

Associated Higgs Production at NNLO+PS with SMEFT contributions

Luc Schnell IMPRS Colloquium July 13, 2023

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1. Introduction

1.1 Associated Higgs Production (Vh)1.2 Theoretical Predictions1.3 SMEFT



1. Introduction

1.1 Associated Higgs Production (Vh)

Sources: ¹<u>ArXiv:1808.08238</u> (ATLAS), ²<u>ArXiv:1808.08242</u> (CMS).





Measure y_b to appreciable precision.

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Measure y_b to appreciable precision. • **Goal:** $h \xrightarrow{y_b} \overline{b}$ SM: $y_b = \frac{\sqrt{2} m_b}{v} \rightarrow$ any deviation therefrom is a clear sign for NP.

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Source: <u>nytimes.com</u>





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Measure y_b to appreciable precision.

SM:
$$y_b = \frac{\sqrt{2} m_b}{v} \rightarrow v$$

down to $\pm 0.05^{1,2}$.

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Currently: $\mu_{h \rightarrow b\bar{b}} = 1.02$ down to $\pm 0.05^{1,2}$.

• Choose your fighter:

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Source: opendata.atlas.cern.





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Candidate Event: $pp \rightarrow H(\rightarrow bb) + Z(\rightarrow ee)$ Run: 337215 Event: 1906922941 2017-10-05 07:55:20 CEST





















Sources: ¹DOI:10.1007/BF01679868 (G. Kramer et al.), ²DOI:10.1016/0550-3213(91)90064-5 (R. Hamberg et al.), ³ArXiv:1112.1531 (T. Gehrmann et al.), ⁴ArXiv:2112.04168 (S. Zanoli et al.).







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double virtual^{1,2}







• Quick side remark on the **double virtual**...



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A COMPLETE CALCULATION OF THE ORDER α_s^2 CORRECTION TO THE DRELL-YAN K-FACTOR

R. HAMBERG and W.L. van NEERVEN*

Instituut-Lorentz, University of Leiden, P.O.B. 9506, 2300 RA Leiden, The Netherlands

II. Institut für Theoretische Physik, Universität Hamburg, D-2000 Hamburg 50, Germany

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T. MATSUURA**

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FormCalc

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Master Integrals



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virtual-real³















virtual-real³





double real





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UV poles, IR poles





double virtual^{1,2}

virtual-real³



SM





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POWHEG MiNNLO_{PS}⁴ code



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POWHEG MiNNLO_{PS}⁴ code





















































• What about **new effects**?









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$$\int \frac{dk^D}{(2\pi)^D} \frac{\dots}{((p+k)^2 - M_{\phi}^2) \dots}$$







• What about **new effects**?



$$\int \frac{dk^D}{(2\pi)^D} \frac{\dots}{((\not p + k)^2 - M_{\phi}^2) \dots}$$







• What about **new effects**?



LHC scale Λ_{LHC}

$$\int \frac{dk^D}{(2\pi)^D} \frac{\dots}{((p+k)^2 - M_{\phi}^2) \dots}$$



independent of kinematics





• What about **new effects**?



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$$\int \frac{dk^D}{(2\pi)^D} \frac{\dots}{((p+k)^2 - M_{\phi}^2) \dots}$$











• What about **new effects**?









(SMEFT)

















2. SMEFT Effects in Vh at NNLO+PS

2.1 QCD and Higgs Operators2.2 EW Operators



2.1 QCD and Higgs Operators

Sources: ¹ArXiv:2204.00663 (U. Haisch et al.).



Effects from QCD and Higgs operators have been implemented¹.

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Effects from QCD and Higgs operators have been implemented¹.

$$Q_{H\Box} = (H^{\dagger}H) \Box (H^{\dagger}H), \qquad Q_{HD} = (H^{\dagger}D_{\mu}H)^{*} (H^{\dagger}H)$$

$$Q_{bH} = y_{b}(H^{\dagger}H) \bar{q}_{L}b_{R}H, \qquad Q_{bG} = \frac{g_{s}^{3}}{(4\pi)^{2}} y_{b} \bar{q}_{L}\sigma_{\mu\nu}$$

$$Q_{HG} = \frac{g_{s}^{2}}{(4\pi)^{2}} (H^{\dagger}H) G_{\mu\nu}^{a} G^{a,\mu\nu}, \qquad Q_{3G} = \frac{g_{s}^{3}}{(4\pi)^{2}} f^{abc} G_{\mu}^{a},$$

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 $H^{\dagger}D^{\mu}H),$

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Effects from QCD and Higgs operators have been implemented¹.

$$Q_{H\Box} = (H^{\dagger}H) \Box (H^{\dagger}H), \qquad Q_{HD} = (H^{\dagger}D_{\mu}H)^{*} (H^{\dagger}),$$

$$Q_{bH} = y_{b}(H^{\dagger}H) \bar{q}_{L}b_{R}H, \qquad Q_{bG} = \frac{g_{s}^{3}}{(4\pi)^{2}} y_{b} \bar{q}_{L}\sigma_{\mu}$$

$$Q_{HC} = \frac{g_{s}^{2}}{(4\pi)^{2}} (H^{\dagger}H) G_{\mu\nu}^{a} G^{a,\mu\nu}, \qquad Q_{3G} = \frac{g_{s}^{3}}{(4\pi)^{2}} f^{abc} G_{\mu}^{a}$$

Source: <u>ArXiv:2204.00663</u> (U. Haisch et al.).

• **Higgs operators** factorize:

$$\Gamma(h \to b\bar{b})_{\rm SMEFT}^{\rm NNLO, fac} = (1 + 2c_{\rm fac}) \Gamma(h \to b\bar{b})_{\rm SM}^{\rm NNLO} ,$$

$$\sigma(pp \to Zh)_{\rm SMEFT}^{\rm NNLO} = (1 + 2c_{\rm kin}) \sigma(pp \to Zh)_{\rm SM}^{\rm NNLO},$$

Sources: ¹ArXiv:2204.00663 (U. Haisch et al.).

 $H^{\dagger}D^{\mu}H),$

 $_{\mu\nu}T^{a}b_{R}HG^{a,\mu\nu}$,

 $^{\nu}_{L},^{\nu}G^{b,\sigma}_{\nu}G^{c,\mu}_{\sigma},$











Effects from QCD and Higgs operators have been implemented¹.



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QCD operators do not









2.1 QCD and Higgs Operators

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• The effects from the QCD operators lead to non-trivial shape changes that depend on the jet clustering parameter R^1 .

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Source: ArXiv:0802.1189 (M. Cacciari et al.).


































• We would like to extend the existing implementation to include electroweak operators as well.



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$Q_{HW} = H^{\dagger} H \ W^{I}_{\mu\nu} W^{I,\mu\nu} ,$	$Q_{HB} = H^{\dagger} H B_{\mu\nu} B^{\mu\nu} ,$
$Q_{HWB} = H^{\dagger} \tau^{I} H \ W^{I}_{\mu\nu} B^{\mu\nu} ,$	
$Q_{Hl}^{(1)} = \left(H^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} H \right) \left(\bar{l} \gamma^{\mu} l \right) ,$	$Q_{Hl}^{(3)} = \left(H^{\dagger} i \overset{\leftrightarrow}{D_{\mu}^{I}} H \right) \left(\bar{l} \tau^{I} \gamma^{\mu} l \right) ,$
$Q^{(1)}_{Hq} = \left(H^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} H \right) \left(\bar{q} \gamma^{\mu} q \right) ,$	$Q_{Hq}^{(3)} = \left(H^{\dagger} i \overset{\leftrightarrow}{D_{\mu}^{I}} H \right) \left(\bar{q} \tau^{I} \gamma^{\mu} q \right) ,$
$Q_{Hu} = \left(H^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} H \right) \left(\bar{u} \gamma^{\mu} u \right) ,$	$Q_{Hd} = \left(H^{\dagger} i \overset{\leftrightarrow}{D}_{\mu} H \right) \left(\bar{d} \gamma^{\mu} d \right) ,$
$Q_{He} = \left(H^{\dagger}i \overset{\leftrightarrow}{D}_{\mu} H\right) \left(\bar{e} \gamma^{\mu} e\right) ,$	

Source: ArXiv:1008.4884 (B. Grzadkowski et al.).









• We would like to extend the existing implementation to include **electroweak operators** as well.

$$\begin{split} Q_{HW} &= H^{\dagger}H \; W_{\mu\nu}^{I} W^{I,\mu\nu} , \qquad Q_{HB} = H^{\dagger}H \; B_{\mu\nu}B^{\mu\nu} , \\ Q_{HWB} &= H^{\dagger}\tau^{I}H \; W_{\mu\nu}^{I}B^{\mu\nu} , \\ Q_{Hl}^{(1)} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right) \left(\bar{l}\gamma^{\mu}l\right) , \qquad Q_{Hl}^{(3)} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H\right) \left(\bar{l}\tau^{I}\gamma^{\mu}l\right) , \\ Q_{Hq}^{(1)} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right) \left(\bar{q}\gamma^{\mu}q\right) , \qquad Q_{Hq}^{(3)} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H\right) \left(\bar{q}\tau^{I}\gamma^{\mu}q\right) , \\ Q_{Hu} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right) \left(\bar{u}\gamma^{\mu}u\right) , \qquad Q_{Hd} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right) \left(\bar{d}\gamma^{\mu}d\right) , \\ Q_{He} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right) \left(\bar{e}\gamma^{\mu}e\right) , \end{split}$$

Source: ArXiv:1008.4884 (B. Grzadkowski et al.).

• Fermion operators factorize:











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 $\left(\ldots\gamma^{\mu}P_{L}+\ldots\gamma^{\mu}P_{R}\right)$

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• Fermion operators factorize:





• Gauge operators do not factorize:







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 $\left(\ldots\gamma^{\mu}P_{L}+\ldots\gamma^{\mu}P_{R}\right)$

$$\begin{split} Q_{HW} &= H^{\dagger}H \; W_{\mu\nu}^{I}W^{I,\mu\nu} , \qquad Q_{HB} = H^{\dagger}H \; B_{\mu\nu}B^{\mu\nu} , \\ Q_{HWB} &= H^{\dagger}\tau^{I}H \; W_{\mu\nu}^{I}B^{\mu\nu} , \\ Q_{Hl}^{(1)} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right)\left(\bar{l}\gamma^{\mu}l\right) , \qquad Q_{Hl}^{(3)} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H\right)\left(\bar{l}\tau^{I}\gamma^{\mu}l\right) , \\ Q_{Hq}^{(1)} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right)\left(\bar{q}\gamma^{\mu}q\right) , \qquad Q_{Hq}^{(3)} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H\right)\left(\bar{q}\tau^{I}\gamma^{\mu}q\right) , \\ Q_{Hu} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right)\left(\bar{u}\gamma^{\mu}u\right) , \qquad Q_{Hd} = \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right)\left(\bar{d}\gamma^{\mu}d\right) , \\ Q_{He} &= \left(H^{\dagger}i\overleftrightarrow{D}_{\mu}H\right)\left(\bar{e}\gamma^{\mu}e\right) , \end{split}$$

Source: ArXiv:1008.4884 (B. Grzadkowski et al.).

• Fermion operators factorize:





• Gauge operators do not factorize:







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$$\frac{ig_{q_L}^Z g_{l_L}^Z}{\left(s_{12} - M_Z^2\right) \left(s_{34} - M_Z^2\right)} \left(\frac{gM_Z}{c_w}\right) \langle 23 \rangle [41]$$

















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$$\langle 23 \rangle [41] = \frac{1}{2} \langle 3 | \gamma^{\mu} | 4] \langle 1 | \gamma_{\mu} | 2]$$

















 $\langle 1 | \gamma_{\mu} | 2]$

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 $\langle 3 \left| \gamma^{\mu} q \right| 3 \rangle [34] + \langle 34 \rangle [4 \left| q \gamma^{\mu} \right| 4]$









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 $(-\langle 12 \rangle \langle 34 \rangle [41]^2 + \langle 23 \rangle^2 [21][43])$ $(s_{12} - M_7^2)$







• Fortran implementation...





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* pow(K%Zb(i4,i1),2) & + pow(K%Za(i2,i3),2) & * K%Zb(i2,i1)*K%Zb(i4,i3))





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	Checks Virtual		
Ratio Vi	rtual (ours/existing)):	0.99999999999971267
Chec	ks correlated Born		
i, j:	1	1	
Ratio:	1.000000008927112		
i, j:	1	2	
Ratio:	1.000000008927112		
i, j:	1	6	
Ratio:	1.000000008927112		
mu, nu:	1		1
Ratio:	1.0000000347233580		
mu, nu:	1		2
Ratio:	0.99999999037223986		
mu, nu:	2		1
Ratio:	0.99999999037223986		
mu, nu:	2		2
Ratio:	1.0000000309149535		
i, j:	2	1	
Ratio:	1.000000008927112		
i, j:	2	2	
Ratio:	1.000000008927110		
i, j:	2	6	
Ratio:	1.000000008927112		
i, j:	6	1	
Ratio:	1.000000008927112		
i, j:	6	2	
Ratio:	1.000000008927112		
i, j:	6	6	
Ratio:	1.000000008927110		
	Checks Born		
Ratio Bo	rn (ours/existing):		1.000000008927110





• Fortran implementation...



• **Next step:** generate phenomenology plots.

* pow(K%Zb(i4,i1),2) & + pow(K%Za(i2,i3),2) & * K%Zb(i2,i1)*K%Zb(i4,i3))

	Checks Virtual		
Ratio Vi	rtual (ours/existing)):	0.99999999999971267
Chec	ks correlated Born		
i, j:	1	1	
Ratio:	1.000000008927112		
i, j:	1	2	
Ratio:	1.000000008927112		
i, j:	1	6	
Ratio:	1.000000008927112		
mu, nu:	1		1
Ratio:	1.0000000347233580		
mu, nu:	1		2
Ratio:	0.99999999037223986		
mu, nu:	2		1
Ratio:	0.99999999037223986		
mu, nu:	2		2
Ratio:	1.0000000309149535		
i, j:	2	1	
Ratio:	1.000000008927112		
i, j:	2	2	
Ratio:	1.000000008927110		
i, j:	2	6	
Ratio:	1.000000008927112		
i, j:	6	1	
Ratio:	1.000000008927112		
i, j:	6	2	
Ratio:	1.000000008927112		
i, j:	6	6	
Ratio:	1.000000008927110		
	Checks Born		
Ratio Bo	rn (ours/existing):		1.000000008927110





3. Conclusions



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• The associated Higgs production (Vh) channel is interesting phenomenologically, since it allows to measure the Higgs couplings precisely.

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- Indirect effects of heavy new physics at the LHC can be parameterised in a modelindependent way using the Standard Model Effective Field Theory (SMEFT).
 - The effects of **QCD and Higgs operators** in Vh have been implemented at NNLO+PS in the MiNNLO_{PS} framework.
 - We are extending this to include the effects of **electroweak operators** as well.

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 - The effects of QCD and Higgs operators in Vh have been implemented at NNLO+PS in the MiNNLO_{PS} framework.
 - We are extending this to include the effects of **electroweak operators** as well.
- Using spinor-helicity amplitudes, the electroweak SM helicity structure can be replaced by new helicity structures without having to recompute the QCD corrections.

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Thank you for your attention!