

# The Anomalous Magnetic Moment of the Muon

Providing an Explanation using a Visible QCD Axion

Joachim Weiss

Master's Thesis with U. Nierste, T. Schwetz-Mangold, R. Ziegler at KIT

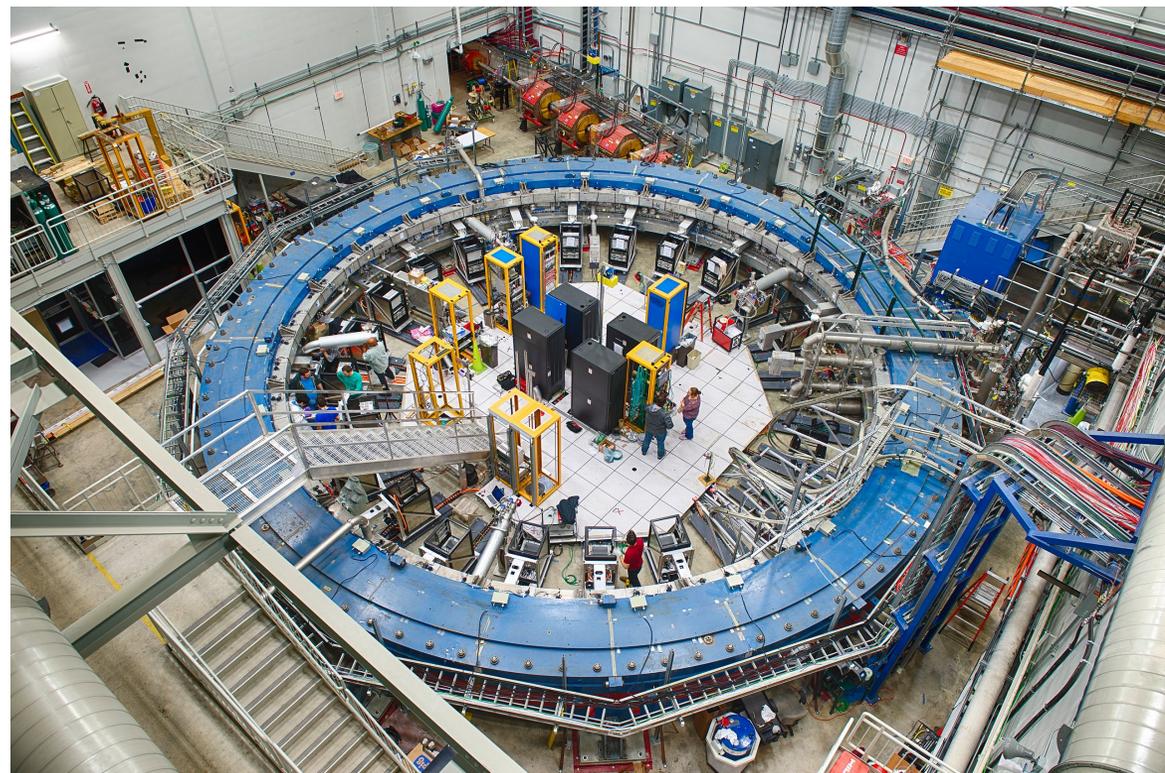
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- (2) Theoretical Calculation of  $(g - 2)_\mu$
- (3) Explaining the Anomaly using a QCD Axion
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- (6) The Model
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# (1) Experiments at BNL and FNAL

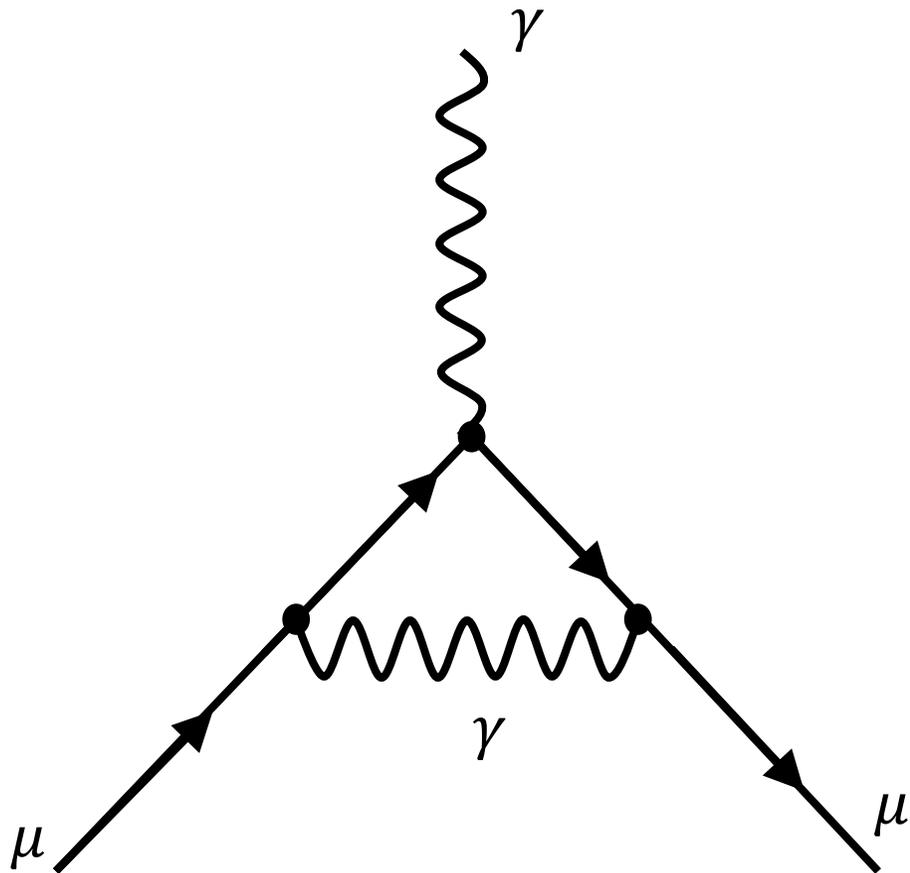
- Magnetic moment  $\vec{\mu}$  of the muon
  - $\vec{\mu} = -g \frac{q}{2m} \vec{L}$
- Measured at Brookhaven National Laboratory (BNL) and Fermi National Accelerator Laboratory (FNAL)
- $a_{\mu} = \left( \frac{g-2}{2} \right)_{\mu}$

From Reidar Hahn - <https://vms.fnal.gov/asset/detail?recid=1950114>, CC BY-SA 4.0, <https://commons.wikimedia.org/w/index.php?curid=58326280>

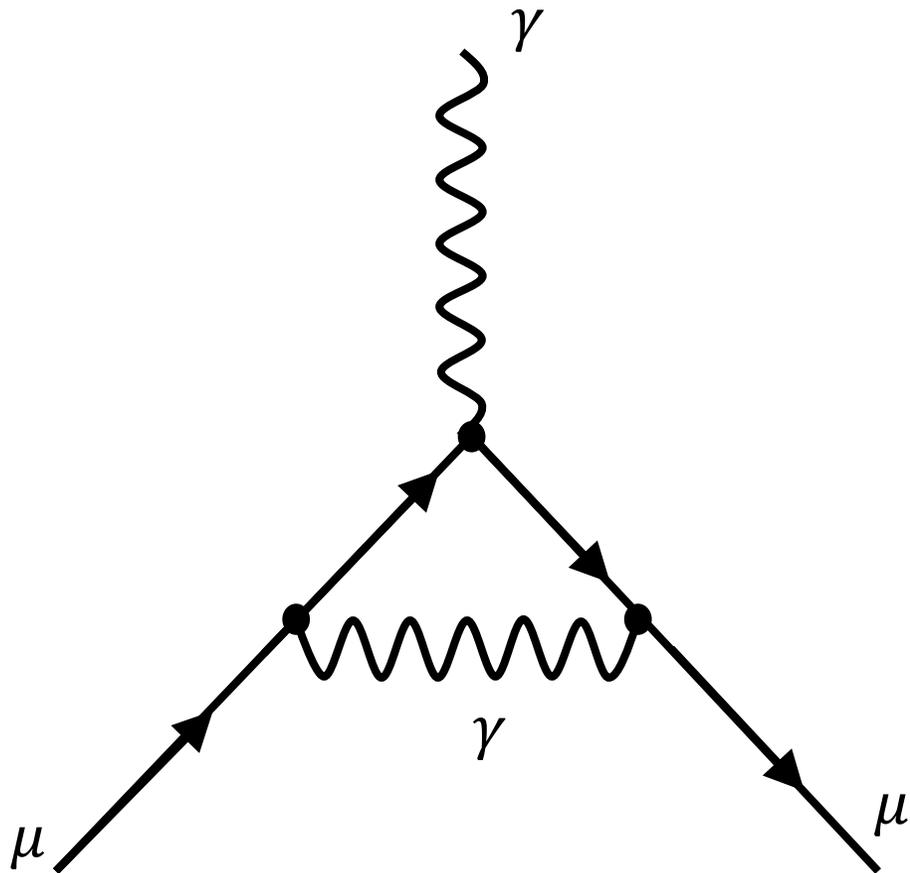


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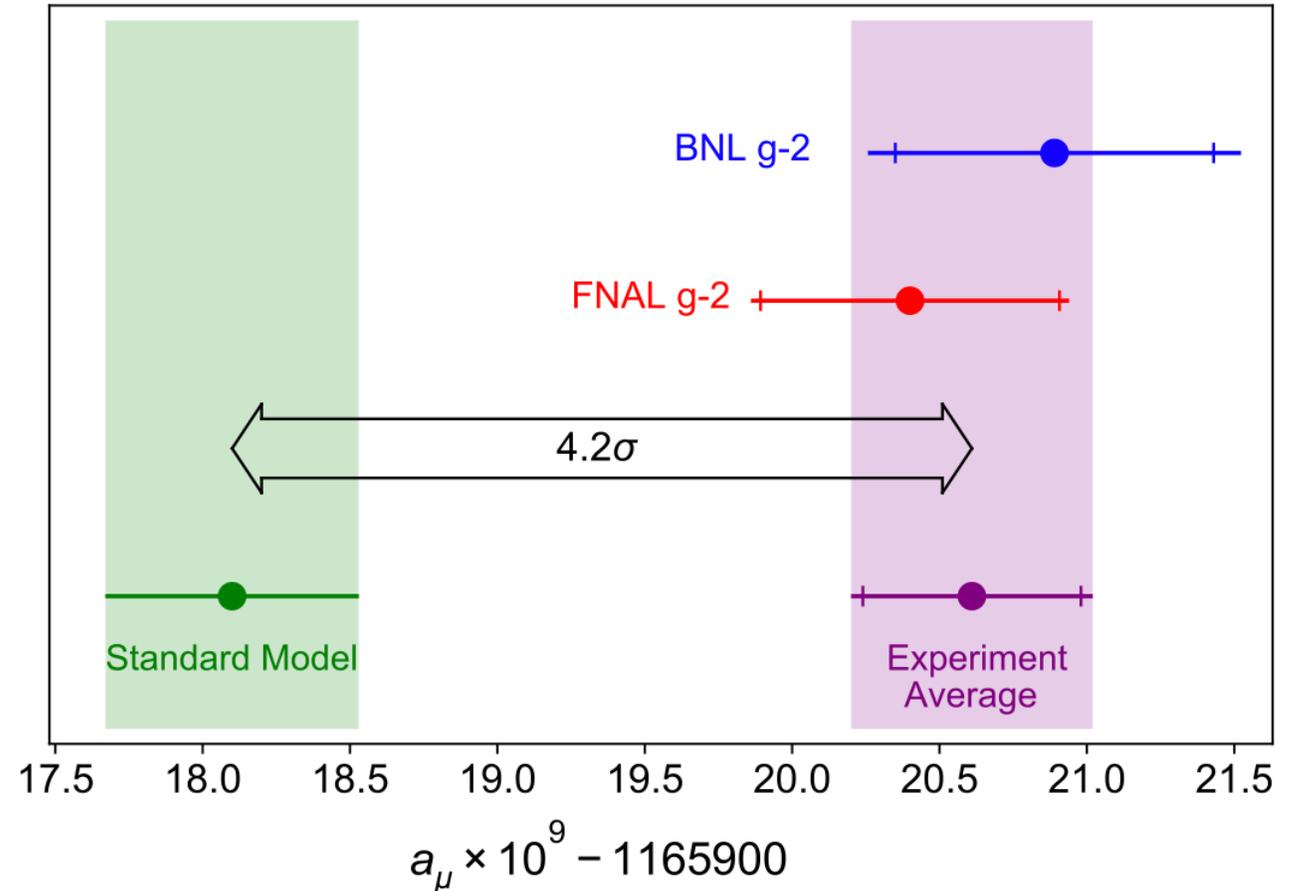
## (2) Theoretical Calculation of $(g - 2)_\mu$



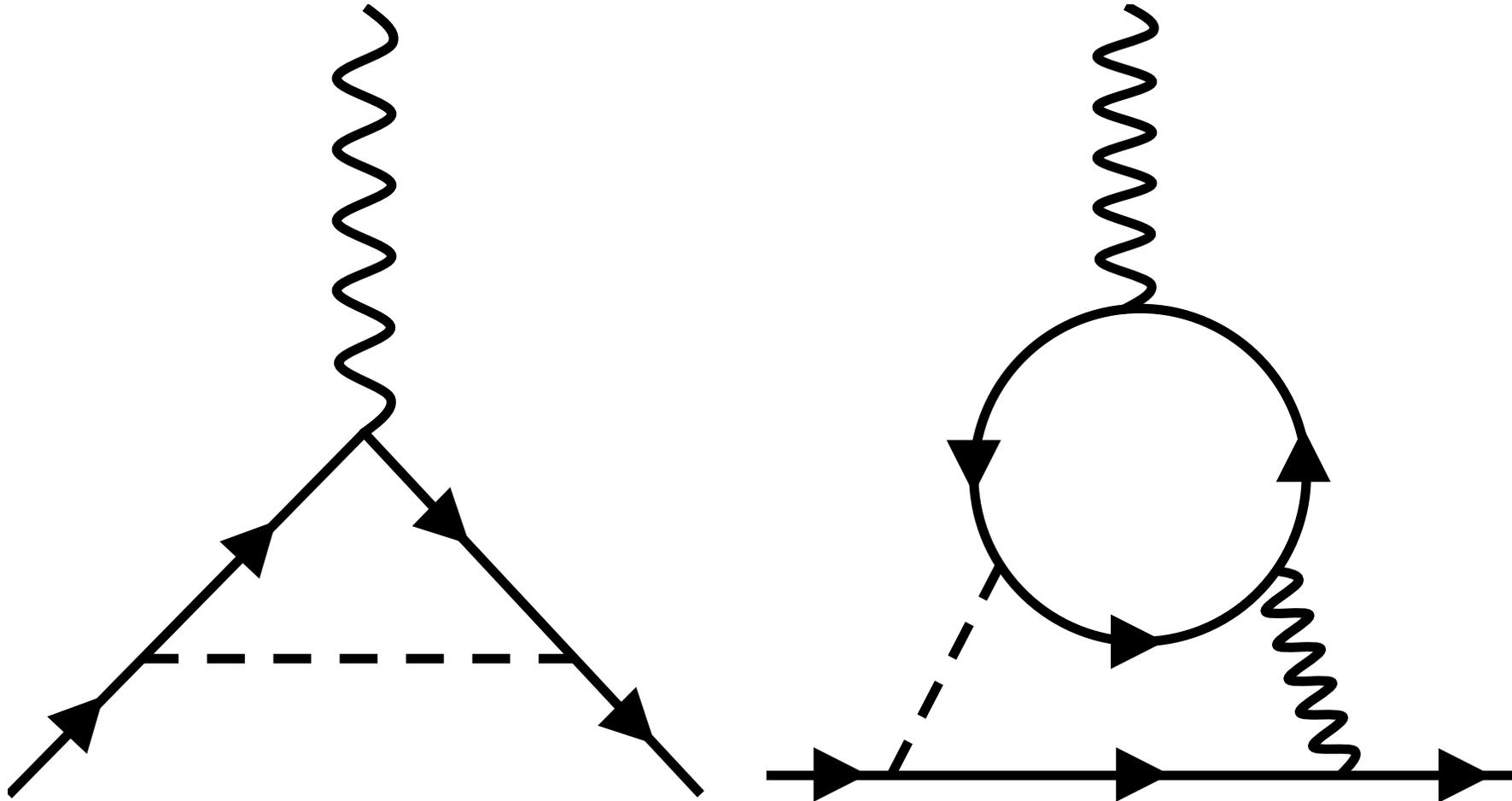
# (2) Theoretical Calculation of $(g - 2)_\mu$



B. Abi et al.: arXiv:2104.03281



# (3) Explaining the Anomaly



# (4) The QCD Axion

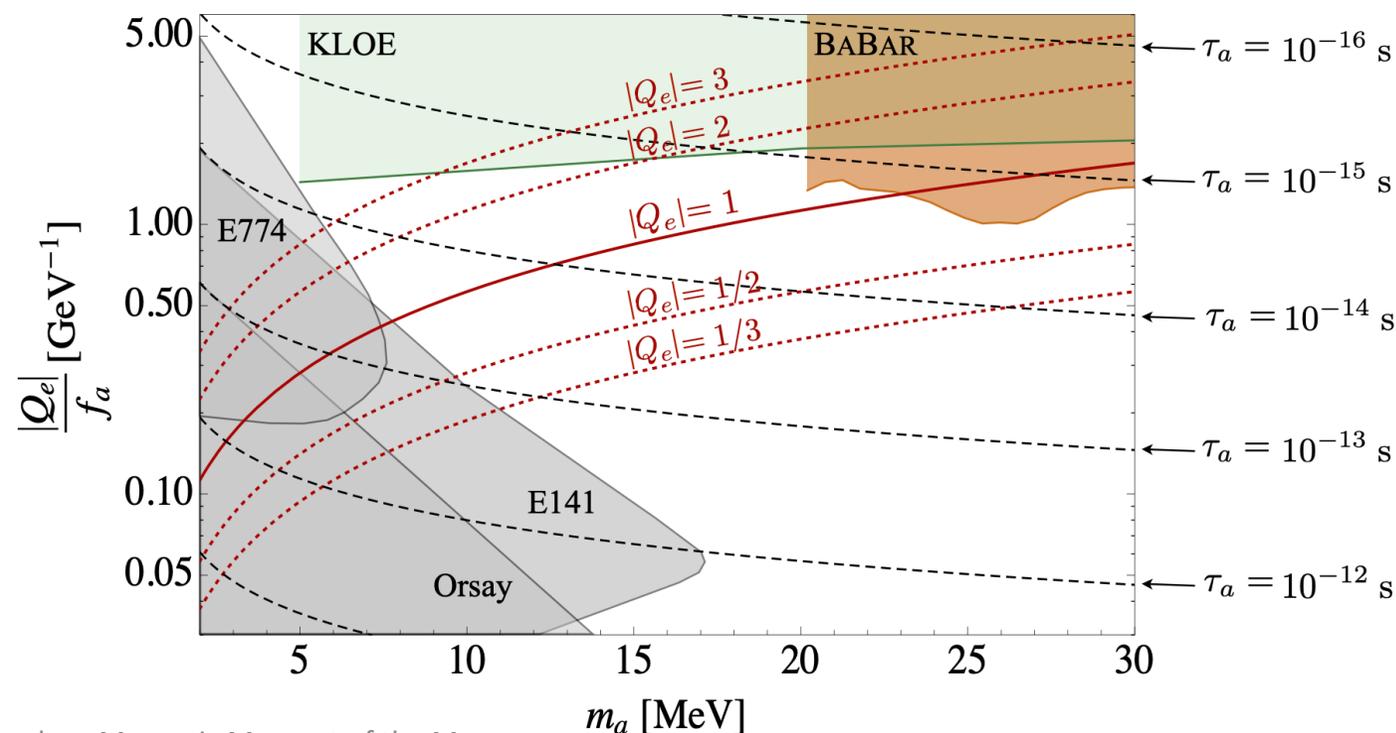
- Pseudoscalar particle
- Coupling to fermions, effective coupling to photons
- Broken  $U(1)$  symmetry to solve the “strong CP puzzle” (R. Peccei, H. Quinn)

$$\mathcal{L}_{QCD} \supset \theta \frac{g^2}{32\pi^2} F_a^{\mu\nu} \tilde{F}_{a,\mu\nu}, \quad \bar{\theta} < 10^{-10}$$

# (4) The Visible QCD Axion

- Decays (promptly) into SM fermions (“visible”)
- D. Alves, N. Weiner (arXiv:1710.03764):

- $m_a \sim \mathcal{O}(\text{MeV})$
- Couples to first generation
  - “Pionphobic”
  - “Nucleophobic”
- Avoids constraints
  - Beam-dump experiments
  - Collider experiments



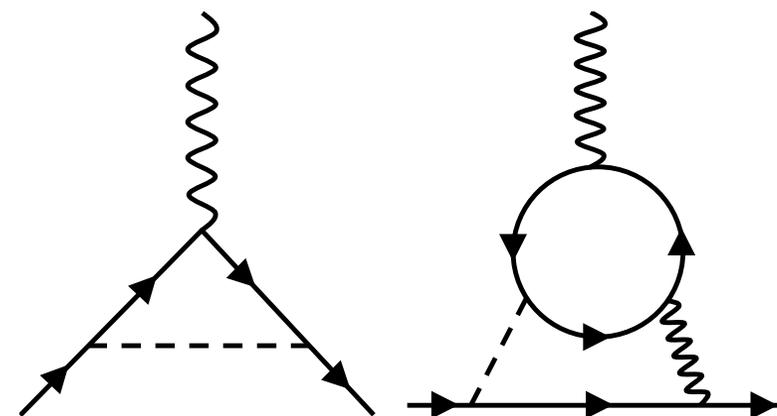
# (4) Goals of the new Model

- Address  $(g - 2)_\mu$ 
  - $\Delta a_\mu = (25.1 \pm 5.9) \cdot 10^{-10} (4,2\sigma)$
  - Use 1-loop and 2-loop contributions
  - Fulfil constraints by beam-dump and collider experiments
- Address strong CP puzzle
  - Identify pseudoscalar mediator  $a$  as visible QCD axion
  - Pionphobic, nucleophobic, promptly decaying, ...

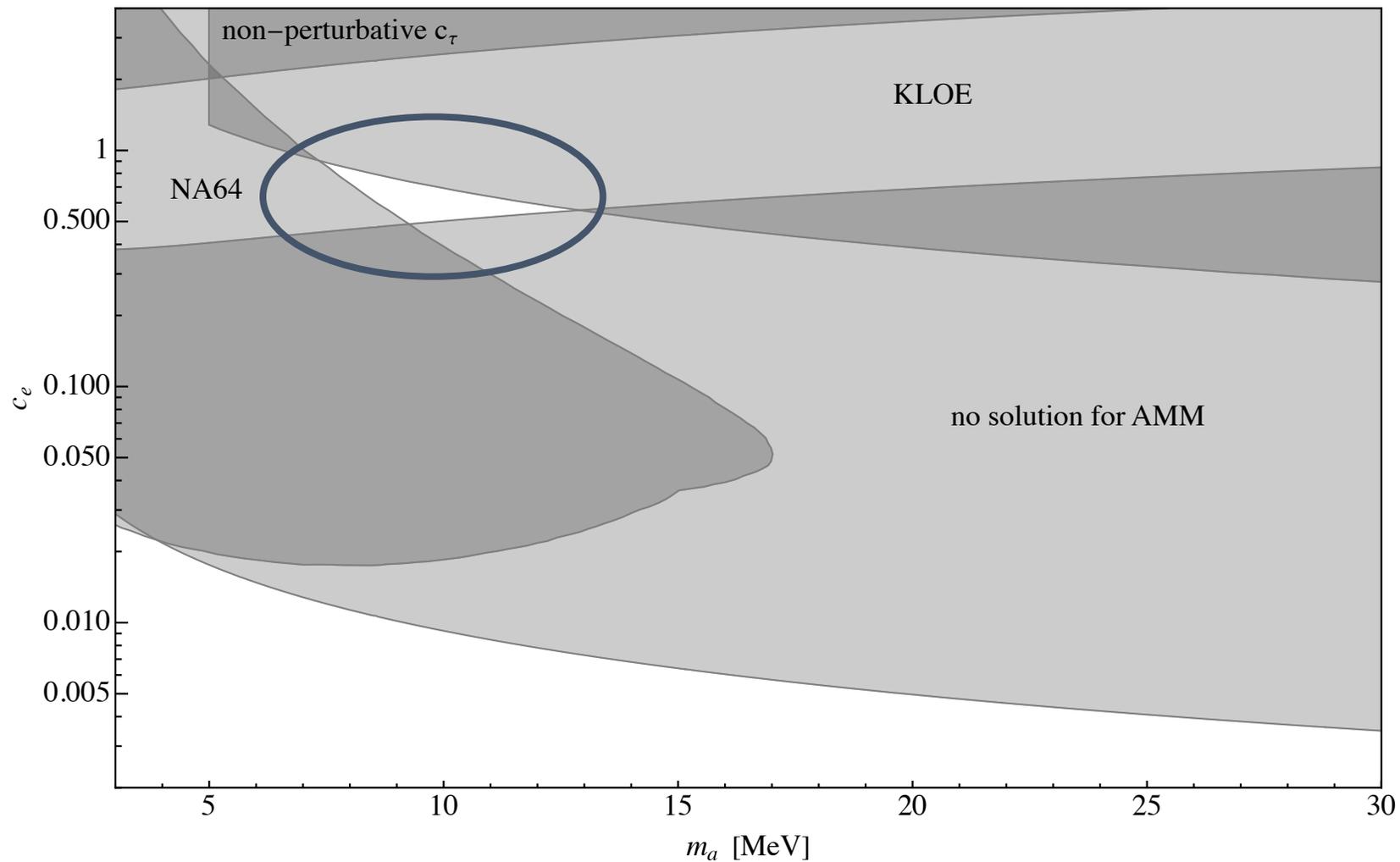
# (5) The Model

$$\mathcal{L}_{\text{eff}} = \frac{\partial^\mu a}{2f_a} c_i \bar{\psi}_i \gamma_\mu \gamma_5 \psi_i$$

- $c_u = \frac{2}{3}$
  - $c_d = \frac{1}{3}$
- } pionphobic, nucleophobic
- $c_s = c_c = c_b = c_t = 0$
  - Hierarchy:  $c_e \simeq c_\tau \gg c_\mu$



# (5) The Model



# (6) Conclusion

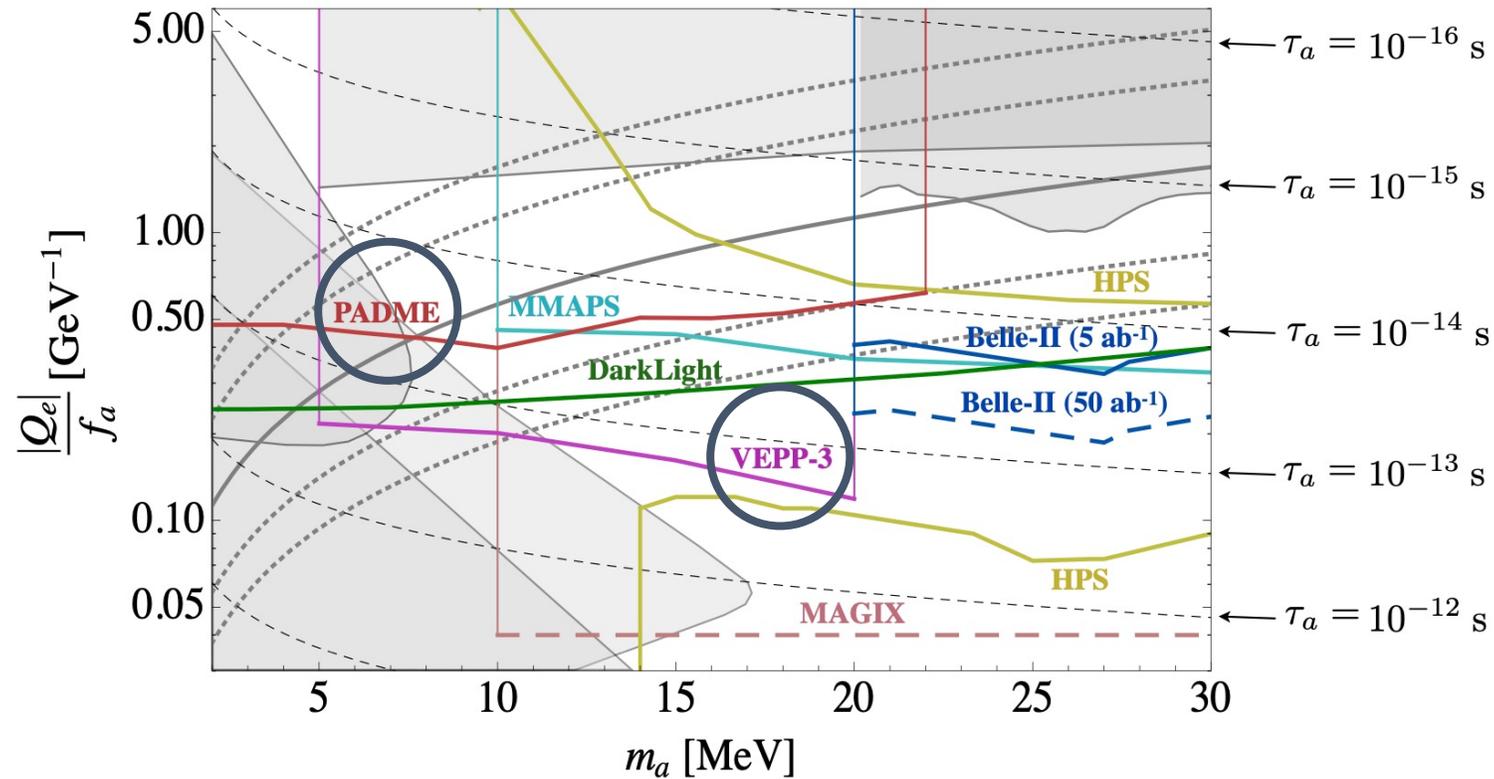
- Complete model addressing the strong CP puzzle and  $(g - 2)_\mu$
- Able to address further issues (XENON1T excess,  $(g - 2)_e$ )
- Fulfills given constraints (pionphobic, beam-dump experiments, ...)
- Future experiments will investigate interesting region (PADME, VEPP-3, ...)
  
- Proof of concept



# Backup

# Backup: Outlook

D. Alves et al.: arXiv:1710.03764



# Backup: Benchmark Model

$$m_a = 5.7 \mu\text{eV} \left( \frac{10^{12} \text{ GeV}}{f_a} \right)$$

$$f_a = \frac{V_{\text{PQ}}}{2N} \in [0.66, 0.75] \text{ GeV}$$
$$c_c = c_s = c_b = c_t \approx 2.5 \cdot 10^{-5}$$

Prompt decay to  $e^+e^-$  :  $\tau \approx 7 \cdot 10^{-15} \text{ s}$

$$\text{BR}(a \rightarrow \gamma\gamma) \approx 4 \cdot 10^{-5}$$

Decay of SM Higgs boson:

$$\text{BR}(h_{\text{SM}} \rightarrow aa) \approx 10^{-2}$$

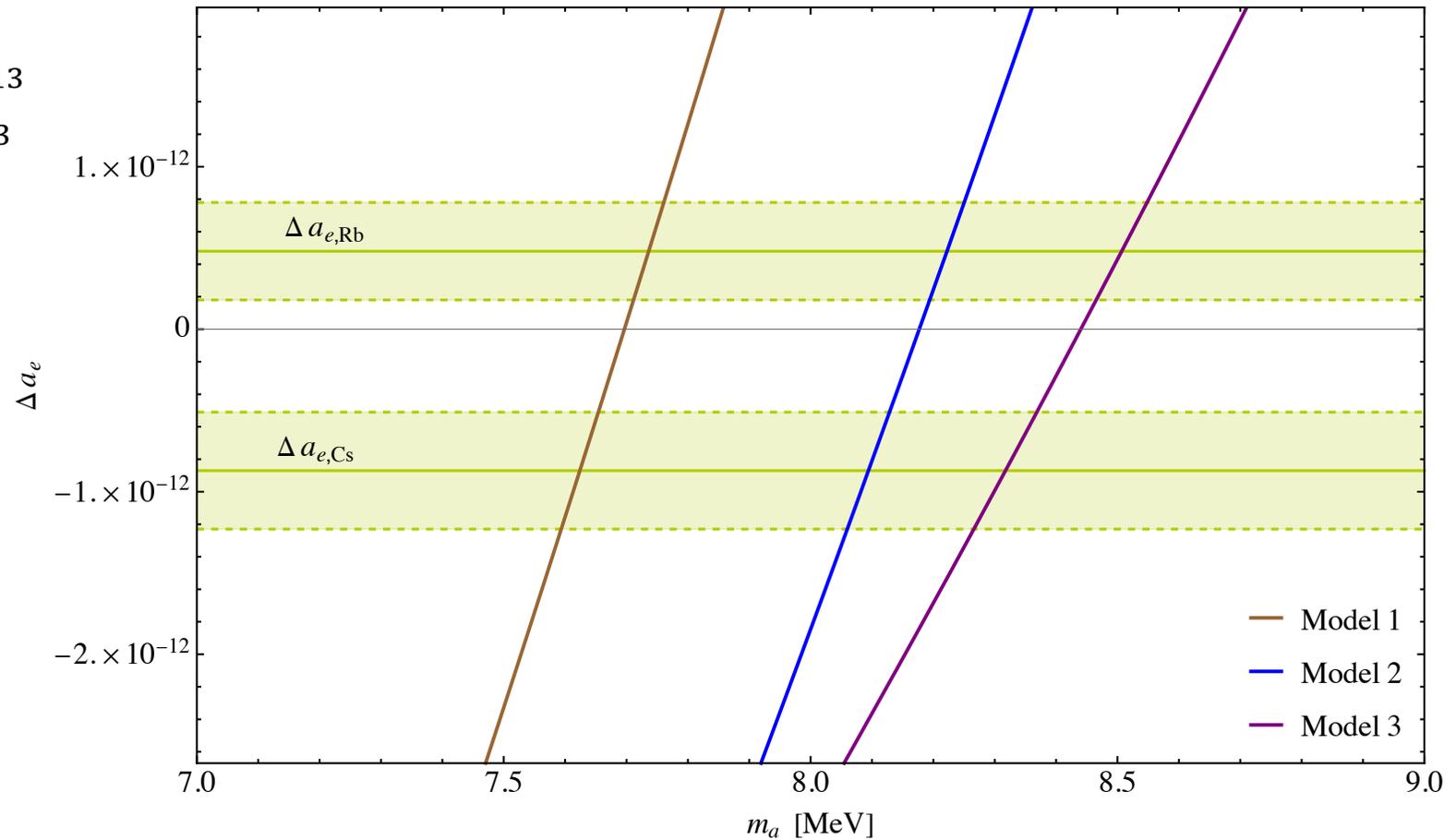
# Backup: Axion Lagrangian

$$\mathcal{L} \supset \frac{a}{f_a} \frac{g_S^2}{32\pi^2} G_{\mu\nu}^A \tilde{G}^{A,\mu\nu} + \frac{a}{f_a} \frac{e^2}{32\pi^2} \frac{E}{N} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{\partial_\mu a}{2f_a} c_i \bar{\psi}_i \gamma^\mu \gamma_5 \psi_i$$

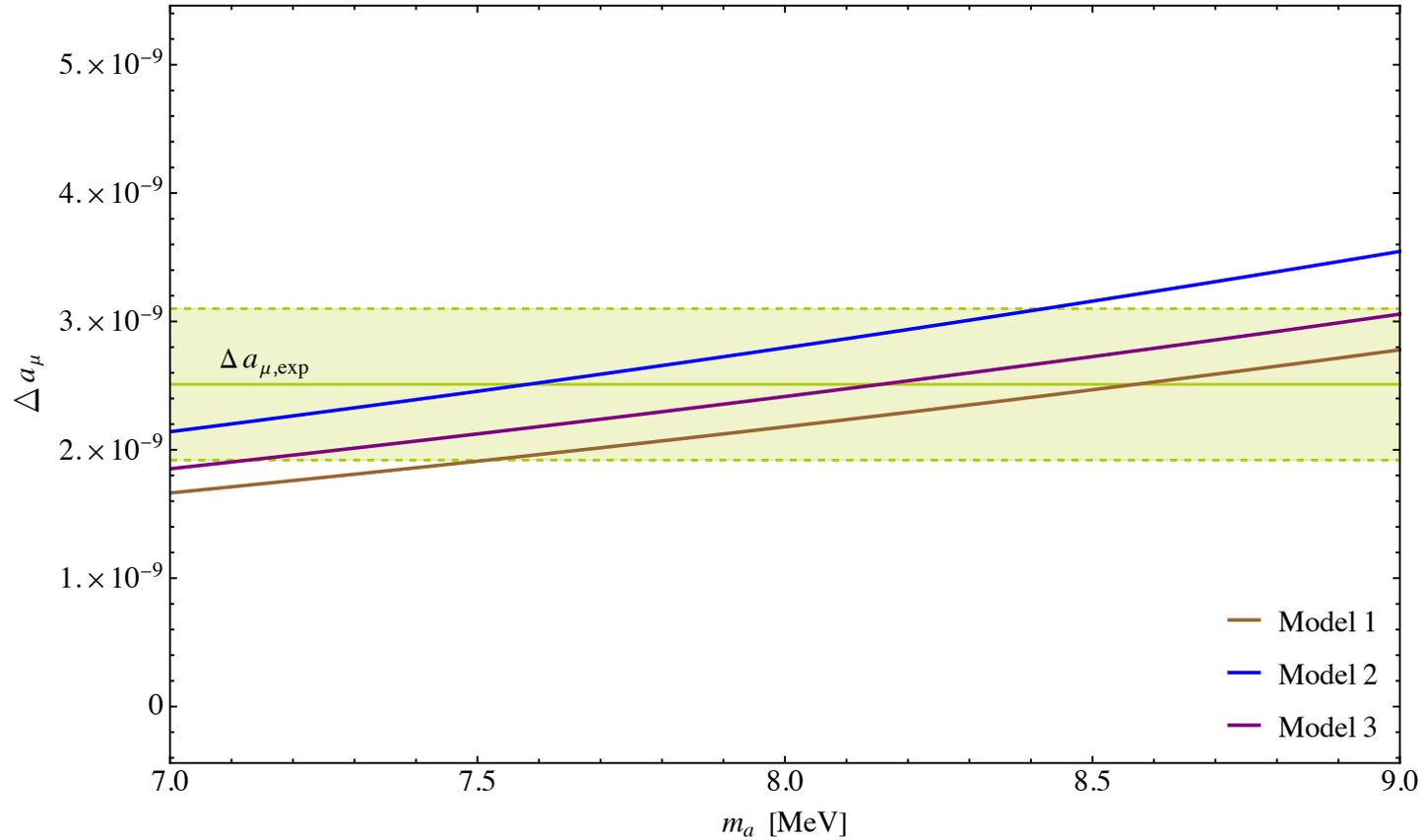
# Backup: Reproducing $\Delta a_e$

$$\Delta a_{e,Cs} = (-8.8 \pm 3.6) \cdot 10^{-13}$$

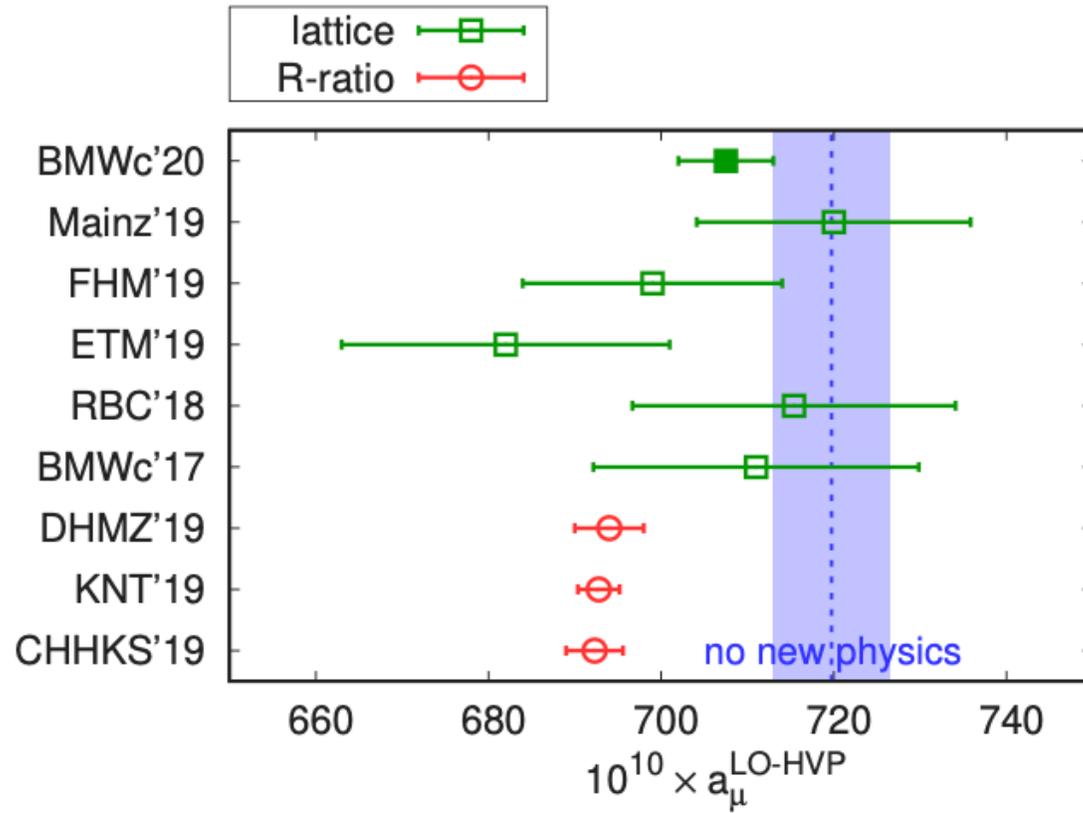
$$\Delta a_{e,Rb} = (4.8 \pm 3.0) \cdot 10^{-13}$$



# Backup: Reproducing $\Delta a_\mu$

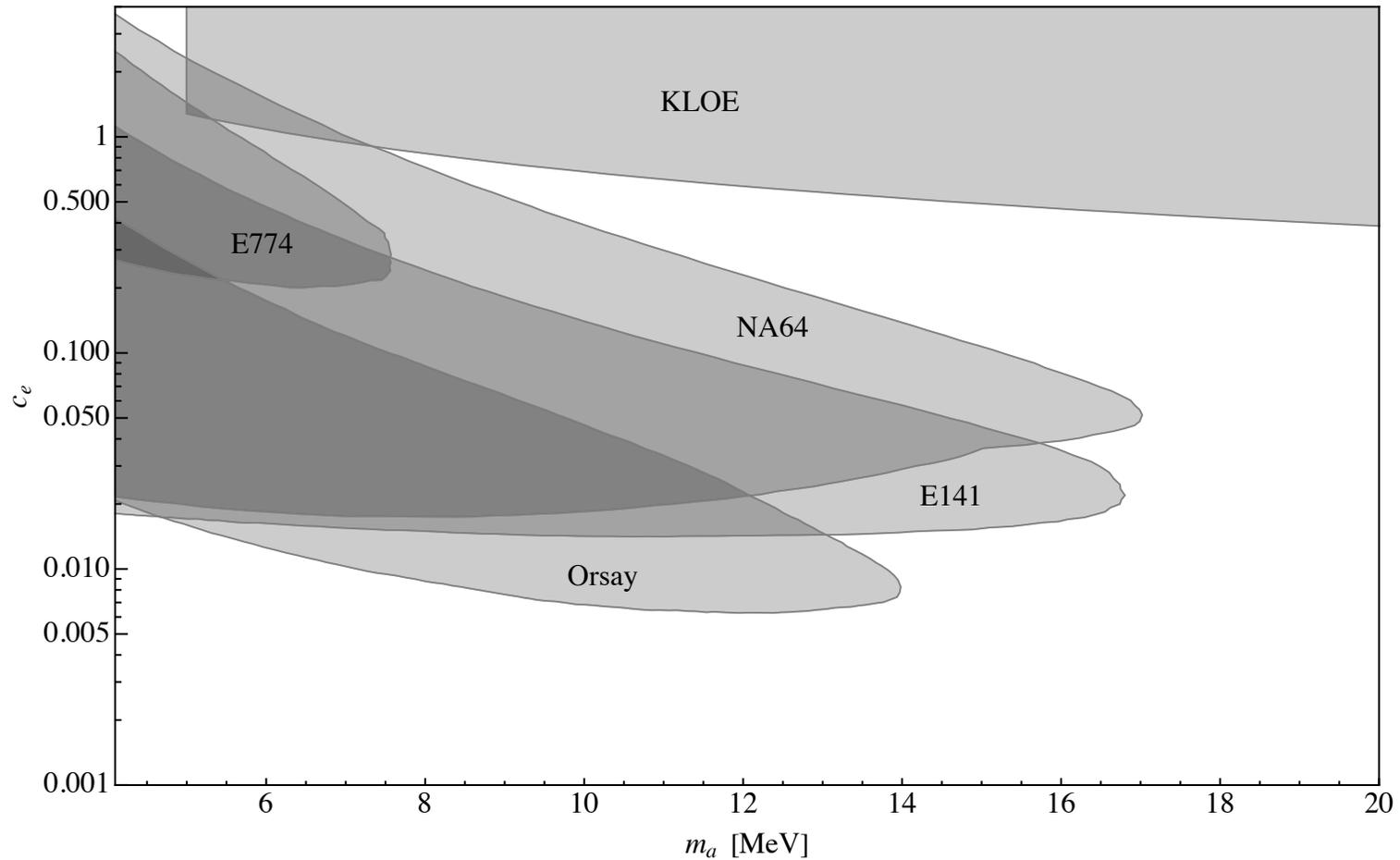


# Backup: Lattice-QCD

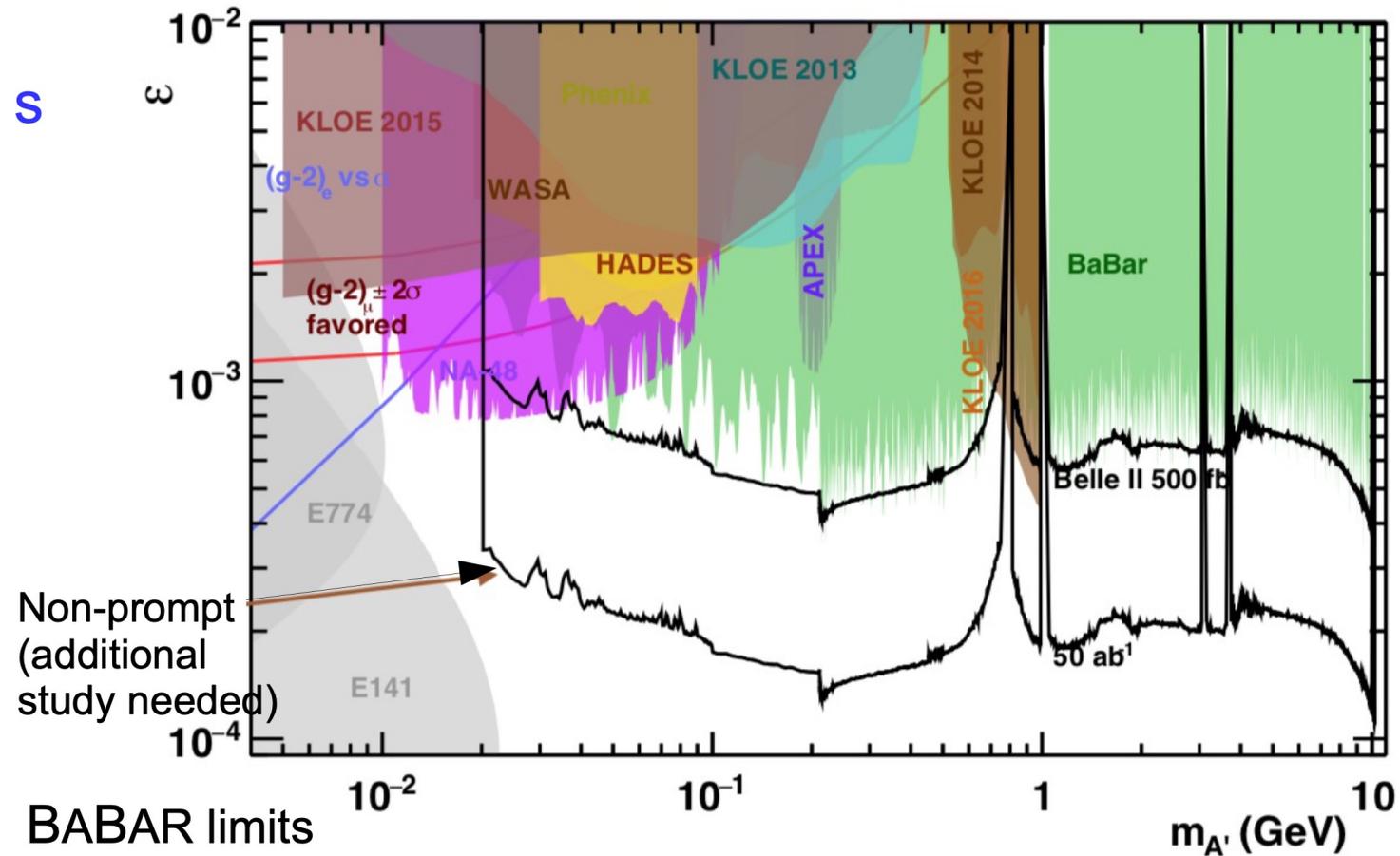


arXiv:2002.12347

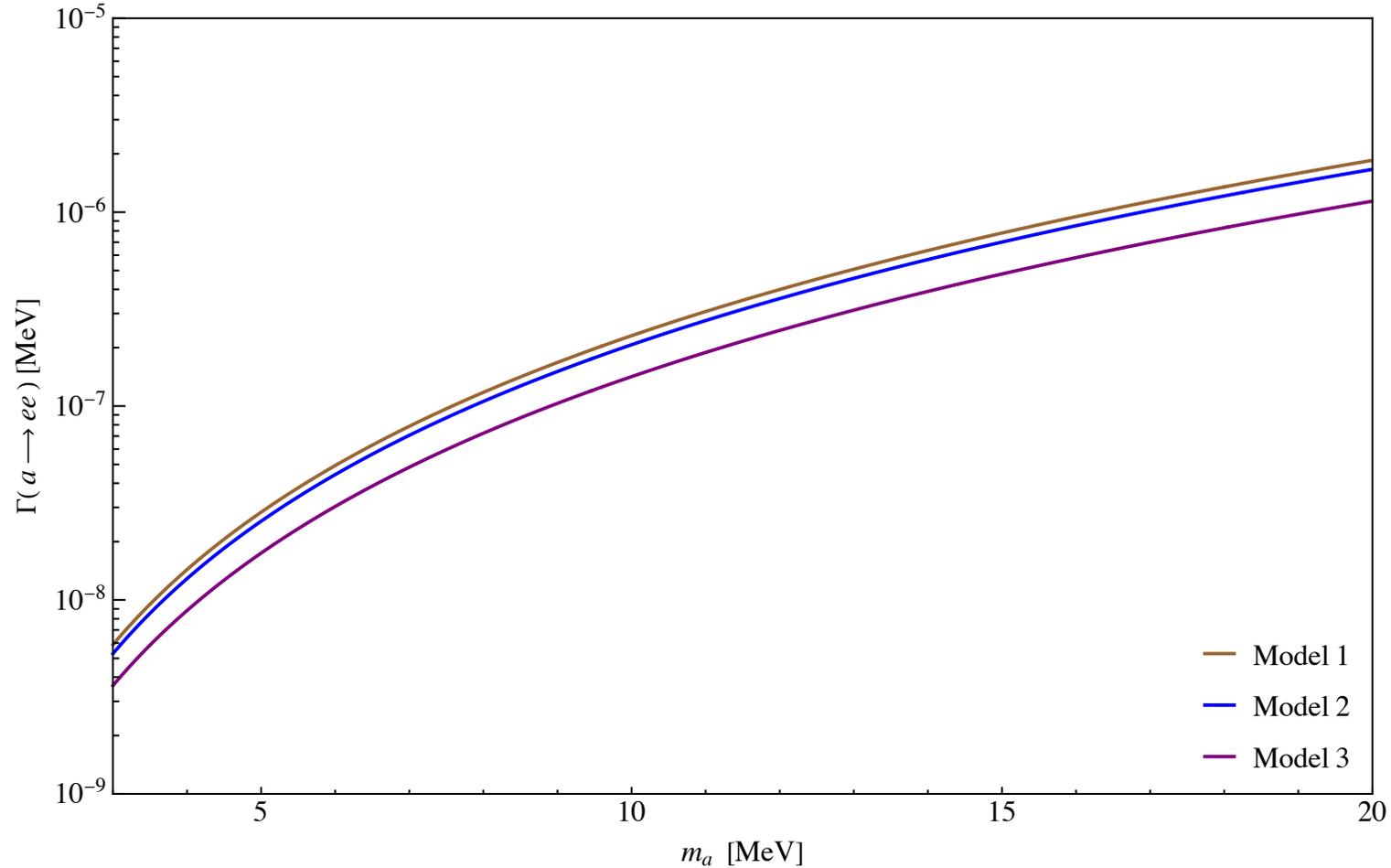
# Backup: Experimental constraints



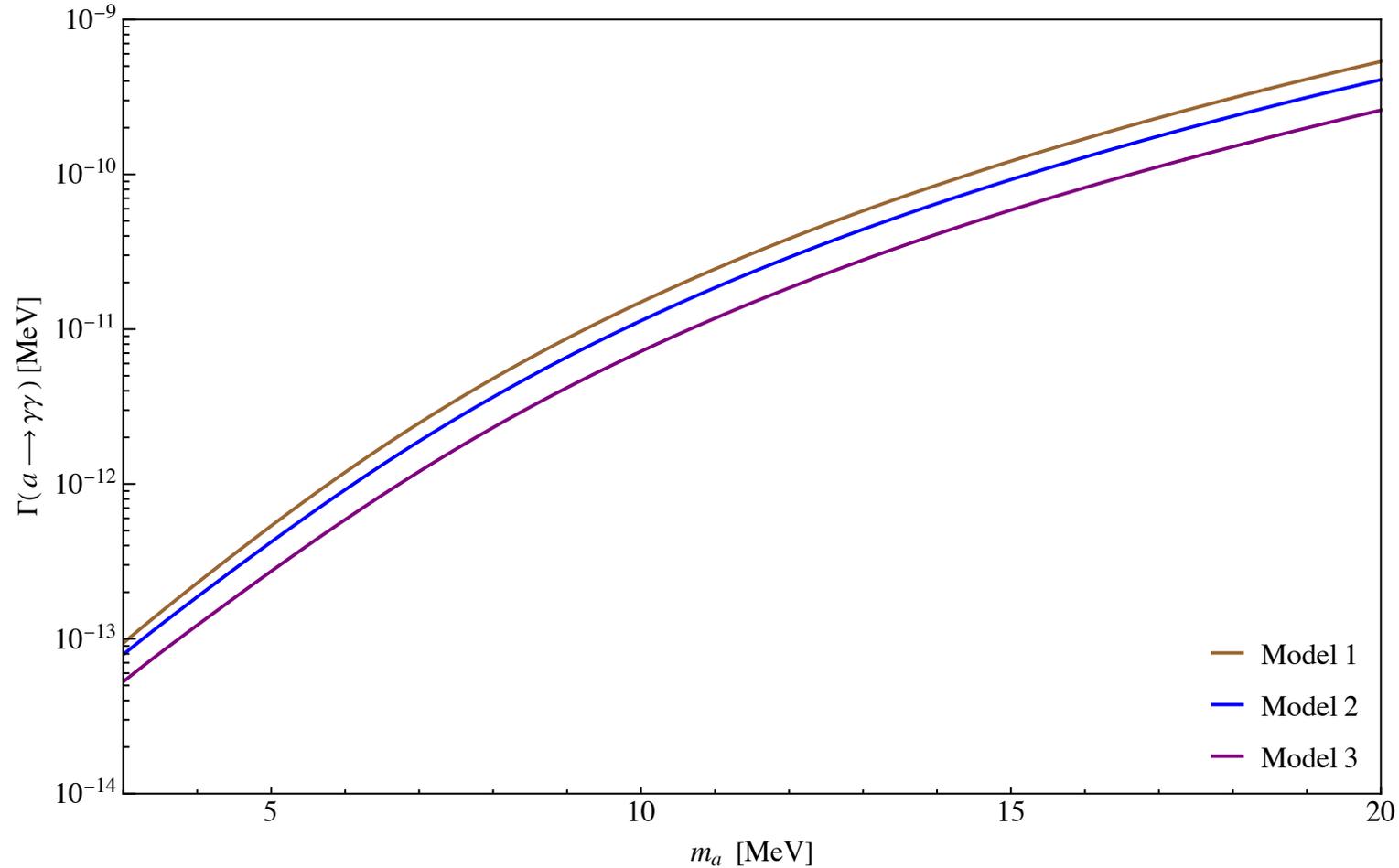
# Backup: BaBar, Belle2



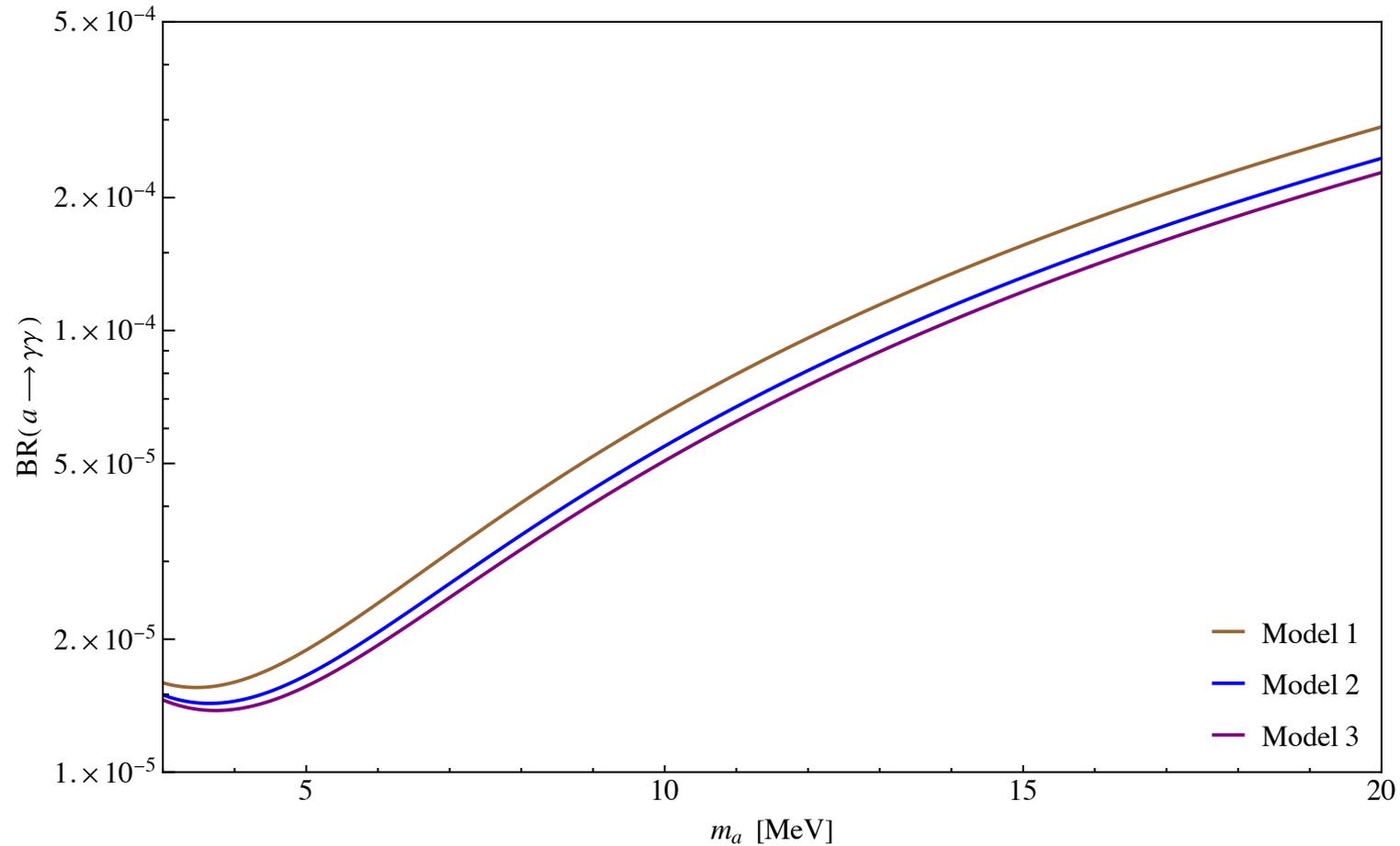
# Backup: Axion decay into electrons



# Backup: Axion decay into photons



# Backup: Branching ratio axion into photons



# Backup: Calculation

1-loop contribution:

$$\Delta a_i^1 = -\frac{m_i^2}{16\pi^2 f_a^2} |c_i|^2 f_1\left(\frac{m_a^2}{m_i^2}\right)$$

2-loop contribution (exact,  $j \in \{s, c, b, t, e, \mu, \tau\}$ ):

$$\Delta a_i^{2,j} = N_j \frac{m_i^2 \alpha}{8\pi^3 f_a^2} c_i c_j f_3\left(\frac{m_a^2}{m_i^2}, \frac{m_a^2}{m_j^2}\right)$$

[R. Ziegler et al: arXiv:2011.08919](#)

2-loop contribution (effective,  $\{u, d\}$ ):

$$\Delta a_i^{2,u/d} = \frac{2\alpha m_i^2}{\pi f_a^2} c_i C_{\gamma\gamma}^{\text{eff}} \left( f_2\left(\frac{m_a^2}{m_i^2}\right) - \ln\left(\frac{16\pi^2 f_\pi^2}{m_i^2}\right) \right)$$

[M. Neubert et al.: arXiv:1708.00443](#)

# Backup: UV Completion/Mass Spectra

**Table 5.1.:** Fermionic part of the particle content. Shown are the PQ charges and hypercharges of all SM fermions.

Particle	$q_{L,i}$	$u_{R,1}$	$u_{R,2}$	$u_{R,3}$	$d_{R,1}$	$d_{R,2}$	$d_{R,3}$	$l_{L,i}$	$e_{R,1}$	$e_{R,2}$	$e_{R,3}$
PQ charge	0	$X_u$	$X_c$	$X_t$	$X_d$	$X_s$	$X_b$	0	$X_e$	$X_\mu$	$X_\tau$
Hypercharge	$\frac{1}{6}$	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{2}$	-1	-1	-1

**Table 5.2.:** Scalar doublet part of the particle content. Shown are their PQ charges, hypercharges and VEV.

Particle	$h_0$	$h_1$	$h_2$	$h_3$	$h_4$	$h_5$
PQ charge	$X_{h0}$	$X_{h1}$	$X_{h2}$	$X_{h3}$	$X_{h4}$	$X_{h5}$
Hypercharge	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
VEV	$v_{h0}$	$v_{h1}$	$v_{h2}$	$v_{h3}$	$v_{h4}$	$v_{h5}$

**Table 5.3.:** Scalar singlet part of the particle content. Shown are their PQ charges, hypercharges and VEV.

Particle	$\phi_0$	$\phi_1$	$\phi_2$	$\phi_3$	$\phi_4$	$\phi_5$	$\phi_6$	$\phi_7$
PQ charge	1	$X_{\phi_1}$	$X_{\phi_2}$	$X_{\phi_3}$	$X_{\phi_4}$	$X_{\phi_5}$	$X_{\phi_6}$	$X_{\phi_7}$
Hypercharge	0	0	0	0	0	0	0	0
VEV	$v_{\phi_0}$	$v_{\phi_1}$	$v_{\phi_2}$	$v_{\phi_3}$	$v_{\phi_4}$	$v_{\phi_5}$	$v_{\phi_6}$	$v_{\phi_7}$

**Table F.2.:** The masses of all CP-even scalars  $m_S$  and CP-odd pseudoscalars  $m_{PS}$  of each model are displayed, sorted numerically. The CP-even mass of 125 GeV corresponds to the SM Higgs boson  $H_{SM}$  and the two massless CP-odd particles represent the two Goldstone bosons of the unbroken symmetries. For every model, the individual estimates of the required parameters are used.

Model 1		Model 2		Model 3	
$m_S$ in GeV	$m_{PS}$ in GeV	$m_S$ in GeV	$m_{PS}$ in GeV	$m_S$ in GeV	$m_{PS}$ in GeV
260	260	180	180	130	130
260	260	180	180	130	130
260	260	180	180	130	130
125	3.8	125	16	130	130
4.4	3.6	16	2.8	125	2.2
3.8	1.7	4.7	1.7	5.9	1.8
3.5	1.5	2.8	1.6	2.3	1.6
1.8	1.4	1.9	1.2	1.8	1.3
1.6	1.1	1.7	0.97	1.6	1.1
1.6	0.79	1.4	0.79	1.5	1.0
1.5	0.56	1.3	0.54	1.4	0.76
1.3	0.32	1.3	0.36	1.2	0.15
0.90	0	1.1	0	0.98	0
0.57	0	0.60	0	0.64	0

# Backup: UV Lagrangian

$$\mathcal{L}_{\text{kin}} = \sum_{i=0}^5 (\partial_\mu h_i)^\dagger \partial^\mu h_i + \sum_{i=0}^7 (\partial_\mu \phi_i)^* \partial^\mu \phi_i .$$

$$h_i = \frac{(v_{hi} + H_i)}{\sqrt{2}} e^{i \frac{X_{hi}}{V_{PQ}} a} \begin{pmatrix} 0 \\ 1 \end{pmatrix} ,$$

$$\phi_i = \frac{(v_{\phi i} + \Phi_i)}{\sqrt{2}} e^{i \frac{X_{\phi i}}{V_{PQ}} a} .$$

$$V_1 = \sum_{i=1}^3 \left( \bar{q}_i Y_{i1}^u \tilde{h}_4 u_{R,1} + \bar{q}_i Y_{i2}^u \tilde{h}_0 u_{R,2} + \bar{q}_i Y_{i3}^u \tilde{h}_0 u_{R,3} \right. \\ \left. + \bar{q}_i Y_{i1}^d h_2 d_{R,1} + \bar{q}_i Y_{i2}^d h_0 d_{R,2} + \bar{q}_i Y_{i3}^d h_0 d_{R,3} \right. \\ \left. + \bar{l}_i Y_{i1}^e h_5 e_{R,1} + \bar{l}_i Y_{i2}^e h_1 e_{R,2} + \bar{l}_i Y_{i3}^e h_3 e_{R,3} \right) + h.c.$$

$$V_2 = \sum_{i=0}^5 \left( -\mu_{hi}^2 h_i^\dagger h_i + \lambda_{hi} (h_i^\dagger h_i)^2 \right) + \sum_{i=0}^7 \left( -\mu_{\phi i}^2 \phi_i^* \phi_i + \lambda_{\phi i} (\phi_i^* \phi_i)^2 \right) .$$

$$V_3^1 = - \left( A_0 \phi_1^* \phi_0^2 + A_1 \phi_2^* \phi_1^2 + A_2 \phi_3^* \phi_2 \phi_0 + A_3 \phi_4^* \phi_3^2 + A_4 \phi_5^* \phi_4^2 + A_5 \phi_6^* \phi_5^2 \right. \\ \left. + A_6 \phi_7^* \phi_4 \phi_6 \right) + h.c. ,$$

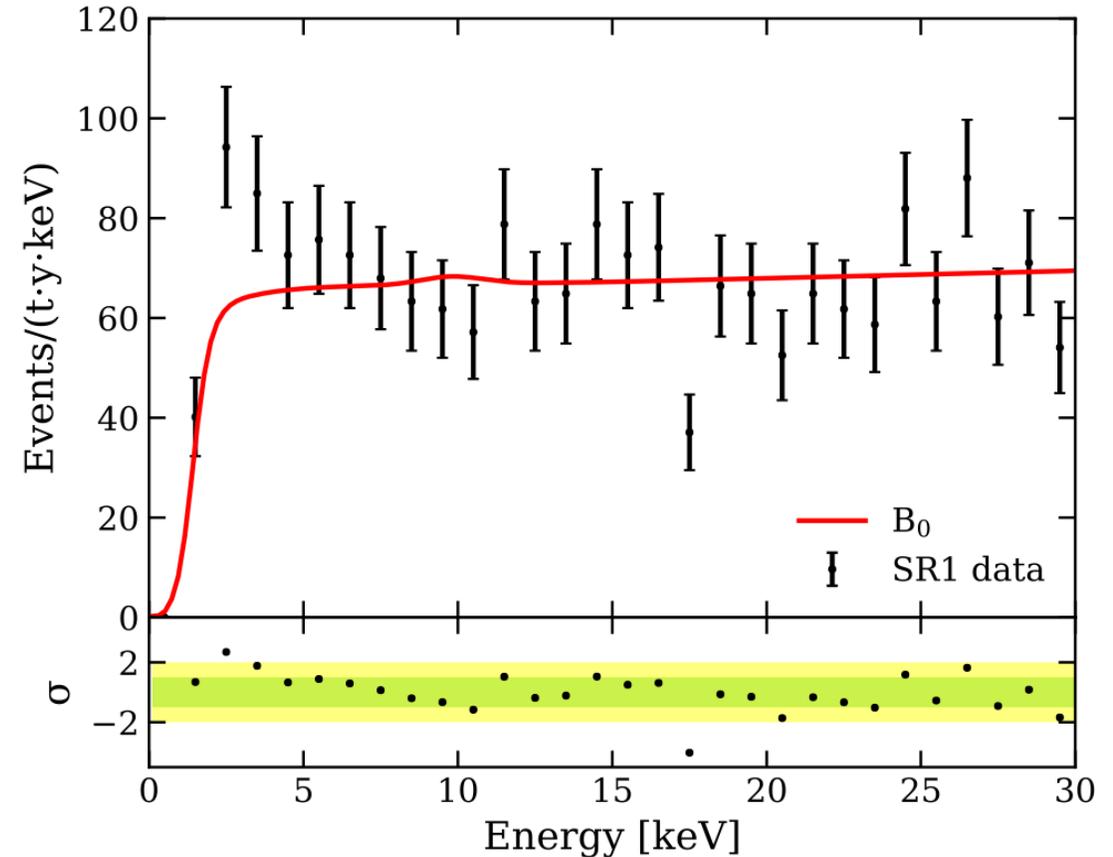
$$V_4^1 = - \left( B_0 \phi_0^* h_0^\dagger h_1 + B_1 \phi_5^* h_0^\dagger h_2 + B_2 \phi_3^* h_2^\dagger h_3 + B_3 \phi_6 h_0^\dagger h_4 + B_4 \phi_7^* h_1^\dagger h_5 \right) + h.c.$$

# Backup: SM Higgs decay into axions

$$\begin{aligned}\mathcal{L}_{haa} &= \frac{\gamma_1}{\sqrt{2}V_{PQ}^2 m_H^2} (\partial_\mu a)^2 H_{\text{SM}}, \\ \gamma_1^{\text{exp}} &= \left(X_{\phi 0}^2 + X_{h 1}^2\right) B_0 v_{\phi 0} v_{h 1} + \left(X_{\phi 5}^2 + X_{h 2}^2\right) B_1 v_{\phi 5} v_{h 2} \\ &\quad + \left(X_{\phi 6}^2 + X_{h 4}^2\right) B_3 v_{\phi 6} v_{h 4} + \alpha_0^2 \sqrt{2} m_H^2 X_{h 0}^2 v_{h 0}, \\ &\approx 2 \left(X_{\phi 0}^2 + X_{\phi 5}^2 + X_{\phi 6}^2\right) B v_{h 1} v_\phi + \sqrt{2} m_H^2 X_{h 0}^2 v.\end{aligned}$$

# Backup: Addressing the XENON1T excess

E. Aprile et al.: arXiv:2006.09721



# Backup: Addressing the XENON1T excess

