

The role of the Higgs radial mode as a UV completion of the non linear sigma model.

Master thesis project

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Introduction

NGB EFT Lagrangian: Weakly or strongly coupled?

Framework:

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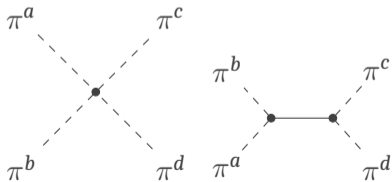


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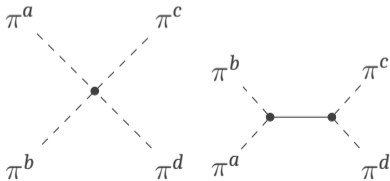


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- **How can we make the theory weakly coupled?**
 - Answer: Adding by hand new degrees of freedom: $\text{NL}\Sigma\text{M} \rightarrow \text{L}\Sigma\text{M}$;





Preliminary concepts

Coset construction!

- Let's take for example the $N_f = 2$ QCD-like theory with SSB $SU(2)_L \times SU(2)_R / SU(2)_V$:

$$\Sigma = \exp\left(\frac{2i}{f}\pi^a X^a\right), \quad \Sigma \xrightarrow{\mathcal{G}} L\Sigma R^\dagger, \quad \Sigma\Sigma^\dagger = \mathbb{I}, \quad \det \Sigma = 1$$

$$\mathcal{L} = \frac{f^2}{4}(\partial_\mu \Sigma^\dagger \partial^\mu \Sigma) \quad \Rightarrow \quad \mathcal{L} = \frac{1}{2}(\partial_\mu \pi^a)^2 + \frac{1}{f^2}\pi^a \partial_\mu \pi^b \pi^c \partial^\mu \pi^d f^{ab\gamma} f^{cd\gamma} + \dots$$

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- $\exists \Lambda^* \sim 4\pi f$ above which the EFT approach fails.
- Σ has not enough d.o.f. to provide a linear irreducible representation of \mathcal{G} .



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allow us to construct a linear irreps of \mathcal{G} just adding a singlet extra radial mode, "curing" the energy growing behaviour of the scattering amplitudes and making the theory weakly coupled at high energies.



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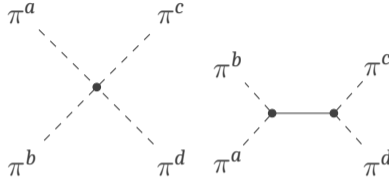
- **Notice the link:**

Linear irreps of $\mathcal{G} \iff$ Renormalizability \iff Weakly coupled theory.



Examples

Standard Model and two flavours QCD-like theory.



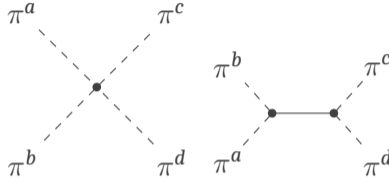
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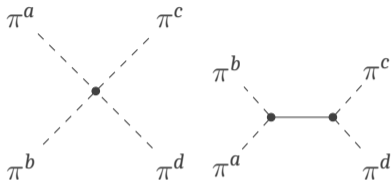
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 - Non Linear Sigma Model just as an EFT (ex: QCD, no other resonances intervene to restore perturbative unitarity) \Rightarrow The theory remain strongly coupled.
 - $SO(4)/SO(3)$ custodial symmetry for the EW sector in the SM, the Higgs boson radial mode *unitarize* the SM amplitudes as the WW-scattering \Rightarrow The theory became weakly coupled.



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What about the more general case?

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 - Determine the minimal number of degrees of freedom needed for a linear representation of \mathcal{G} .
 - Finding the minimal number of degrees of freedom necessary to *unitarize* the theory.



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Non linear sigma model and coset Construction

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- $\tilde{\Sigma} \xrightarrow{\mathcal{G}} UL\tilde{\Sigma}R^\dagger, \quad U \in U(1), L \in SU(N)_L, R \in SU(N)_R$
- We perform the chiral perturbation theory of the Non Linear Sigma Model

$$\mathcal{L}_{nl\sigma m} = \frac{1}{2}(\partial_\mu \pi^a)^2 + \frac{1}{2}(\partial_\mu \eta)^2 + \frac{\tilde{\mathcal{C}}_\pi}{6f^2} [(\pi^a \partial_\mu \pi^a)^2 - \pi^a \pi^a \partial_\mu \pi^b \partial_\mu \pi^b] + \dots$$

$$\mathcal{A}(ab \rightarrow cd) = i\tilde{\mathcal{C}}_\pi \left(\delta^{ab} \delta^{cd} \frac{s}{f^2} + \delta^{ac} \delta^{bd} \frac{t}{f^2} + \delta^{ad} \delta^{bc} \frac{u}{f^2} \right). \quad \boxed{\tilde{\mathcal{C}}_\pi = \frac{N}{N^2 - 2}}$$



Linear Sigma Model

2 Linear and non linear sigma model: More general case

- Without the pseudo-reality condition, \mathcal{H} contains more than one singlet radial mode and the more general linear irreps could be found in $\mathcal{H} \sim \text{Mat}(N, \mathbb{C}) \Rightarrow 2N_{\mathbb{R}}^2$ d.o.f.

$$2N_{\mathbb{R}}^2 = \underbrace{N^2 - 1}_{\pi^a(x)} + \underbrace{1}_{\eta(x)} + \underbrace{N^2 - 1}_{H^a(x)} + \underbrace{1}_{h(x)}$$



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- Using polar decomposition one can find an useful parametrization of \mathcal{H} :

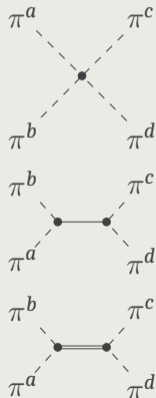
$$\mathcal{H} = \tilde{\Sigma}(x)\phi(x) = e^{\frac{2i}{f}\pi^a X^a} e^{\frac{i}{f}\eta} \left(1 + \frac{H^a(x)}{f} X^a \right) (f + h(x)).$$



Linear sigma model framework

2 Linear and non linear sigma model: More general case

Feynmann Diagrams



Lagrangian

$$\mathcal{L} = \underbrace{\frac{1}{4} \text{Tr}[\partial_\mu \mathcal{H}^\dagger \partial^\mu \mathcal{H}]}_{(1)} - \underbrace{V(\mathcal{H} \mathcal{H}^\dagger)}_{(2)} + \underbrace{\text{Pol}(\det \mathcal{H})}_{(3)}$$

$$(1) \Rightarrow \mathcal{L}_{4\pi} + \mathcal{L}_{h\pi} + \mathcal{L}_{H\pi}$$

$$(2) \Rightarrow \text{Gives information on } m_h^2 \text{ and } m_H^2.$$

$$(3) \Rightarrow \text{Explicitly breaking of the } U(1) \text{ abelian symmetry} \Rightarrow \text{mass contribution to the } \eta \text{ field.}$$



Restoring perturbative unitarity in $\pi\pi$ scattering

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- Looking at the two to two $\pi\pi$ scattering we notice that

$$\mathcal{A}_{tot} = \mathcal{A}_{4\pi} + \mathcal{A}_h + \mathcal{A}_H$$

$$\mathcal{A}^{\text{tot}}(\pi^a \pi^b \rightarrow \pi^c \pi^d) = \frac{i}{f^2} (\tilde{\mathcal{C}}_\pi - \tilde{\mathcal{C}}_h - \tilde{\mathcal{C}}_H) \left(\delta^{ab} \delta^{cd} s + \delta^{ac} \delta^{bd} t + \delta^{ad} \delta^{bc} u \right)$$

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$$\tilde{\mathcal{C}}_h = \frac{2}{N}, \quad \tilde{\mathcal{C}}_H = \frac{N}{N^2 - 2} - \frac{2}{N}$$

- Finally we arrive to the main result:

$$\tilde{\mathcal{C}}_\pi + \tilde{\mathcal{C}}_h + \tilde{\mathcal{C}}_H = 0$$

ALL THE HIGGS RADIAL MODES ARE INVOLVED IN RESTORING PERTURBATIVE UNITARITY.



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- In the $N_f = 2$ case it is sufficient just a singlet radial mode, in order to make the theory weakly coupled. However we proved that, in the general case, we need more extra radial modes (*"Adjoint higgses"*).



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 - *Warning:* Those theories manifest, as well as the Standard Model, the hierarchy problem.
- Our studies provides *new model building tools* to extend the UV completion through the LΣM procedure for more general SSB patterns.



The role of the Higgs radial mode as a
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model. *Thank you for listening!*

Any questions?