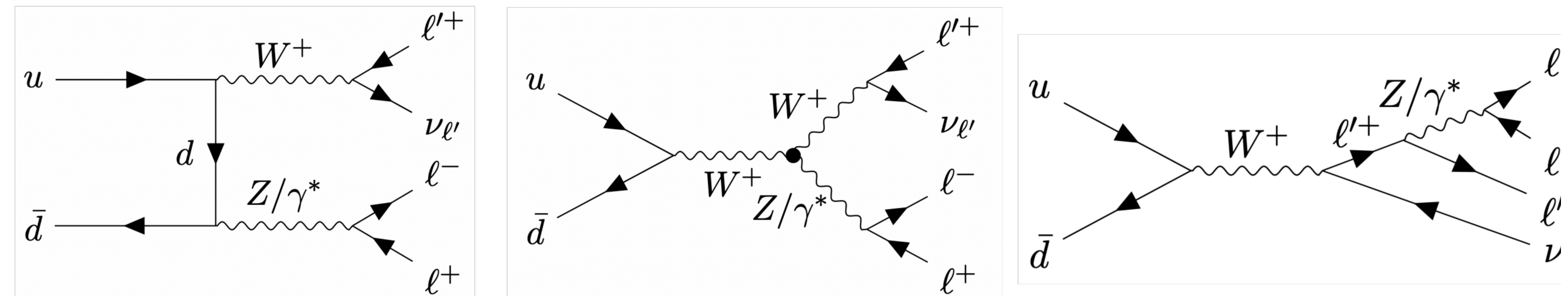
[Grazzini, Kallweit, Lindert, Pozzorini, Wiesemann (2020)]

THIS TALK:

- NNLO+PS (QCD) calculation using MiNNLO_{PS}
- Combination of NNLO+PS (QCD) with NLO+PS (EW) computations

- 1) How to obtain NNLO+PS (QCD) and NLO+PS (EW) results for WZ production
 - a. Generation of NNLO+PS (QCD) results using MiNNLO_{PS}
 - b. Generation of NLO+PS (EW) results using POWHEG
 - c. Combination of NNLO+PS (QCD) and NLO+PS (EW) results (different matching schemes)
 - d. Treatment of the shower (PY8) and implementation of a veto procedure
- 2) Phenomenological analysis and comparison with data
- 3) Summary and outlooks

$$pp \rightarrow l'^{\pm} \nu_{l'} l^+ l^- + X$$



1) Generation of NNLO+PS (QCD) results using MiNNLO_{PS}

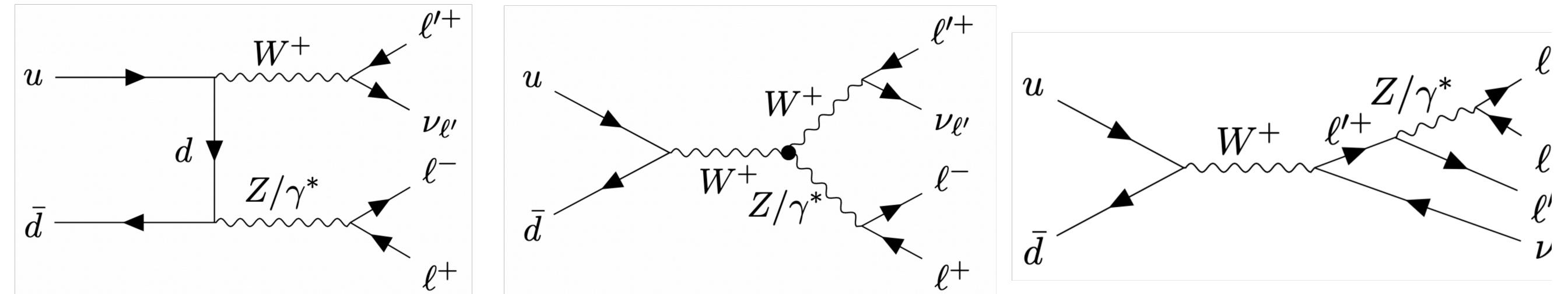
$$d\sigma_F^{\text{MiNNLO}_{\text{PS}}} = d\Phi_{\text{FJ}} \bar{\mathbf{B}}^{\text{MiNNLO}_{\text{PS}}} \times \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(\mathbf{p}_{\text{T,rad}}) \frac{\mathbf{R}_{\text{FJ}}}{\mathbf{B}_{\text{FJ}}} \right\}$$

$$\bar{\mathbf{B}}^{\text{MiNNLO}_{\text{PS}}} \sim e^{-S} \left\{ d\sigma_{\text{FJ}}^{(1)} (1 + S^{(1)}) + d\sigma_{\text{FJ}}^{(2)} + (D - D^{(1)} - D^{(2)}) \right\}$$

Very simplified notation!

- Order $\mathcal{O}(\alpha^4 \alpha_s^2)$.
- No loop-induced gluon-fusion contributions.
- Important NNLO corrections (10-15%), due to radiation zero effect at LO (= vanishing of the leading helicity amplitudes in some kinematic regions).

$$pp \rightarrow l'^{\pm} \nu_{l'} l^+ l^- + X$$



2) Generation of NLO+PS (EW) results using POWHEG

$$d\sigma_F^{\text{pwg}} = d\Phi_F \bar{B}^{\text{pwg}} \times \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + \int d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{T,\text{rad}}) \frac{R_F}{B_F} \right\}$$

$$\bar{B}^{\text{pwg}} = d\sigma_F^{(1)} + d\sigma_F^{(2)}$$

Very simplified notation!

- Order $\mathcal{O}(\alpha^5)$.
- Real radiation corresponds to photon radiation.
- No photon-photon contribution at this order.
- Photon-quark contributions are not considered (formally, they are $\mathcal{O}(\alpha^6 L)$).



3) Combination of NNLO_{QCD}+PS and NLO_{EW}+PS: MATCHING SCHEMES

ADDITIVE SCHEME

$$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}} - \text{LO}$$

$$\mathcal{O}(\alpha^4), \mathcal{O}(\alpha^4\alpha_s), \mathcal{O}(\alpha^4\alpha_s^2), \mathcal{O}(\alpha^5)$$

MULTIPLICATIVE SCHEME

$$\text{NNLO}_{\text{QCD}} \times \text{NLO}_{\text{EW}}/\text{LO}$$

$$\mathcal{O}(\alpha^4), \mathcal{O}(\alpha^4\alpha_s), \mathcal{O}(\alpha^4\alpha_s^2), \mathcal{O}(\alpha^5), \mathcal{O}(\alpha^5\alpha_s), \mathcal{O}(\alpha^5\alpha_s^2)$$

- The **multiplicative scheme is preferable** in the high energy limit, where EW Sudakov-logs are dominant and dominant QCD effects arise at scales below the hard scale. —> **QCD factorizes**.
- This assumption is **violated when giant K-factors are present** (= hard vector-boson+jet topologies, with a soft second vector boson).
- The **average** of the two schemes can give a **pragmatic estimate** in these regions.

3) Combination of NNLO_{QCD}+PS and NLO_{EW}+PS: TREATMENT OF THE SHOWER

1. The **formal accuracy** of the calculation **must not be spoiled**.
2. We must **avoid double counting**.

We let the QCD and/or QED showers radiate in whole the phase space and then we apply the following **veto procedure**:

NNLO_{QCD}+PS :

- **QCD** shower is **restricted** by the transverse momentum of the hardest QCD emission generated at Les Houches level (as commonly done in POWHEG).

- **QED** shower is **unconstrained**.

NLO_{EW}+PS :

- **QCD** shower is **unconstrained**.
- **QED** shower is **restricted** by the transverse momentum of the hardest QED emission generated at Les Houches level (POWHEG multiple-radiation scheme → **three different starting scales** for ISR, FSR from W decay, FSR from Z decay).

ADDITIVE: 1. $\text{NNLO}_{\text{QCD}}^{(\text{QCD}, \text{QED})_{\text{PS}}} + \text{NLO}_{\text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}} - \text{LO}^{(\text{QCD}, \text{QED})_{\text{PS}}} = \text{NNLO}_{\text{QCD+EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}}$

2. $\text{NNLO}_{\text{QCD}}^{(\text{QCD}, \text{QED})_{\text{PS}}} + \text{NLO}_{\text{EW}}^{(\text{QED})_{\text{PS}}} - \text{LO}^{(\text{QED})_{\text{PS}}}$

3. $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}} + \text{NLO}_{\text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}} - \text{LO}^{(\text{QCD})_{\text{PS}}}$

MULTIPLICATIVE: 4. $\text{NNLO}_{\text{QCD}}^{(\text{QCD}, \text{QED})_{\text{PS}}} \times \text{NLO}_{\text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}} / \text{LO}^{(\text{QCD}, \text{QED})_{\text{PS}}} = \text{NNLO}_{\text{QCD} \times \text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}}$

5. $\text{NNLO}_{\text{QCD}}^{(\text{QCD}, \text{QED})_{\text{PS}}} \times \text{NLO}_{\text{EW}}^{(\text{QED})_{\text{PS}}} / \text{LO}^{(\text{QED})_{\text{PS}}}$

6. $\text{NLO}_{\text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}} \times \text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}} / \text{LO}^{(\text{QCD})_{\text{PS}}}$

7. $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}} \times \text{NLO}_{\text{EW}}^{\text{f.o.}} / \text{LO}^{\text{f.o.}}$

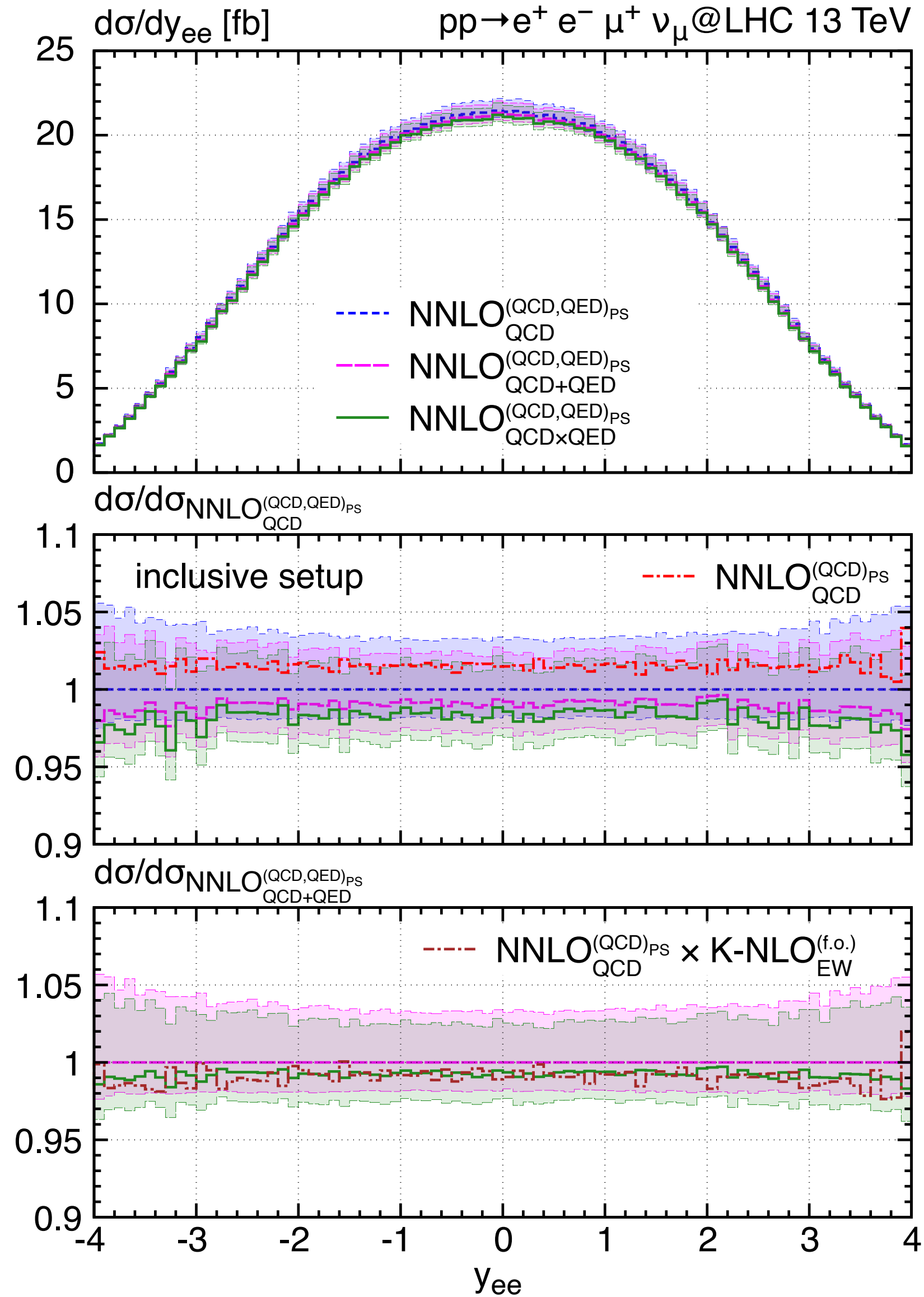
NOTATION:

$$(\text{N})\text{NLO}_{\text{X}}^{(\text{Y})_{\text{PS}}}$$

X = QCD,EW calculation

Y = QCD,QED showers (PY8)

Rapidity of the Z boson - inclusive setup

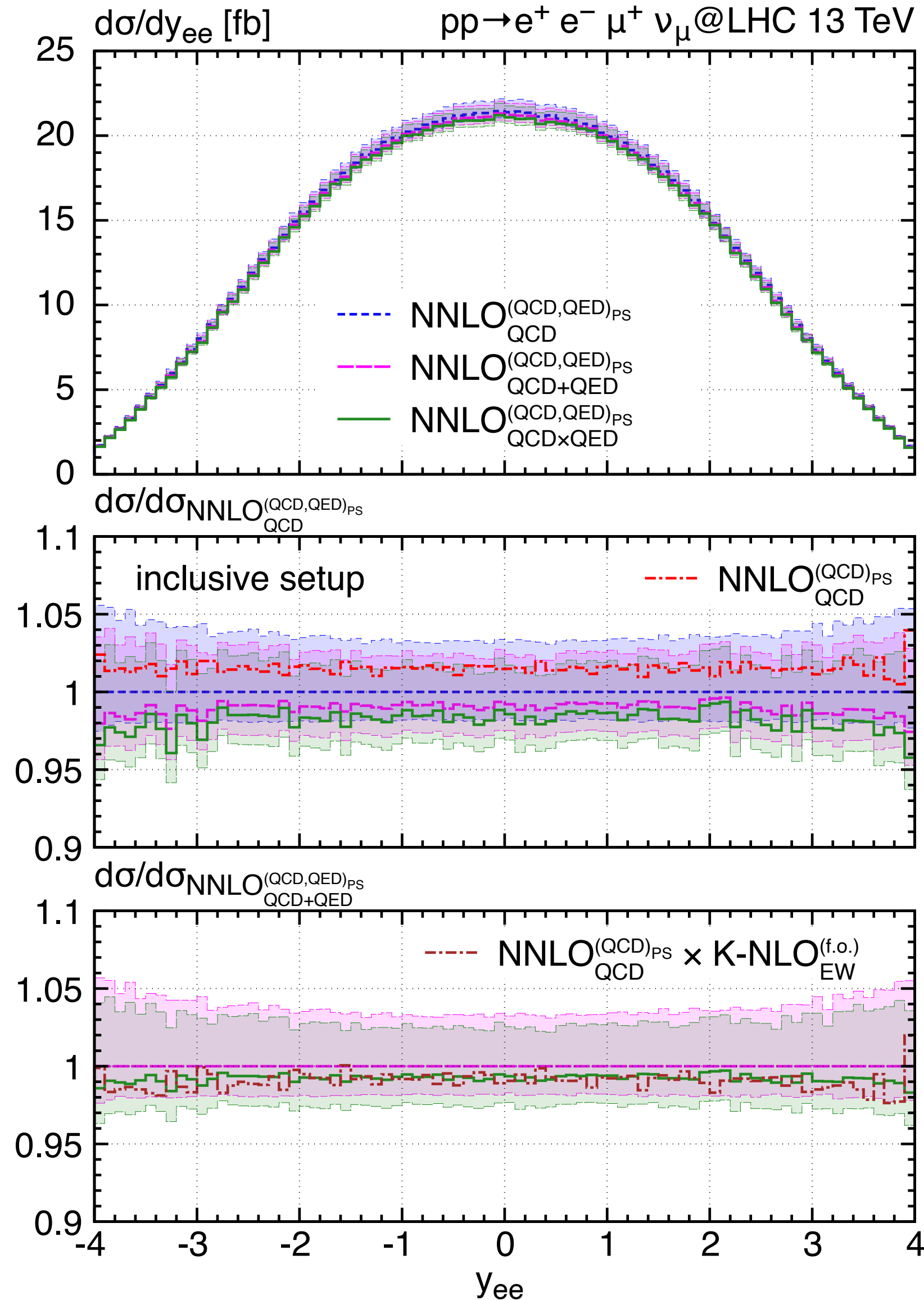


LEGEND:

- $\text{NNLO}_{\text{QCD}}^{(\text{QCD},\text{QED})_{\text{PS}}}$
- $\text{NNLO}_{\text{QCD+EW}}^{(\text{QCD},\text{QED})_{\text{PS}}}$
- $\text{NNLO}_{\text{QCD}\times\text{EW}}^{(\text{QCD},\text{QED})_{\text{PS}}}$
- $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}}$
- $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}} \times K_{\text{EW}}^{\text{f.o.}}$

Rapidity of the Z boson - inclusive setup

- Almost no shape effect
- EW corrections are 2-3%
- Additive ● and multiplicative ● schemes are almost identical
- Fixed-order K-factor ● is in excellent agreement \rightarrow effects of secondary photon emission are negligible for this observable



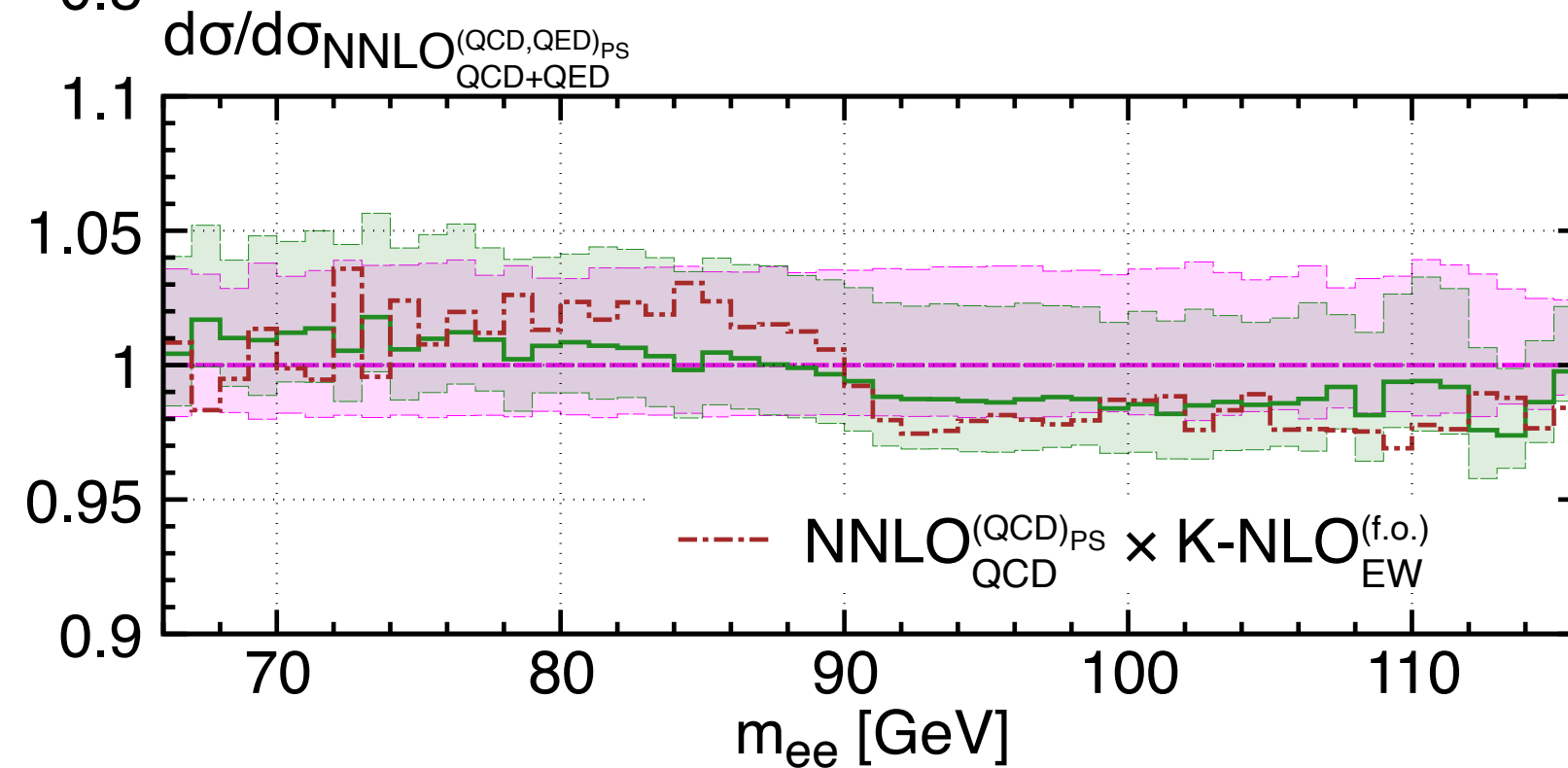
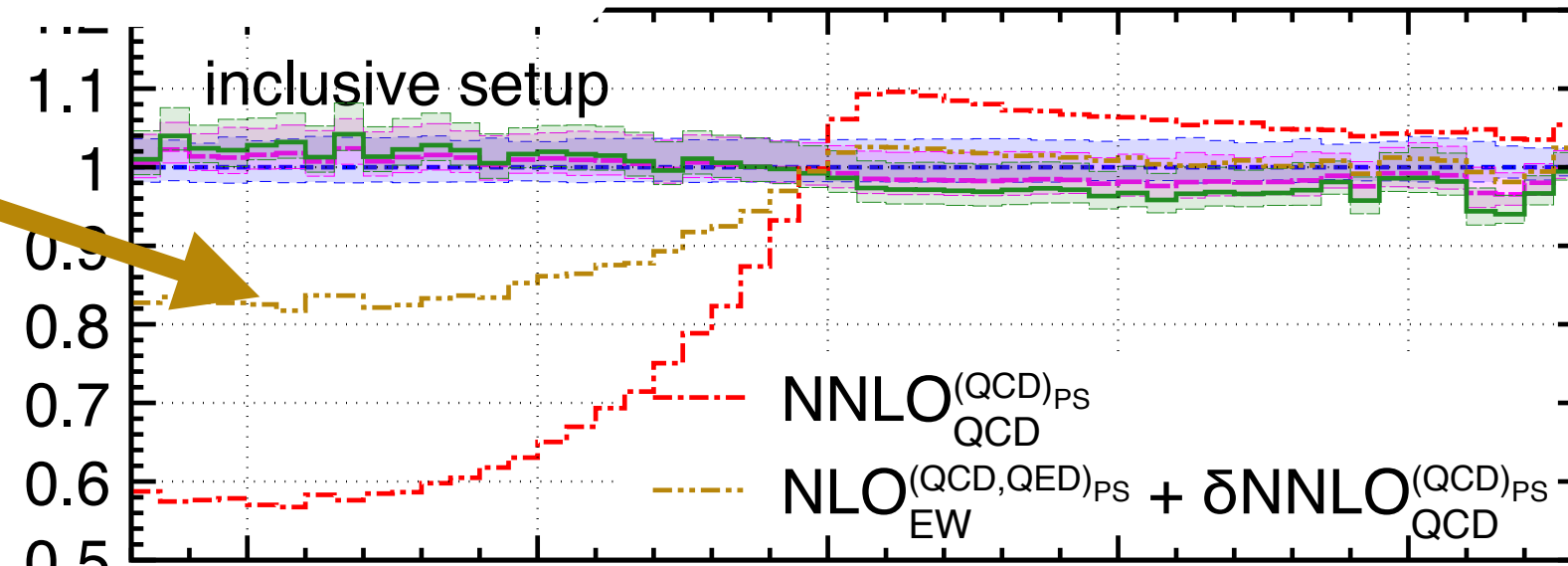
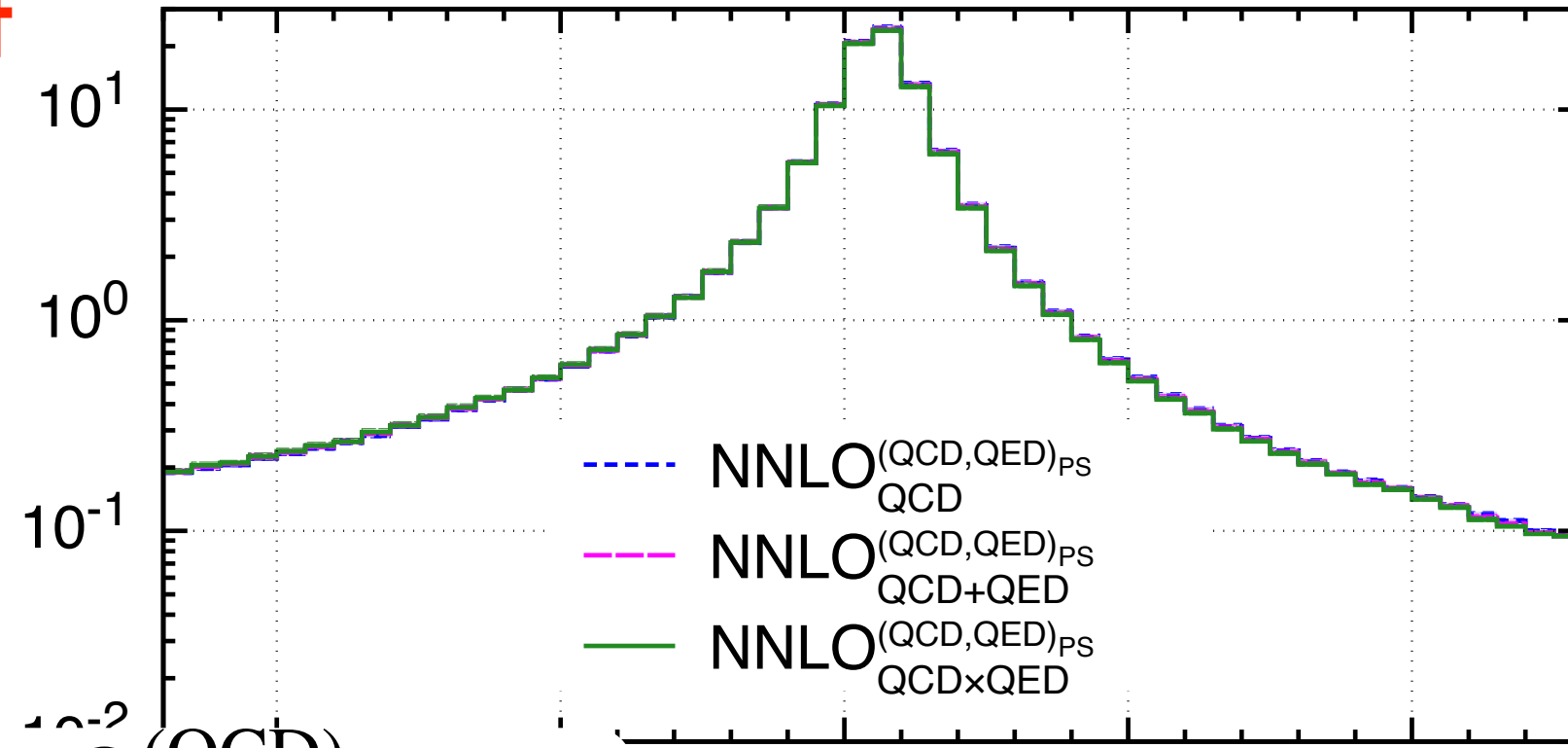
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- $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}}$
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Invariant mass of the Z boson - inclusive setup

● $\text{NNLO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}} + \text{NLO}_{\text{EW}}^{(\text{QCD}, \text{QED})_{\text{PS}}} - \text{LO}_{\text{QCD}}^{(\text{QCD})_{\text{PS}}}$

$d\sigma/dm_{ee}$ [fb/GeV] $pp \rightarrow e^+ e^- \mu^+ \nu_\mu$ @LHC 13 TeV



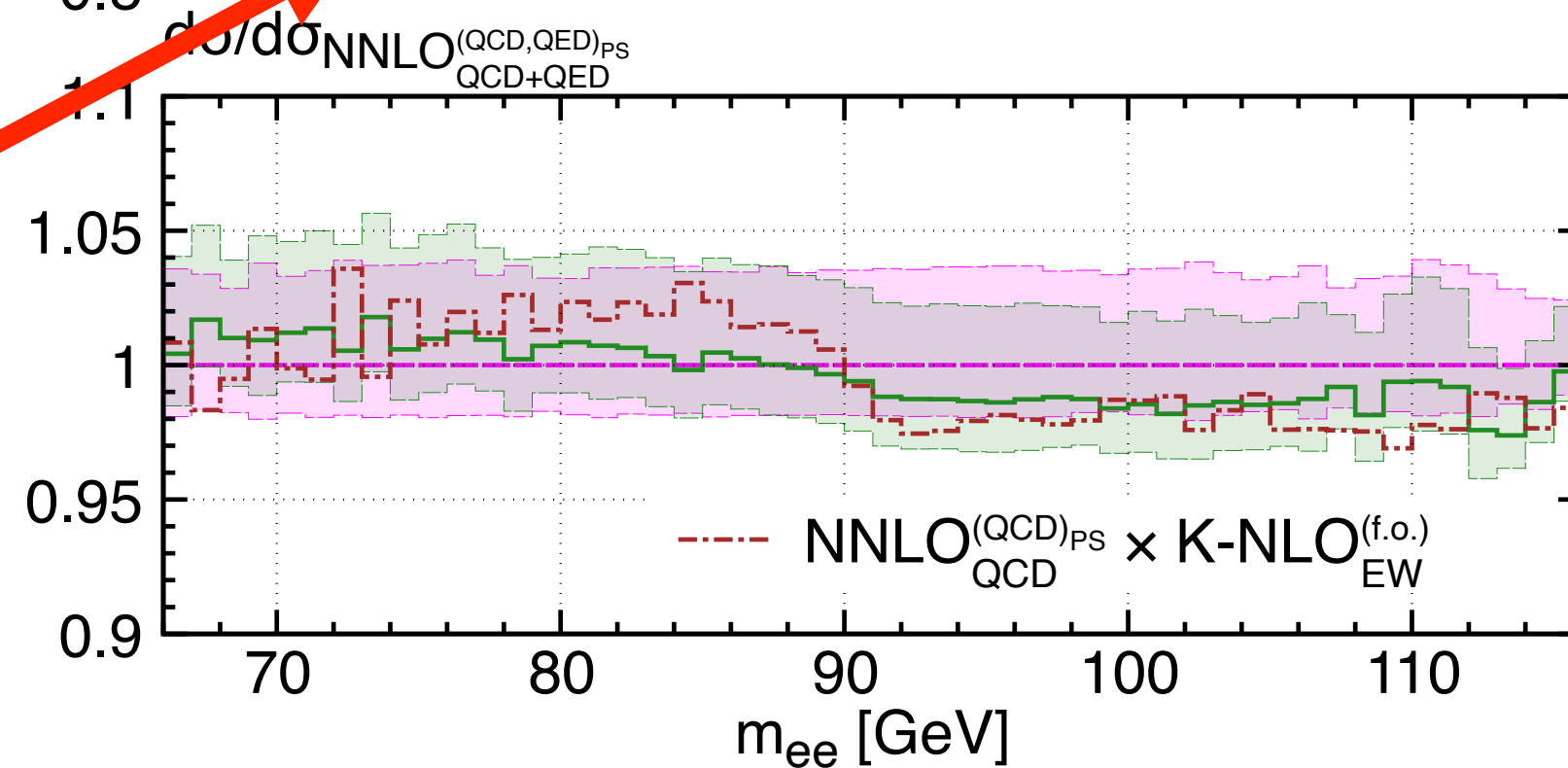
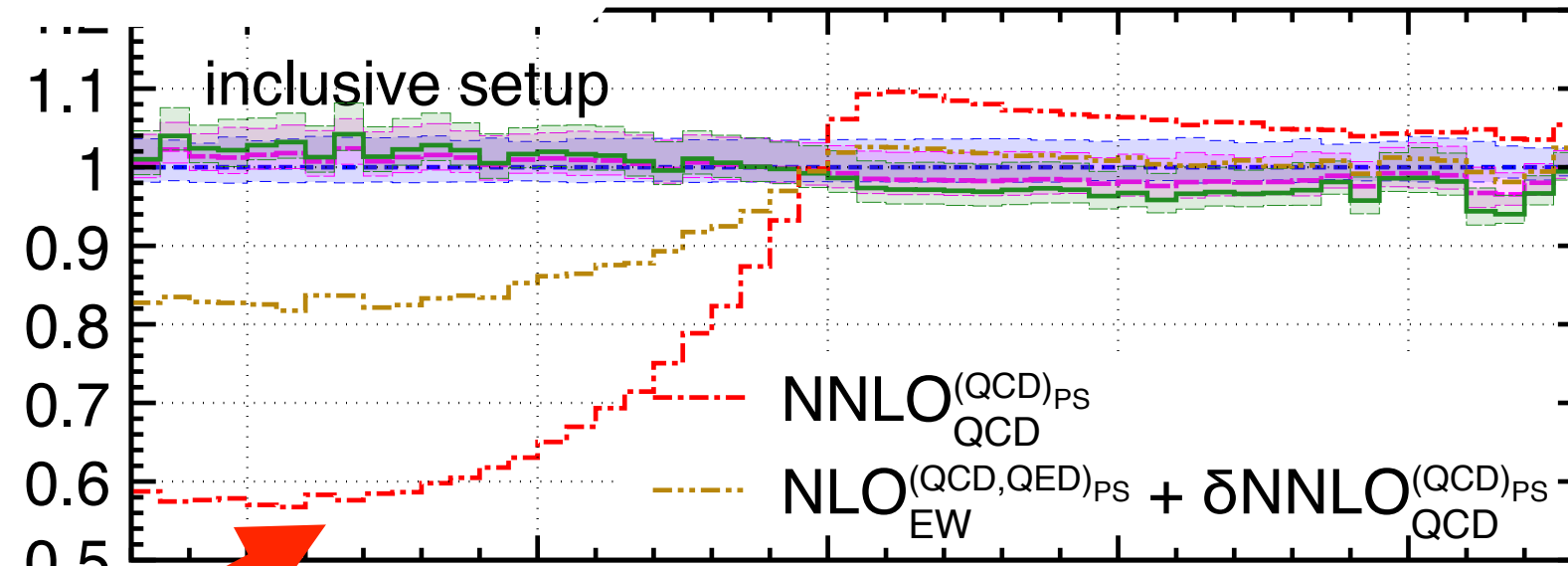
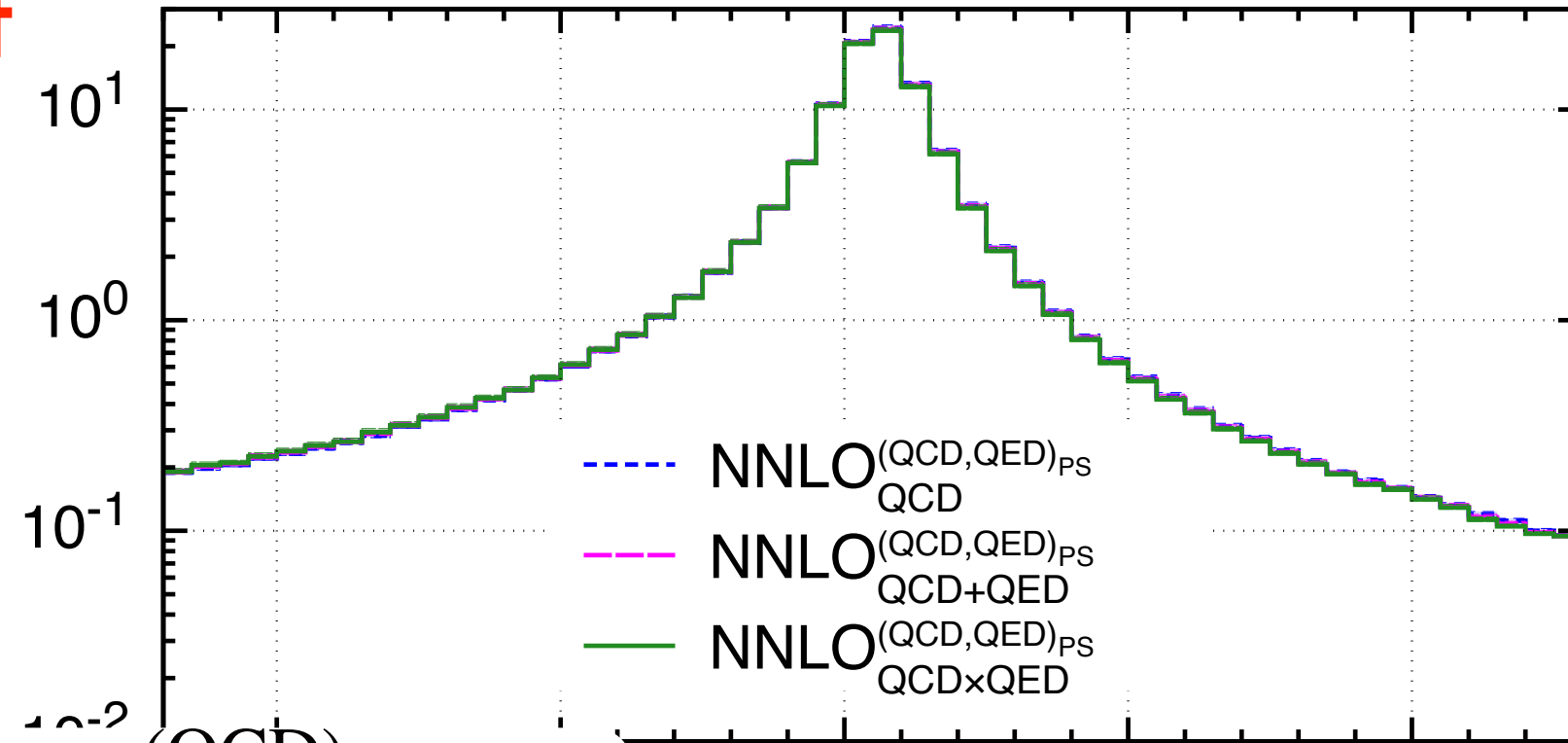
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$d\sigma/dm_{ee}$ [fb/GeV] $pp \rightarrow e^+ e^- \mu^+ \nu_\mu$ @LHC 13 TeV



Large effects from collinear QED radiations (~40%), which are absent in ●

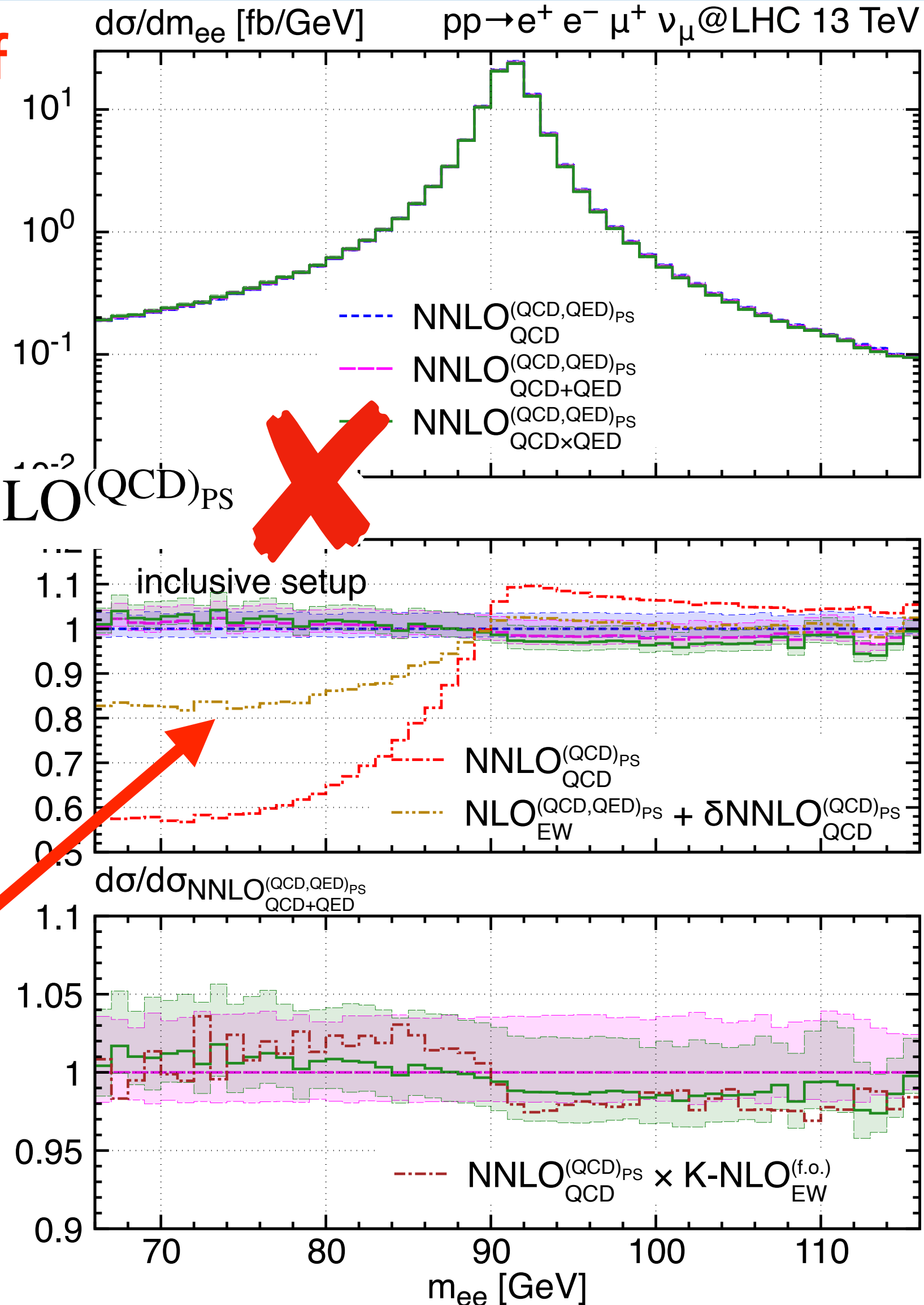
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Invariant mass of the Z boson - inclusive setup

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● misses important QED-QCD effects originating from QED emissions on top of the NNLO calculation —> **DISCARDED**



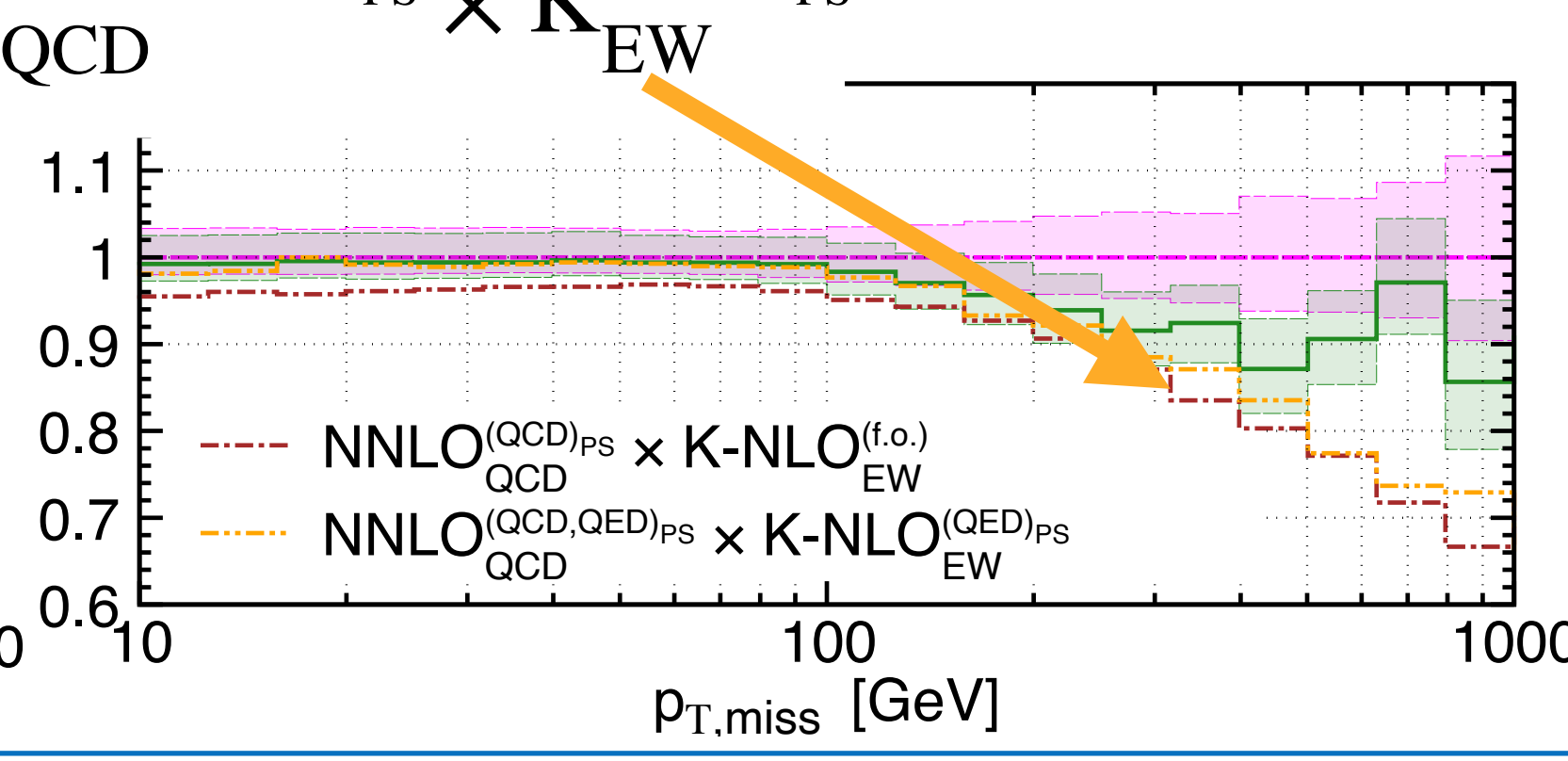
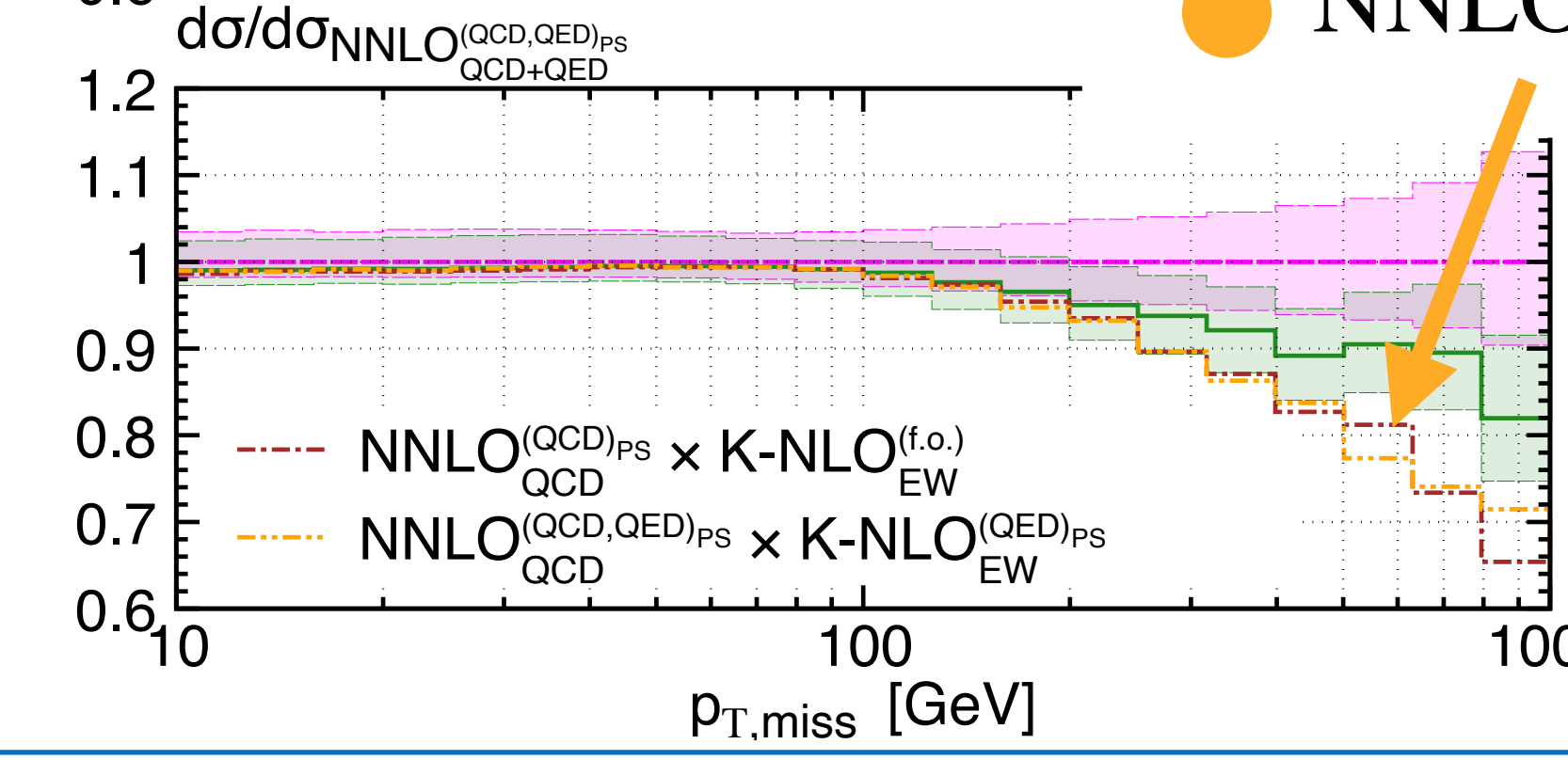
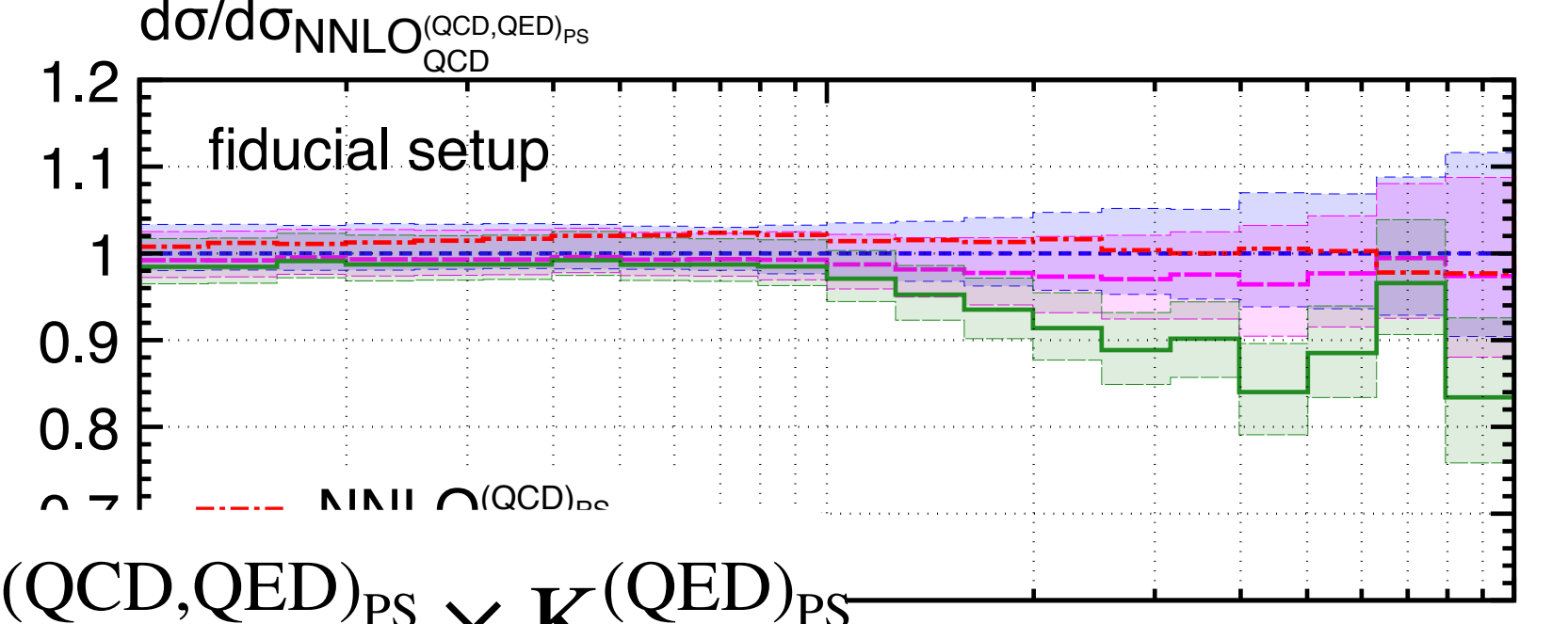
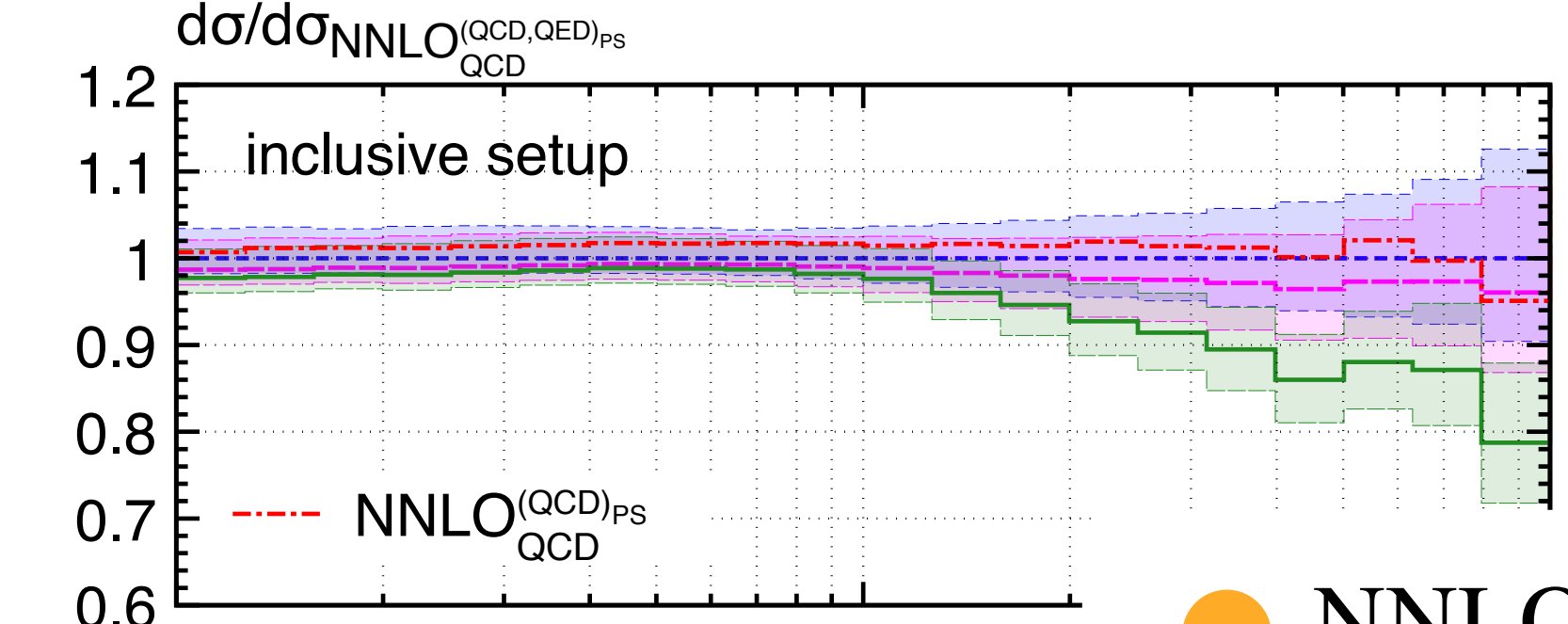
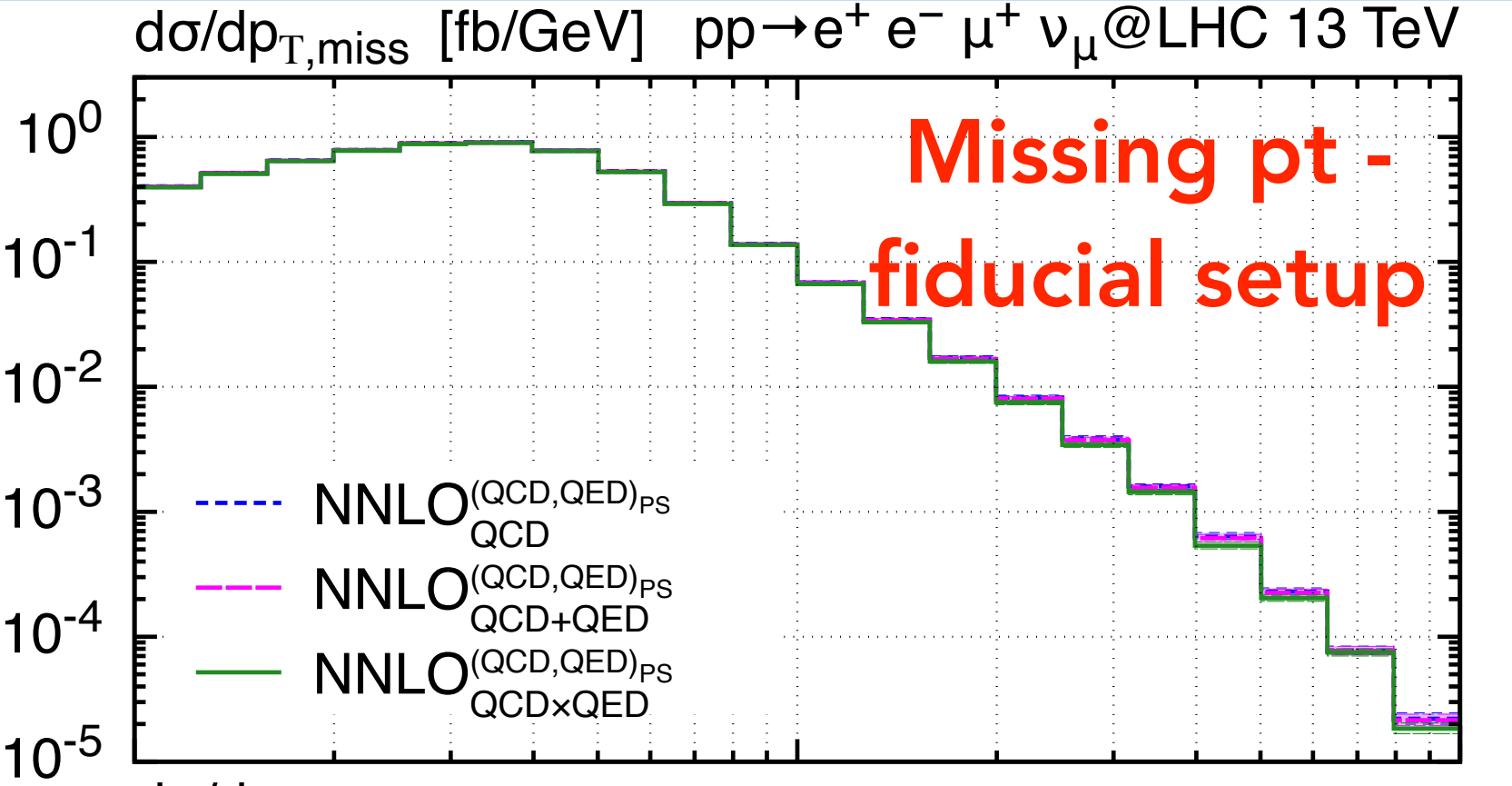
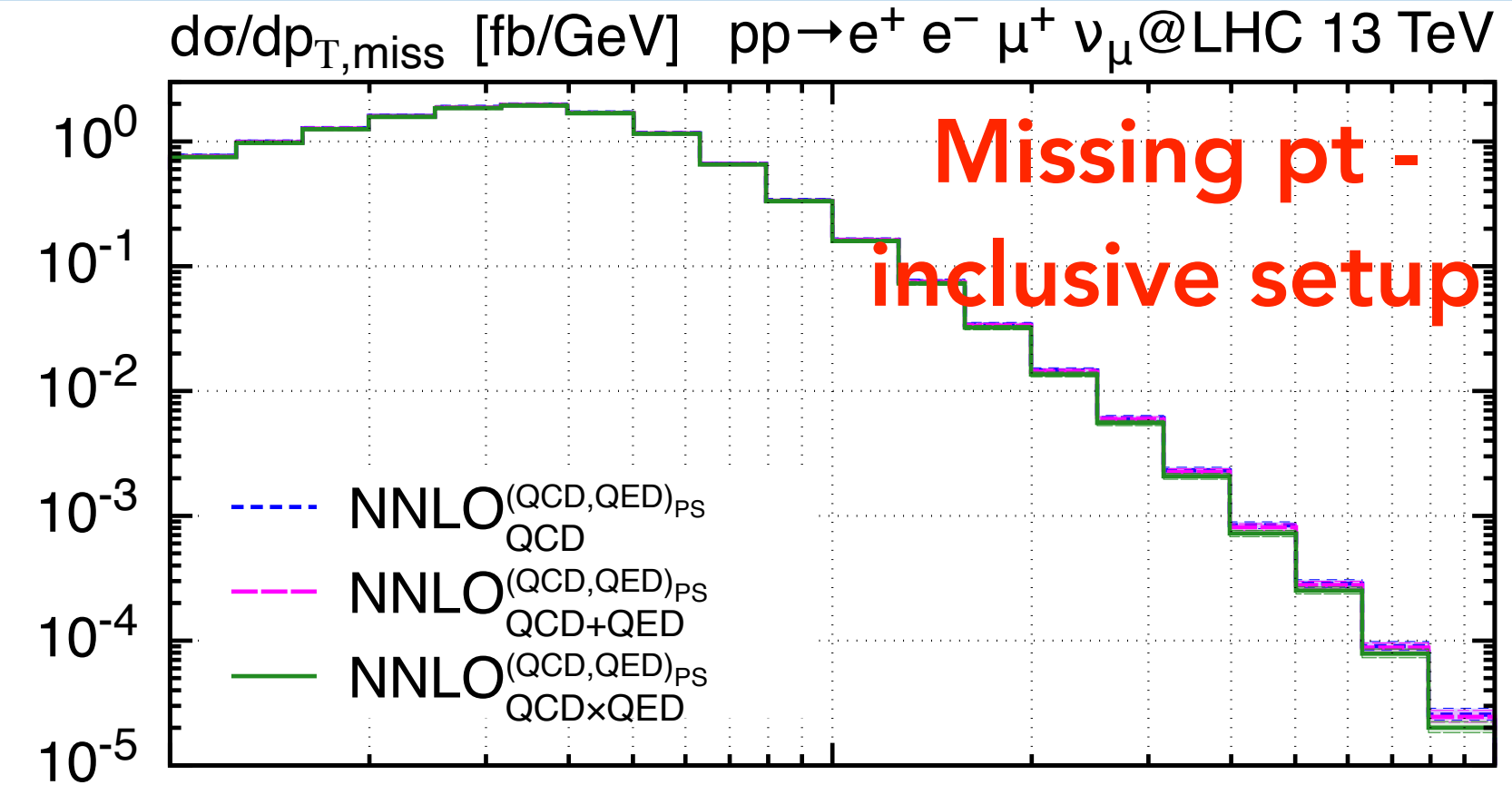
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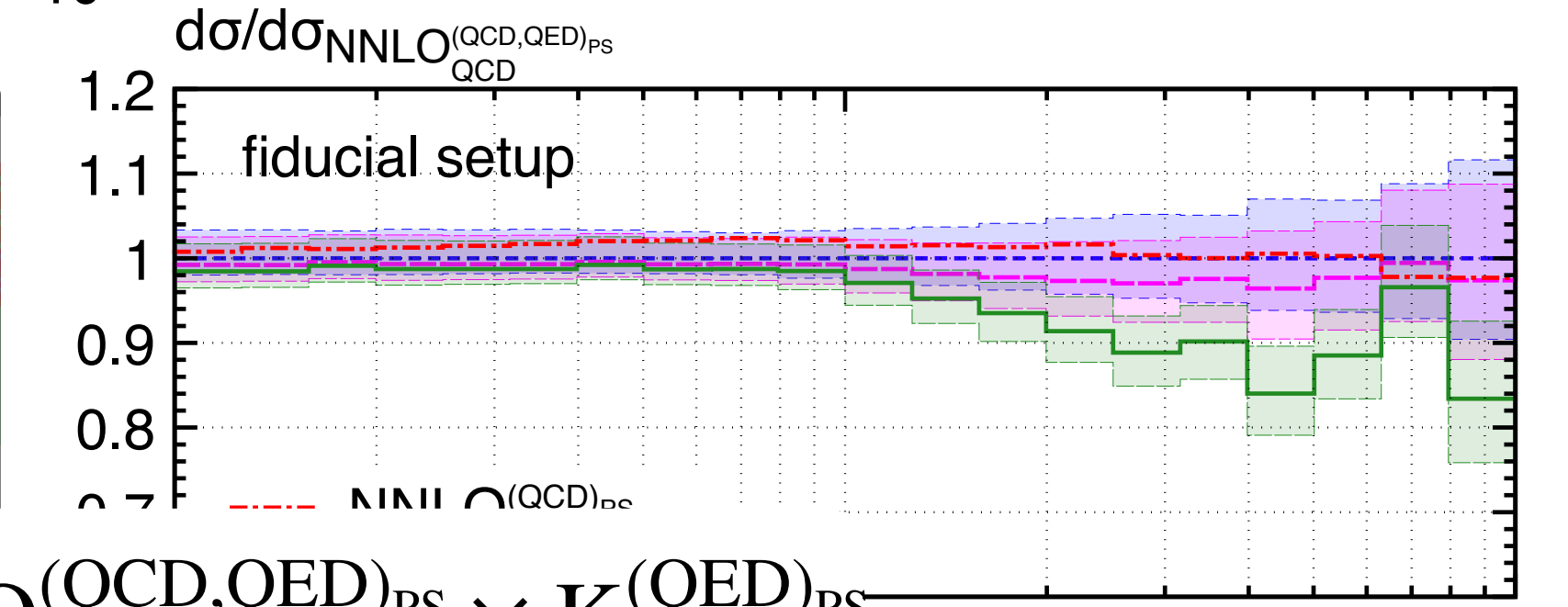
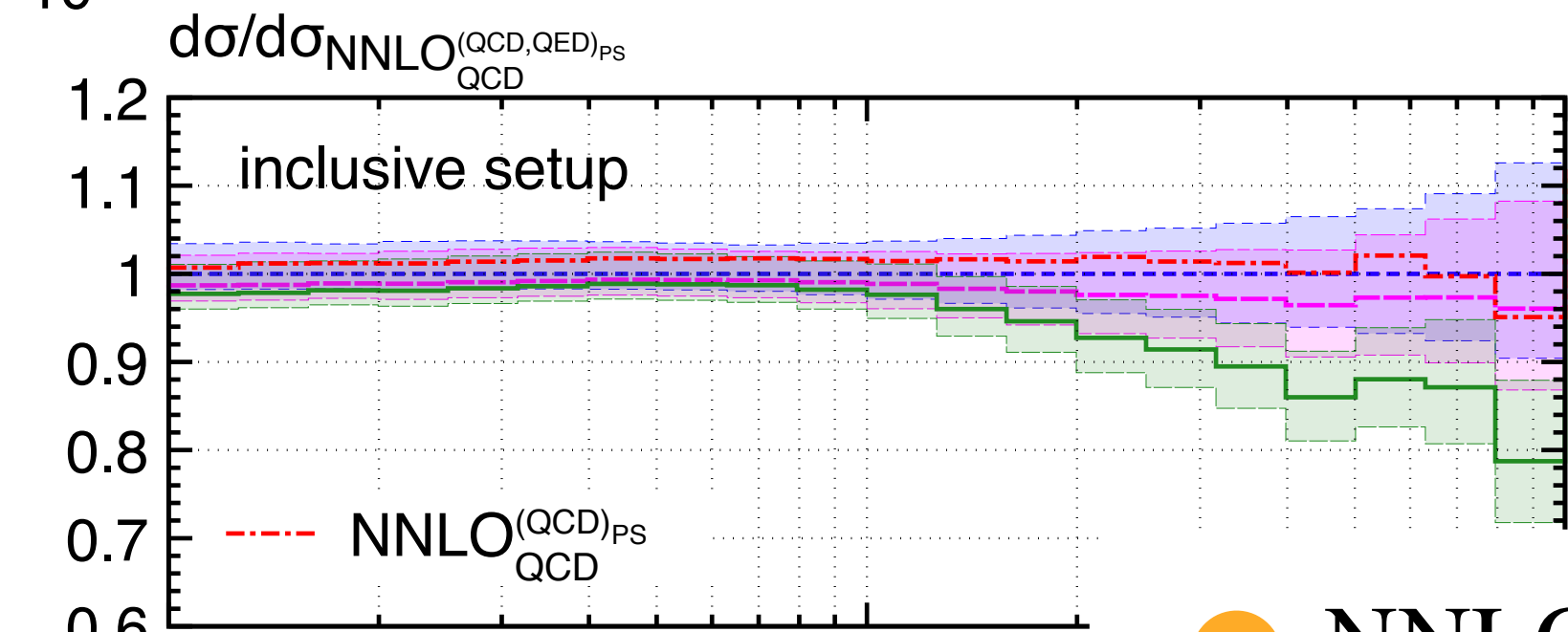
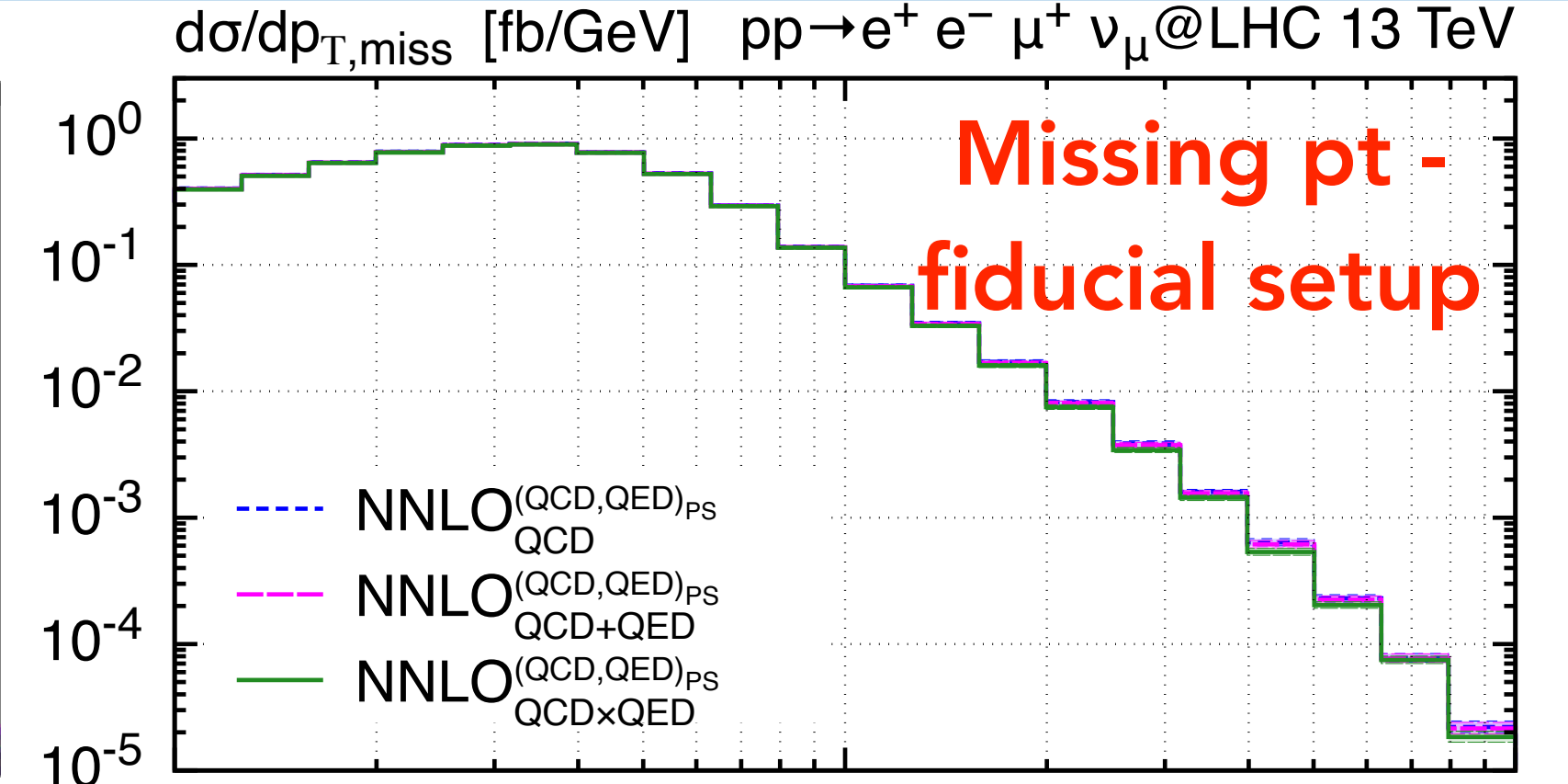
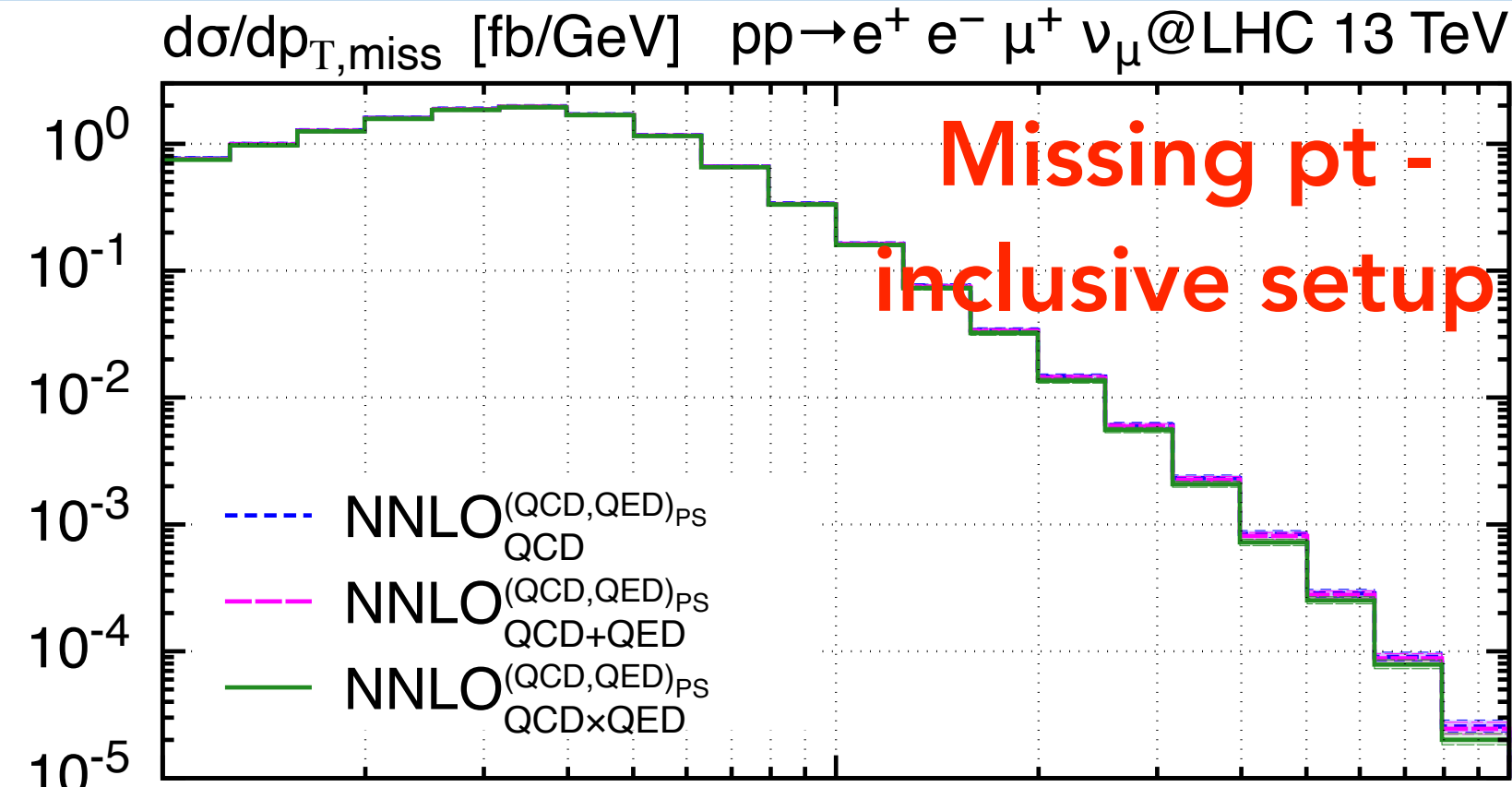




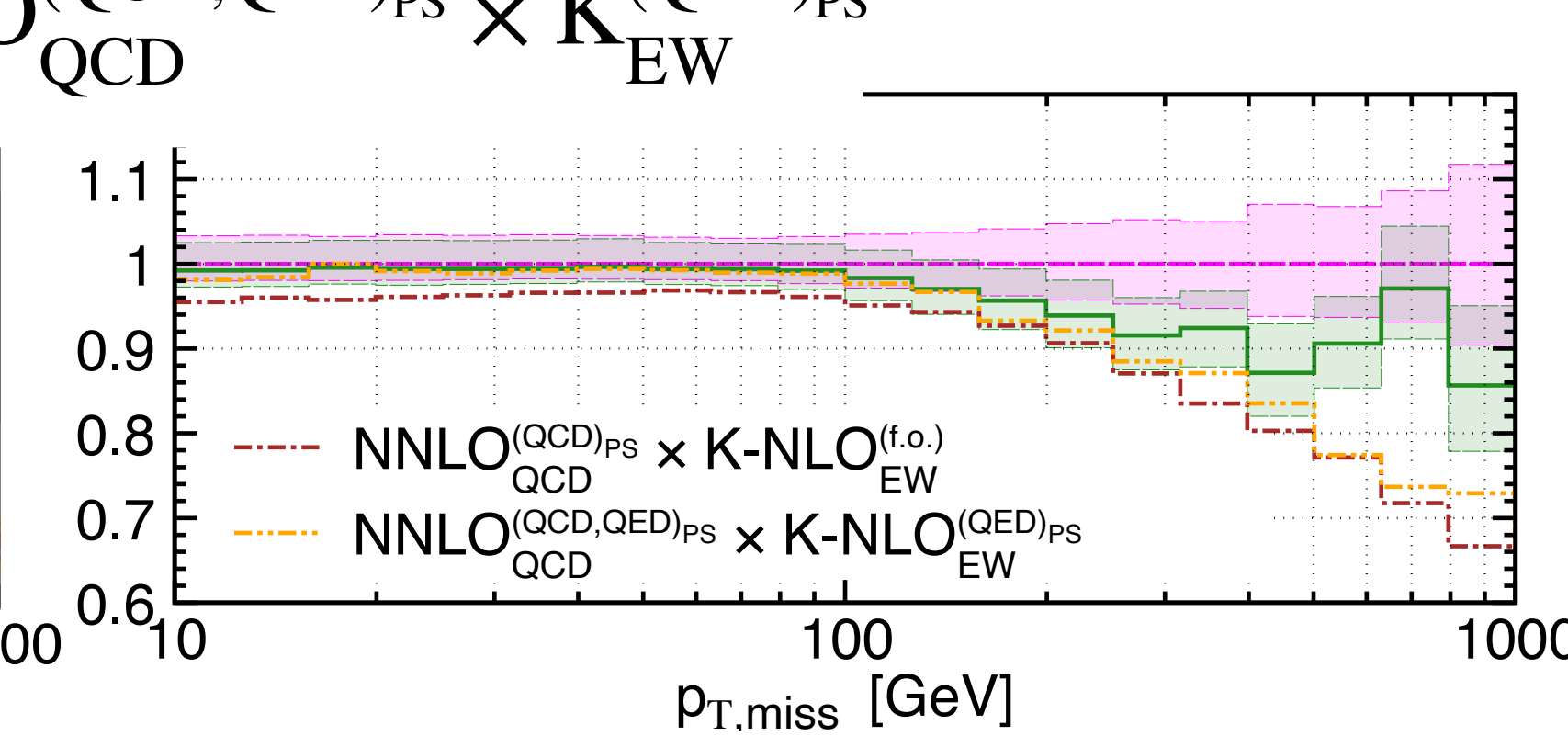
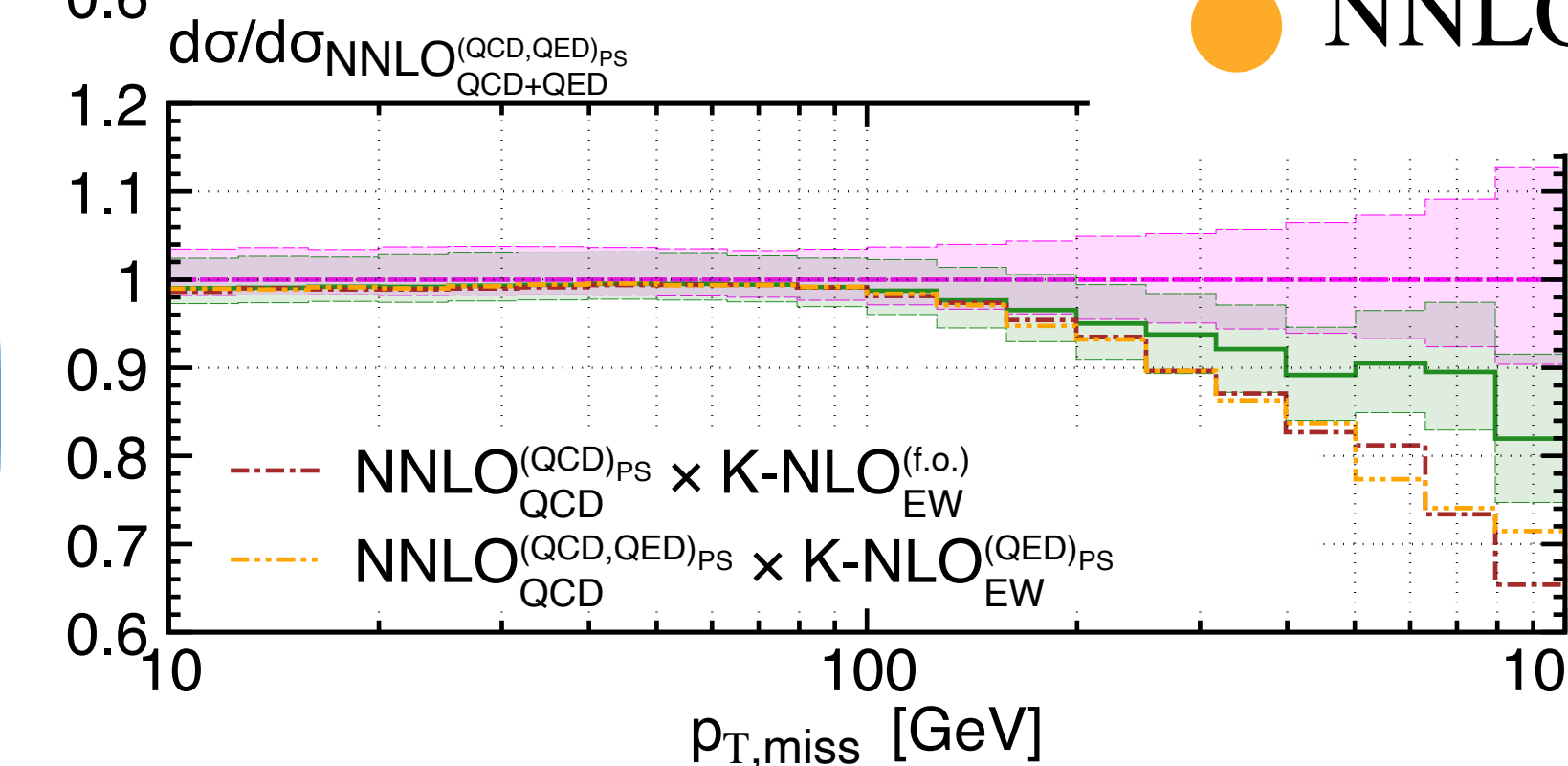
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● is affected by giant K-factors



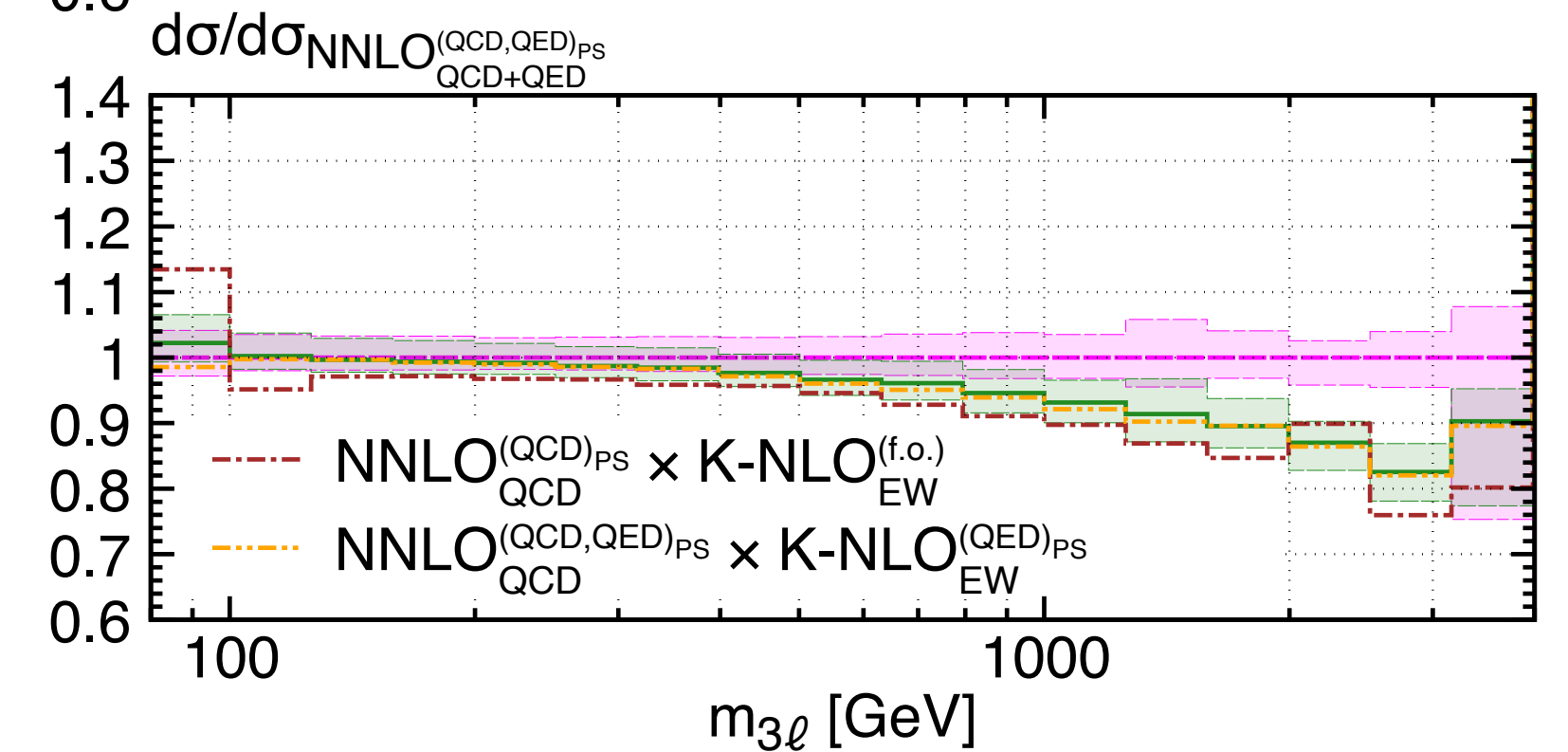
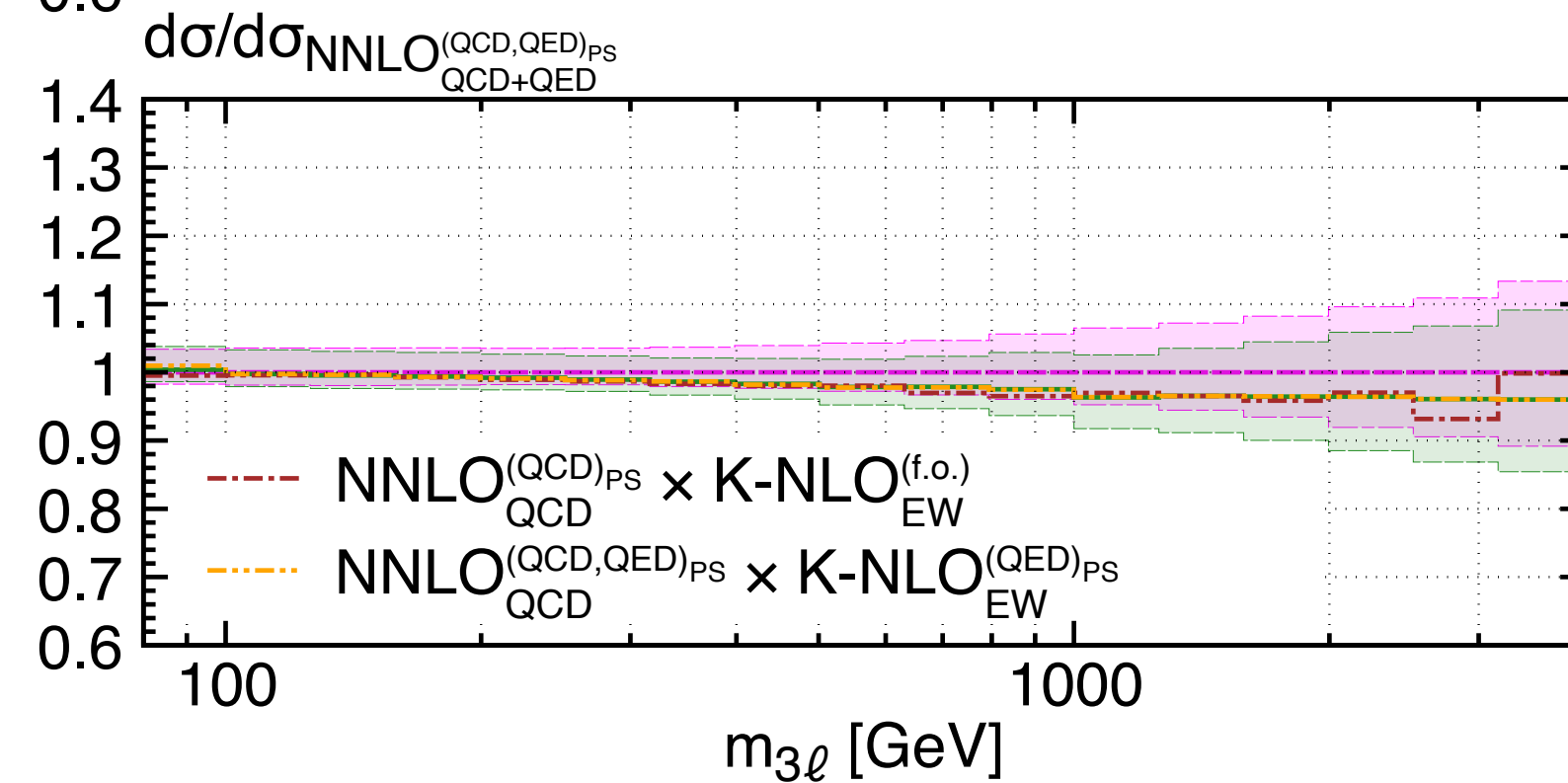
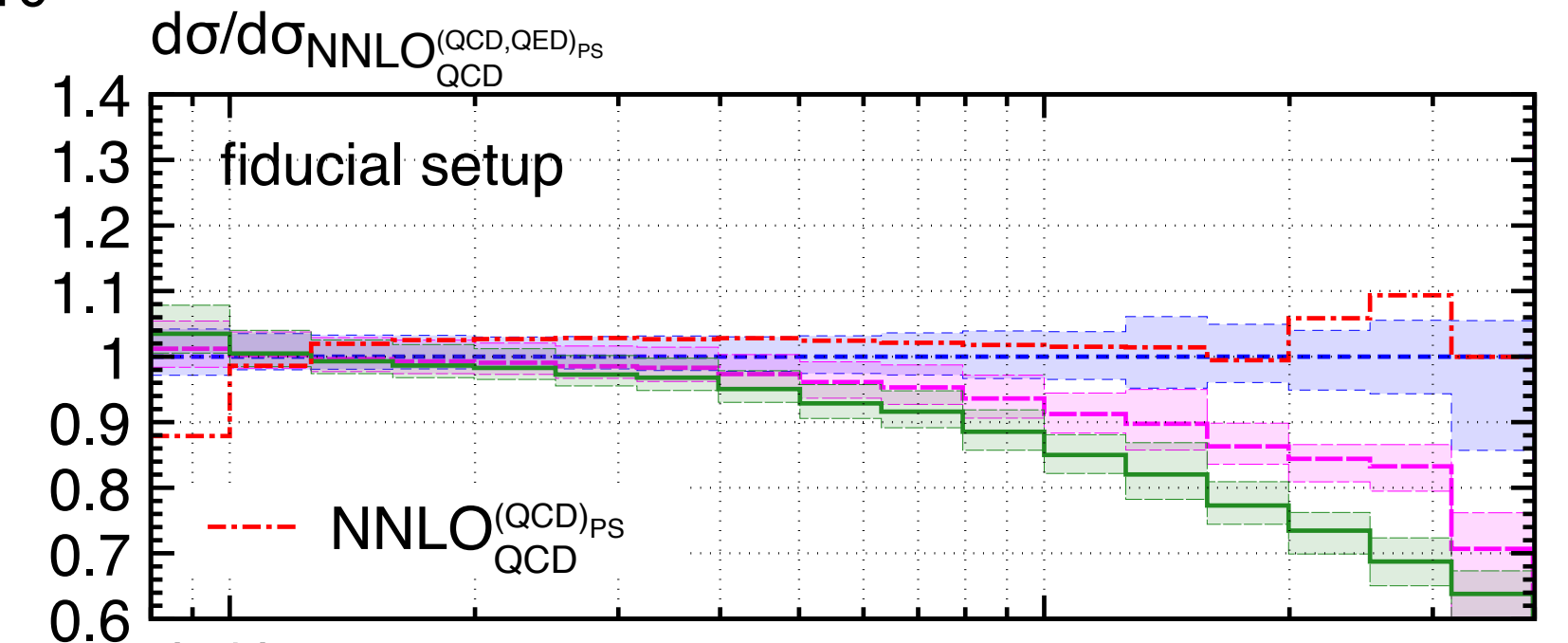
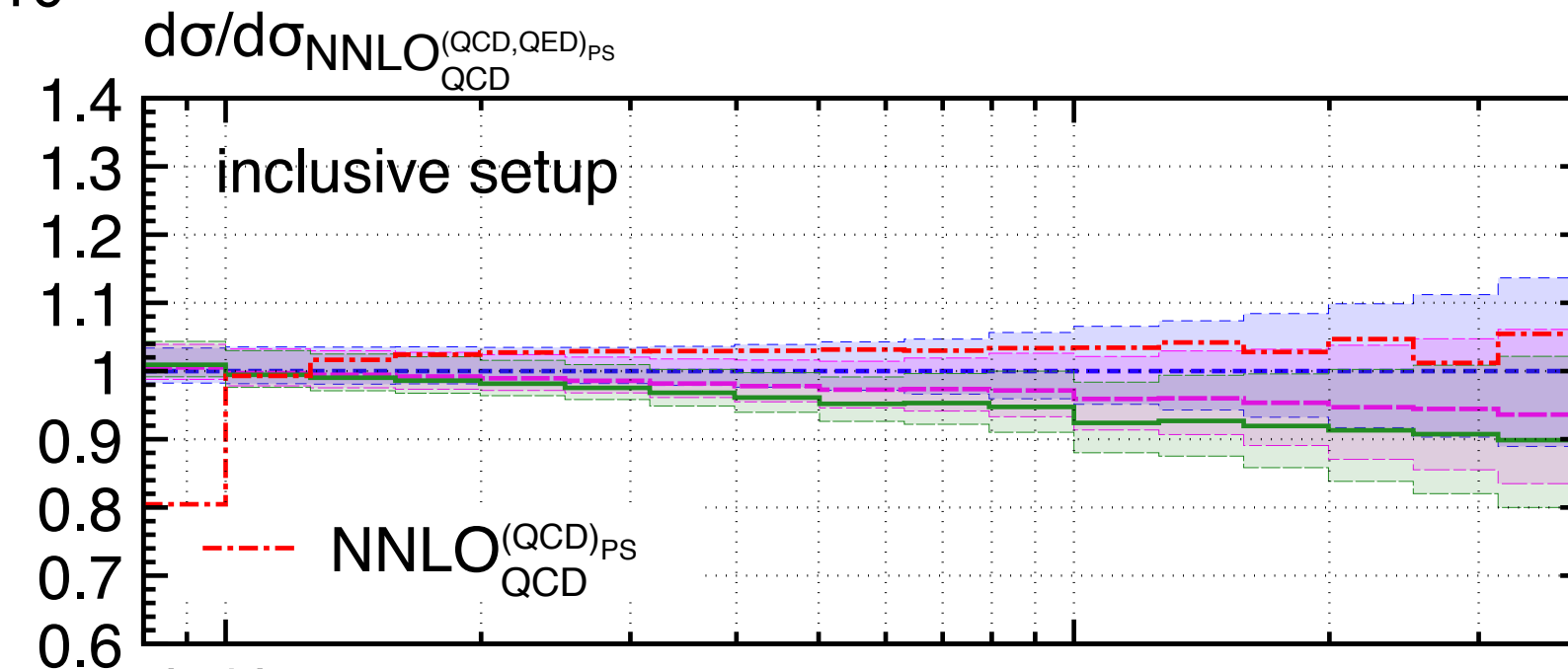
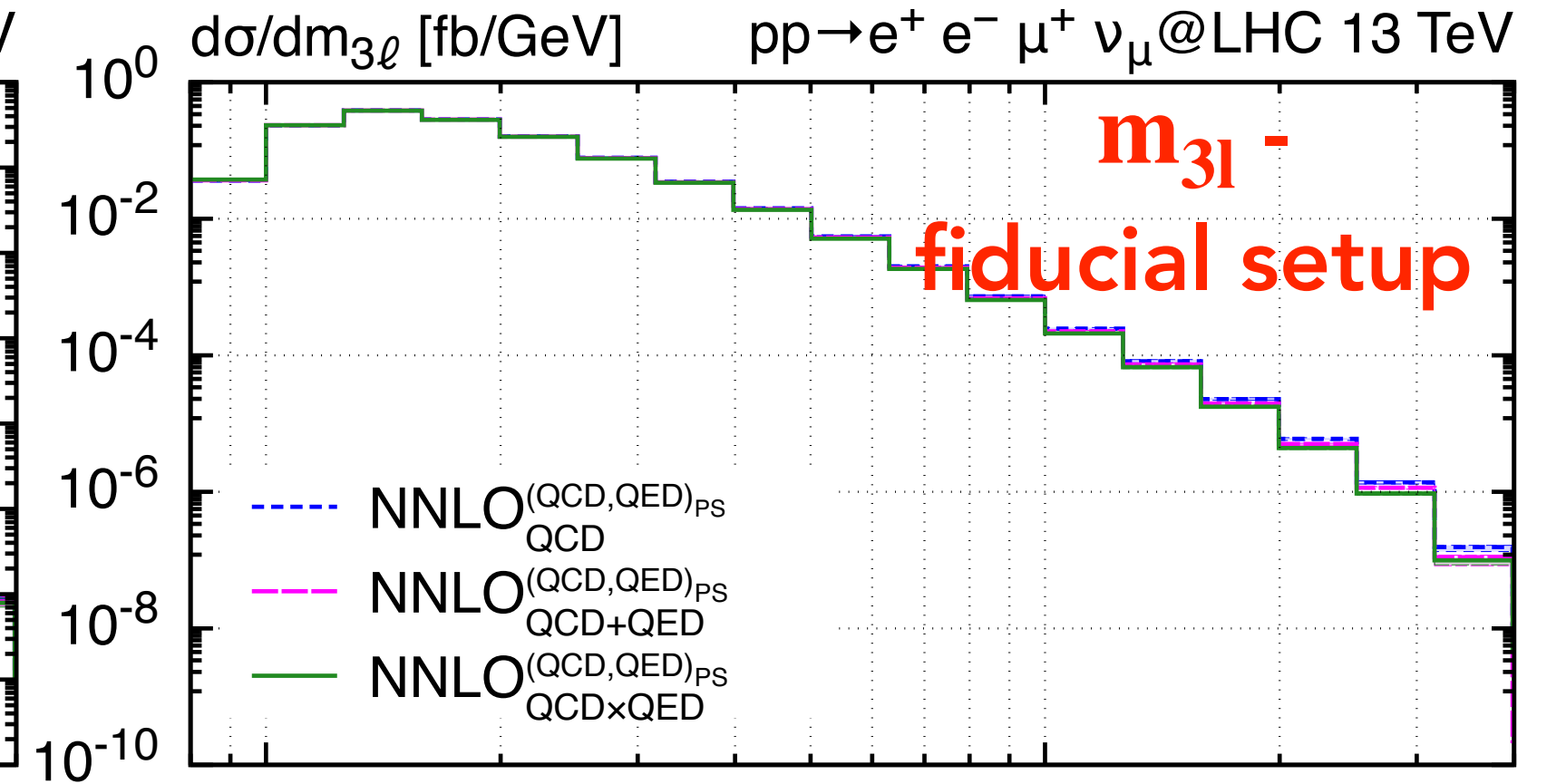
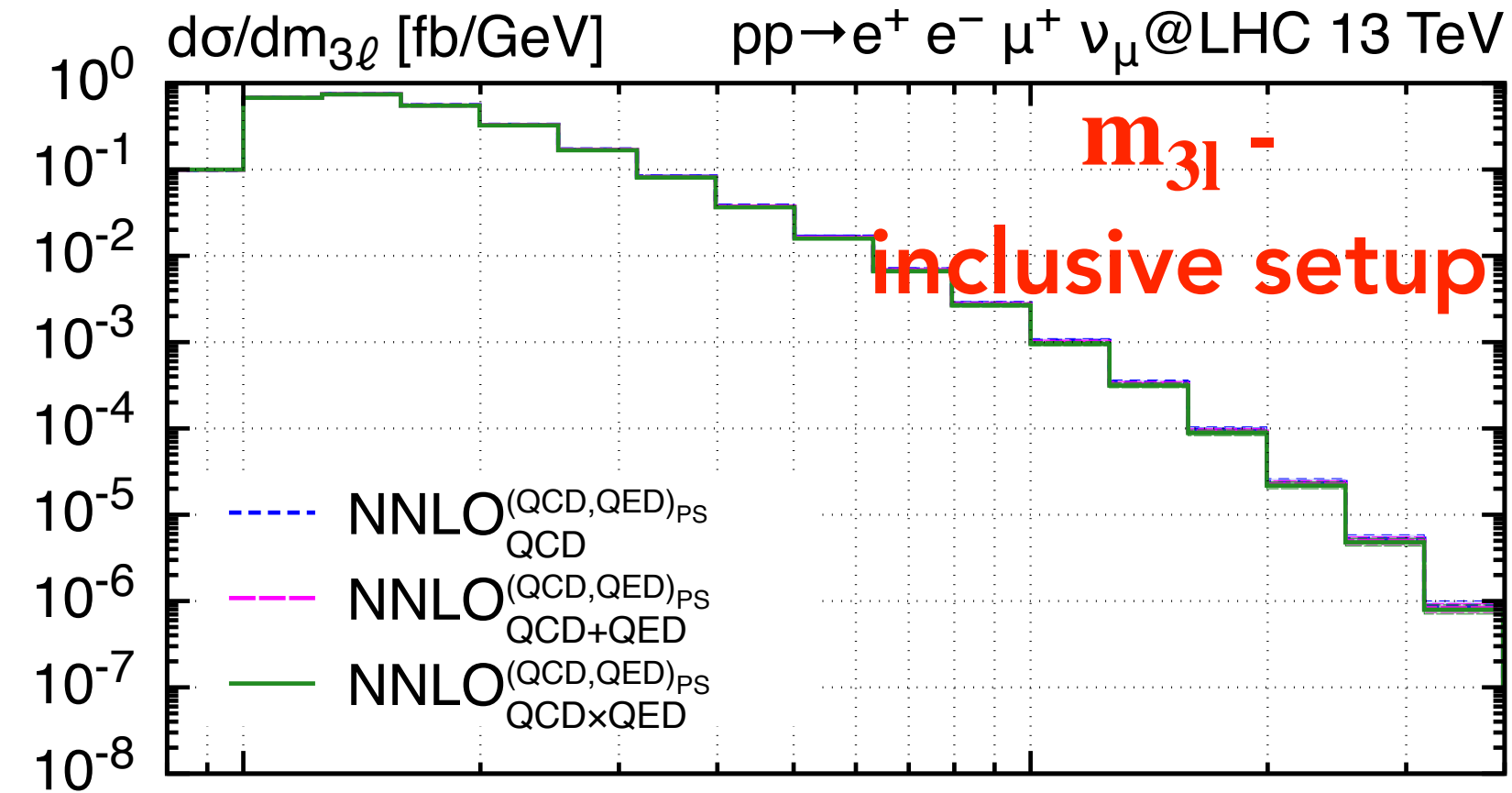
● $\text{NNLO}_{\text{QCD}}^{(\text{QCD}, \text{QED})_{\text{PS}}} \times K_{\text{EW}}^{(\text{QED})_{\text{PS}}}$





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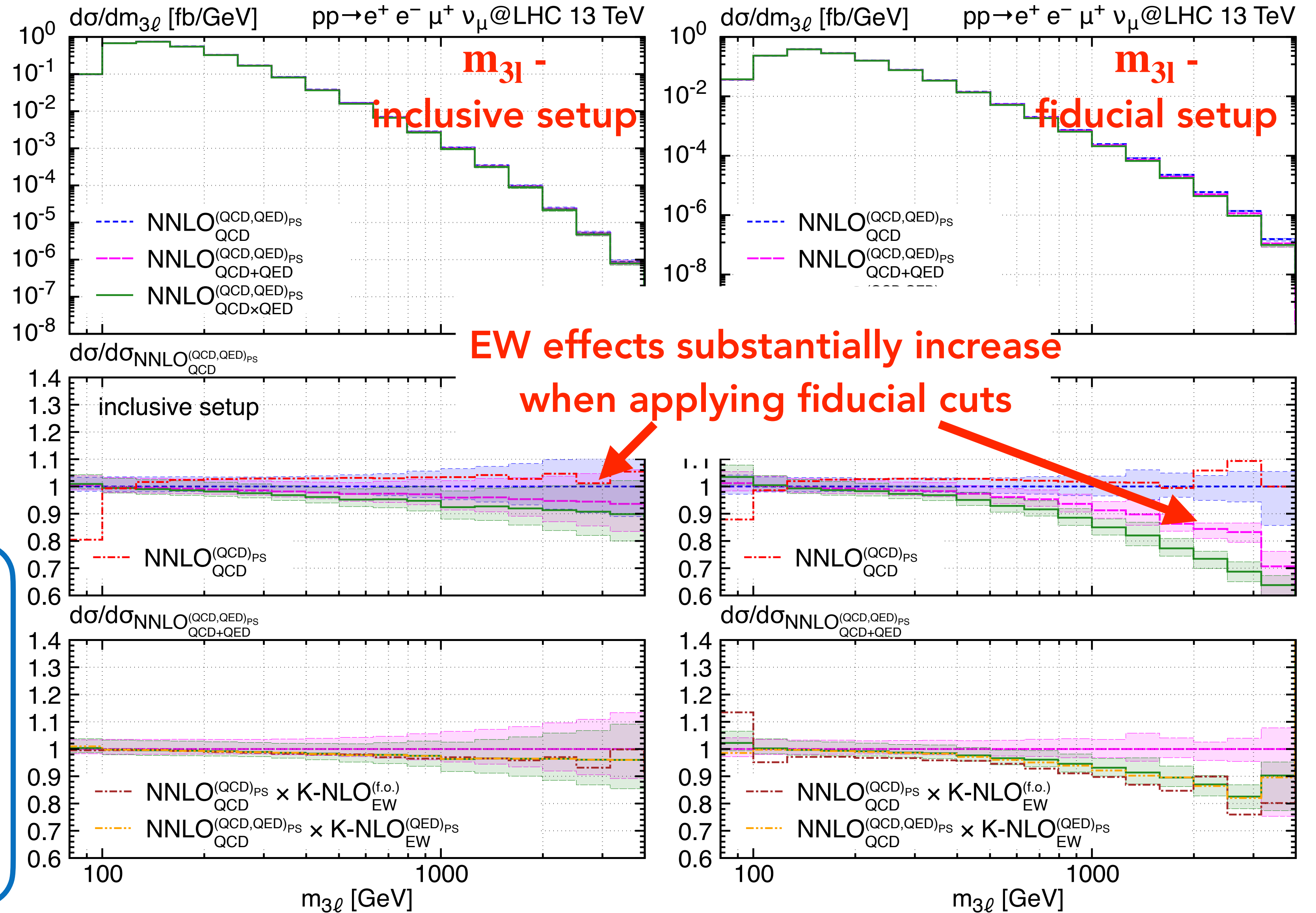




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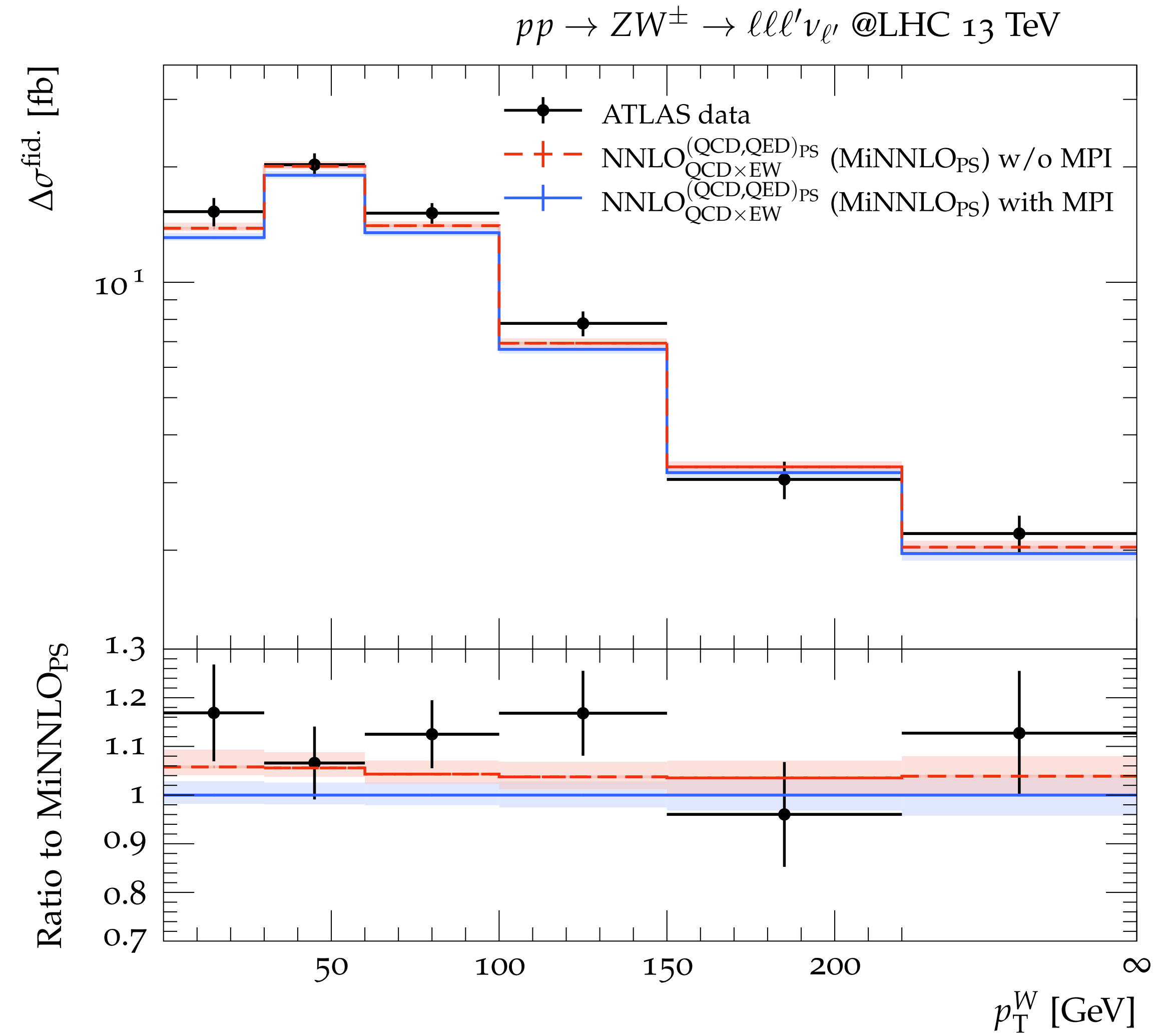
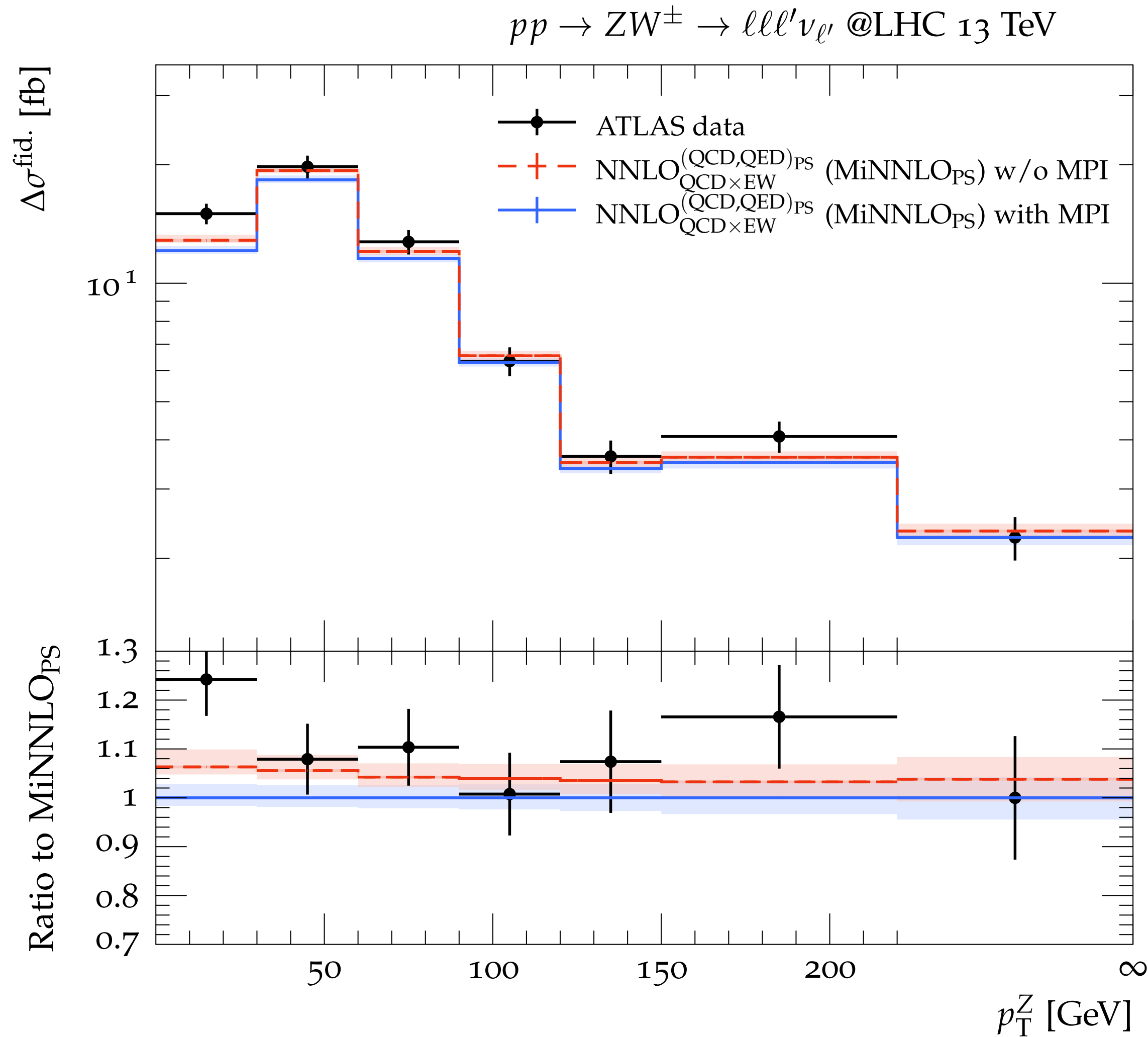
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In the inclusive case, Sudakov-logs are suppressed because not all the Mandelstam invariants are large in the very forward regime. These regions are removed when applying fiducial cuts.



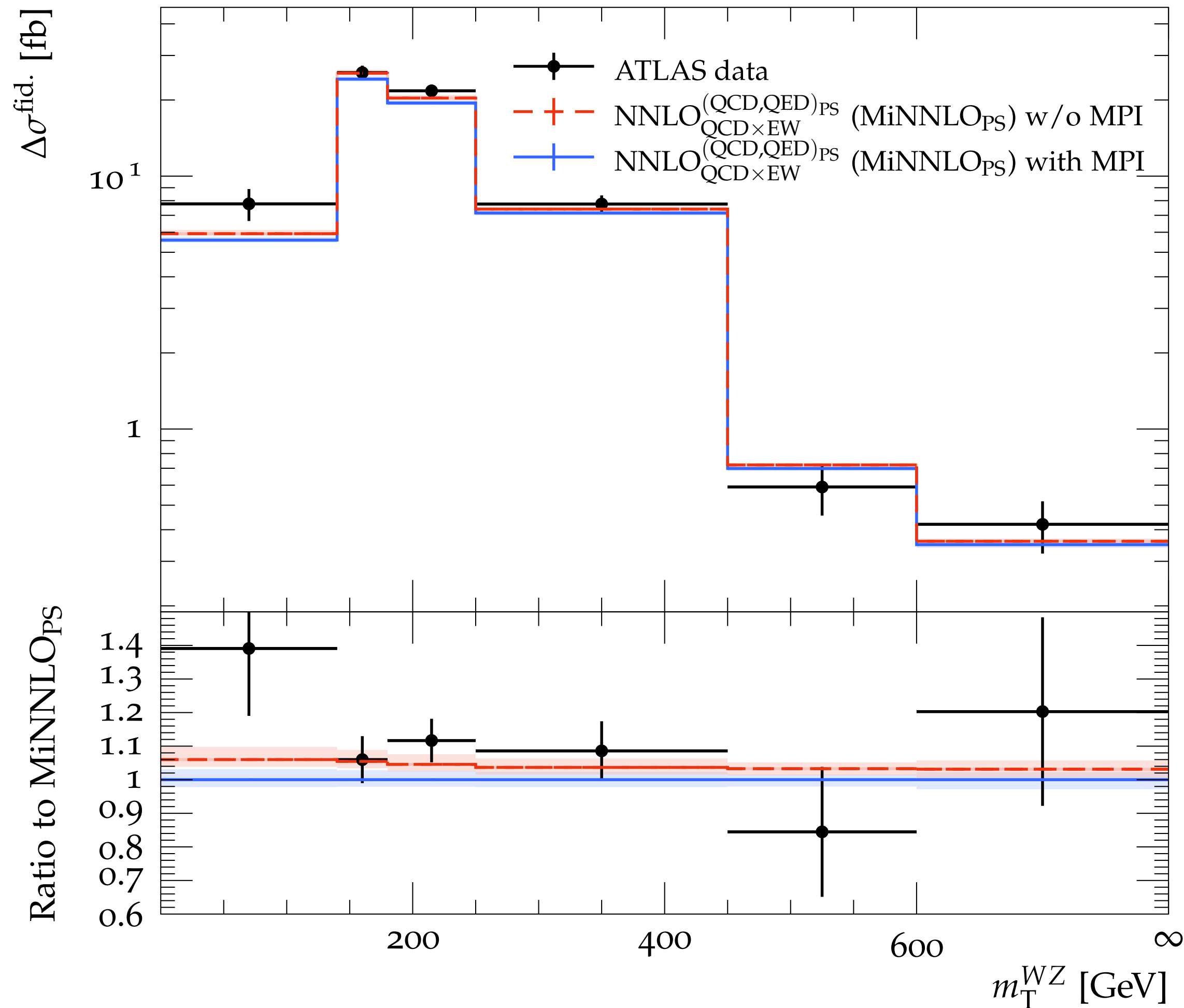


ATLAS data from Eur. Phys. J. C 79 (2019)

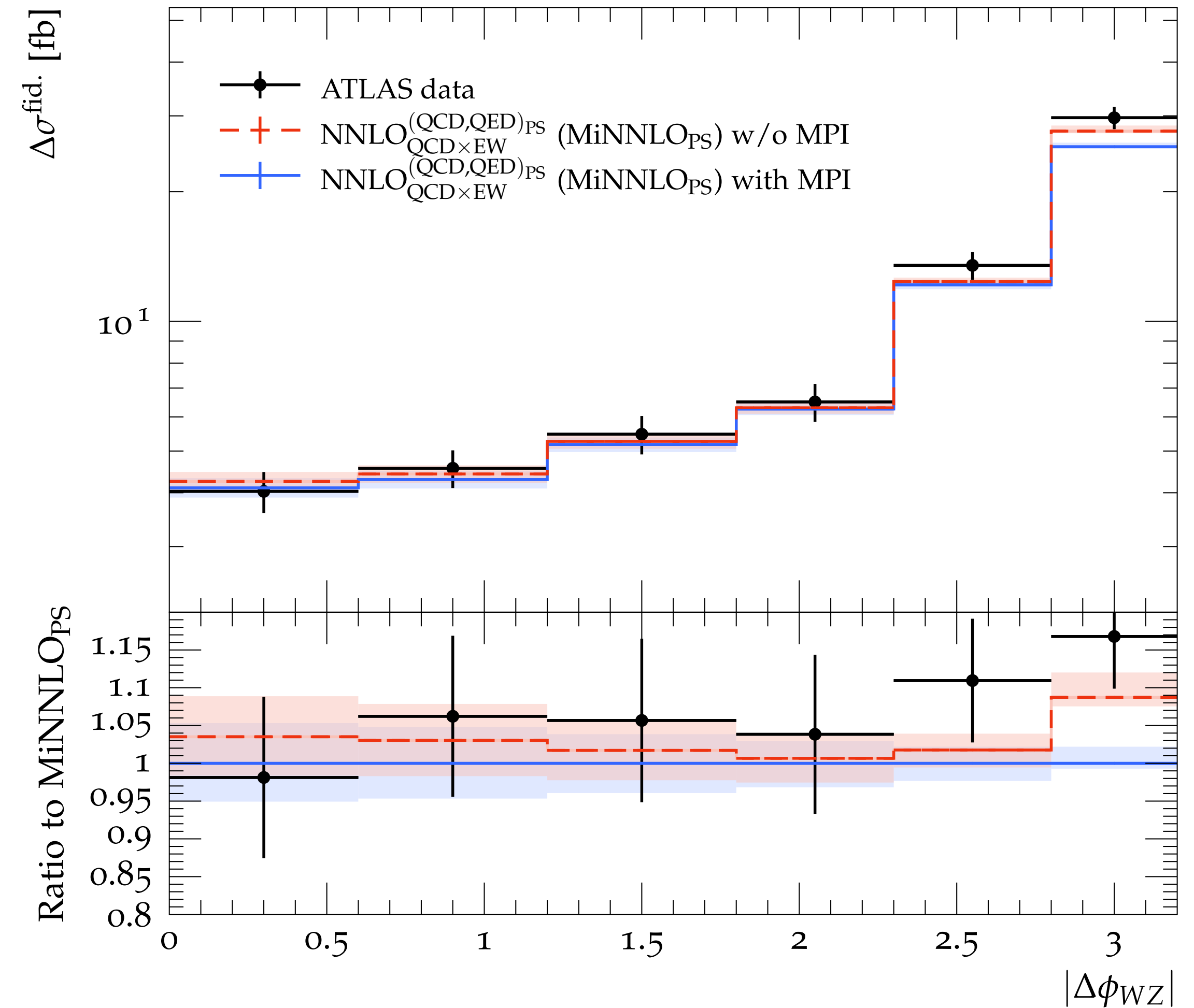


ATLAS data from Eur. Phys. J. C 79 (2019)

$pp \rightarrow ZW^\pm \rightarrow lll'v_{\ell'} @LHC 13 \text{ TeV}$



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- **NNLO+PS (QCD)** predictions are strongly needed for a realistic description of LHC events.
- The **MiNNLO_{PS} method** is a powerful tool for reaching this accuracy.
- In the context of precision physics, the inclusion of **NLO EW corrections** on top of the NNLO calculations is particularly important.
- I showed and discuss results for **WZ production** at NNLO (QCD) and NLO (EW) accuracy matched to parton showers for 13 TeV LHC collisions.
- The natural next step is the implementation of the combined generation of NNLO QCD and NLO EW accurate events, rather than an a posteriori recombination.

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Thank you!