

# Multi-loop results for the anomalous magnetic moment of the muon

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- I. Introduction & Motivation
- II. Outline of the calculation
- III. Results & Discussion
- IV. Summary & Conclusion

In collaboration with: [P.A. Baikov](#) and [K.G. Chetyrkin](#)  
Nucl. Phys. **B183** (Proc.Suppl.) + work in preparation ...



# Introduction

Generalities; Anomalous magnetic moment of a lepton:  $a_\ell$

- Magnetic moment of any system:

I.) Motion of el. charges



II.) Intrinsic mag. moments of elementary particles

- Dirac theory predicts for a lepton  $\ell = e, \mu, \tau$ :

$$\vec{\mu}_\ell = g_\ell \left( \frac{e}{2m_\ell} \right) \vec{S}, \quad g_\ell = 2 \quad (\text{free, non-interacting})$$

- **Quantum fluctuations:**  $\rightsquigarrow$  deviation from  $g_\ell = 2$ :  
parametrized by

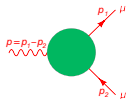
$$g_\ell = 2(1 + a_\ell) \quad \rightsquigarrow \text{precises test of QFT}$$



# Introduction

Generalities; Anomalous magnetic moment of the muon:  $a_\mu$

- More formally:



$$= \bar{u}(p_1) \left[ \gamma^\mu F_E(p^2) + i \frac{\sigma^{\mu\nu} p_\nu}{2m_\mu} F_M(p^2) \right] u(p_2)$$

- In the static limit ( $p^2 \rightarrow 0$ ):  $F_E(0) = 1$ ,  $F_M(0) = a_\mu$
- Muon very interesting:

Quantum fluctuations due to heavier particles  $M$ :

$$\frac{\delta a_\ell}{a_\ell} \propto \frac{m_\ell^2}{M^2} \quad (M \gg m_\ell)$$

$M$  heavy SM or BSM particle

ratio:  $m_\mu / m_e \sim 200$

Higher sensitivity to physics beyond SM

→ Consider anomalous magnetic moment of muon

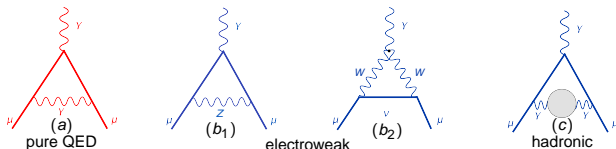
- Loop calculations not only mathematical task to increase precisions  $\rightsquigarrow$  also allow to access energy regimes not yet reachable by collider experiments through virtual particles



# Introduction

## Theory: Higher order corrections

- Higher order corrections are classified into 3 classes:



$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{EW}} + a_{\mu}^{\text{Had}}$$

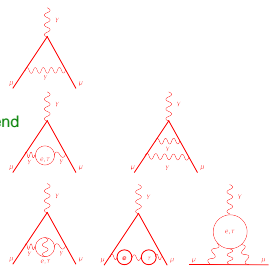
- The QED part is known to 4-loops (and leading terms in 5 loops!)  
(next slides)
- The EW part is known to 2-loops  
R. Jackiw, S. Weinberg; G. Altarelli et al.; I. Bars, M. Yoshimura; A. Czarnecki et al.
- The hadronic part is known but with limited accuracy  
Bouchiat, et al.; M. Gourdin, et al.; Brodsky, de Rafael; Hagiwara et al.; Alemany et al.; Davier et al.; Passera et al.  
Dominant theoretical uncertainties to the muon anomaly  
Effects at present not from first principles

# Introduction

Theory: QED contributions to  $a_\mu$

$$\begin{aligned} a_\mu^{\text{QED}} &= \left(\frac{\alpha}{\pi}\right) 0.5 \quad \text{Schwinger} \\ &+ \left(\frac{\alpha}{\pi}\right)^2 0.765857410(27) \quad \text{Sommerfield; Petermann; Suura & Wichmann; Elend} \\ &+ \left(\frac{\alpha}{\pi}\right)^3 24.05050964(87) \quad \text{Barbieri, Laporta, Remiddi et al.; Kinoshita et al.;} \\ &\quad \text{Czarnecki, Skrzypek; Friot, Greynat, de Rafael} \\ &+ \left(\frac{\alpha}{\pi}\right)^4 130.8055(80) \quad \text{Kinoshita, Lindquist; Kinoshita, Nio; Kinoshita, Nizic,} \\ &\quad \text{Okamoto; Aoyama, Hayakawa, Kinoshita, Nio; Lautrup, de Rafael} \\ &+ \left(\frac{\alpha}{\pi}\right)^5 663(20) \quad \text{In progress} \\ &\quad \text{Kinoshita et al.; Kataev; Laporta; Baikov et al.} \end{aligned}$$

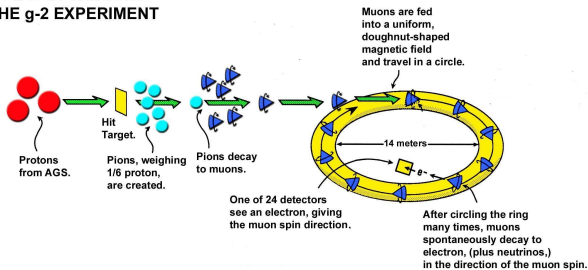
growing coefficients



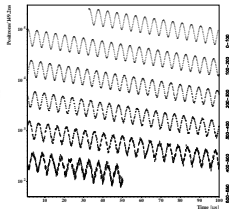
# Introduction

## Experiment

### LIFE OF A MUON: THE g-2 EXPERIMENT



<http://www.g-2.bnl.gov/>



(g-2) Collaboration (H.N. Brown et al.)

**Quantum Fluctuations!**

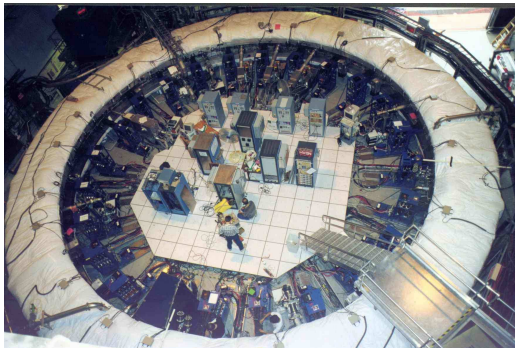
$$\omega_a = \omega_s - \omega_c \propto \frac{a_\mu B}{m}$$

$\omega_s$ : spin precession frequency,  $\omega_c$ : cyclotron frequency

positron time spectrum:  $N_0(E)e^{t/\gamma\tau} [1 + A(E) \cos(\omega_a t + \phi(E))]$

# Introduction

## Experiment



<http://www.g-2.bnl.gov/>



muon momentum: 3.094 GeV,    Radius: 7.112 m,    field: 1.45 T

# Introduction

## The value

The present experimental value is terrifically accurate!:

$$a_{\mu}^{\text{exp}} = 116592080(63) \cdot 10^{-11}$$

E821: Final Report: PRD73 (2006)

with statistical error( $54 \cdot 10^{-11}$ ) and systematic error( $33 \cdot 10^{-11}$ )

↪ Improvement of a factor of 14 compared to the classic CERN experiment

- QED four-loop contributes as much as  $380.8 \cdot 10^{-11}$   
(compared to the exp. uncertainty of  $\sim 60 - 70 \cdot 10^{-11}$ )

- <sup>1</sup>  $a_{\mu}^{\text{SM}} = 116591834(2)(41)(26) \cdot 10^{-11}$

The current theory prediction shows an

“interesting but not yet conclusive discrepancy” of  $\sim 3.2\sigma$

(or  $\sim 1.9\sigma$  if one uses  $\tau$ -data to describe the hadronic effects)

- <sup>2</sup> Update, new data (KLOE, BABAR):  $e^+e^-$  based:  $3.6\sigma$ ,  $\tau$ -data:  $2.4\sigma$

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<sup>1</sup> A. Höcker, W. Marciano, PDG

<sup>2</sup> Davier et al.

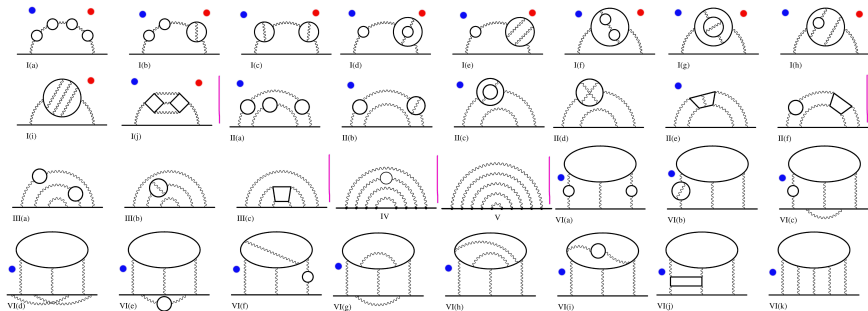




# Calculation

...back to theory, some diagrams...

- 12672 diagrams
- 6 classes (I - VI), 32 gauge invariant subsets
- (For simplicity external photon omitted)



diagrams from M. Nio, T. Aoyama, M. Hayakawa, T. Kinoshita

•: in this talk    •: Kinoshita et al.

# Calculation

## logarithmically enhanced contributions

There are 2 **sources** of numerically leading enhanced logs of large ratio  $\frac{M_\mu}{m_e} = 206.7682838$  pure QED contributions:

$$\begin{aligned} \text{LBL: light by light scattering} & \sim \left(\frac{\alpha}{\pi}\right)^3 \left(\frac{2\pi^2}{3} \ln\left(\frac{M_\mu}{m_e}\right) + \dots\right), \\ \text{VP: vacuum polarization} & \sim \left(\frac{\alpha}{\pi}\right)^2 \left(\frac{1}{3} \ln\left(\frac{M_\mu}{m_e}\right) + \dots\right) \end{aligned}$$

→ We will consider mixed VP insertions: photon propagator composed from electron loops and photon exchanges only!

$$a_\mu = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \left[ d_R \left( \frac{-x^2 M_\mu}{1-x m_e}, \alpha \right) - 1 \right] \quad \text{B. Lautrup, de Rafael}$$

with  $d_R(q^2/m^2, \alpha) = 1/(1 + \alpha \Pi^{\text{OS}}(q^2/m_e^2, \alpha^{\text{OS}}))$   
 $\Pi^{\text{OS}}$  proper photon VP in **OS-scheme**

$$\Pi(q^2) \sim \text{diagram} \dots$$

# Calculation

## Vacuum polarization function

Required ingredients for the calculation:

$$\Pi(q^2) \sim \text{diagram}$$

### ■ Vacuum polarization function

- High-energy limit  $\rightarrow$  massless propagators

The value of  $\overline{\Pi}(Q^2, m = 0, \overline{\alpha})$  is known at 4-loops with **Baicer** (P. Baikov, 2000-2008) in  **$\overline{\text{MS}}$ -scheme**

P.A. Baikov, K.G. Chetyrkin, J.H. Kühn

Traditionally in calculations of  $a_\ell$  everybody uses the classical **on-shell scheme**:

$\alpha$  and all **lepton masses** are **on-shell (OS)** and:

$$\Pi^{\text{OS}}(Q = 0, m, \alpha^{\text{OS}}) = 0$$

$\rightsquigarrow$  **Transform massless propagator from  $\overline{\text{MS}}$   $\rightarrow$  on-shell scheme**

### ■ **Need:** $\overline{\text{MS}}$ $\leftrightarrow$ on-shell relation at 4-loop

for **fine structure constant** conversion:  $\overline{\alpha} \rightarrow \alpha^{\text{OS}}$

# Calculation

On-shell  $\leftrightarrow$   $\overline{\text{MS}}$ -scheme for  $\alpha, m_\ell$

- In QED  $\overline{\text{MS}}$  in OS-schemes are related through:

$$\frac{\alpha^{\text{OS}}}{1 + \Pi^{\text{OS}}(\mathbf{Q}, m, \alpha^{\text{OS}})} = \frac{\bar{\alpha}}{1 + \bar{\Pi}(\mathbf{Q}, \bar{m}, \bar{\alpha})} \quad \begin{array}{l} \text{scheme invariant} \\ \text{concept of the} \\ \text{invariant charge} \end{array}$$

Required:  $\bar{\Pi}(\mathbf{Q} = 0, \bar{m}, \bar{\alpha})$

$$\bar{\Pi}(q^2) \sim \begin{array}{c} m_\ell \\ \circlearrowleft \\ \text{---} q \text{---} \end{array}$$

$$\text{Inversion} \rightsquigarrow \bar{\alpha} = \alpha^{\text{OS}} \left( 1 + \sum_{i \geq 1} c_{\bar{\alpha}\alpha}^{(i)} \left( \frac{\alpha^{\text{OS}}}{\pi} \right)^i \right)$$

- Need:  $\overline{\text{MS}} \leftrightarrow$  On-shell relation for conversion of lepton mass  $m_\ell$  at 3-loop

Steinhauser, Chetyrkin; Melnikov, van Ritbergen

# Result

## $\overline{MS}$ -OS-relation for conversion of fine structure constant

$$\overline{\alpha} = \alpha^{\text{OS}} \left( 1 + \sum_{i \geq 1} C_{\overline{\alpha}\alpha}^{(i)} \left( \frac{\alpha^{\text{OS}}}{\pi} \right)^i \right)$$

$$\begin{aligned} C_{\overline{\alpha}\alpha}^{(4)} = & \frac{14327767}{9331200} + \frac{8791}{3240} \pi^2 + \frac{204631}{259200} \pi^4 - \frac{175949}{4800} \zeta_3 + \frac{1}{24} \pi^2 \zeta_3 + \frac{9887}{480} \zeta_5 - \frac{595}{108} \pi^2 \ln 2 \\ & - \frac{106}{675} \pi^4 \ln 2 + \frac{6121}{2160} \pi^2 \ln^2 2 - \frac{32}{135} \pi^2 \ln^3 2 - \frac{6121}{2160} \ln^4 2 + \frac{32}{225} \ln^5 2 - \frac{6121}{90} a_4 - \frac{256}{15} a_5 \\ & + \ell_{\mu m} \left[ -\frac{383}{31104} + \frac{23}{108} \pi^2 - \frac{41}{144} \zeta_3 - \frac{2}{9} \pi^2 \ln 2 \right] + \frac{43}{144} \ell_{\mu m}^2 + \frac{13}{108} \ell_{\mu m}^3 + \frac{1}{81} \ell_{\mu m}^4, \\ \ell_{\mu m} = & \ln \frac{\mu}{m}, \quad a_n = \text{Li}_n \left( \frac{1}{2} \right) \end{aligned}$$

**Important:**  $\Pi(q^2/m^2, \alpha) = \Pi^\infty(q^2/m^2, \alpha) + \mathcal{O}(m^2/q^2)$   
then the resulting error in  $a_\mu$  will be:

$$a_\mu = \frac{\alpha}{\pi} \int_0^1 dx (1-x) \left[ d_R^\infty \left( \frac{-x^2 M_\mu}{1-x m_e}, \alpha \right) - 1 \right] + \mathcal{O} \left( \frac{m_e}{M_\mu} \right)$$

of order  $m_e/M_\mu$  with  $d_R^\infty = 1/(1 + \alpha\Pi^\infty)$



# Result

## Analytical 5-loop contributions to $a_\mu$

The resulting contributions to  $a_\mu$  coming from 4-loop terms in the photon propagator read :

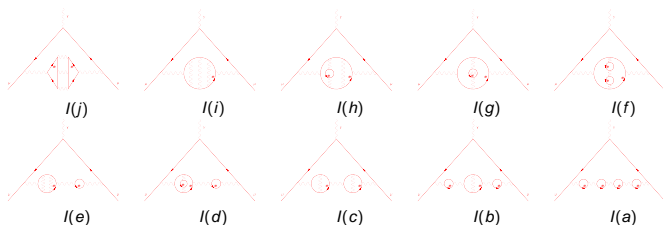
$$a_\mu^{\text{asyp.}} = \sum_{i \geq 2} a_\mu^{\text{asyp.},(i)} \left(\frac{\alpha}{\pi}\right)^i$$

$$\begin{aligned} a_\mu^{\text{asyp.},(5)} = & -\frac{296496193}{41990400} + \frac{45709}{58320} \pi^2 + \frac{212701}{518400} \pi^4 - \frac{4488523}{259200} \zeta_3 + \frac{35}{144} \pi^2 \zeta_3 + \frac{4}{3} \zeta_3^2 + \frac{10909}{720} \zeta_5 \\ & + \frac{35}{8} \zeta_7 - \frac{55}{24} \pi^2 \ln 2 - \frac{53}{675} \pi^4 \ln 2 + \frac{6121}{4320} \pi^2 \ln^2 2 - \frac{16}{135} \pi^2 \ln^3 2 - \frac{6121}{4320} \ln^4 2 \\ & + \frac{16}{225} \ln^5 2 - \frac{6121}{180} a_4 - \frac{128}{15} a_5 + \ell_{\mu e} \left[ \frac{1416095}{279936} + \frac{41}{972} \pi^2 - \frac{1855}{432} \zeta_3 - \frac{10}{3} \zeta_5 - \frac{2}{9} \pi^2 \ln 2 \right] \\ & + \ell_{\mu e}^2 \left[ -\frac{1507}{1944} + \frac{8}{81} \pi^2 + \frac{4}{3} \zeta_3 \right] - \frac{83}{243} \ell_{\mu e}^3 + \frac{8}{81} \ell_{\mu e}^4 + \mathcal{O}\left(\frac{m_e}{M_\mu}\right), \quad \ell_{\mu e} = \ln \frac{M_\mu}{m_e}, \quad a_n = \text{Li}_n\left(\frac{1}{2}\right) \end{aligned}$$

Numerically:  $\left(\frac{\alpha}{\pi}\right)^5 a_\mu^{\text{asyp.},(5)} = \left(\frac{\alpha}{\pi}\right)^5 62.2667 = 0.42105 \cdot 10^{-11}$   
(compared to  $\left(\frac{\alpha}{\pi}\right)^5 663(20)$ )

# Comparison

Analytical  $\longleftrightarrow$  Numerical results



Subset	analytical	numerical	
$I(j)$	$-1.21429 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	-1.24726(12)	✓
$I(i)$	$+0.25237 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	-	
$I(g) + I(h)$	$+1.50112 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+1.56070(64)	✓
$I(f)$	$+2.89019 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+2.88598(9)	✓
$I(e)$	$-1.33141 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	-1.20841(70)	✓
$I(d)$	$+7.44918 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+7.45270(88)	✓
$I(c)$	$+4.81759 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+4.74212(14)	✓
$I(b)$	$+27.7188 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+27.69038(30)	✓
$I(a)$	$+20.1832 + \mathcal{O}\left(\frac{m_e}{m_\mu}\right)$	+20.14293(23)	✓

-Numerics from: T. Aoyama, M. Hayakawa, T. Kinoshita, M. Nio, N. Watanabe (2008) M. Nio, T. Aoyama, M. Hayakawa, T. Kinoshita (2007) and T. Kinoshita, M. Nio (2006)

-Remaining differences should come from power suppressed corrections to the asymptotic result of  $\mathcal{O}(m_e/M_\mu)$

-Agreement with Kataev, where available

# Summary & Conclusion

- Analytical and numerical methods successfully help each other in computing the QED contribution to  $a_\mu$
- The conversions formula for  $\alpha^{\text{OS}}/\alpha^{\overline{\text{MS}}}$  is evaluated to 4-loops
  - ⇒ one could reexpress any QED order  $\alpha^5$  result in terms of the running  $\alpha^{\overline{\text{MS}}}$  or vice versa
- The asymptotic contribution to the vacuum polarization part of  $a_\mu$  in order  $\alpha^5$  (5-loop) is computed
  - ⇒ completely supports the numerical result of the Kinoshita group

