

Edge states and the Luttinger liquid in M-theory

Constantin Greubel

Max-Planck-Institut für Physik (Werner Heisenberg Institut)

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- Motivation
- Introduction
 - What is M-theory?
 - What is the ABJM model?
- Model and Calculations
 - Reduction to Type IIA SUGRA
 - Brane embedding
 - Finite chemical potential
- Results
 - Spectral function
 - Phase diagram
- Condensed Matter

Lower dimensional CMT with String Theory?

Condensed matter systems

- Compute quantities in condensed matter systems with strong coupling

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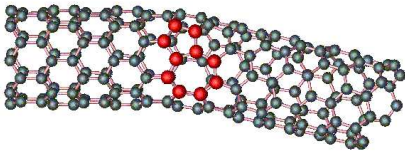
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- Nanotechnology makes manipulations on atomic scale possible
- physics on lower dimensional hypersurfaces:
charges and currents confined, but fields propagate all dim's

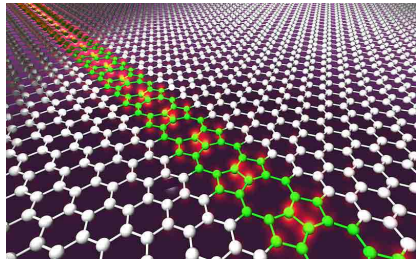
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[V. H. Crespi]



[Y. Lin, USF]

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- no complete background-independent formulation yet!
- UV-completion of SUGRA
- 11-dimensional theory
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From string theory we can uplift to M-theory

ABJM

- dubbed after Aharony, Bergman, Jafferis and Maldacena [0806.1218]
- holographic principle, gauge/gravity duality
- N_c M2-branes probe a $\mathbb{C}^4/\mathbb{Z}_k$ singularity

(Super-) Gravity side

limit for large N_c where 't Hooft coupling $\lambda = N_c/k$ is kept fixed:

$AdS_4 \times S^7/\mathbb{Z}_k$ in SUGRA

limit for $k \rightarrow \infty$ type IIA string theory on $AdS_4 \times \mathbb{C}P^3$

Gauge theory side

1 + 2 dimensional theory with $\mathcal{N} = 6$ SUSY

Conformal field theory i. e. no scales

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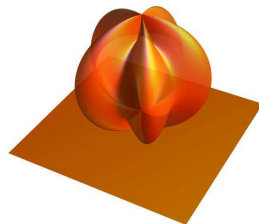
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Conformal field theory i. e. no scales

Chern-Simons terms \rightarrow often occur in condensed matter systems!

M-theory content

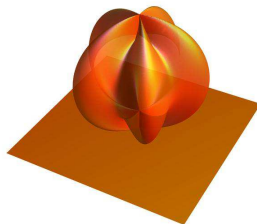
Coincident Stack of N_c M2-branes
transverse space is $\mathbb{C}^4/\mathbb{Z}_k$ orbifold



[J. Davey]

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Dimensional reduction to type IIA

- Zoom to near-horizon $\Rightarrow AdS_4 \times S^7/\mathbb{Z}_k$ SUGRA
- Radius of compact direction $R/k\ell_p = 2^{5/6}\pi^{1/3}(N_c k)^{1/6}/k$
- sending the Chern-Simons level $k \rightarrow \infty$ gives type IIA string theory
- Branes become: M2 \rightarrow D2

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From now, working in the type IIA background

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String theory: quark $\hat{=}$ string between D6 and D2 branes

String theory: quark mass $\hat{=}$ length of this string

Holographic principle: length at boundary

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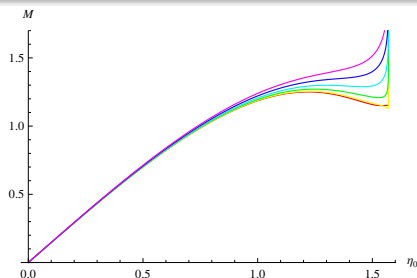
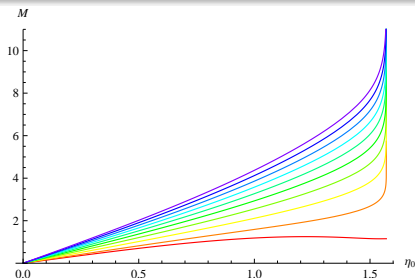
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Brane embedding

Introduce $U(1)$ gauge field A_μ living on D6-brane

Gauge field A_μ gauged to have A_t the only nonvanishing component

Generalized action

action: Minimal area law

$$S \propto \int d\sigma^{1+6} \sqrt{-\det(\mathcal{P}[g])}$$

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$$S_{\text{D6}} \propto \int d\chi \chi^{-4} \left[2 \cos \eta \sqrt{1 + \eta'^2 \chi^2 (1 - \chi^3) - \chi^4 (\partial_\chi A_t)^2} - \sin^2 \eta \right]$$

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Before embedding can be computed: Solve for $\partial_\chi A_t$

Chemical potential

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Asymptotics of gauge field

At the boundary of AdS ($\chi \rightarrow 0$)

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EOM for gauge field

Solving the equation of motion for the gauge field

$$\partial_\chi A_t = d \sqrt{\frac{1 + \eta'^2 \chi^2 (1 - \chi^3)}{4 \cos^2 \eta + d^2 \chi^4}} \quad \mu = -d \int_0^1 d\chi \sqrt{\frac{1 + \eta'^2 \chi^2 (1 - \chi^3)}{4 \cos^2 \eta + d^2 \chi^4}}$$

Scalar excitations

Scalar fluctuations

Fluctuations of the brane in one component $\eta(\chi)$
scalar mode $\eta(\chi) \rightarrow \eta(\chi) + \delta\eta(\sigma^i)$

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vector mode $A_\mu(\chi) \rightarrow A_\mu(\chi) + \delta A_\mu(\sigma^i)$

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Fluctuations of the brane in one component scalar mode $\eta(\chi) \rightarrow \eta(\chi) + \delta\eta(\sigma^i)$

Vector fluctuations

Fluctuations of the gauge field vector mode $A_\mu(\chi) \rightarrow A_\mu(\chi) + \delta A_\mu(\sigma^i)$

Equations of motion for fluctuations

Plug into DBI-action and expand to 2nd order in the fluctuations
Equations of Motion are partial differential equations

Fourier trafo $\delta\eta(\chi, \vec{x}) = \int \frac{d^4k}{(2\pi)^4} e^{i\vec{k}\vec{x}} \delta\eta(\chi, \vec{k})$ with $\vec{k} = (\omega, 0, q)$

Linear eom $\partial_\chi^2(\delta\eta) + \mathfrak{C}_1\partial_\chi(\delta\eta) + \mathfrak{C}_2(\delta\eta) = 0$

The spectral function

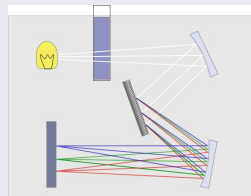
Definition

The spectral function is related to the retarded correlator G^R by

$$\Im \mathfrak{R} = -2\text{Im} G^R \quad \text{where}$$

$$G^R(k) = -i \int d^4x \, e^{ikx} \theta(x^0) \langle [J(x), J(0)] \rangle$$

with $G^R = \frac{\delta^2 \mathcal{S}_{\text{SUGRA}}}{\delta \tilde{A}^2} \Big|_{\text{boundary}}$ [\[Son, Starinets, '02\]](#)



The spectral function

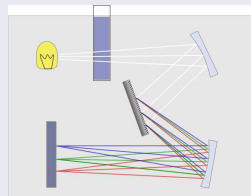
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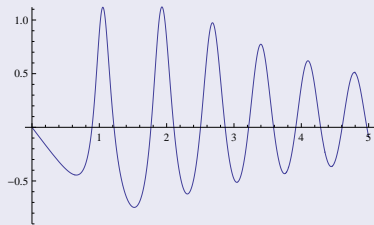
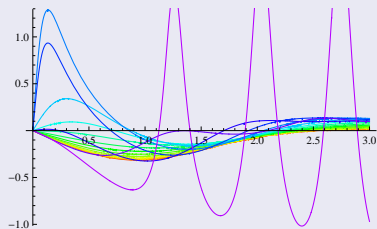


In our coordinates

$$\Re(\omega, \vec{q} = 0) \propto \text{Im} \frac{\partial_x \eta}{\eta} \Big|_{\text{boundary}}$$

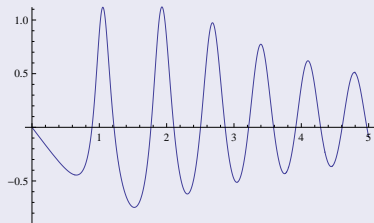
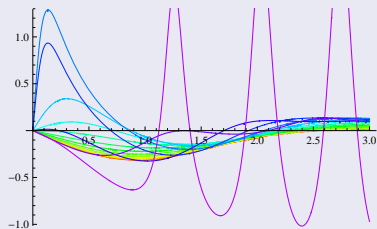
Numerical Calculation

Spectral function for vanishing density



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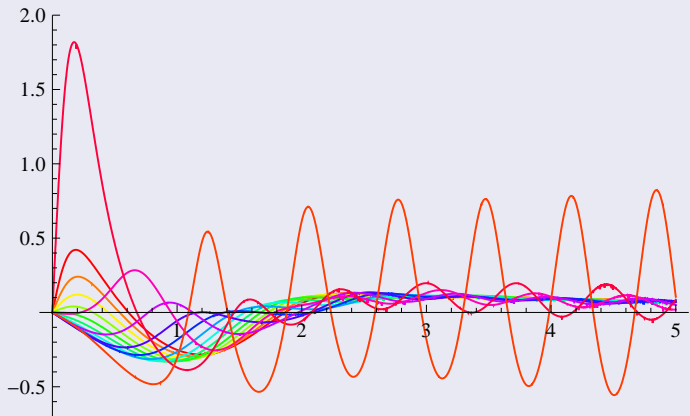
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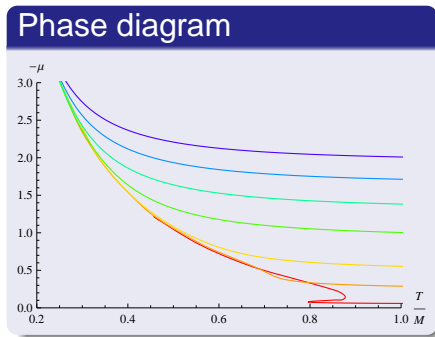
- higher quark mass gives very pronounced peaks
- expected from earlier studies in D3/D7

Numerical Calculation

Spectral function for non-vanishing density

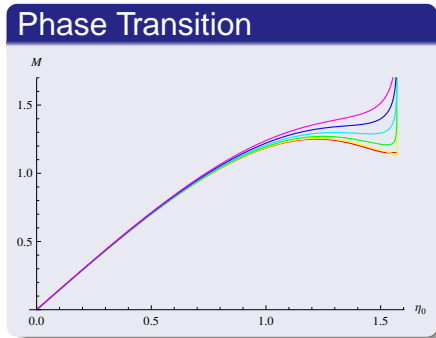
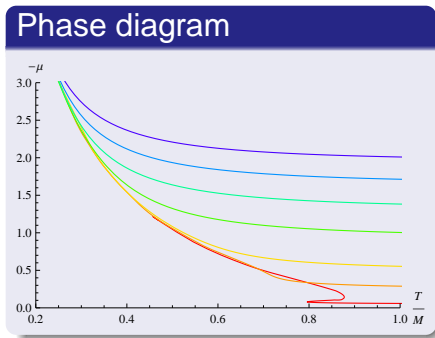


Phase diagram



PROBLEM:
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SOLUTION:

comparison with the embeddings indicates that instability lies at the phase transition

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Main idea

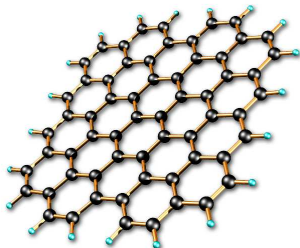
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But no description with Fermi-Landau possible!
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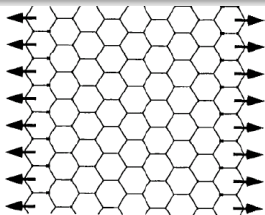


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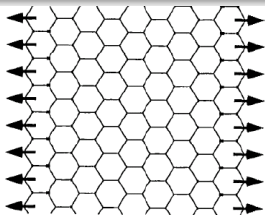


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- AdS_4 : codimension 0,1 or 2 defect
- here: codimension 0 defect
- calculated in [\[Erdmenger et al., 0909.3845\]](#)

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- Embedding with phase transition
- Chemical potential
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Outlook

- Connection to condensed matter systems
- Compute further observables: quark condensate, . . .
- Build in chirality

Conclusion

Thank you for your attention!