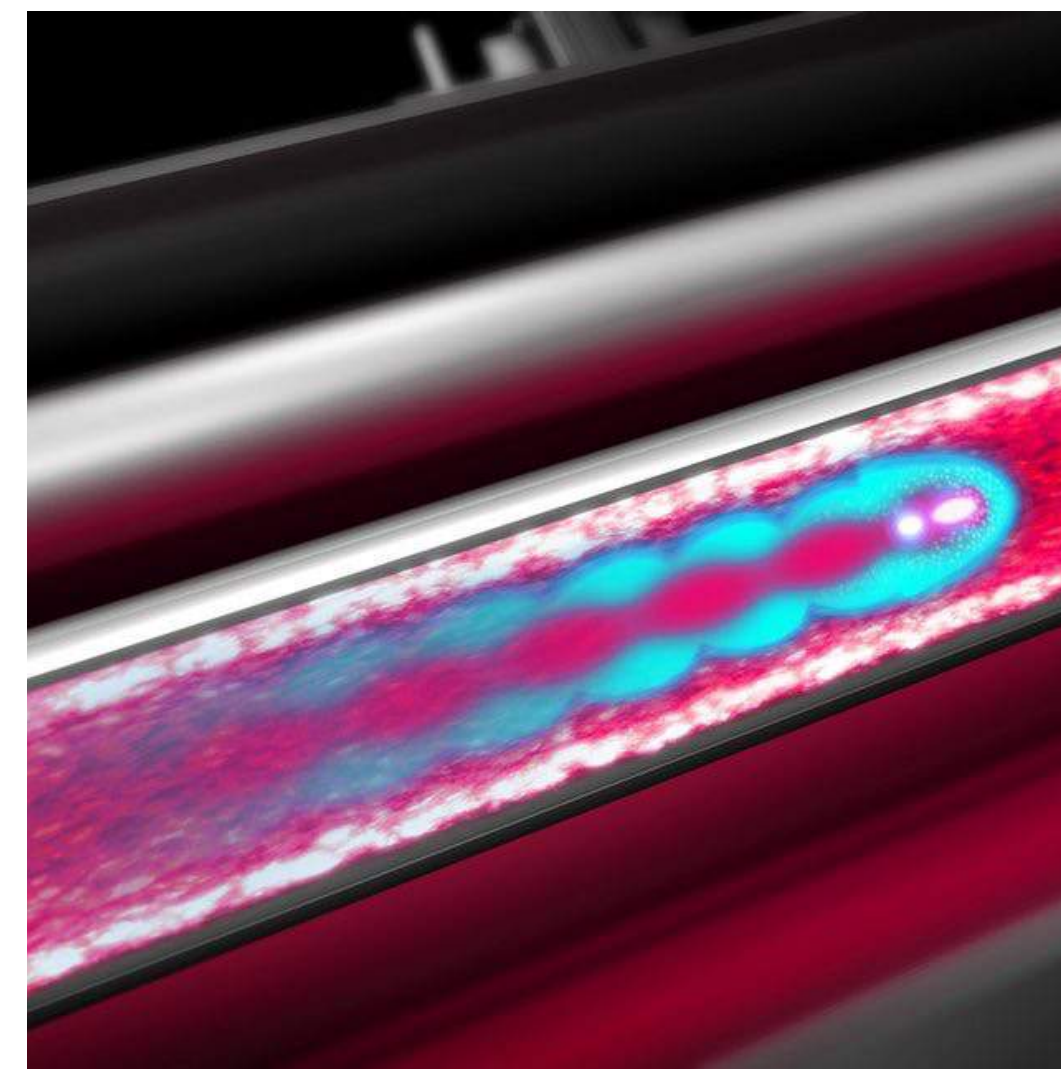
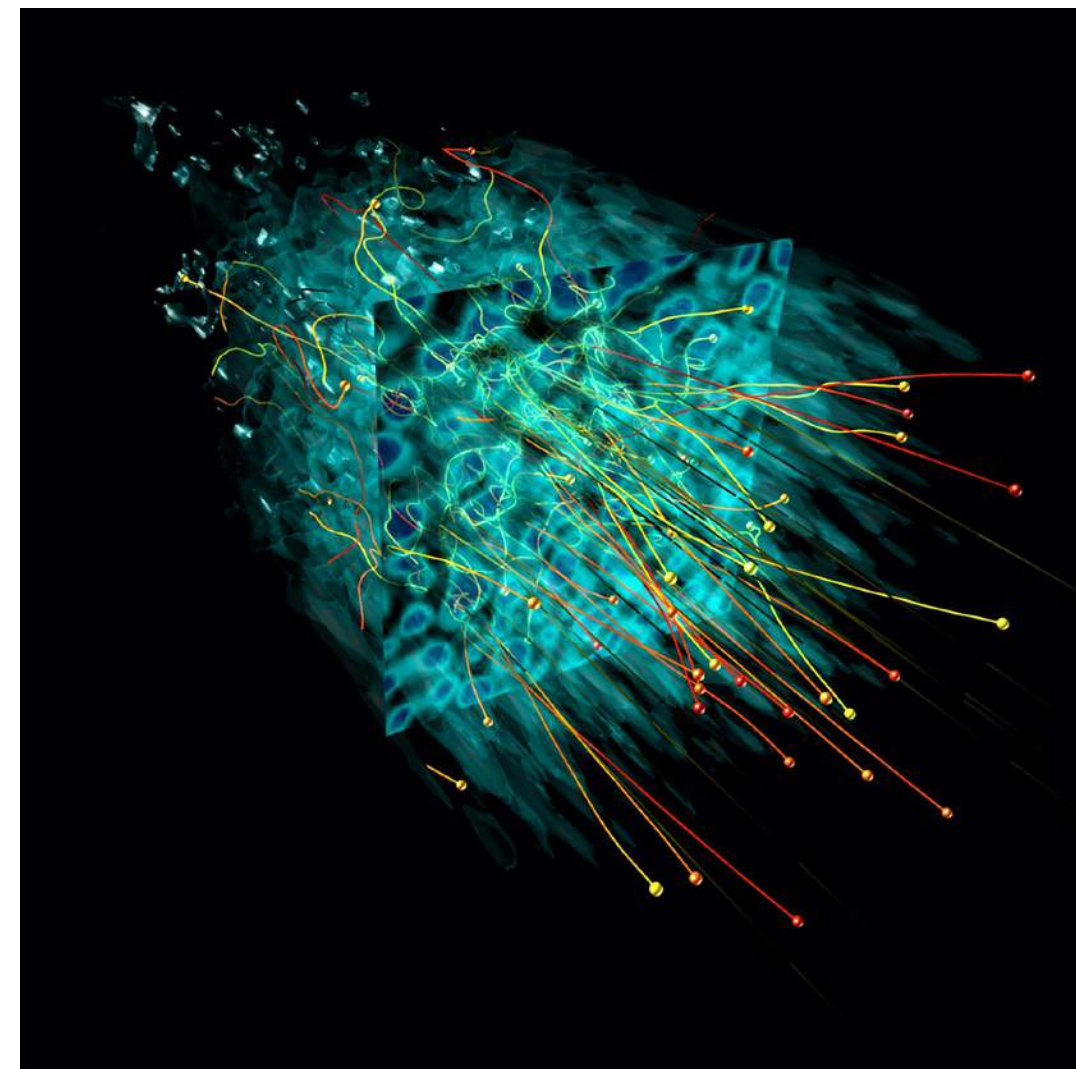


Cosmic-Ray Driven Instabilities

Frederico Fiuza

fiuza@slac.stanford.edu



Acknowledgements



SLAC/Stanford:

J. R. Peterson

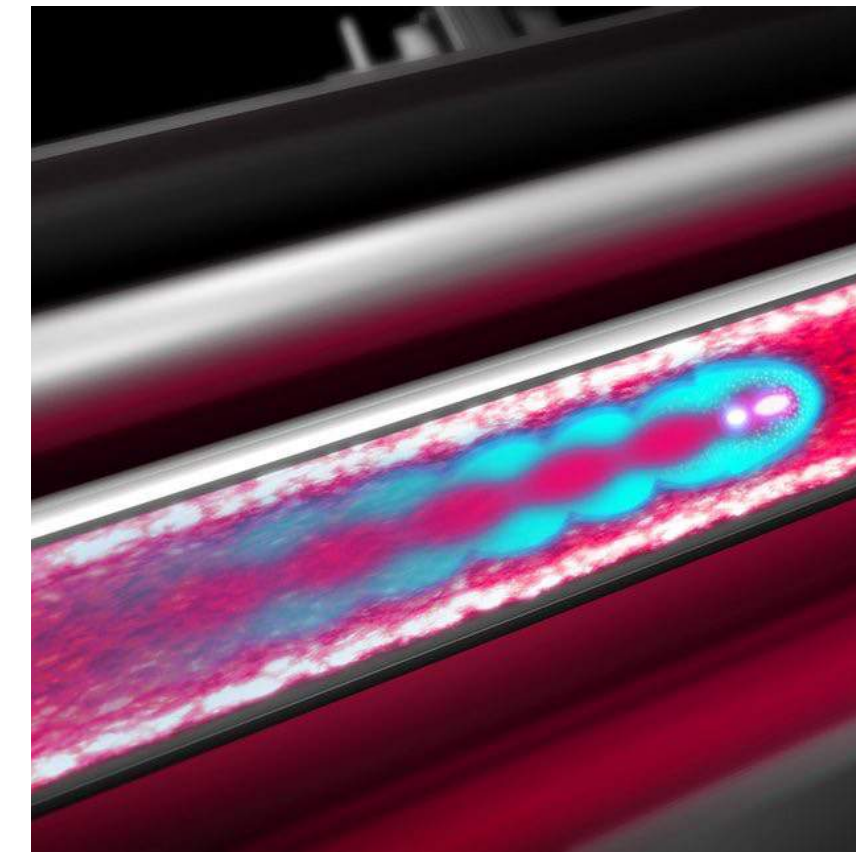
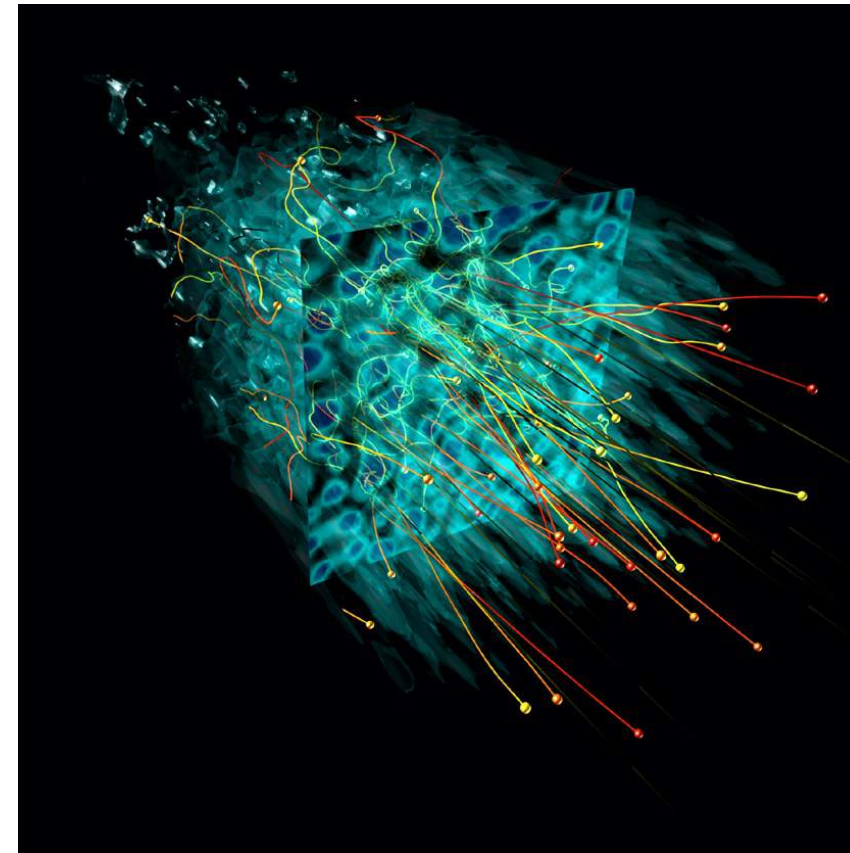
A. Vanthieghem

C. Ruyer

M. Hogan

V. Yakimenko

S. Glenzer



Collaborators:

P. Claveria, S. Corde (E. Polytechnique)

K. Marsh, C. Joshi (UCLA)

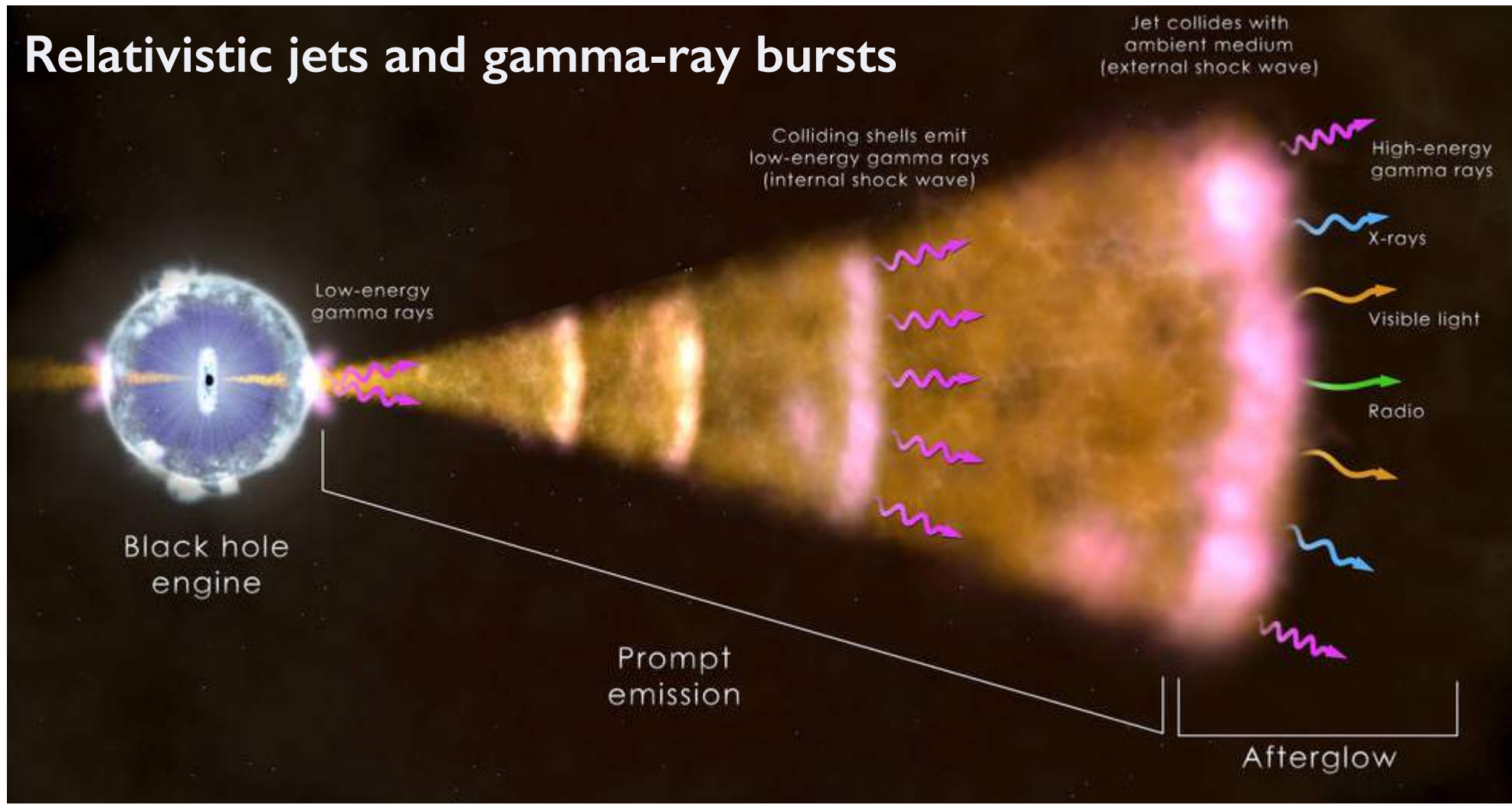
M. Tamburini (MPIK)

OSIRIS code provided by the OSIRIS Consortium (IST/UCLA)

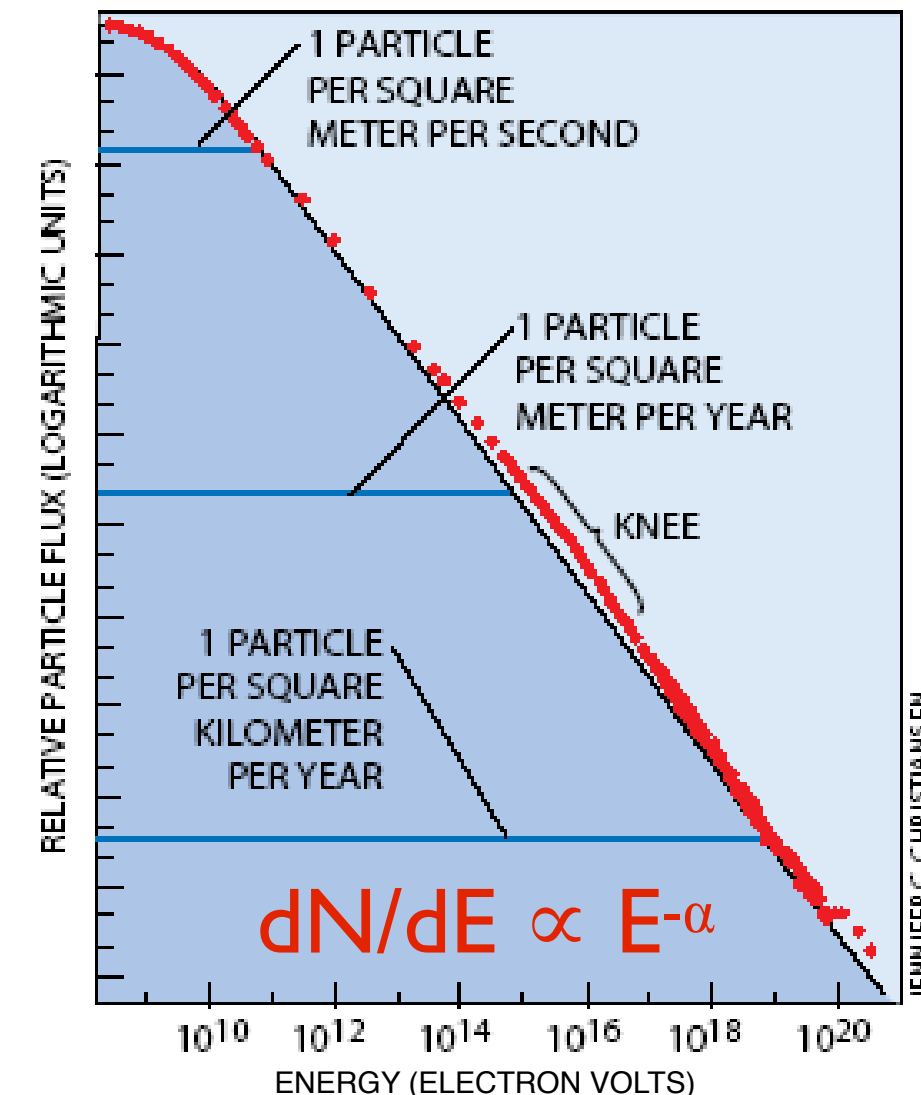
HPC resources at Quartz (LLNL), Theta (ANL), and Cori (NERSC)

Financial support from the DOE Early Career Research Program,
DOE Fusion Energy Sciences, and National Science Foundation

Astrophysical plasmas are extraordinary particle accelerators



Scientific American, (c) 1998

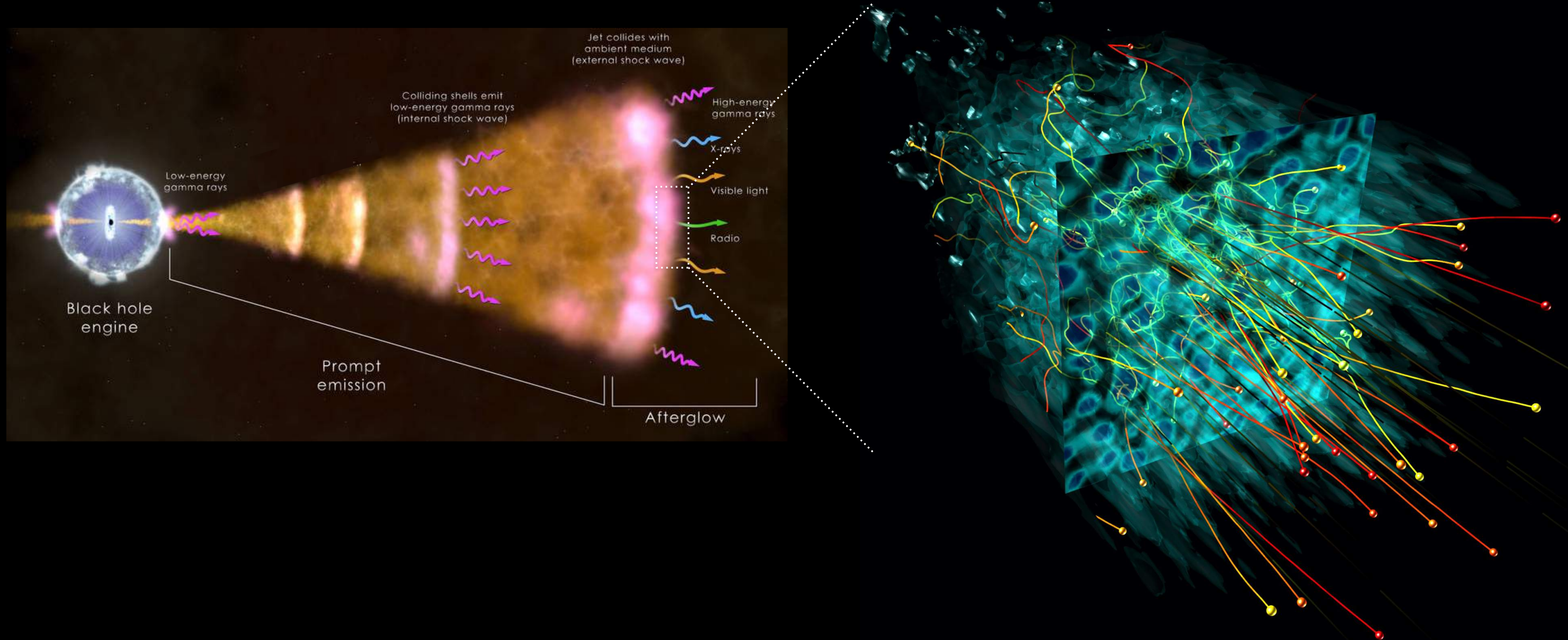


- What do astrophysical shocks manage to be such efficient accelerators?
- How are ultrahigh-energy cosmic rays (with up to 10^{20} eV!) accelerated?
- What controls the EM and particle signatures of compact objects and their mergers?
- Can we study these processes in the lab and harness them to produce more efficient accelerators and light sources?

The answers are tied to kinetic plasma processes that control energy partition in these extreme astrophysical environments

Advances in numerical tools and accelerator facilities/ light sources are now enabling studies of relevant plasma processes to help unveil these mysteries

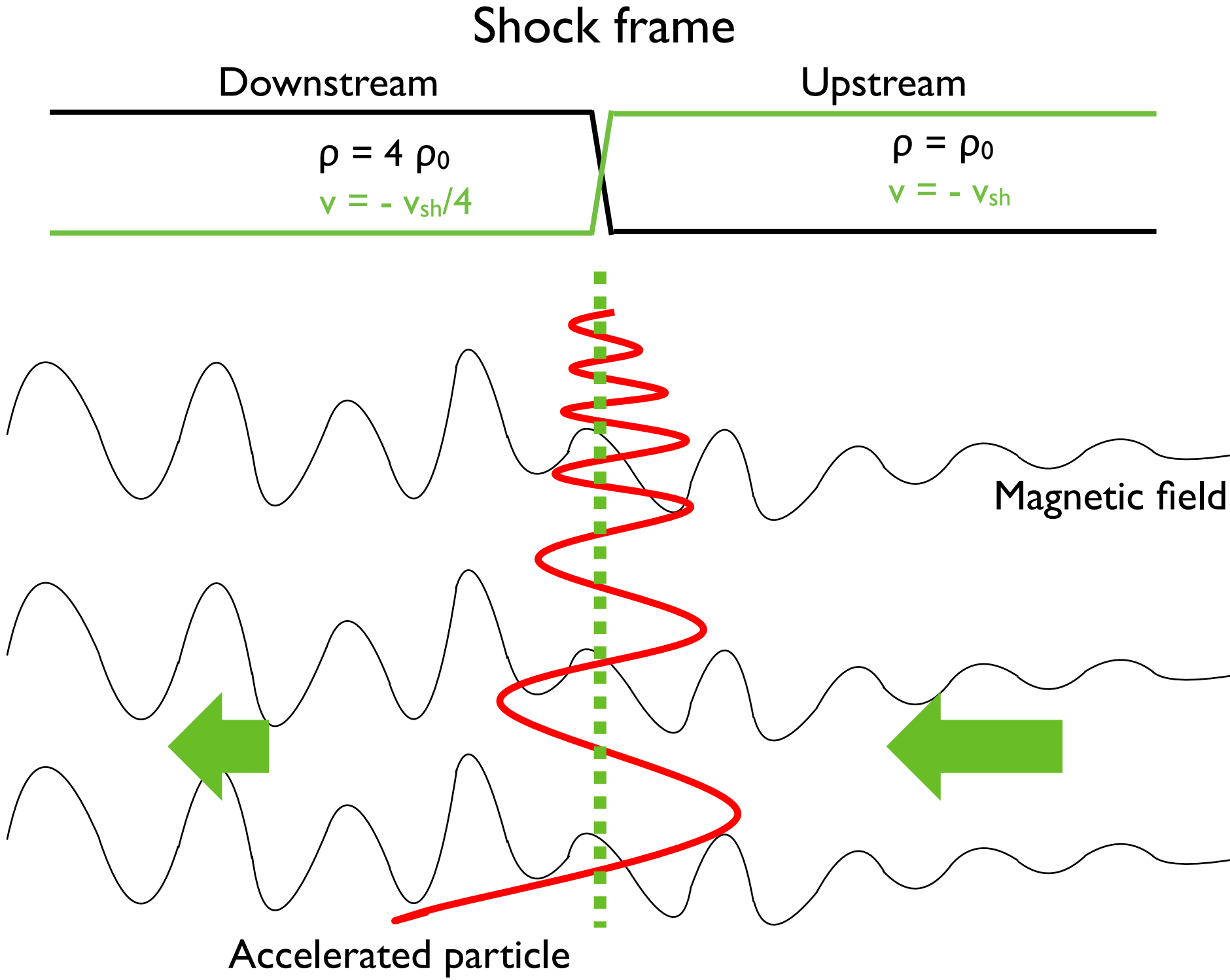
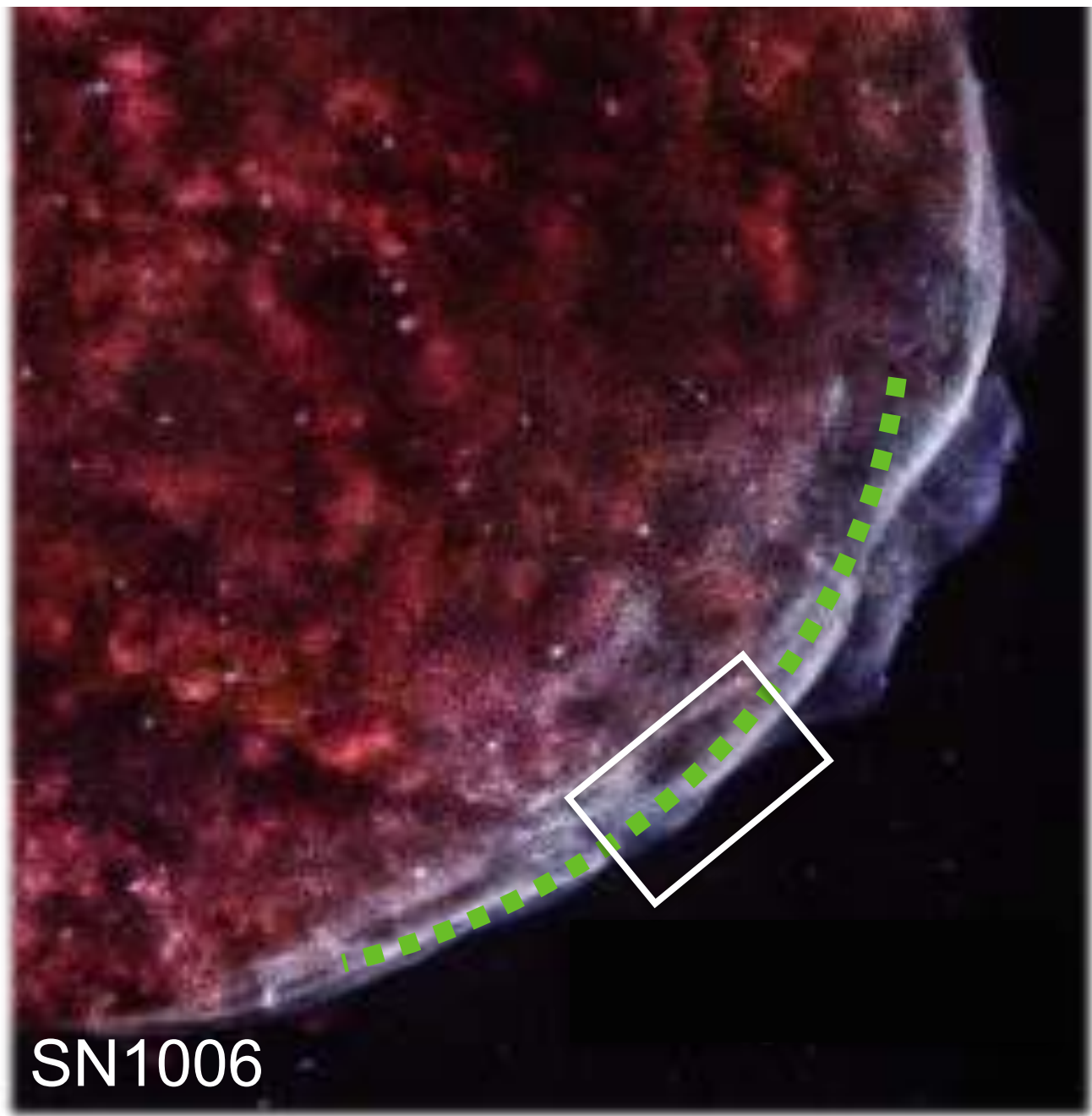
Cosmic accelerators and the importance of CR-driven instabilities



Collisionless shocks are thought to be dominant source of energetic cosmic rays



SNR collisionless shock



1st order Fermi mechanism in shocks*
(a.k.a. diffusive shock acceleration)

Fractional energy gain per shock crossing:

$$\frac{\Delta \epsilon}{\epsilon} = \frac{v_{sh}}{c}$$

Fractional particle loss per shock crossing:

$$\frac{dn}{n} = -\frac{v_{sh}}{c}$$

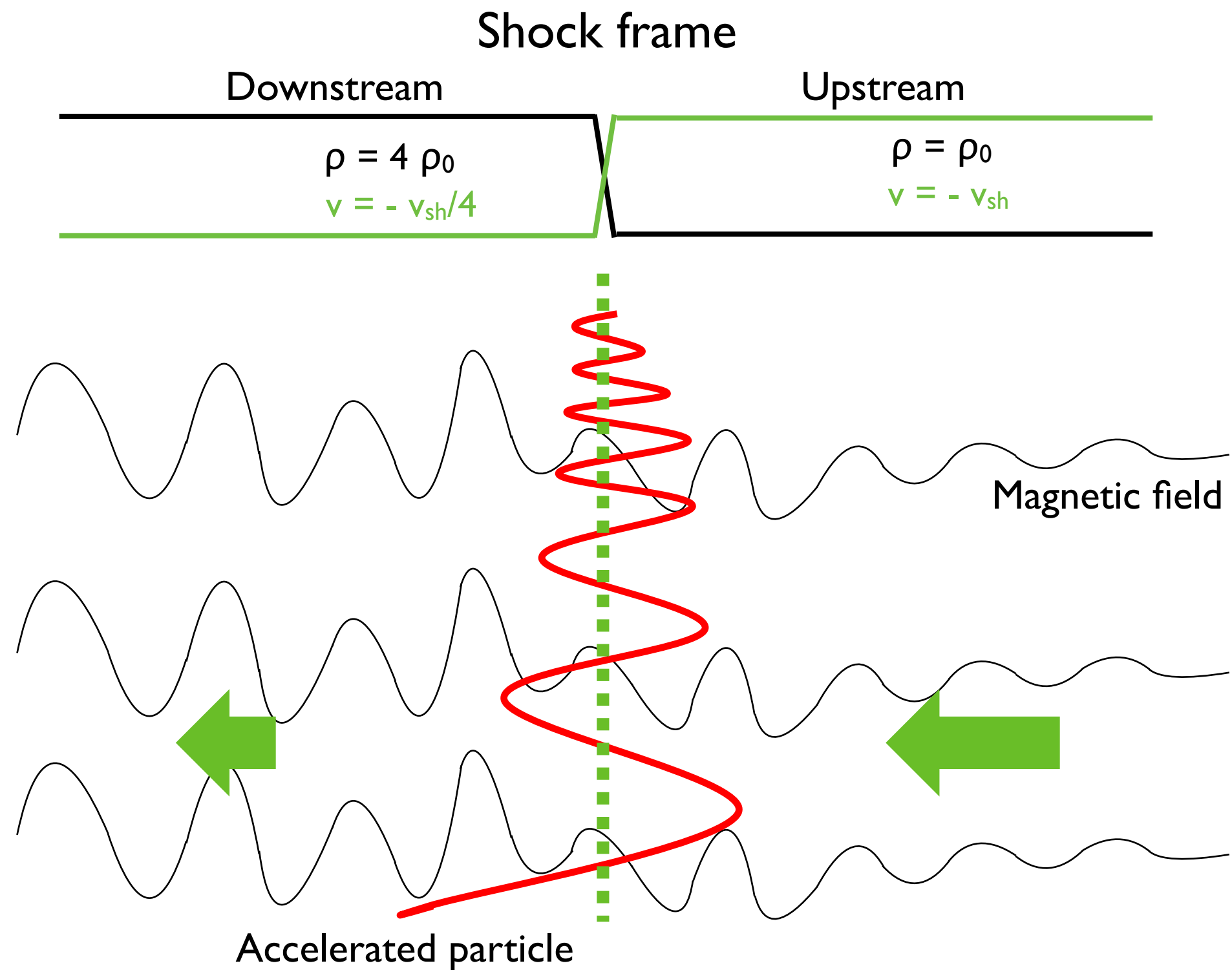
Power-law spectrum:

$$\frac{dN}{d\epsilon} \propto \epsilon^{-2}$$

The microphysics that controls the particle acceleration efficiency at the shock and magnetic field amplification is not yet well understood

* Krymskii 1977, Axford, Leer, Skadron 1977, Bell 1978, Blandford & Ostriker 1978

Which processes control magnetic field amplification?



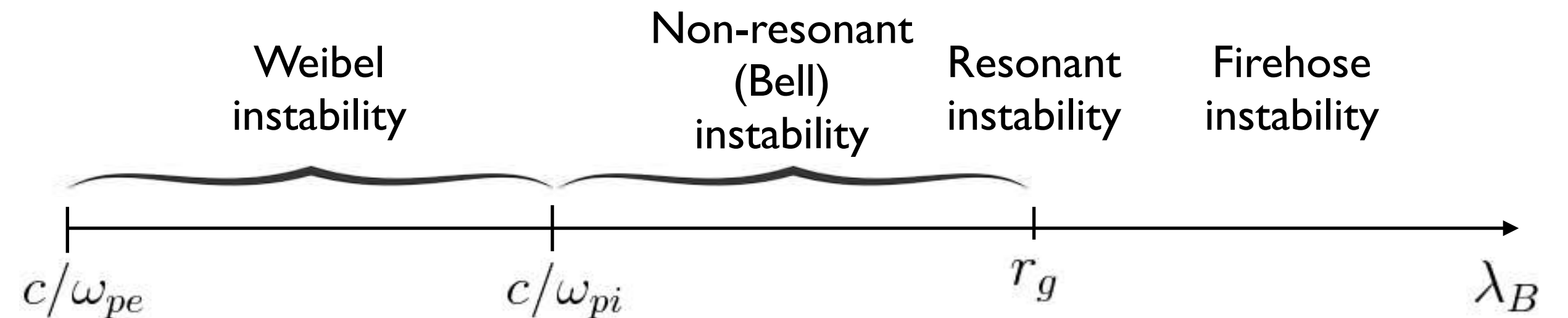
Efficient particle acceleration requires **strong magnetic fluctuations**

Fastest acceleration if scale of fluctuations comparable to Larmor radius of particles (Bohm diffusion)

$$\delta B \gg B$$

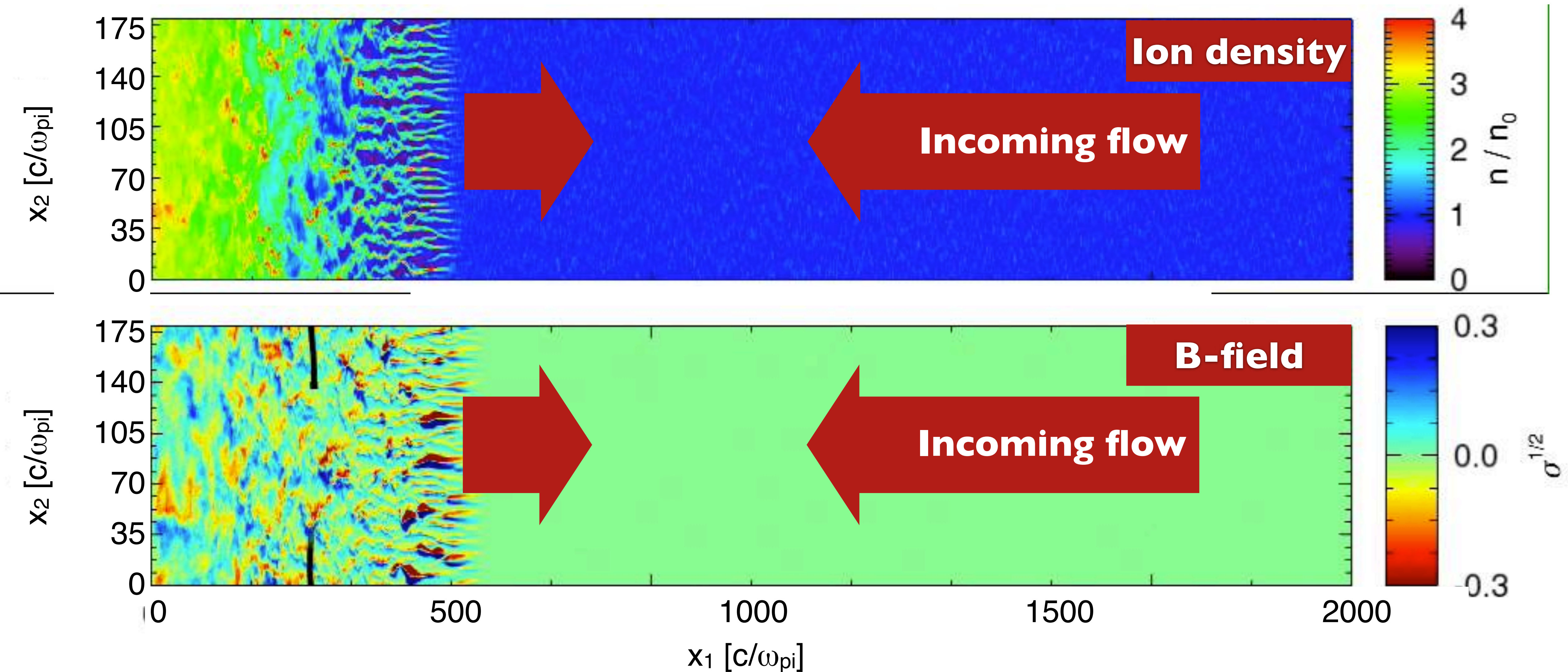
+

$$\lambda_B \sim r_g$$

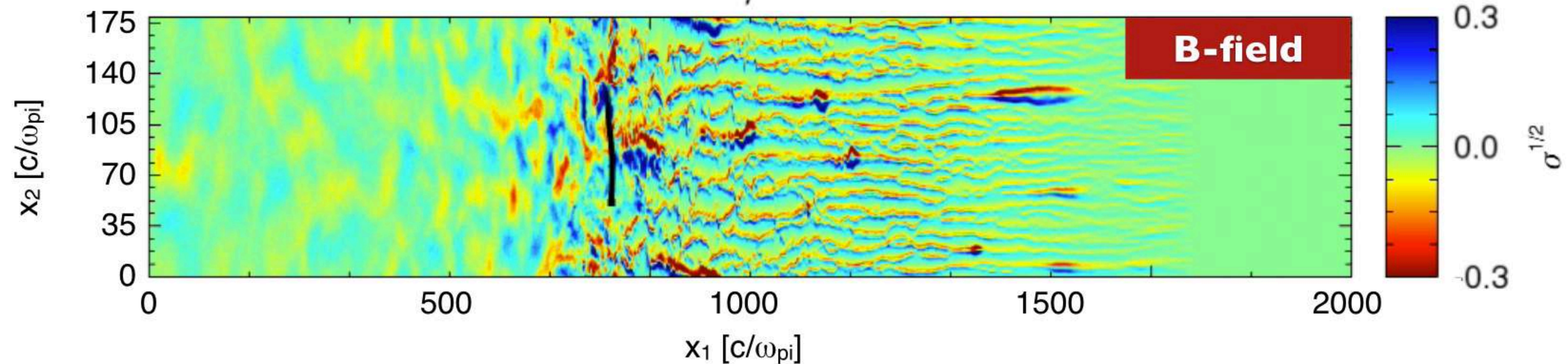
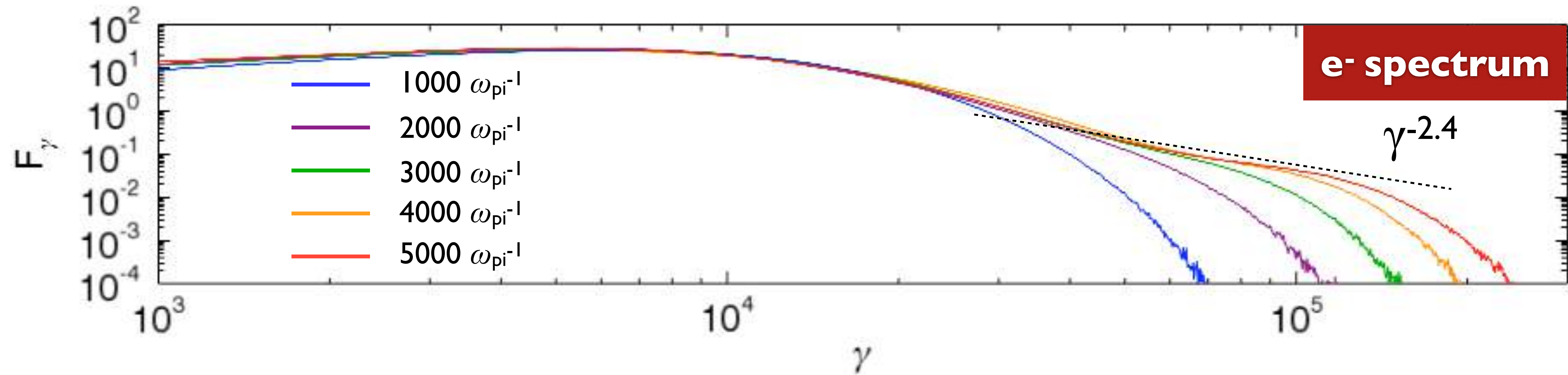


Nonlinear interplay between different instabilities is critical for particle acceleration but it is not clear how it depends on the shock and ambient plasma conditions

Relativistic collisionless shocks in weakly magnetized plasmas

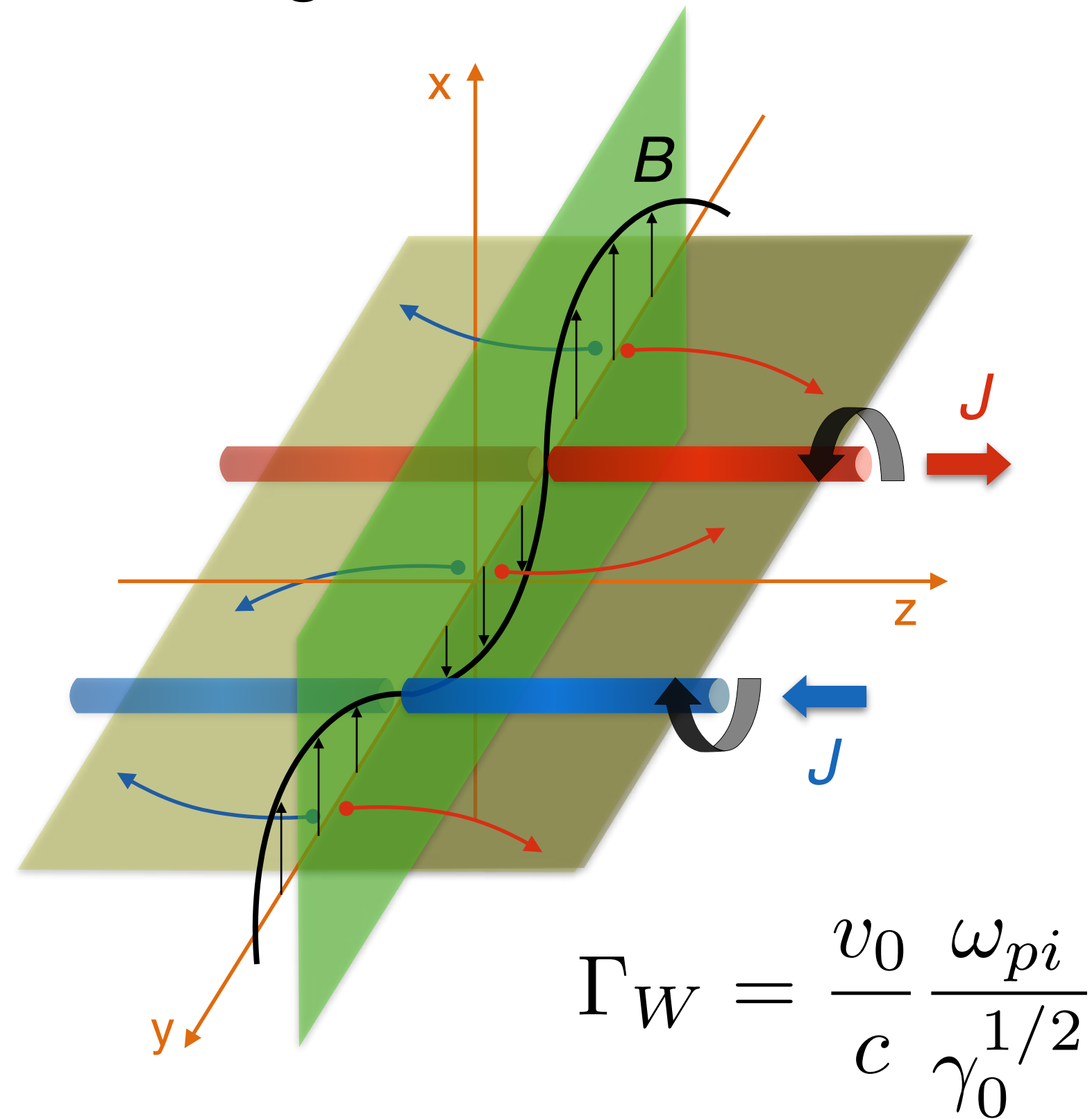


Efficient development of high-energy power-law distribution

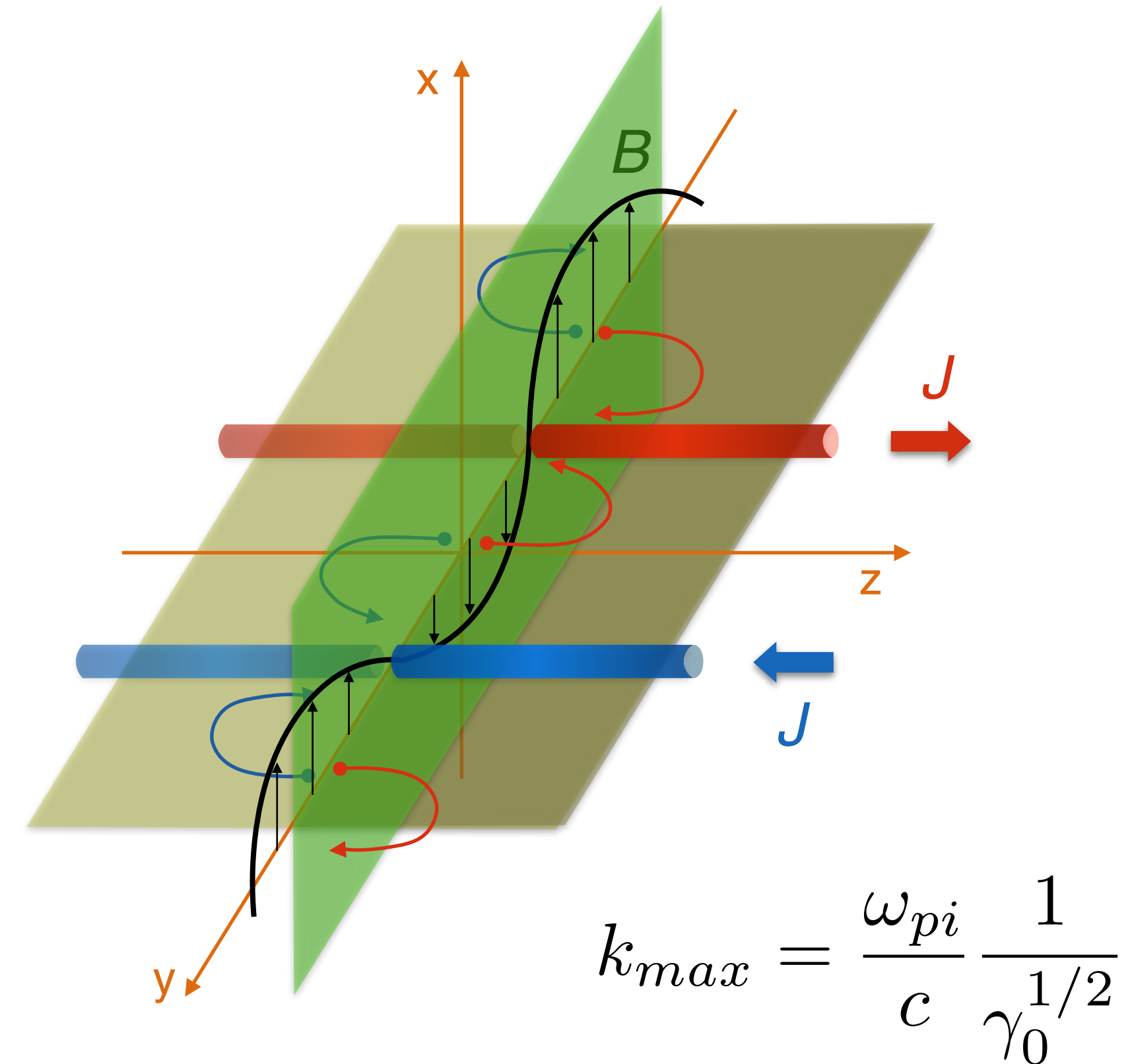


Weibel instability amplifies B-fields at kinetic scales

Linear regime

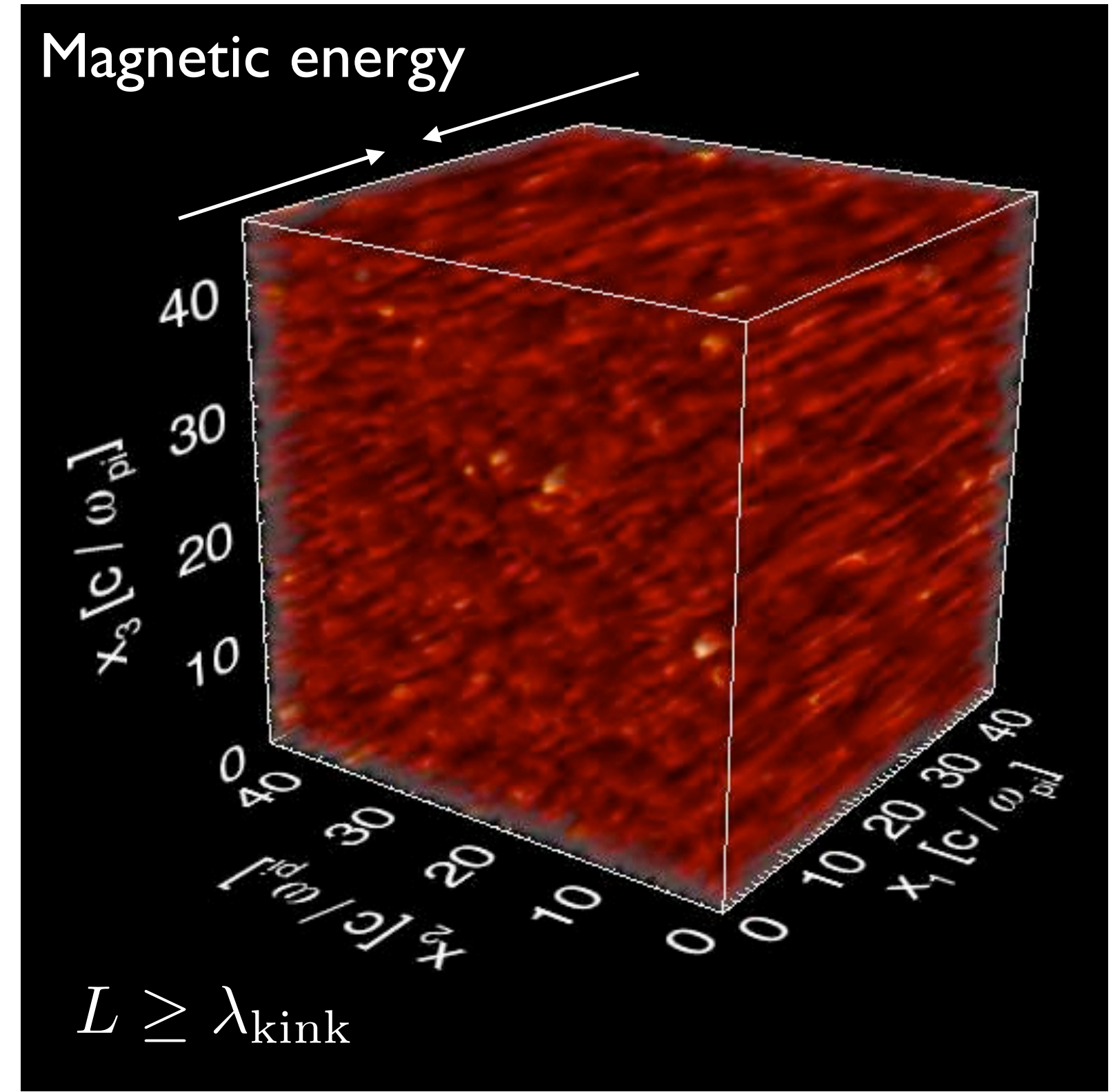
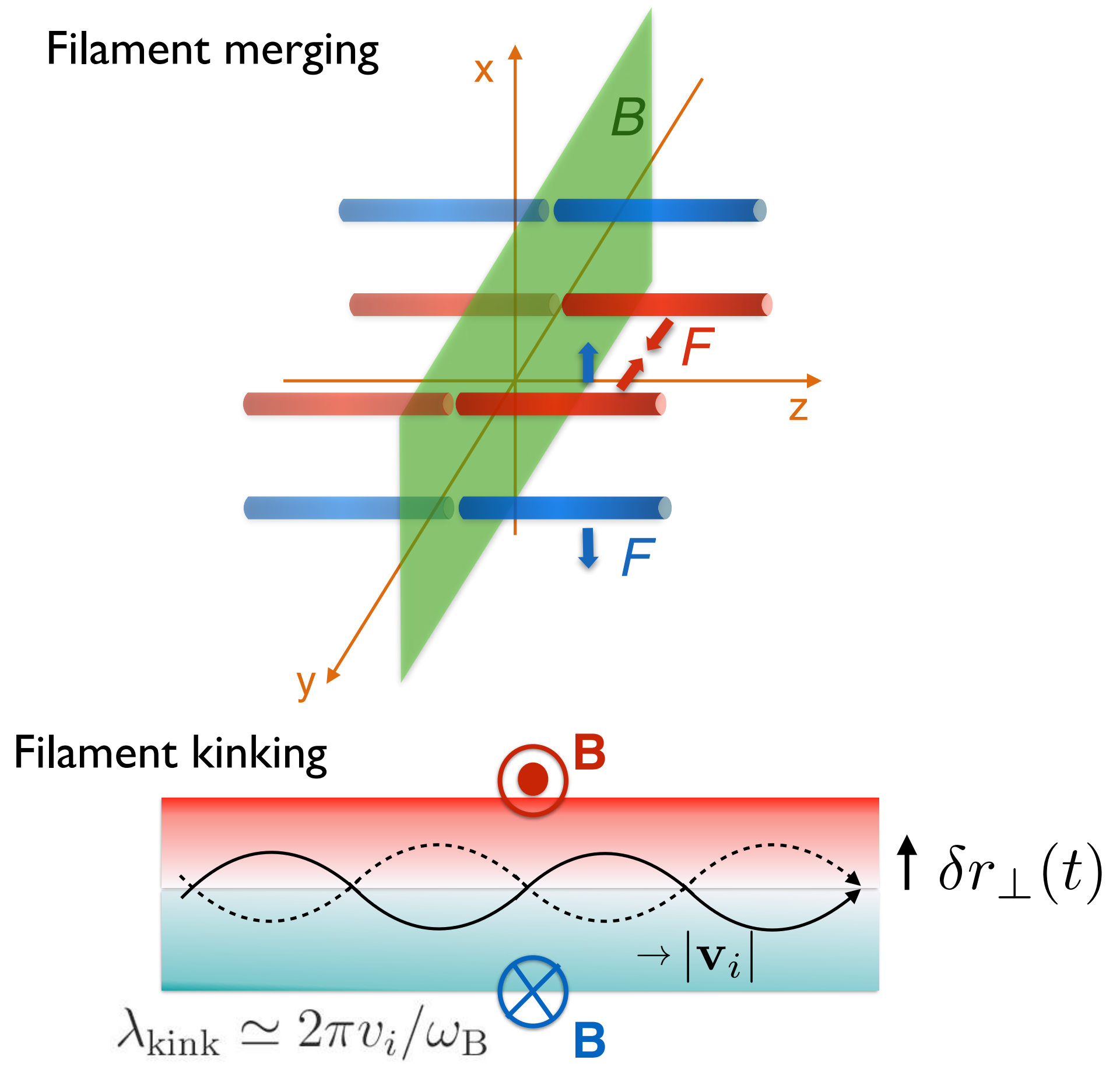


Saturation



Instability can transfer 1-10% of kinetic energy of plasma flows into magnetic energy

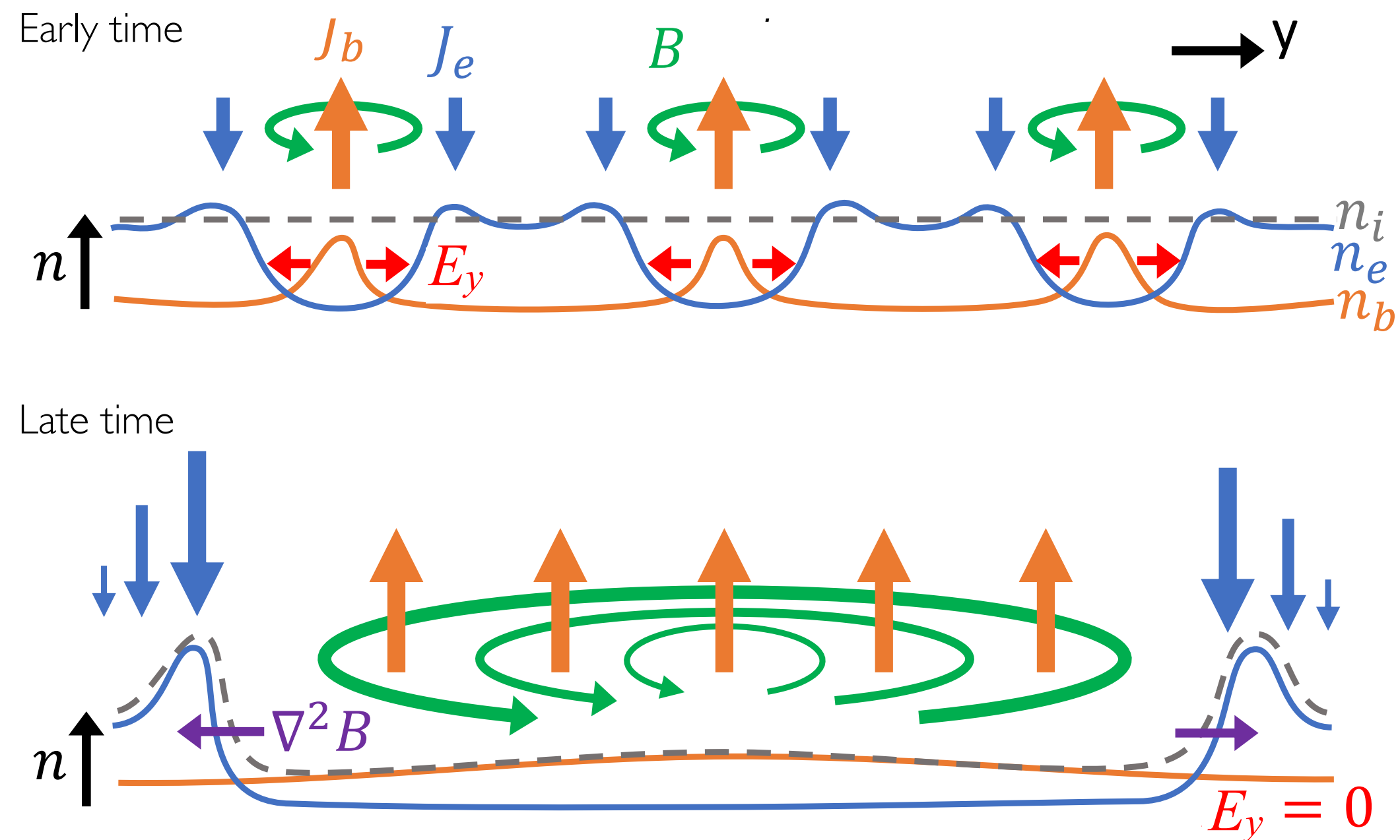
Nonlinear evolution critical for onset of magnetic turbulence



Competition between merging and kink-mode controls shock formation

Secondary instabilities dictate late time evolution

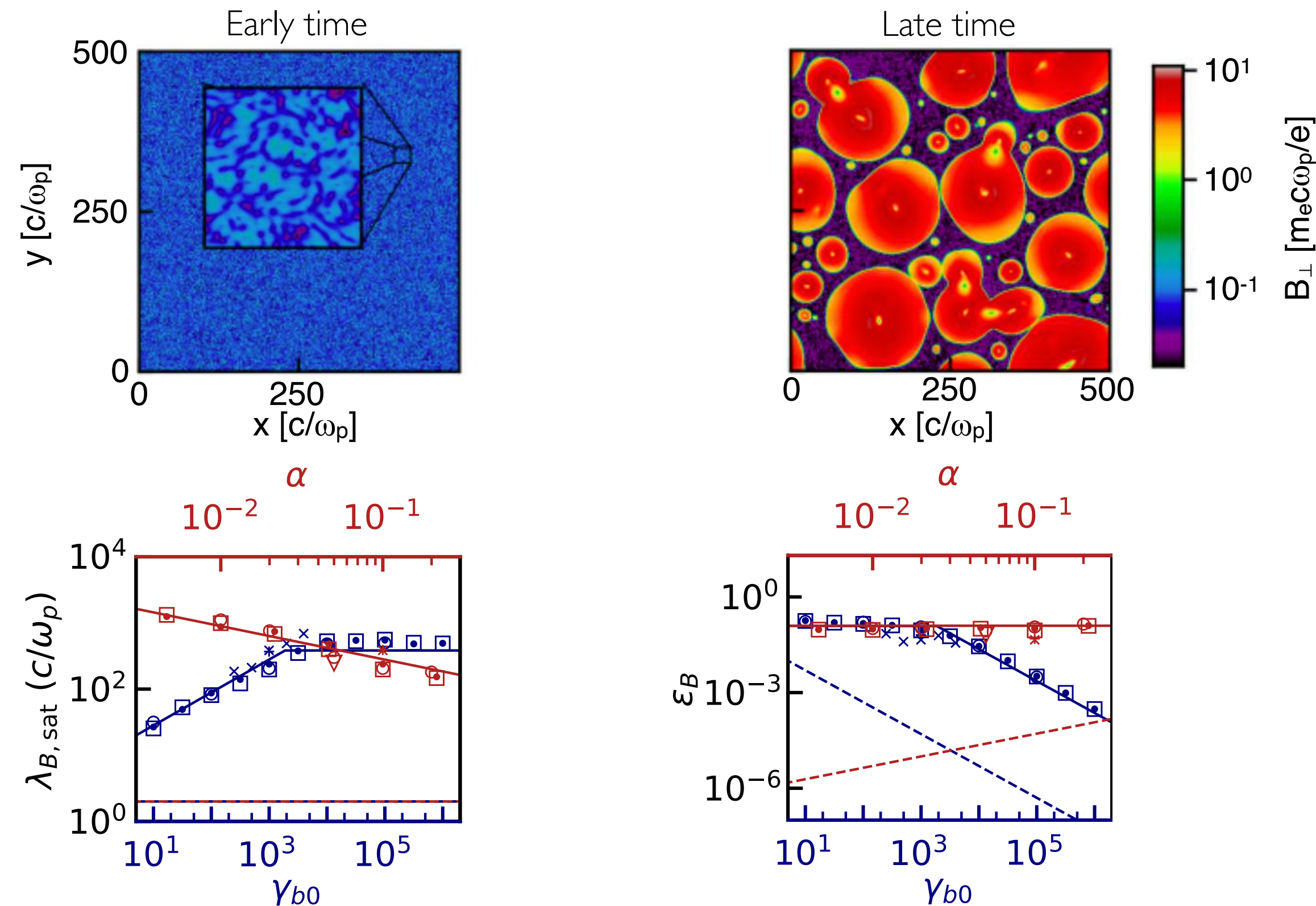
Magnetic field amplification by new plasma cavitation instability



Magnetic coherence scale \gg plasma kinetic scale critical to explain GRB polarization and high-energy particle acceleration

Late time evolution is also dictated by secondary instabilities

Magnetic field amplification by new plasma cavitation instability

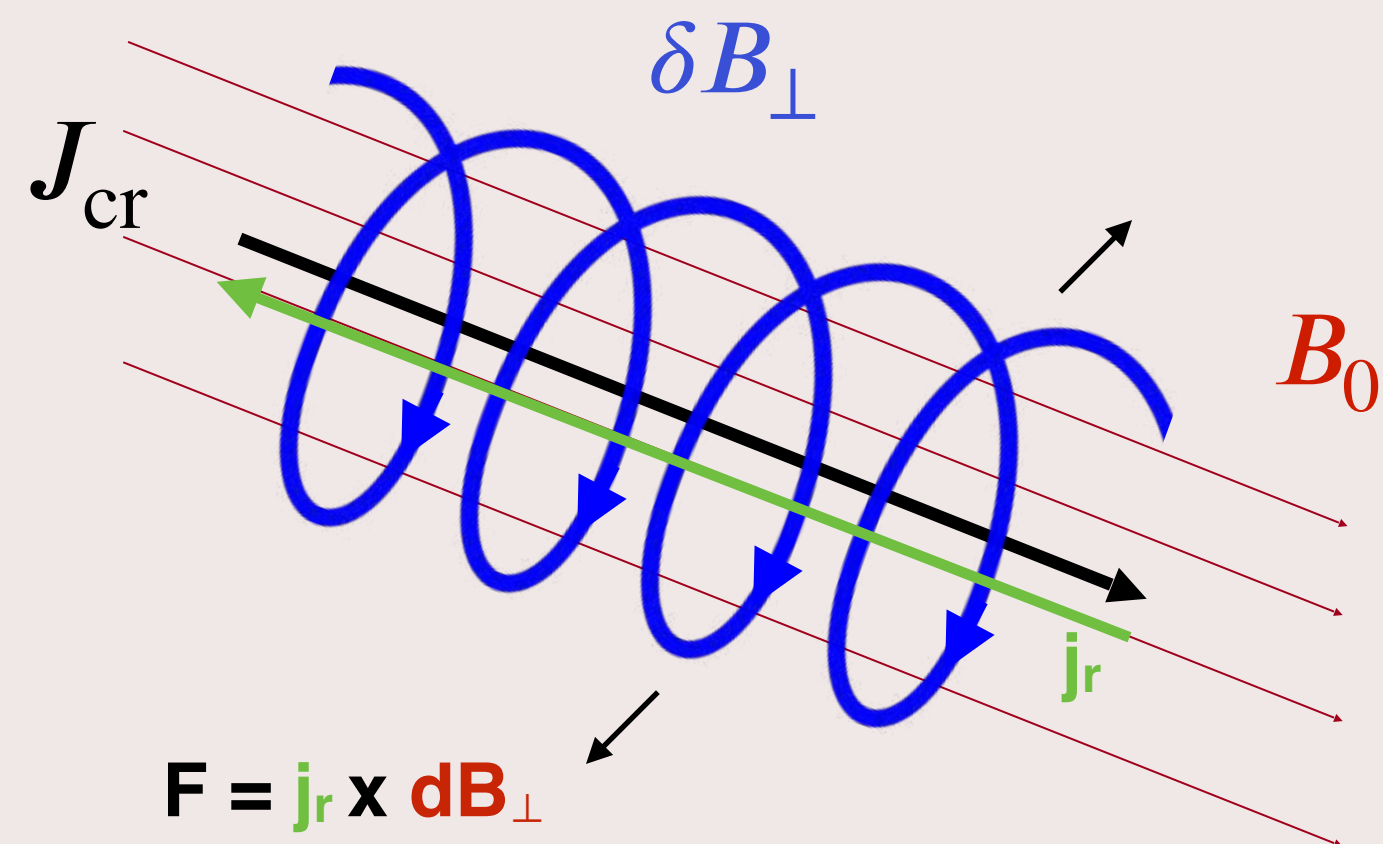


Magnetic coherence scale \gg plasma kinetic scale critical to explain GRB polarization and high-energy particle acceleration

Nonresonant (Bell) and firehose instabilities can generate large-scale modes

Nonresonant (Bell's) current-driven instability*

Driven by charged beam drifting along the external magnetic field

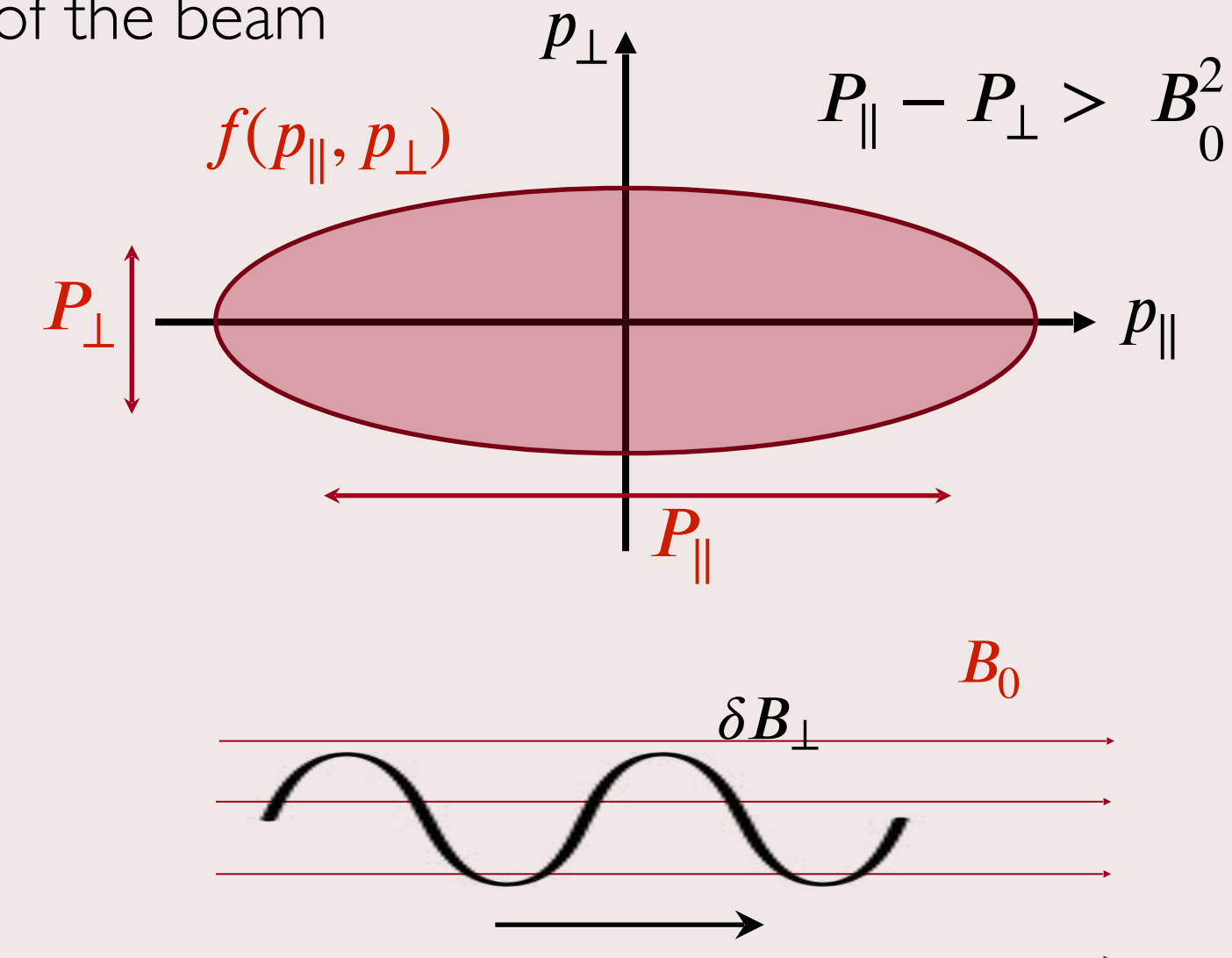


What is the dominant saturation mechanism?

- Equipartition (Bell 2004)
- Resonant pitch-angle scattering (Caprioli & Spitkovsky 2014)
- Background plasma acceleration (Riquelme & Spitkovsky 2009, Gargate et al. 2010)

Firehose instability**

Driven by (ram or thermal) pressure anisotropy of the beam

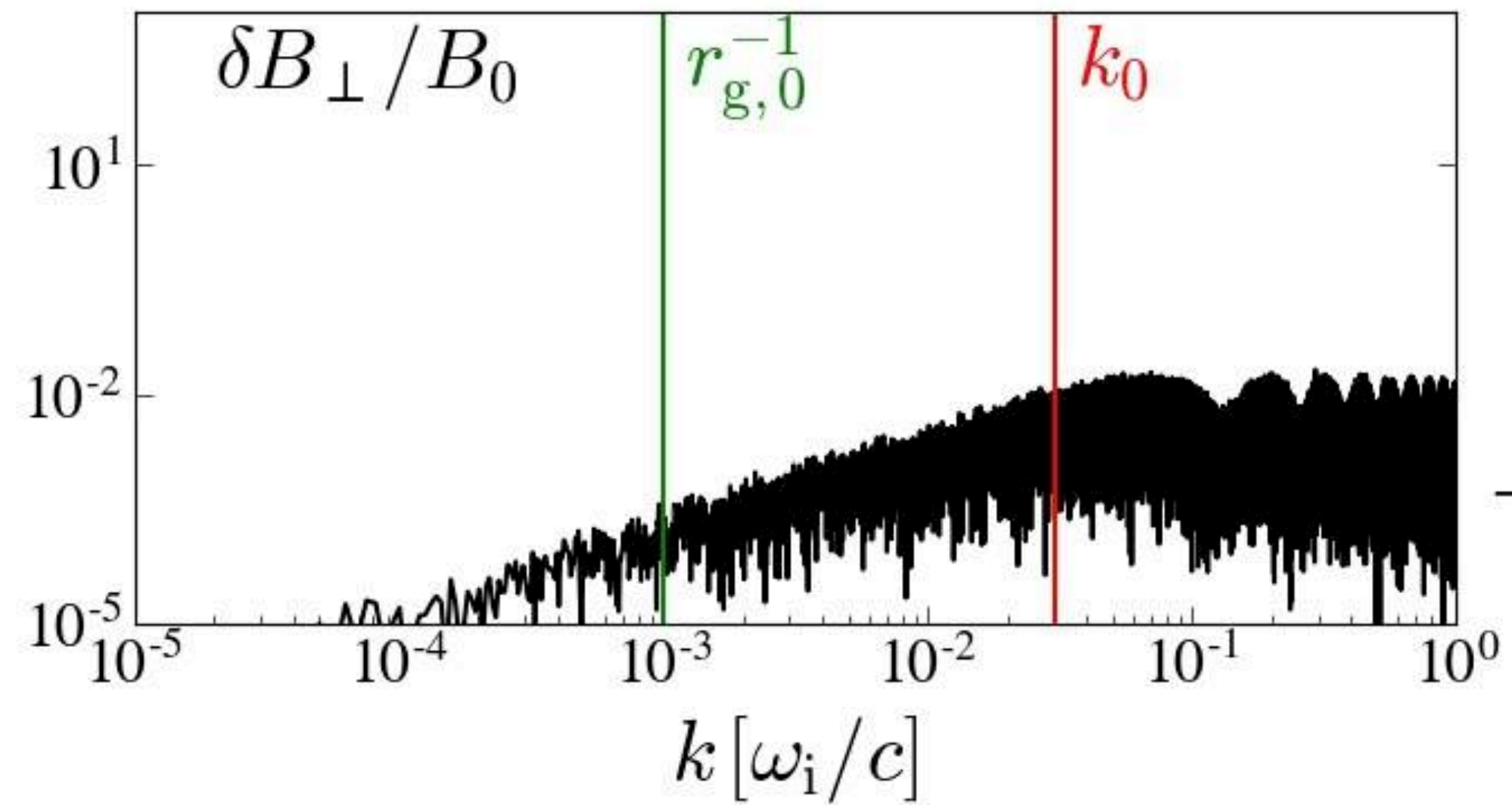


Dominant saturation mechanism?

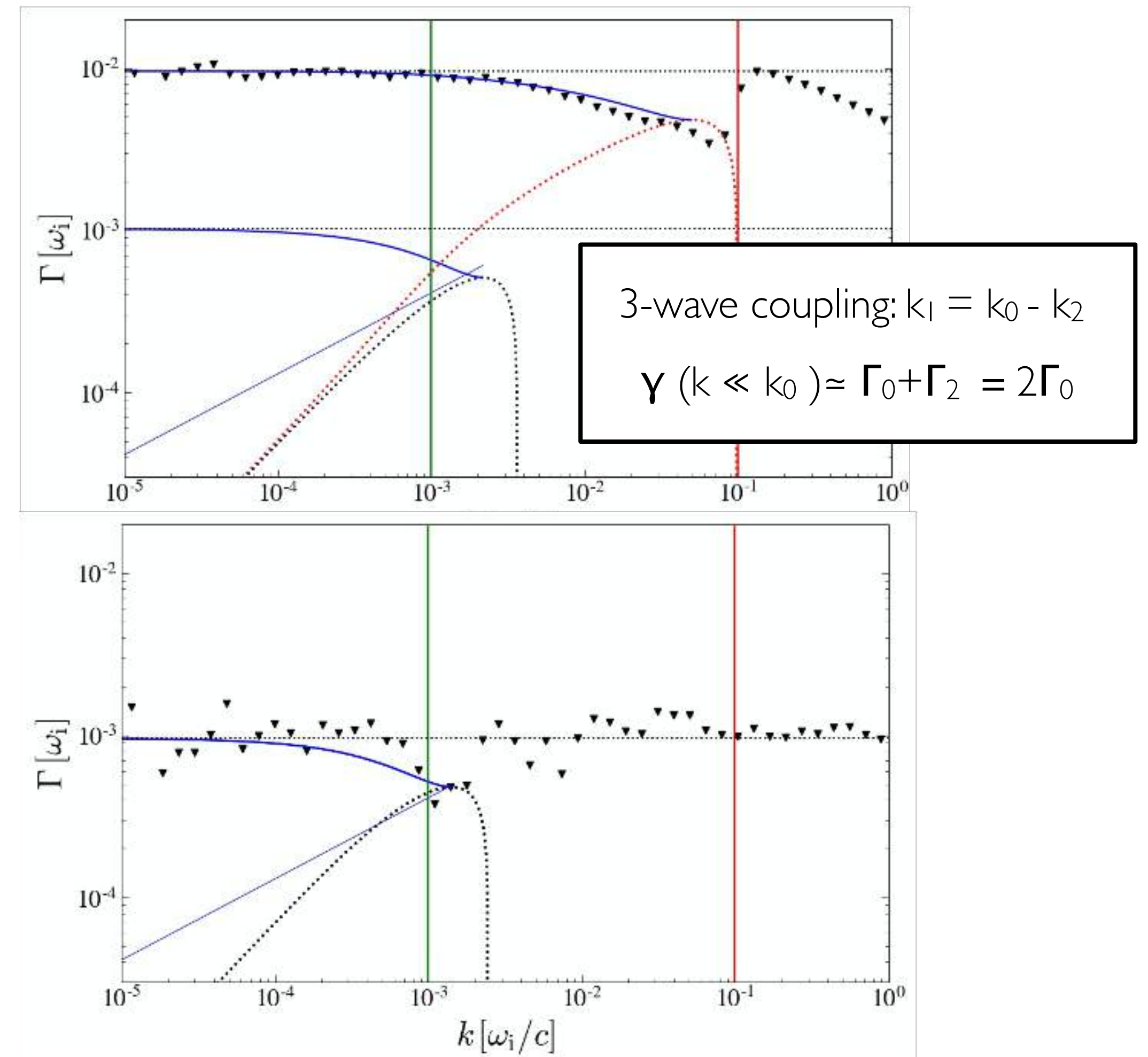
- Thermal + magnetic pressure balance

Nonlinear mode coupling important for B-field amplification at large scales

Fully-kinetic PIC simulations with OSIRIS



Nonlinear mode coupling



- short-scale limit (Bell)

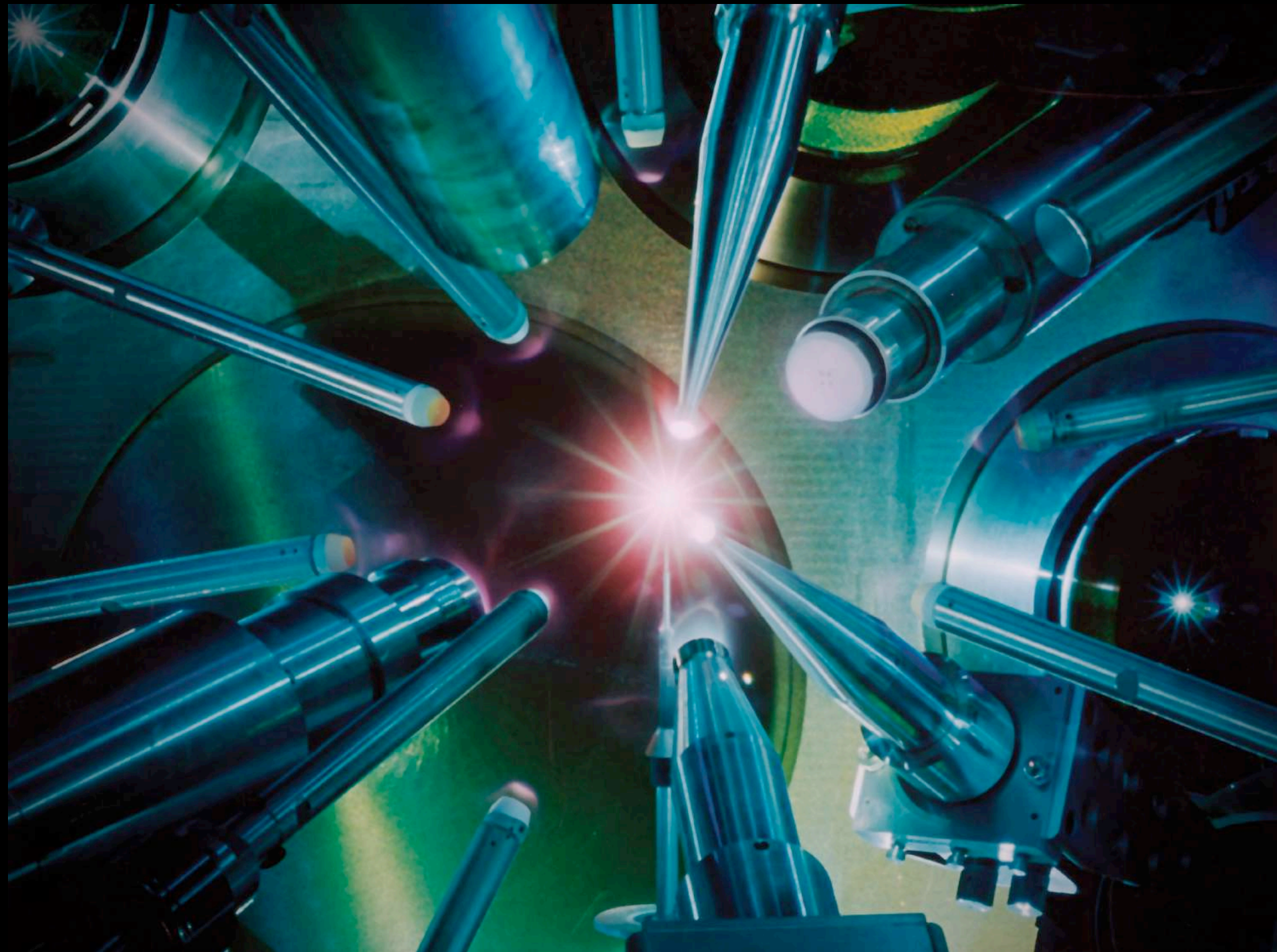
$$\omega = \beta_A (k^2 - k k_0)^{1/2}, k_0 = J_{cr} / B_0$$

- Large-scale limit (firehose)

$$\omega = k \beta_A \left(1 - \frac{P_{ram}}{B_0^2} \right)^{1/2}$$

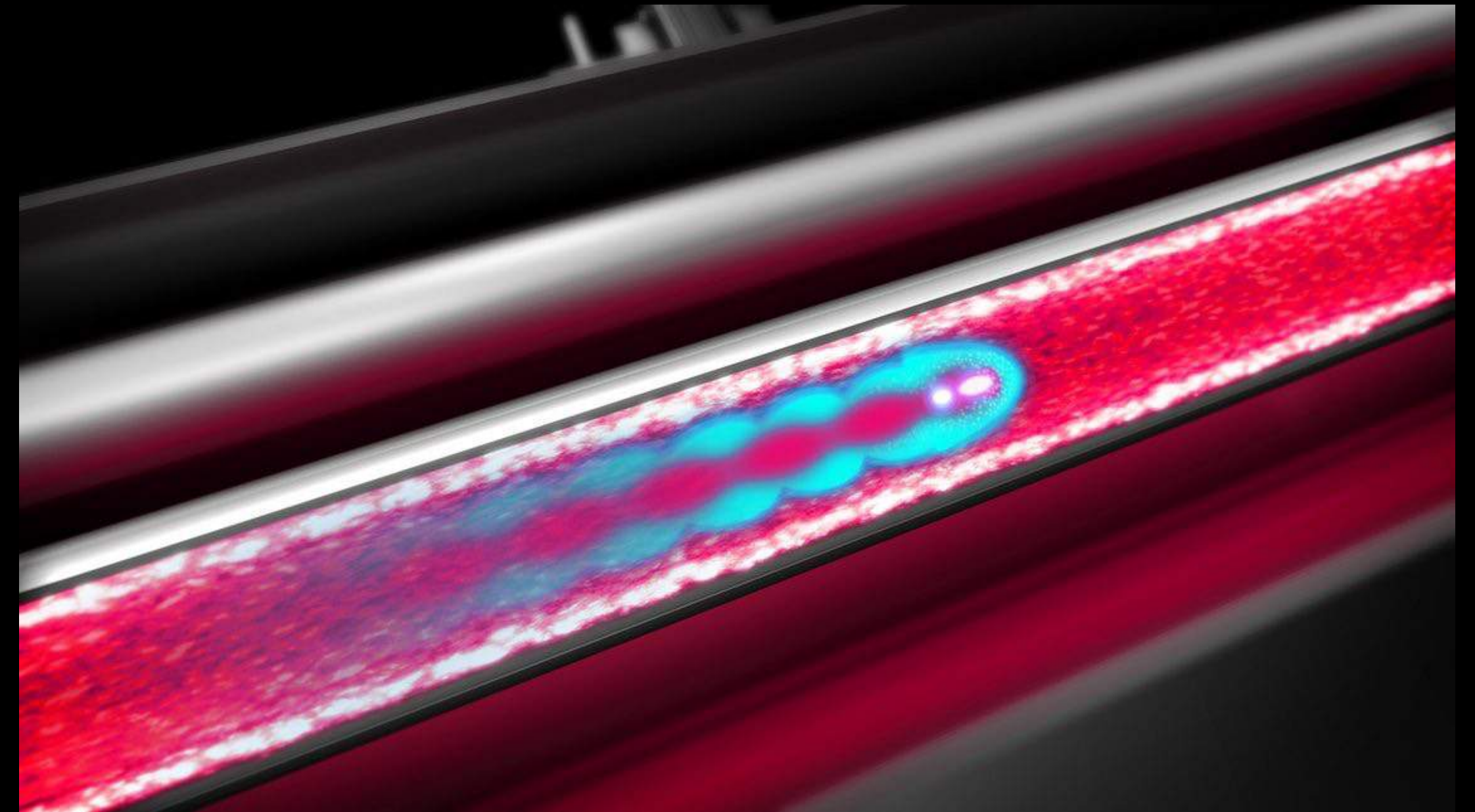
Can laboratory experiments help unveil the nonlinear evolution of CR-driven instabilities?

High-intensity and high-energy lasers



OMEGA EP @ LLE

High-intensity and high energy particle beams



FACET-II @ SLAC

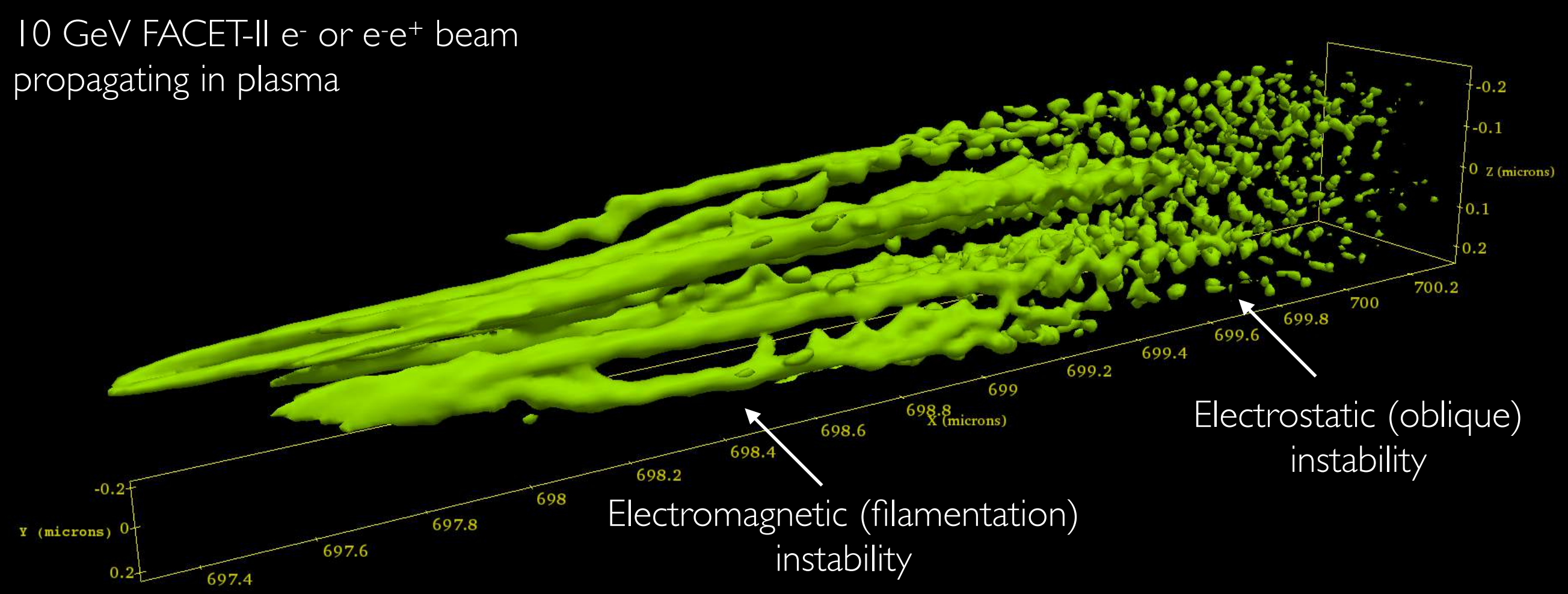
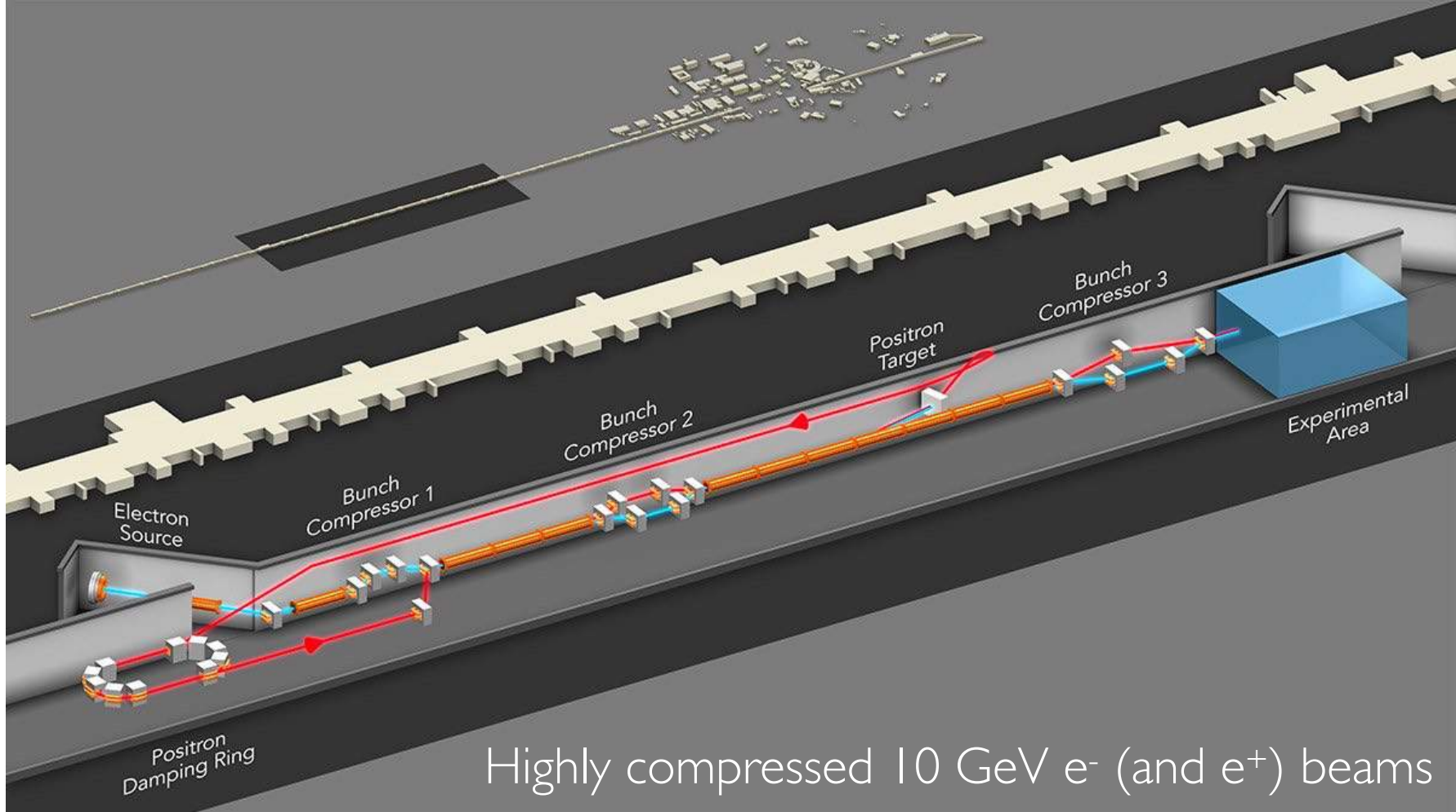
We are not going to put a GRB in a bottle!
... but can we test critical theoretical aspects and validate models/codes?

Studying filamentation instability with FACET-II



FACET-II could explore nonlinear competition between cosmic-ray driven instabilities

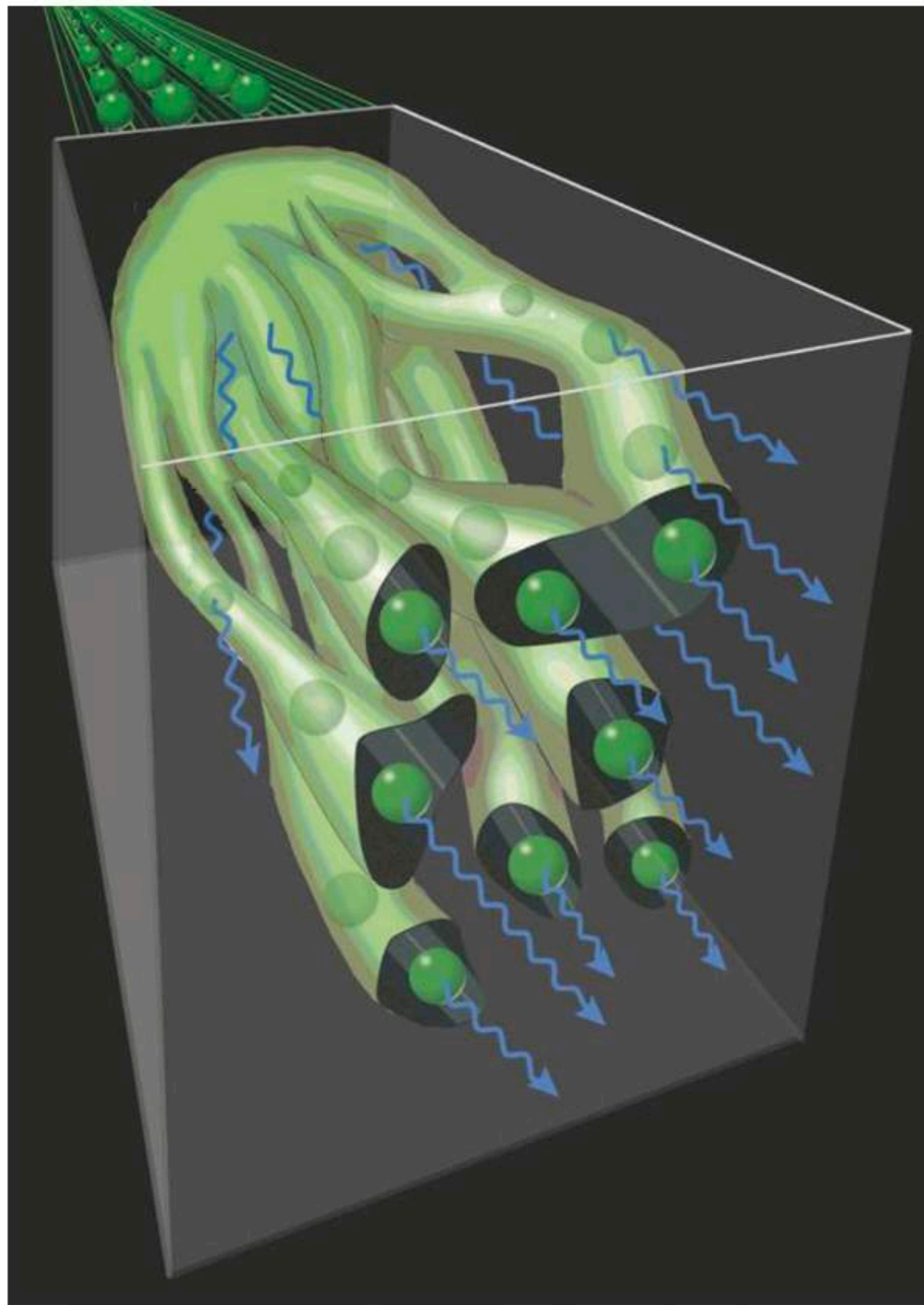
Approved FACET-II experiment E-305: beam filamentation and bright gamma-ray bursts



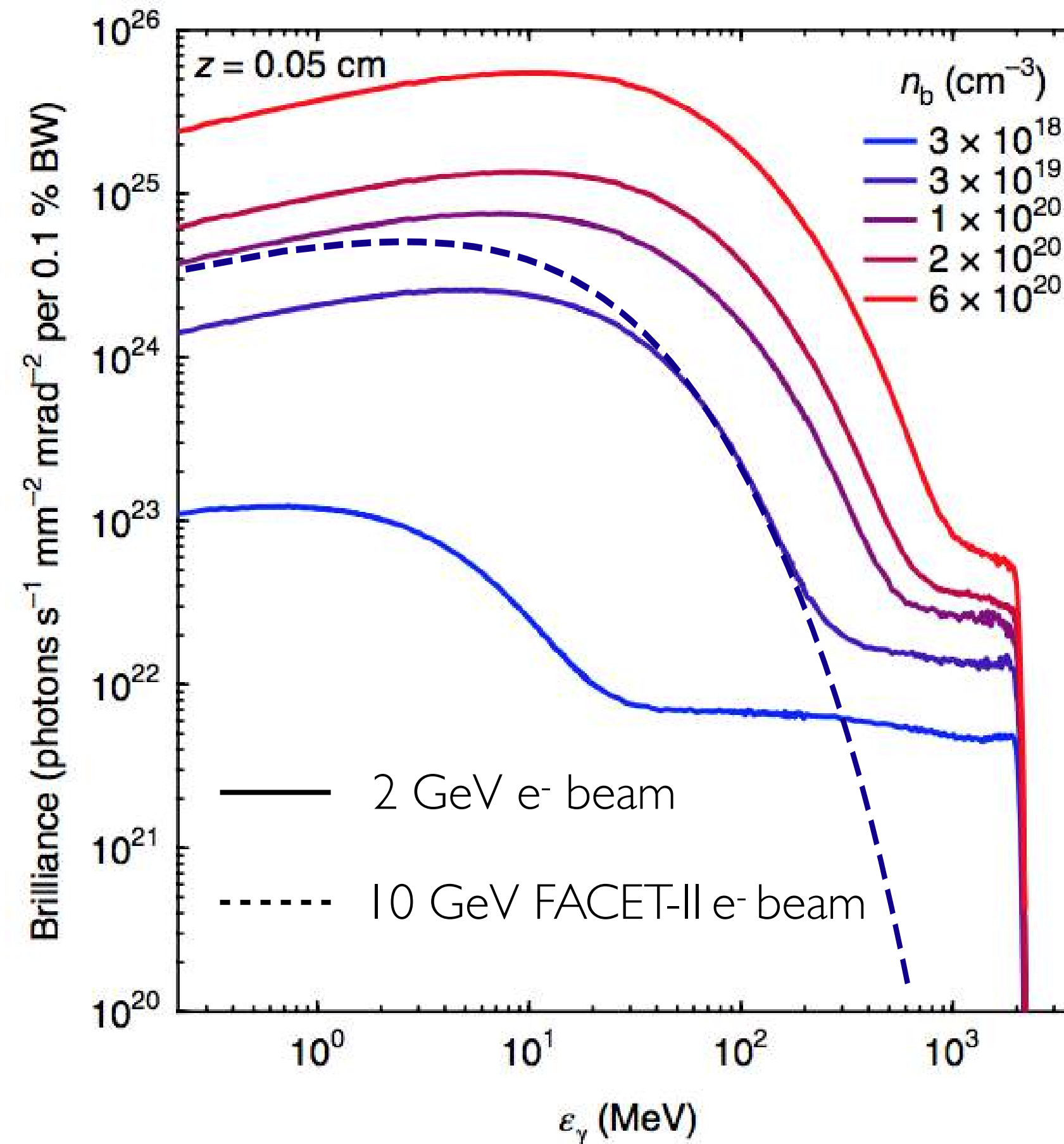
A. Sampath et al. PRL (2021), P. Claveria et al, PRR (2022)

High brilliance gamma-ray flashes

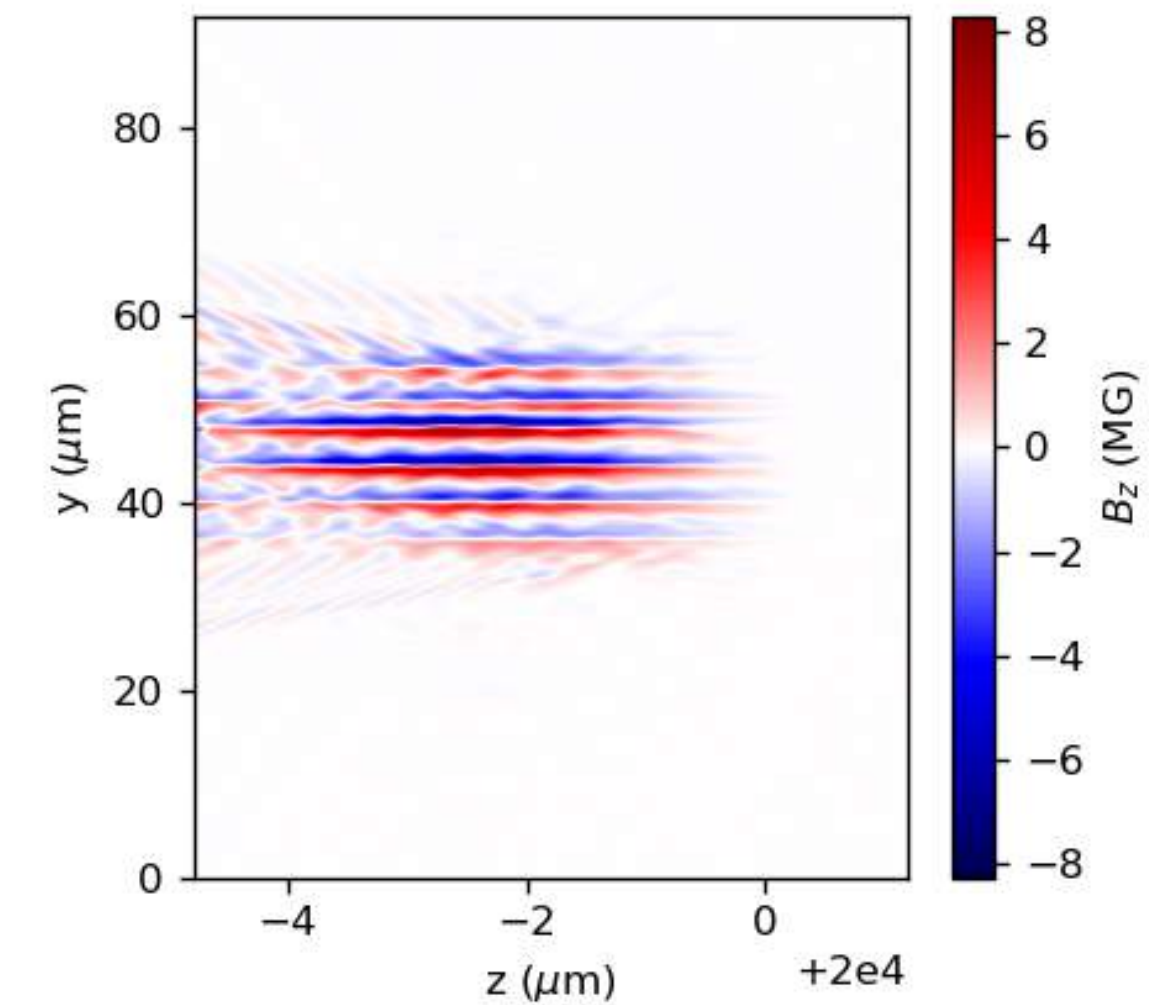
Synchrotron emission is self-amplified B-fields



Gamma-ray source with unprecedented high brilliance



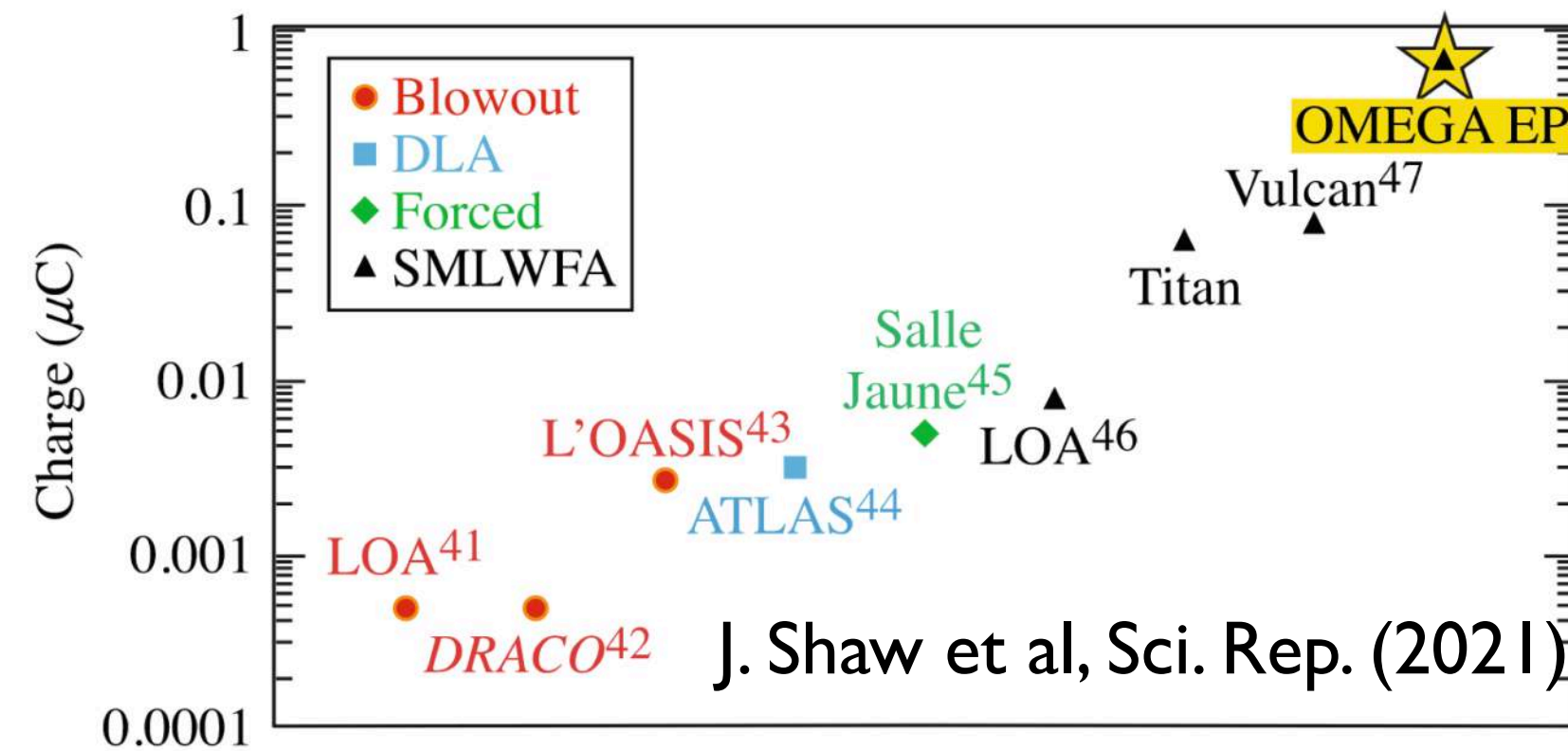
B-fields for propagation over 2 cm gas jet



Cavitation instability driven by μC laser-driven electron beams

Self-modulated LWFA regime explored at OMEGA EP

1 kJ, 1 ps laser \rightarrow 1 μC , 200 MeV ($\eta \sim 10\%$)

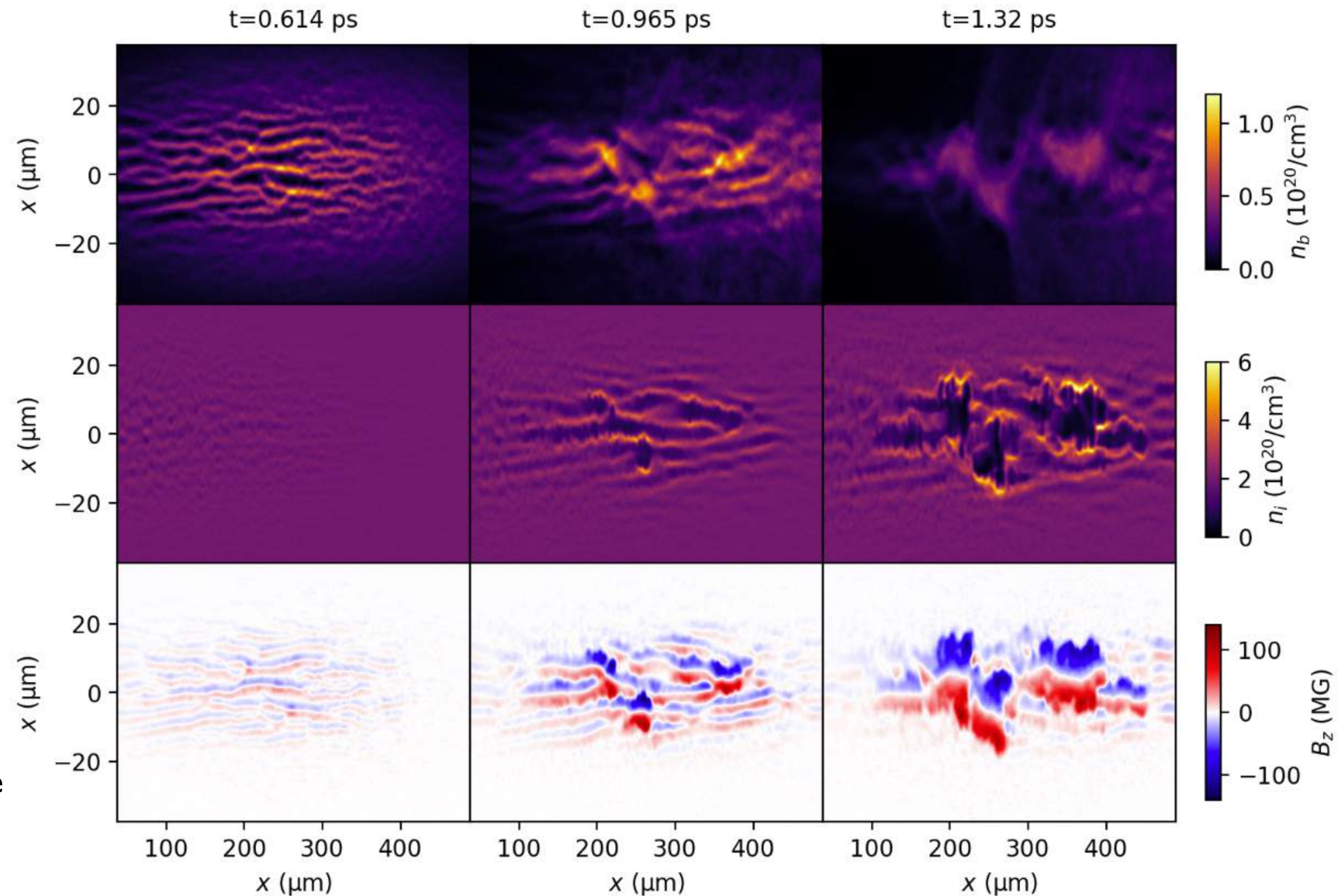


With NIF ARC or PETAL (assuming similar efficiency)

10 kJ laser \rightarrow 10 μC , 50 MeV

Cavitation instability could be driven to nonlinear regime

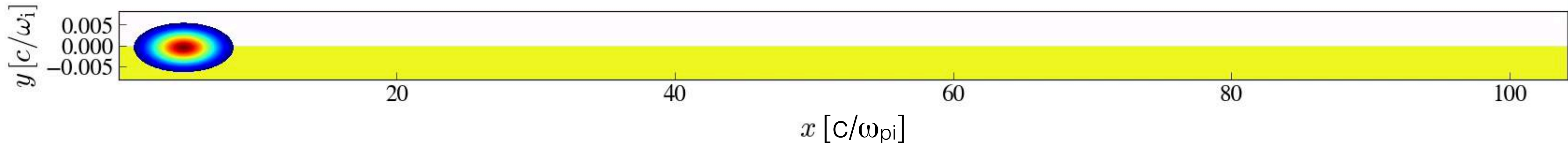
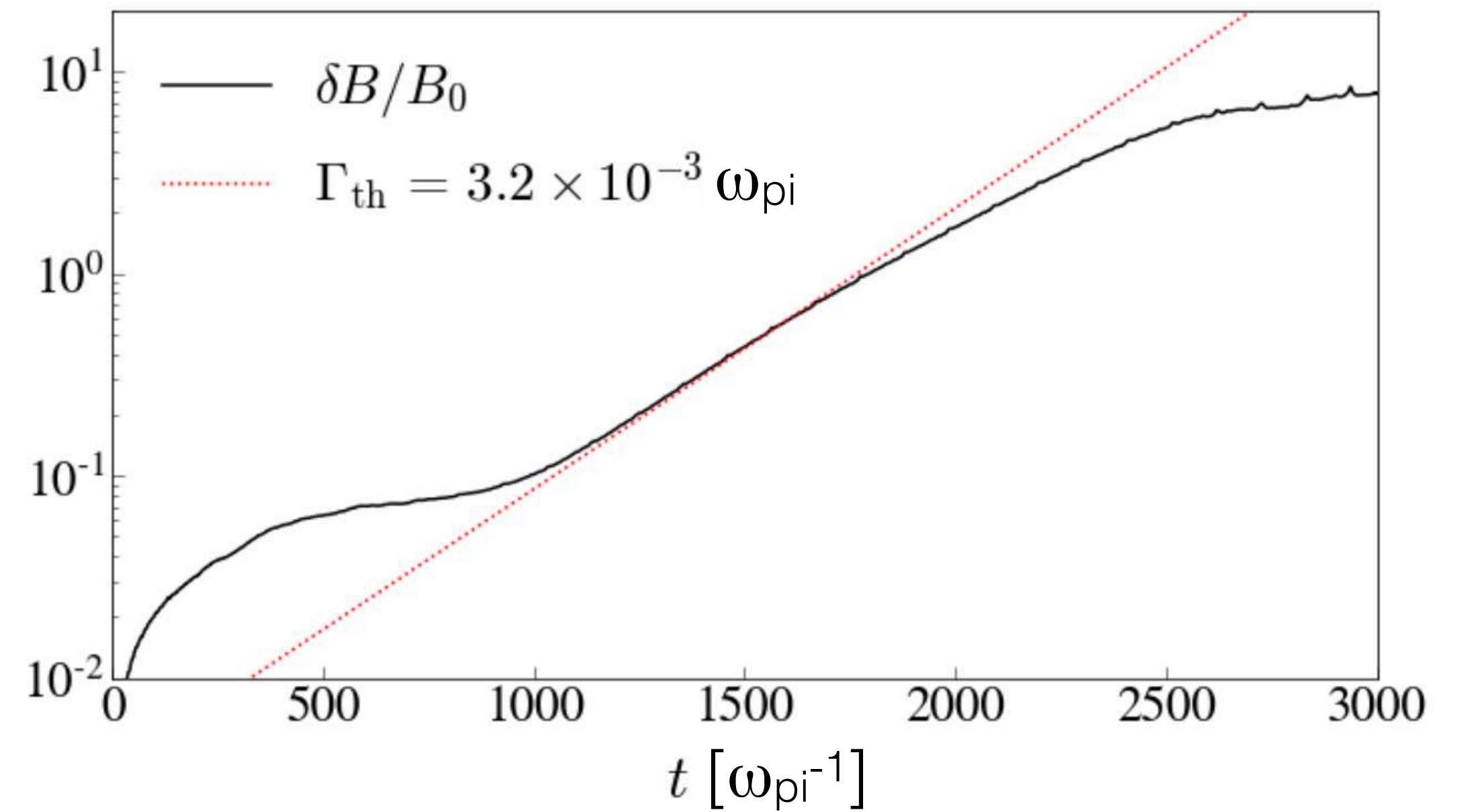
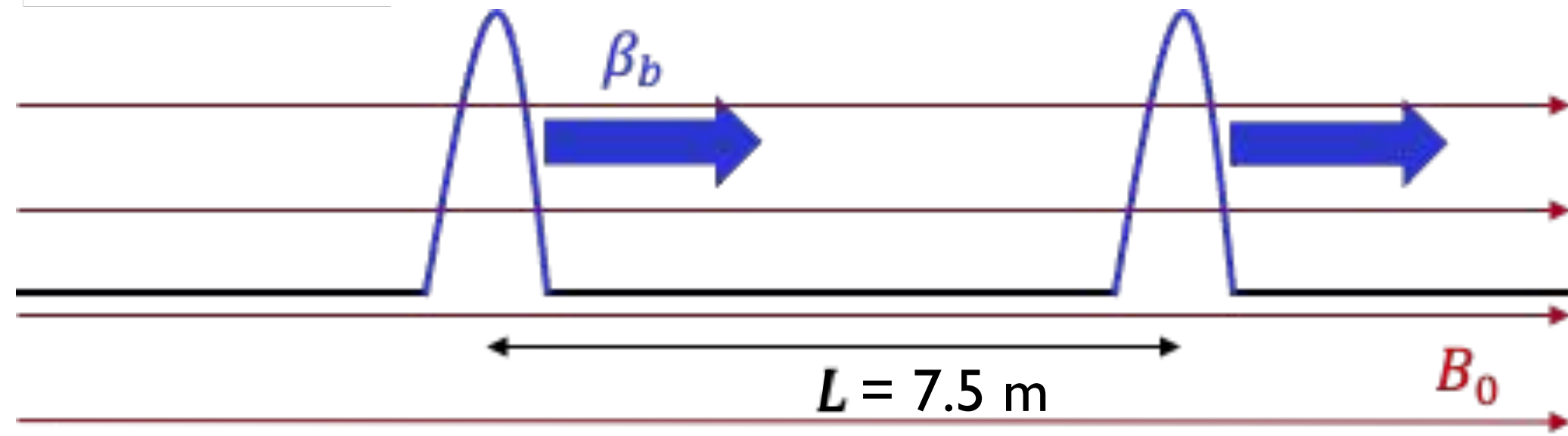
$\epsilon_b = 40 \text{ MeV}$, $\tau_b = 1.5 \text{ ps}$, $\sigma_{\perp} = 25 \mu\text{m}$, $n_b/n_0 = 0.2$



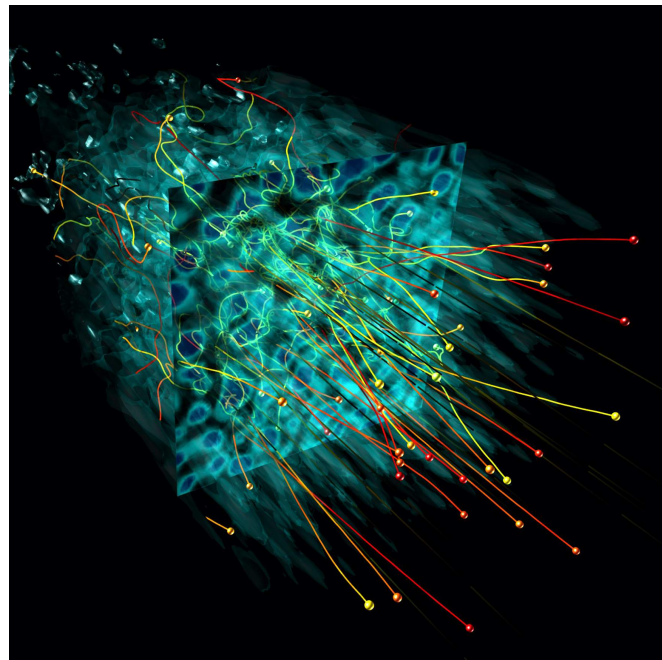
Bell instability driven by proton bunch train from SPS at CERN



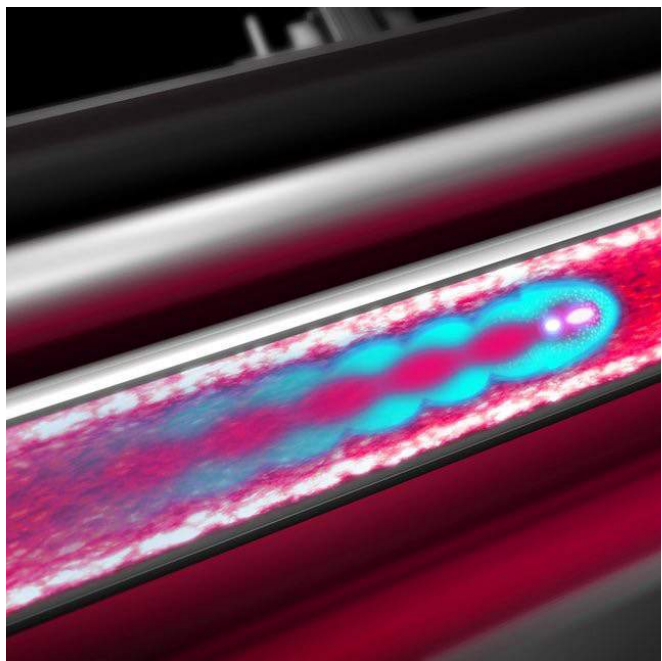
- Bunch density: $4 \times 10^{12} \text{ cm}^{-3}$
- $\sigma_z = 12 \text{ cm}$
- $\sigma_r = 0.02 \text{ cm}$
- Proton energy: 400 GeV
- $n_0 \simeq 10^{13} \text{ cm}^{-3}$
- $B_0 \simeq 0.2 \text{ T}$



Opening new windows into cosmic accelerators



Astrophysical observations reveal that astrophysical plasmas are **extraordinary particle accelerators** and insist on the need to understand better the plasma physics at these extreme environments



The combination of first-principles simulations and laboratory experiments is **opening new windows** into the physics of cosmic-ray driven instabilities that can **advance our understanding of particle acceleration** and...



... **establish new connections** between **plasma-based accelerators** and **high-energy astrophysics** with far-reaching implications for our understanding of the most spectacular phenomena in the Universe