Master Thesis Presentation "Long-distance contributions in rare semileptonic B decays" Supervisor: Prof. Gino Isidori

Nikolaos Kalntis, MSc Physics ETH Zurich

Physics Division, Theory Group Lawrence Berkeley National Laboratory (LBNL) & UC Berkeley

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- Introduction
- Mathematical Framework

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- One and Two Particle long-distance contributions

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- Projects beyond my Master Thesis

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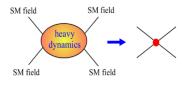
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- Difficult to be calculated: non-perturbative character.

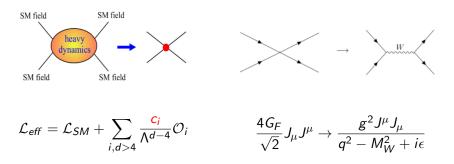
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Theoretical Framework: $B \rightarrow K \ell \ell$

To search for New Physics in $B \rightarrow K\ell\ell$, we use the EFT framework.

$$\mathcal{H}_{eff} = \frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i C_i \mathcal{O}_i \,. \tag{1}$$

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 $\mathcal{H}_{c\bar{c}}(q^2)$: Non - local hadronic contributions.

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We calculate them from data using dispersion relations.

Subtracted dispersion relation for the hadronic contributions

$$\Delta \mathcal{H}_{c\bar{c}}(q^2) \equiv \frac{q^2}{\pi} \int_{s_0}^{\infty} ds \frac{\rho_{c\bar{c}}(s)}{s(s-q^2)}, \qquad \rho_{c\bar{c}}(s) = \frac{1}{\pi} Im[\mathcal{H}_{c\bar{c}}(s)].$$
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We want to have a quantitative estimate of these contributions.

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$$\mathcal{B}(B \to KV_j) \times \mathcal{B}(V_j \to \ell^+ \ell^-) = \tau_B \frac{G_F^2 \alpha^2 |\lambda_t|^2}{128\pi^5} \times \int_{4m_{\ell^2}}^{(m_B - m_K)^2} dq^2 k(q^2)^3 [\beta(q^2) - \frac{1}{3}\beta(q^2)^3] |f_+(q^2)|^2 |\eta_j|^2 \left| \frac{q^2}{m_j^2} A_j^{1P}(q^2) \right|^2.$$
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J/Ψ	$(1.791 \pm 1.231) imes 10^4$	
$\Psi(2s)$	$(1.834 \pm 1.346) imes 10^3$	
Ψ(3770)	4.358 ± 0.389	
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- $\eta_{J/\Psi}$, $\eta_{\Psi(2s)} \gg \mathcal{O}(1)$: Large violations of quark-hadron duality only for these two resonances \rightarrow Perturbative calculations can not be trusted.

Proceeding in a similar way, the 2P contribution is

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Compatible with $\eta_{D,\bar{D}} \lesssim \mathcal{O}(1)$ (perturbative arguments).

Nikolaos Kalntis (LBNL/UC Berkeley)

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- These methods could be used for the predictions of the η values for other similar decays, where the η factors have not been yet calculated in the bibliography, given the lack of reliable perturbative estimates.
- LHC Run 3 coming soon! Hope to collect more data and clarify whether the LFU violations is a statistical fluctuation or not.

During my studies: Co-authored two publications on the search for the QCD critical point $^{1\ 2}$

²N. G. Antoniou, F. K. Diakonos, **N. Kalntis** and A. Kanargias, "Higher cumulants of baryon number in critical QCD." *Nuclear Physics A* 986, (167-174 (2019) = • = • •

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- <u>Future</u>: Integration of CTs and extension to the non-abelian case (under progress).

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- Currently writing a paper with my collaborators from LBNL/ UC Berkeley:

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- One paper so far ³: Proof that the gravitational Schwinger effect prevents the Big-Rip singularity in the far-future of a universe dominated by phantom energy. Analysis done **only** in the context of GR (in contrast to previous analysis assuming modified gravity).
- Currently writing a paper with my collaborators from LBNL/ UC Berkeley: Constraining beyond the ACDM models (DM interactions) using Fisher forecasting.

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Thank you very much for your time!

Additional/Explanatory Slides!

One-Particle and Two-Particle States	η value
J/Ψ	$(1.791 \pm 1.231) imes 10^4$
$\Psi(2s)$	$(1.834 \pm 1.346) imes 10^3$
Ψ(3770)	4.358 ± 0.389
Ψ(4040)	1.538 ± 0.622
DD	$\leq 0.247 \pm 0.216$
DD*	$\leq 1.014 \pm 0.358$

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$$\rho^{1P} \propto \mathcal{M}_1(B \to KV) \mathcal{M}_2(V \to \ell \ell).$$
(9)

$$\rho^{2P} \propto \int \frac{d^4 \ell_1 d^4 \ell_2}{E_1 E_2} \delta^4(s - \ell_1 - \ell_2) \mathcal{M}_1(B \to KMM') \mathcal{M}_2(MM' \to \ell\ell) \,. (10)$$

DD case:

 $\begin{aligned} \mathcal{M}_1(B \to \mathsf{K}\mathsf{D}\mathsf{D}) &= & \mathsf{D}\mathsf{D}^* \text{ case:} \\ \alpha_1 + \beta_1 p \cdot q + \gamma p \cdot \ell_1, & \mathcal{M}_1(B \to \mathsf{K}\mathsf{D}\mathsf{D}^*) = \beta p \cdot \epsilon(\mathsf{D}^*), \\ (-ie)(\ell_1 - \ell_2)^{\mu}: \ \gamma \mathsf{D}\mathsf{D} \text{ vertex (scalar } C \ \epsilon(\mathsf{D}^*)^{\mu}: \ \gamma \mathsf{D}\mathsf{D}^* \text{ vertex.} \\ \\ \mathsf{QED}) \end{aligned}$

In principle, the non-local hadronic contributions takes the following form

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$$i \int d^4 x e^{iq \cdot x} \langle \mathcal{K}(p) | \mathcal{T}[J^{em}_{\mu}(x), C_1 \mathcal{O}^c_1(0) + C_2 \mathcal{O}^c_2(0)] | \mathcal{B}(p+q) \rangle \qquad (11)$$
$$= \mathcal{H}_{c\bar{c}}(q^2)((p \cdot q)q_{\mu} - q^2 p_{\mu}). \qquad (12)$$

The leading contribution to $\rho_{c\bar{c}}(s)$ is given by the one-particle intermediate hadronic states, with $c\bar{c}$ valence quarks (charmonium resonances: J/Ψ and its excited states $\Psi(2s)$, ...).

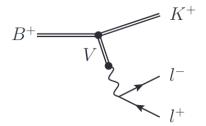


Figure: One-particle intermediate hadronic states.

The next-to-leading contribution to $\rho_{c\bar{c}}(s)$ is given by the two-particle intermediate hadronic states (one-loop).

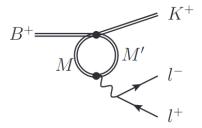


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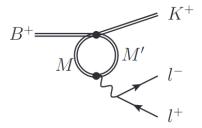


Figure: Two-particle intermediate hadronic states.

Most important contributions are from DD and the D^*D mesons in the loop (D: scalar bosons, D^* : vector bosons).

Nikolaos Kalntis (LBNL/UC Berkeley)

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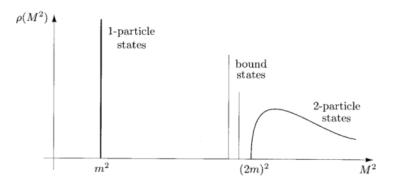
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Integral weights lower energies more heavily and lessens the influence of the high energy region.

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18.07.2022



$$\rho(M^2) = Z\delta(M^2 - m^2) + \sigma(M^2)\theta(M^2 - s_0).$$
(15)

$$\mathcal{H}(s) \propto \frac{1}{\pi} \int_0^\infty ds' \frac{\rho(s')}{s' - s - i\epsilon} = \frac{\rho_0}{m^2 - s - i\epsilon} + \int_{s_0}^\infty ds' \frac{\sigma(s')}{s' - s - i\epsilon} \,. \tag{16}$$