

15min+5min for questions

Emerging from the Swamp

Antonia Paraskevopoulou



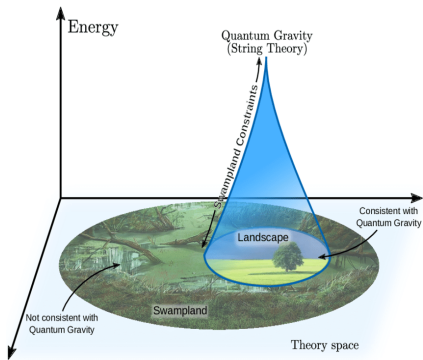
Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

IMPRS Young Scientist Workshop @ Ringberg Castle

November 23rd, 2023

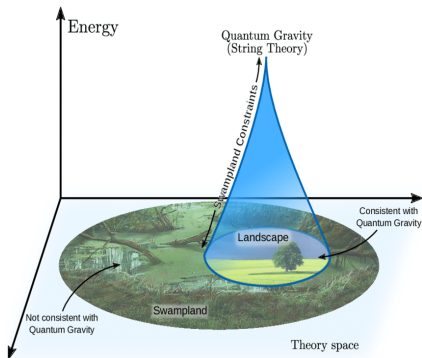
The Swampland Programme

Consistent set of conjectures motivated mainly (but not exclusively) by string theory (see e.g. Palti '19).



The Swampland Programme

Consistent set of conjectures motivated mainly (but not exclusively) by string theory (see e.g. Palti '19).



- No Global Symmetries Conjecture
- Distance Conjecture [Ooguri, Vafa '06]
- Weak Gravity Conjecture
- (A)dS Distance Conjecture
- Gravitino Conjecture
-
- The Emergence Proposal [Palti '19]
- The Emergent String Conjecture [Lee, Lerche, Weigand '18]

What do swampland conjectures look like?

- **Swampland Distance Conjecture:**[Ooguri, Vafa '06] At an infinite distance in moduli space, a tower of exponentially light states appears in our EFT, with masses given by

$$M(p) \sim M(p_0)e^{-\alpha d(p_0,p)}, \quad (1)$$

where α is some positive $\mathcal{O}(1)$ constant.

What do swampland conjectures look like?

- **Swampland Distance Conjecture:**[Ooguri, Vafa '06] At an infinite distance in moduli space, a tower of exponentially light states appears in our EFT, with masses given by

$$M(p) \sim M(p_0)e^{-\alpha d(p_0,p)}, \quad (1)$$

where α is some positive $\mathcal{O}(1)$ constant.

- **Emergent String Conjecture:**[Lee, Lerche, Weigand '18] A quantum gravitational theory in an infinite distance limit either decompactifies, or reduces to an asymptotically tensionless, weakly coupled string theory.

What do swampland conjectures look like?

- **Swampland Distance Conjecture:**[Ooguri, Vafa '06] At an infinite distance in moduli space, a tower of exponentially light states appears in our EFT, with masses given by

$$M(p) \sim M(p_0)e^{-\alpha d(p_0,p)}, \quad (1)$$

where α is some positive $\mathcal{O}(1)$ constant.

- **Emergent String Conjecture:**[Lee, Lerche, Weigand '18] A quantum gravitational theory in an infinite distance limit either decompactifies, or reduces to an asymptotically tensionless, weakly coupled string theory.

→ restriction on possible towers!

What do swampland conjectures look like?

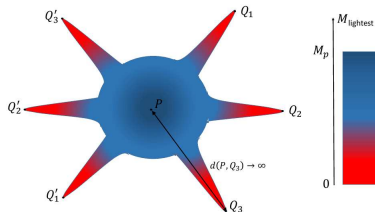
- **Swampland Distance Conjecture:**[Ooguri, Vafa '06] At an infinite distance in moduli space, a tower of exponentially light states appears in our EFT, with masses given by

$$M(p) \sim M(p_0)e^{-\alpha d(p_0,p)}, \quad (1)$$

where α is some positive $\mathcal{O}(1)$ constant.

- **Emergent String Conjecture:**[Lee, Lerche, Weigand '18] A quantum gravitational theory in an infinite distance limit either decompactifies, or reduces to an asymptotically tensionless, weakly coupled string theory.

→ restriction on possible towers!



The Emergence Proposal

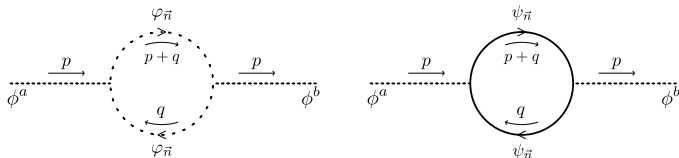
Emergence Proposal: *In a theory of Quantum Gravity all light particles in a perturbative regime have NO kinetic terms in the UV.*

Emergence Proposal: *In a theory of Quantum Gravity all light particles in a perturbative regime have NO kinetic terms in the UV. These terms appear as an IR effect due to loop corrections induced by towers of light states (**strong**). Alternatively, the 1-loop kinetic terms are analogous to tree level ones (**weak**). [e.g. Castellano, Herraez, Ibañez '22]*

The Emergence Proposal

Emergence Proposal: *In a theory of Quantum Gravity all light particles in a perturbative regime have NO kinetic terms in the UV. These terms appear as an IR effect due to loop corrections induced by towers of light states (**strong**). Alternatively, the 1-loop kinetic terms are analogous to tree level ones (**weak**). [e.g. Castellano, Herraez, Ibañez '22]*

Comparison with usual renormalization procedure in QFT:



Integrating out light states with $m_{\vec{n}}(\phi^a) = m_{\vec{n}}(\phi_0^a + \delta\phi^a)$, where ϕ^a are scalars will produce a correction to the propagator matrix

$$D_{ab}(p^2) = \frac{1}{p^2 - \Pi_{ab}(p^2)}, \quad \Pi_{ab}(p^2) = \sum_{\vec{n}} \Pi_{ab, \vec{n}}(p^2). \quad (2)$$

Usual wavefunction renormalization

$$G_{ab}^{(1)} = \sum_{\vec{n}} \frac{\partial \Pi_{ab, \vec{n}}(p^2)}{\partial p^2} \Big|_{p^2=0} \quad (3)$$

Usual wavefunction renormalization

$$G_{ab}^{(1)} = \sum_{\vec{n}} \frac{\partial \Pi_{ab, \vec{n}}(p^2)}{\partial p^2} \Big|_{p^2=0} \quad (3)$$

gives us

$$G_{ab}^{(1)} \simeq \sum_{\vec{n}} \text{deg}_{\vec{n}} \partial_a m_{\vec{n}} \partial_b m_{\vec{n}}(\dots). \quad (4)$$

Usual wavefunction renormalization

$$G_{ab}^{(1)} = \sum_{\vec{n}} \frac{\partial \Pi_{ab, \vec{n}}(p^2)}{\partial p^2} \Big|_{p^2=0} \quad (3)$$

gives us

$$G_{ab}^{(1)} \simeq \sum_{\vec{n}} \text{deg}_{\vec{n}} \partial_a m_{\vec{n}} \partial_b m_{\vec{n}}(\dots). \quad (4)$$

Similarly for the gauge kinetic functions:

$$f_{ab}^{(1)} \simeq \sum_{\vec{n}} \text{deg}_{\vec{n}} q_{a, \vec{n}} q_{b, \vec{n}}(\dots). \quad (5)$$

Usual wavefunction renormalization

$$G_{ab}^{(1)} = \sum_{\vec{n}} \frac{\partial \Pi_{ab, \vec{n}}(p^2)}{\partial p^2} \Big|_{p^2=0} \quad (3)$$

gives us

$$G_{ab}^{(1)} \simeq \sum_{\vec{n}} \text{deg}_{\vec{n}} \partial_a m_{\vec{n}} \partial_b m_{\vec{n}}(\dots). \quad (4)$$

Similarly for the gauge kinetic functions:

$$f_{ab}^{(1)} \simeq \sum_{\vec{n}} \text{deg}_{\vec{n}} q_{a, \vec{n}} q_{b, \vec{n}}(\dots). \quad (5)$$

All sums are restricted!

The Species Scale

The cut-off of our theory is the **species scale**. For a $4D$ theory, that is [Dvali et al. '07]

$$\tilde{\Lambda} \sim \frac{M_{\text{pl}}}{N_{\text{sp}}^{1/2}}, \quad (6)$$

The Species Scale

The cut-off of our theory is the **species scale**. For a 4D theory, that is [Dvali et al. '07]

$$\tilde{\Lambda} \sim \frac{M_{\text{Pl}}}{N_{\text{sp}}^{1/2}}, \quad (6)$$

where $N_{\text{sp}} = \begin{cases} \# \text{particles with } m < \tilde{\Lambda} & (\text{QFT picture}) \\ \text{S of minimum black holes} & (\text{BH picture}) \end{cases}$

- QFT picture: Graviton propagator [Donoghue '93]

$$\pi^{-1}(p^2) = p^2 \left(1 - \frac{N_{\text{sp}} p^2}{120\pi M_{\text{Pl}}^2} \log\left(-\frac{p^2}{\mu^2}\right) + \gamma \sum_{n=1}^{N_{\text{sp}}} \frac{p^2}{M_{\text{Pl}}^2} \frac{m_n}{\sqrt{-p^2}} \right), \quad (7)$$

- BH picture:

$$\tilde{\Lambda} = r_{\text{min}}^{-1}.$$

The Species Scale

The cut-off of our theory is the **species scale**. For a 4D theory, that is [Dvali et al. '07]

$$\tilde{\Lambda} \sim \frac{M_{\text{Pl}}}{N_{\text{sp}}^{1/2}}, \quad (6)$$

where $N_{\text{sp}} = \begin{cases} \# \text{particles with } m < \tilde{\Lambda} & \text{(QFT picture)} \\ \text{S of minimum black holes} & \text{(BH picture)} \end{cases}$

- QFT picture: Graviton propagator [Donoghue '93]

$$\pi^{-1}(p^2) = p^2 \left(1 - \frac{N_{\text{sp}} p^2}{120\pi M_{\text{Pl}}^2} \log\left(-\frac{p^2}{\mu^2}\right) + \gamma \sum_{n=1}^{N_{\text{sp}}} \frac{p^2}{M_{\text{Pl}}^2} \frac{m_n}{\sqrt{-p^2}} \right), \quad (7)$$

- BH picture:

$$\tilde{\Lambda} = r_{\text{min}}^{-1}.$$

Possibly more reliable method: BH picture!

- Object of interest: Prepotential

$$\mathcal{F} = \sum_{g=0}^{\infty} \mathcal{F}_g \mathcal{W}^{2g} \quad (8)$$

- Key advantage: SUSY restricted!

- Object of interest: Prepotential

$$\mathcal{F} = \sum_{g=0}^{\infty} \mathcal{F}_g \mathcal{W}^{2g} \quad (8)$$

- Key advantage: SUSY restricted!
- Tool: Schwinger integrals [Gopakumar, Vafa '95]

$$\mathcal{F}_g = -\frac{(2g-1)B_g}{(2g)!} \sum_{n=0}^{\infty} \int_0^{\infty} ds s^{2g-3} e^{-sZ_n} \quad (9)$$

- Object of interest: Prepotential

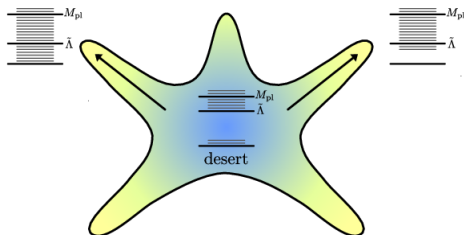
$$\mathcal{F} = \sum_{g=0}^{\infty} \mathcal{F}_g \mathcal{W}^{2g} \quad (8)$$

- Key advantage: SUSY restricted!
- Tool: Schwinger integrals [Gopakumar, Vafa '95]

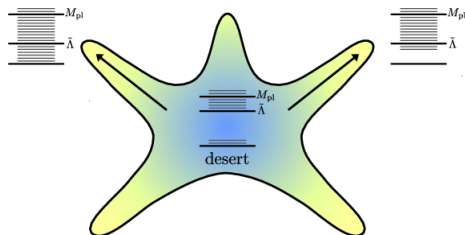
$$\mathcal{F}_g = -\frac{(2g-1)B_g}{(2g)!} \sum_{n=0}^{\infty} \int_0^{\infty} ds s^{2g-3} e^{-sZ_n} \quad (9)$$

- Method: Integrate out the **full, infinite** towers of states with mass scales below the species scale using zeta function regularization.
- At decompactification limit: Full prepotential for the resolved conifold obtained for the first time this way, including perturbative terms \rightarrow exact emergence.

Summary



- Sneak peak into the Swampland Programme
- Emergence Proposal: strong, weak, exact...
- All terms from a one-loop computation at decompactification limit!
- How far can this go? More complicated models, Gopakumar-Vafa invariants...
- Integrating out extended objects?



- Sneak peak into the Swampland Programme
- Emergence Proposal: strong, weak, exact...
- All terms from a one-loop computation at decompactification limit!
- How far can this go? More complicated models, Gopakumar-Vafa invariants...
- Integrating out extended objects?

Thank you for your attention!